

QUANTUM SIMULATION [Computational Quantum Physics]

Exact Diagonalization Approaches

- Interactions
- "N" particles that is large enough
- Interacting Spins
- Interacting Electrons

Solve Schrödinger Equation

$$i \frac{\partial}{\partial t} |\Psi(t)\rangle = \hat{H} |\Psi(t)\rangle$$

(Time Dependent SE)



$$|\Psi(0)\rangle = \text{Initial Condition}$$

Computational techniques

allow to solve time dependent
efficiently in some special cases

$$\hat{H} |\psi_k\rangle = E_k |\psi_k\rangle$$

\rightarrow Eigenkets

$$E_0 \leq E_1 \leq E_2$$



No efficient comp. techniq
is known

Exact Diagonalization Approaches (classical Methods)

$$\hat{H} = \sum_i^N J \hat{\sigma}_i^z \hat{\sigma}_{i+1}^z - g \sum_i^N \hat{\sigma}_i^x \rightarrow \text{transverse field}$$

↓ $2^N \times 2^N$ matrix

$$\hat{\sigma}_i^2 = \underbrace{\hat{I} \otimes \hat{I} \otimes \hat{I} \dots \hat{I}}_{i-1 \text{ identities}} \dots \otimes \underbrace{\hat{I} \dots \hat{I}}_{N-i \text{ identities}}$$

$$\hat{I} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \quad \hat{G}^{-2} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\boxed{U_{N+1}^z = U_1^z}$$

Basis for the Eigenkets of \hat{H}

$$|\Psi\rangle = \sum_{S_1 \dots S_N} C_{S_1 \dots S_N} |S_1 \dots S_N\rangle$$

$$|S_1 S_2 \dots S_N\rangle = |S_1\rangle \otimes |S_2\rangle \otimes \dots \otimes |S_N\rangle$$

$$\hat{S}_i^2 |S_1, \dots, S_N\rangle = S_i |S_1, \dots, S_N\rangle$$

↓ ± 1 (eigenvalues)

Write Explicit Expression for the matrix elements

$f_1 =$

In total N element

$2^N \times 2^N$

$$\sum_i \int \sigma_i^z \sigma_{i+1}^z = \langle S_1, \dots, S_N | \hat{H} | S_1, \dots, S_N \rangle = \sum_{i=1}^N \int S_i S_{i+1}$$

$$\rightarrow = \langle S_1, \dots, S_N | \hat{\sigma}_i^x | S'_1, \dots, S'_N \rangle =$$

$$= \langle S_1 | \hat{\sigma}_1^x \langle S_2 | \hat{\sigma}_2^x \dots \otimes \langle S_N | \hat{\sigma}_N^x | S'_1 \rangle | S'_2 \rangle \dots$$

$$= \langle S_1 | \hat{I}_1 | S'_1 \rangle \langle S_2 | \hat{I}_2 | S'_2 \rangle \dots \langle S_i | \hat{I}_i | S'_i \rangle \dots \langle S_N | \hat{I}_N | S'_N \rangle$$

$$= \delta_{S_1, S'_1} \delta_{S_2, S'_2} \dots \delta_{S_i, -S'_i} \dots \delta_{S_N, S'_N}$$

\downarrow
 due to σ_x^2

$$\hat{\sigma}_i^x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \langle \uparrow | \sigma_x | \downarrow \rangle = 1$$

$$\langle \uparrow | \sigma_x | \uparrow \rangle = 0$$

We have found that SPARSE MATRIX.

