Fast simulation of the CEPC detector with Delphes and its validation *

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Abstract: In this paper, the DELPHES is used to simulate the detector at the Circular Electron-Positron Collider (CEPC). The geometry and performance of the CEPC detector are presented. The fast simulation of CEPC in the DELPHES framework is validated with a series of benchmark processes. The comparisons between DELPHES and the full simulation, which is based on Geant4 & Marlin, show that DEPHES simulates the CEPC detector well. The differences of physics analysis between hadron and lepton collider are also stressed.

Key words: CEPC, DELPHES , fast simulation

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1 Introduction

CEPC[1, 2] is a next generation electron-positron collider proposed by Chinese scientists. The machine is expected to collide electron and positron beams at the center-of-mass energy of 240 - 250 GeV to maximize the Higgs production

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cross section of the $e^+e^- \to ZH$ process, with an instantaneous luminosity of 2×10^{34} cm$^{-2}s^{-1}$. CEPC is designed to deliver a total of 5 ab^{-1} integrated luminosity with two detectors in 10 years, over 10^6 Higgs events will be produced. The large statistics with clean backgrounds will enable CEPC to perform Higgs precision measurements, Standard Model (SM) tests, and searches for potential new physical phenomena. In order to investigate the precison and sensitivity of the CEPC, theorists would get involved to test various new models. To save computing time and to simplify the procedure of full detector simulation and reconstruction, a dedicated fast simulation tool is highly demanded.

Delphes [3] is a fast simulation framework developed in 2009, and the latest prime version was released in 2013, which is widely used for phenomenological studies on the simulation of various detector designs. The Delphes framework simulates the response of a general purpose collider detector, whose components are organized concentrically with a cylindrical symmetry around the beam axis. The energy and momentum are simulated according to the detector resolutions and efficiencies required by users. Eventually, all kinds of particles are reconstructed with algorithm based on PFA[4] philosophy, then clustering into jets with FastJet [16] .

This paper is organized as following. In Sect. **2**, the CEPC detector geometry are briefly introduced. Then in Sect. **3**, the MC generation, full simulation, and reconstruction chain are presented in detail, the data analysis of e^+e^- experiment is also disscused as well. In Sect. **4**, a series of benchmark processes are chosen to the validate the fast simulation on CEPC by comparing it with full simulaiton. In the end the conclusion remarks are presented.

2 The CEPC detector concept

The CEPC detector concept design[1] takes the ILC detector, ILD[5, 6], as a reference and adopts the philosophy of PFA, which benefits from its precision tracking system and high granularity calorimetric system.

The CEPC detector consists of three main sub-detectors and a superconducting solenoid of 3.5 T. The three sub-detectors are, from inner to outer, a hybrid tracking system composed of several silicon based devices and a Time Projection Chamber (TPC), a high granularity calorimetry system, and a Muon detector.

The hybrid tracking system has five parts. A vertex detector (VTX), constructed with high spatial resolution pixel sensor, is placed very close to the interaction point (IP) and the inner radius is only 16 mm. The VTX provides very precision measurements of the impact parameters of tracks, which is used for the $b-/c-$jet flavor tagging and

τ-tagging. A Silicon Inner Tracker (SIT) is just outside and cooperating with the VTX for vertex reconstruction and flavor tagging. A set of Forward Tracking Disks (FTDs) are placed in the forward region to increase the geometric acceptance of tracking system with coverage up to $|\cos\theta| = 0.99$. A Silicon External Tracker (SET) and End-cap Tracking Disks (ETD) are taken as the outermost layer of whole tracker system, which provide precision position measurements of tracks entering the calorimetric system. The TPC, with a 2.35m half-length and 1.8m outer radius, provides over 200 hits per track and 100 μm resolution in $r\phi$ plane, which allows for excellent pattern recognition, track reconstruction efficiency, and potential dE/dx-based particle identification.

A calorimetric system consisting of Electromagnetic Calorimeter (ECAL) and Hadron Calorimeter (HCAL) with very fine granularity is placed inside the solenoid. The system plays an essential role in the Particle-Flow Algorithm (PFA), providing excellent separation of showers from different particles and jet energy resolution of 3-4%.

A superconducting solenoid of 3.5 T is surrounding the calorimetry system. The return yoke is placed outside the solenoid. The CEPC muon system acts as the muon identifier, the solenoid flux return yoke and the support structure for the whole spectrometer. High muon detection efficiency, low hadron mis-identification rate, modest position resolution and large coverage are the main concerns of the design.

3 Monte Carlo samples, detector simulation, reconstruction, and analysis

For the CEPC detector design and optimization, a whole set of $e^+e^- \to ZH$ signal process and Standard Model (SM) backgrounds have been generated [7] with the generic Monte-Carlo generator Whizard 1.95 [8]. To simulate the detector response, a full simulation package, Mokka [9], based on Geant4 [10] and a fast simulation framework of the Delphes [3] are used. The Particle Flow Algorithm (PFA) philosophy is utilized in the reconstruction of both the full and fast simulation.

After full simulation, hits in different sub-detectors are digitized properly and reconstructed with reconstruction software Marlin [11]. A dedicated PFA, Arbor[12], is used for particle reconstruction, and Pandora [13], an alternative one, is taken as reference. All jets are reconstructed with LCFIPlus package[14], where a e^+e^- k_t algorithm [15], often referred to also as Durham algorithm, is used for jet-clustering.

The detector model implemented in the Delphes is same as the one in full simulation with necessary simplifications. For the charged tracks simulated with Delphes, there is a common strategy to smear their momentum.

The angular resolution is assumed to be perfect, so that the smearing only applied to the transversal momentum. In practice, the resolution of their momentum is described as a Gaussian and $\Delta(1/p_t) = a \oplus (bp_t)$, where a and b are taken same as CEPC pre-CDR. The reconstruction efficiency is 100% in the detector coverage since the it is greater than 99% for $p_t > 0.1$ GeV.

For neutral objects which mainly rely on the calorimetric system, some points should be clarified. The first, the fake rate for electrons, muons and photons is not implemented in the current version of DELPHES ; The second, photon conversions into electron-positron pairs are ignored neither. Both the true photons and electrons reaching calorimeter without a reconstructed track are considered as photon in DELPHES . The resolutions of photon and neutral hadrons are taken from pre-CDR, which are consistent those of full simulation.

The overall reconstruction of particles implemented in DELPHES is mainly based on a perfect PFA. For the charged tracks, the reconstruction involves both the tracker system and the calorimetric system. Since the resolution of tracking system is better than this of calorimeter in the CEPC energy region, it can be convenient to use the tracking information within the tracker acceptance for estimating the charged particle momenta. The efficiencies could be parameterized as a function of the energy and pseudo-rapidity. But it should be beared in mind that the detection efficiencies of fast simulaiton is higher than full simulation by about 5% per photon since the photon conversions is not simulated.

In order to get same performance of jet-clustering, e^+e^- k_t algorithm is also used in fast simulation. The exclusive mode of fast jet is used and the input particles are forced into fixed number of jets without any y_{ij} or P_t requirements, which can be applied at analysis stage.

It is worthy to note that the analysis strategy at e^+e^- collider is different from the one of hadron colliders, such as the Atlas and the CMS. The e^+e^- experiments have much less backgrounds than hadron ones and almost pile-up free, and usually almost all final states particles of a event are detected and used in analysis. For instance, all the final physics objects will be used in the analysis of $e^+e^- \to ZH$, with $Z \to \mu^+\mu^-$ and $H \to 2$Jets. As a consequence, the muon pair should be identified firstly with both Pid and kinematic constrains. Then, the remain particles are forced into 2 jets with exclusive mode of e^+e^- k_t jet-clustering algorithm. And any background contaminations of mis-combination can be suppressed in the final event selection with additional requirements.

4 Validation with full simulation

The fast simulation and its reconstruction in DELPHES must be validated by comparing the relative distributions of all types of physics objects to those of full simulation, as well as the efficiencies. Several benchmark processes are selected, which cover all the physics objects of CEPC experiments, such as leptons, photons, and jets.

4.1 $e^+e^- \to \mu^+\mu^- H$

The $e^+e^- \to \mu^+\mu^- H$ process is one of the most important analyes on at e^+e^- Higgs factory, which enables the model-independent measurement of Higgs mass and production rate. The di-muons come from Z boson decay and Higgs is tagged from the recoiling side, since the initial four-momentum is precisely defined. So the study on the recoil mass distribution produces important properties of the Higgs but without detecting any Higgs decay products. Fig. 1 shows the comparisons of invariant mass of and recoil mass against $\mu^+\mu^-$ pair between fast and full simulaitons. It can be found that these two are consistent very well.