Most of its matrix elements are 0.

$$2^N \times 2^N \text{ MATRIX ELEMENTS} \rightarrow 2^N + N2^N \sim 2^N$$

In software (eg Scipy. sparse)

Store the matrix elements


(i, j, M_{ij})

→ M_{ij} ≠ 0

matrix elements
are not zero

Memory Usage

N	Memory ($ \Psi\rangle$)
20	8 Megabytes
30	8 Gigabytes
40	8 Terabytes

$$\text{Memory} = 2^N \times 8 \text{ BYTES}$$

Exploit Sparsity (only for Time-Dependent Sch. Eq.)

$$i \frac{d}{dt} C_{S_1, \dots, S_N}(t) = \sum_{S'_1 \dots S'_N} H(S_1, \dots, S_N; S'_1, \dots, S'_N) C_{S'_1, \dots, S'_N}(t)$$

$$i \frac{d}{dt} \vec{C}(t) = \hat{H} \vec{C}(t) \quad O(2^N)$$

$$\vec{C}(0) \longrightarrow \vec{C}(t) \quad O(2^N)$$

What if I want the Ground State?

$$e^{-\tau \hat{H}} |\Psi(0)\rangle = |\Psi(\tau)\rangle \quad \text{Imaginary Time, } \tau \in \mathbb{R}$$

$$e^{-it\hat{H}} |\Psi(0)\rangle = |\Psi(t)\rangle \quad \text{Real Time, } t \in \mathbb{R}$$

$$|\Psi(z)\rangle = \sum_k e^{-\frac{\hat{H}_k}{2}} c_k |\tilde{E}_k\rangle, \quad c_k = \langle \tilde{E}_k | \Psi_{(0)} \rangle$$

$$\hat{H} = \tilde{E}_k |\tilde{E}_k\rangle$$

$$\sum_k c_k e^{-\frac{\tilde{E}_k}{2}} |\tilde{E}_k\rangle = c_0 e^{-\frac{\tilde{E}_0}{2}} |\tilde{E}_0\rangle + \sum_{k \geq 1}^N c_k e^{-\frac{\tilde{E}_k}{2}} |\tilde{E}_k\rangle$$

$$= e^{-\frac{\tilde{E}_0}{2}} \left[c_0 |\tilde{E}_0\rangle + \sum_{k \geq 1} c_k e^{-\frac{\tilde{E}_k - \tilde{E}_0}{2}} |\tilde{E}_k\rangle \right]$$

$$|\Psi(z)\rangle \propto c_0 |\tilde{E}_0\rangle + \sum_{k \geq 1} c_k e^{-\frac{\Delta_k}{2}} |\tilde{E}_k\rangle$$

proportional

$$\Delta_k = \tilde{E}_k - \tilde{E}_0 \geq 0$$

$\tau \rightarrow \infty$ $c_0 |\tilde{E}_0\rangle$ will survive.

$$\frac{1}{\Delta_1} \ll 2, \quad |c_0| \neq 0$$

Algorithm

$$\frac{\partial}{\partial \epsilon} |\Psi(z)\rangle = - \hat{H} |\Psi(z)\rangle$$

$$\frac{\partial}{\partial \epsilon} c_{s_1 \dots s_N} = - \sum_{s'_1 \dots s'_N} H(s_1 \dots s_N; s'_1 \dots s'_N) c_{s'_1 \dots s'_N}$$