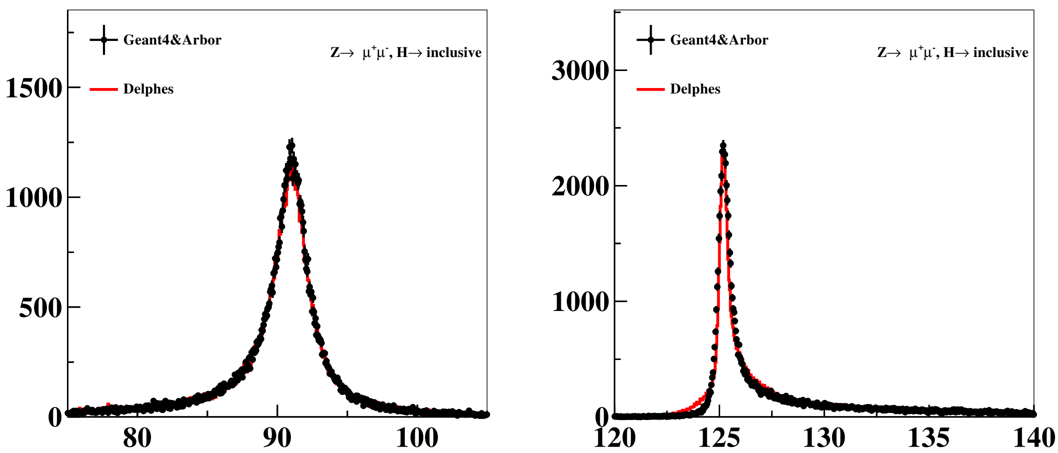

Fig. 1. The invariant mass (left) of and recoil mass (right) against $\mu^+\mu^-$ distributions in $e^+e^- \to ZH$, $Z \to \mu^+\mu^-$ and $H \to$ inclusive. The dot with error bar and red histogram represent full and fast simulaiton, respectively (same in the following plots).

4.2 $e^+e^- \to \nu\bar{\nu}H$, $H \to WW^*$

The $e^+e^- \to \nu\bar{\nu}H$, $H \to WW^*$ process includes the Higgs-bremstruhlung and WW fusion contributions. All the visible final state particles come from WW^* system except the initial state radiation photon(s) and eventually from

Higgs, which contains all types of physics objects at the CEPC, such as jets, leptons, and photons and is very useful to validate the performance of fast simulation. The plots in Fig. 2 show distributions of the invariant masses of and recoil mass against all visible decay products of the process, which represent Higgs and Z (here it includes the WW fusion contribution), respectively.

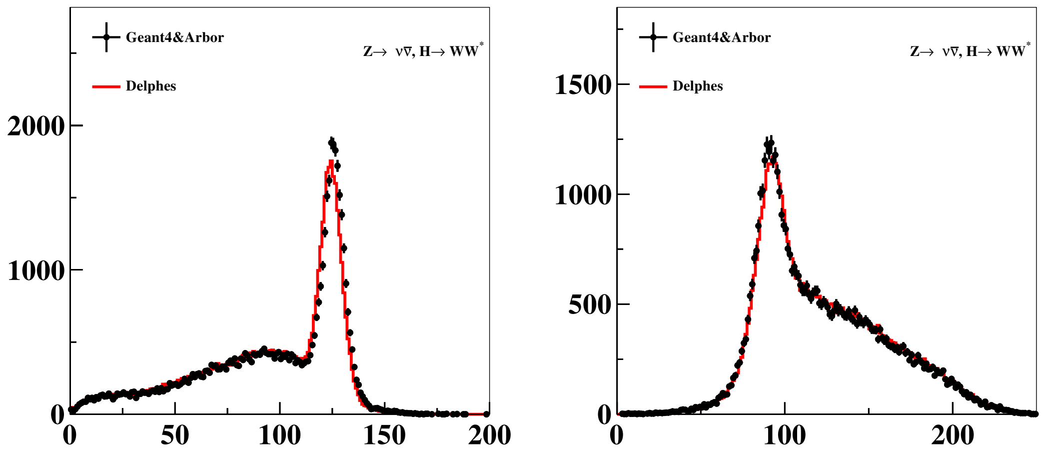

Fig. 2. The distributions of invariant mass(left) of and recoil mass (right) against WW^* system in $e^+e^- \to \nu\bar{\nu}H$, $H \to WW^*$.

4.3 $e^+e^- \to ZH \to 2(q\bar{q})$

The $e^+e^- \to ZH \to 2(q\bar{q})$ process contains four jets in the final states, which is very useful to validate the performances related jets, such as jet-clustering, jet energy resolution, and jet pairing, etc. In both full and fast simulaiton, all the final states particles are forced into four jets with ee_{kt} algorithm. The the four jets are paired by minimize $\chi^2 = (M_{j_1,j_2} - M_Z)^2 + (M_{j_3,j_4} - M_H)^2$, where M_Z and M_H are the world averages of Z and Higgs boson masses. The invariant mass distribution of jet pairs are shown in Fig.3. From Fig. 3, it can be seen that both of the distributions are well consistent. In Fig. 4, the scattering plots of M_{j_1,j_2} vs. M_{j_3,j_4} demonstate similar patterns in the available kinematic region.

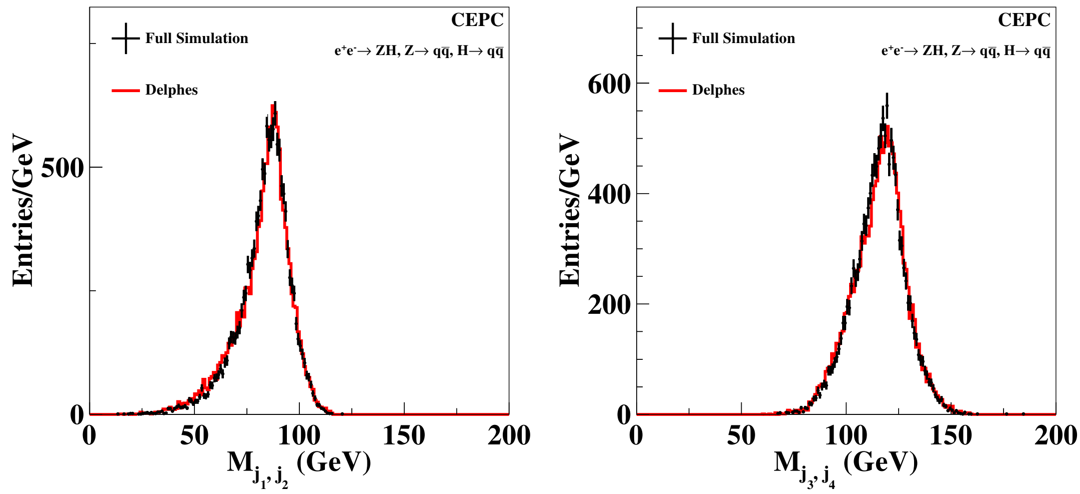

Fig. 3. The invariant mass distributions of the jet pairs, which are peaking at Z and Higgs masses.

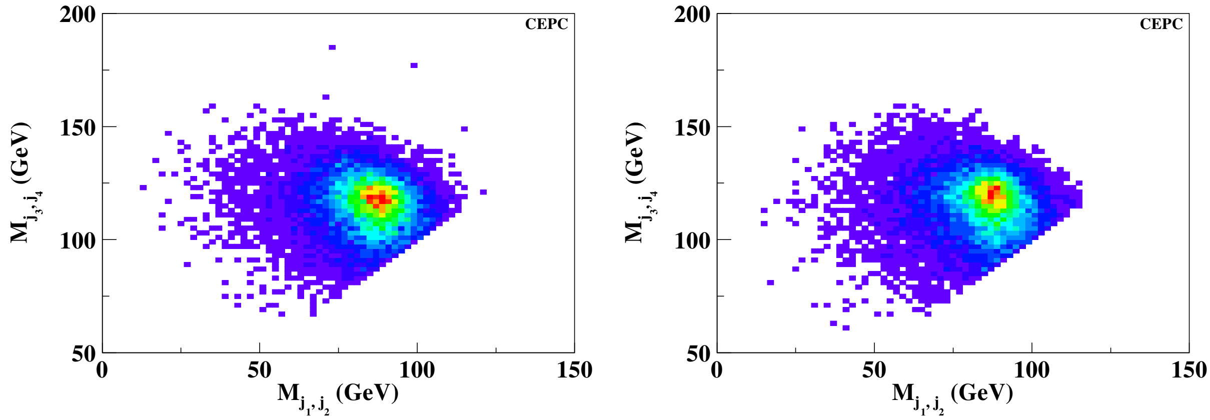

Fig. 4. The Scattering plots of M_{j_1,j_2} vs. M_{j_3,j_4} for full and fast simulaitons, respectively.

5 Conclusion

A dedicated fast simulaiton tool is ensential for the CEPC phenomenological study. The fast simulaiton frame work of the CEPC detector based Delphes is introduced. To validate the performance, comprehensive comparisons between DELPHES fast simulaiton and full simulaiton are performed. The results show well agreement between fast and full simulaitons on the CEPC.

On the anlysis in high energy e^+e^- experiments, it worthy to note that exclusive jet-clustering algorithm is favored in most cases, since it uses all final state particles and doesn't loose any available information. It also should be noted that the flavor tagging module are improved based on the machine learning results of full simulaiton, which is very useful since it provides more realistic experimental condiation.

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