$$\frac{\partial}{\partial \epsilon} \vec{C}(z) = - \hat{H} \cdot \vec{C}(z)$$

complexity $O\left(\frac{1}{\Delta_1} 2^N\right)$

Use a quantum computer

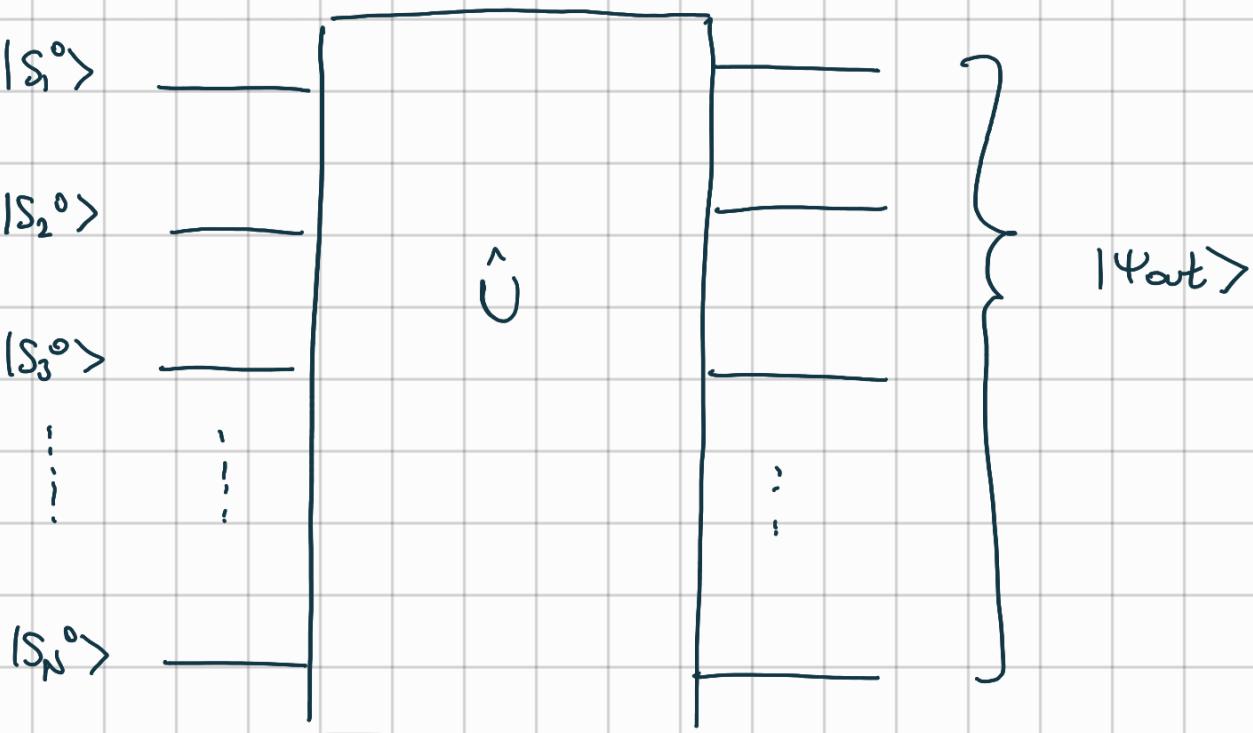
$$|\Psi(+)\rangle = e^{-i\hat{H}t} |\Psi_0\rangle, \text{ Only if } |\Psi_0\rangle \text{ is a product state.}$$

$$|\Psi_0\rangle = |\phi_1\rangle \otimes |\phi_2\rangle \otimes \dots \otimes |\phi_n\rangle$$

Single-Spin wave function
(no entanglement)

$$c_{s_1^0 \dots s_n^0} \neq 0$$

Quantum Circuits



$\hat{U} = \hat{U}_L \dots \hat{U}_2 \hat{U}_1$, \hat{U}_i are k -local operations

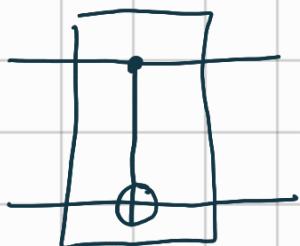
1-local Gates

$$\boxed{X} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\boxed{-I} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

2 qubit gates

CNOT



Trotter Decomposition of Unitary Dynamics

$$\hat{H} = \hat{H}_A + \hat{H}_B, \quad [\hat{H}_A, \hat{H}_B] \neq 0 \text{ (not commute)}$$

$$e^{-i \delta_t \hat{H}} = e^{-i \delta_t \hat{H}_A} e^{-i \delta_t \hat{H}_B} + O(\delta_t^2)$$

First Order
Trotter-Suzuki
Formula

$$e^{-i \delta_t \hat{H}} = e^{-i \frac{\delta_t}{2} \hat{H}_A} e^{-i \delta_t \hat{H}_B} e^{-i \frac{\delta_t}{2} \hat{H}_A} + O(\delta_t^3)$$

Second Order
Formula

$$\hat{H} = \sum_k \hat{k}_k, \quad \hat{k}_k \text{ are } k\text{-local}$$

"act only on few qubits"

$$e^{-i\delta_t \hat{H}} = e^{\underbrace{-i\delta_t \hat{k}_1}_{U_1}} \cdot e^{\underbrace{-i\delta_t \hat{k}_2}_{U_2}} \cdots e^{\underbrace{-i\delta_t \hat{k}_n}_{U_n}} + O(\delta_t^2)$$

$e^{-it\hat{H}}$ time-independent

$$e^{-it\hat{H}} |\psi(0)\rangle = e^{-i\delta_t \hat{H}} \cdot e^{-i\delta_t \hat{H}} \cdots e^{-i\delta_t \hat{H}} |\psi(0)\rangle$$

divide to smaller time spaces

where $t = p \times \delta_t$

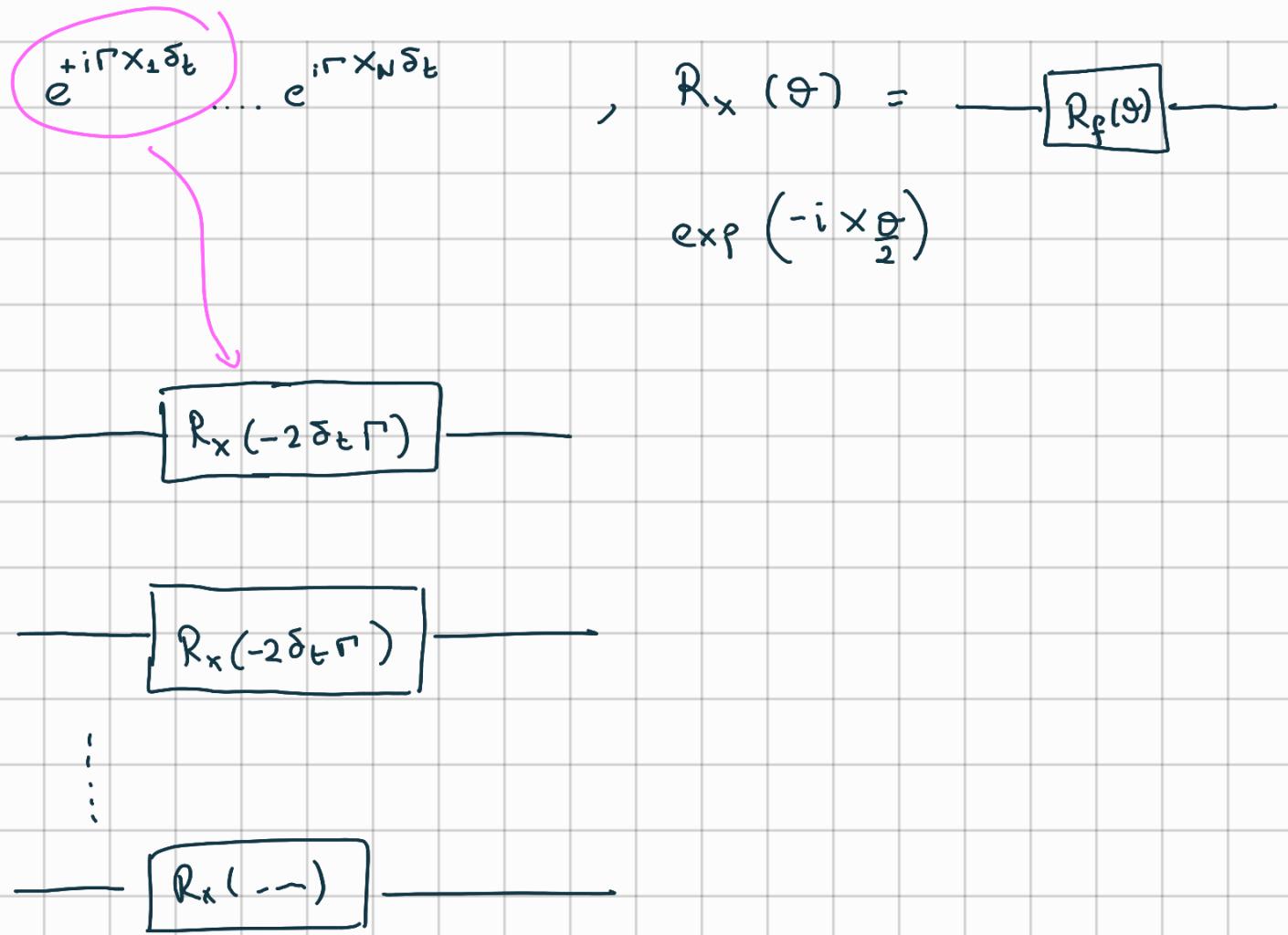
Total Error $O(\delta_t^2 \times p) = O(\delta_t \times t)$

Example: The T.F. Ising Model

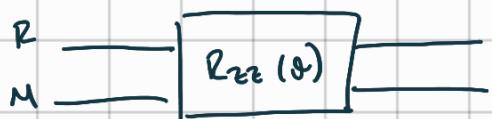
$$\hat{H} = -J \sum_i \hat{x}_i + J \sum_i \hat{z}_i \hat{z}_{i+1}, \quad \hat{z}_{N+1} = \hat{z}_1$$

$\hat{k}_k \rightarrow \textcircled{I} \quad \hat{x}_1, \hat{x}_2, \dots, \hat{x}_N \rightarrow 1 \text{ qubit gates}$

$\textcircled{II} \quad \hat{z}_1, \hat{z}_2, \dots \rightarrow 2 \text{ qubit gates}$

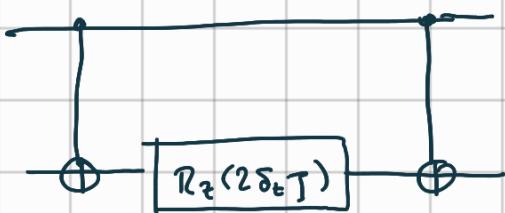


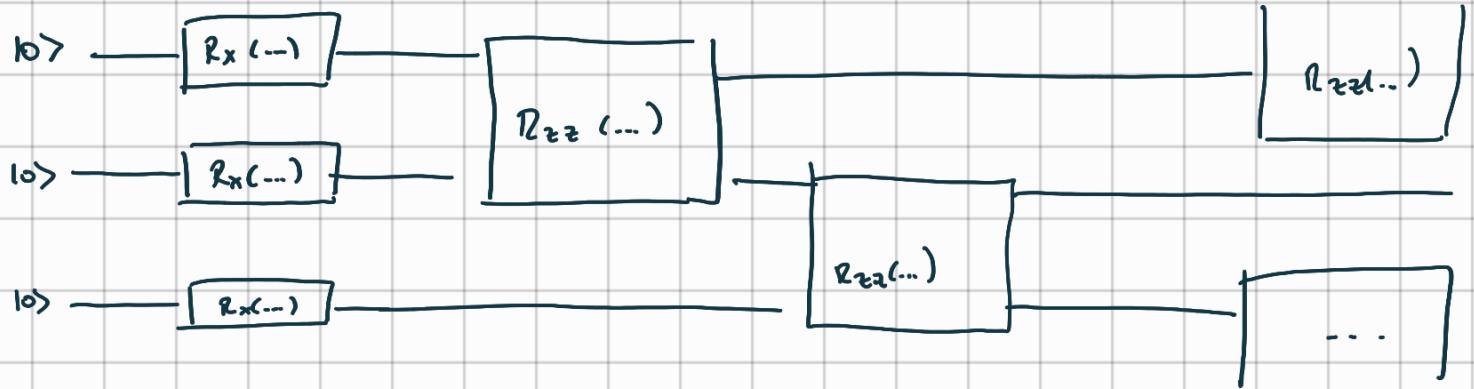
$$e^{-i\delta_t \hat{z}_1 \hat{z}_2} \quad e^{-i\delta_t \hat{z}_2 \hat{z}_3} \quad \dots \quad e^{-i\delta_t \hat{z}_N \hat{z}_1}$$



$$\hat{R}_{zz}(\theta) = e^{-i\frac{\theta}{2} \hat{z}_R \hat{z}_M}$$

$$R_z(\theta) = e^{-i\frac{\theta}{2} \hat{z}_R^2}$$





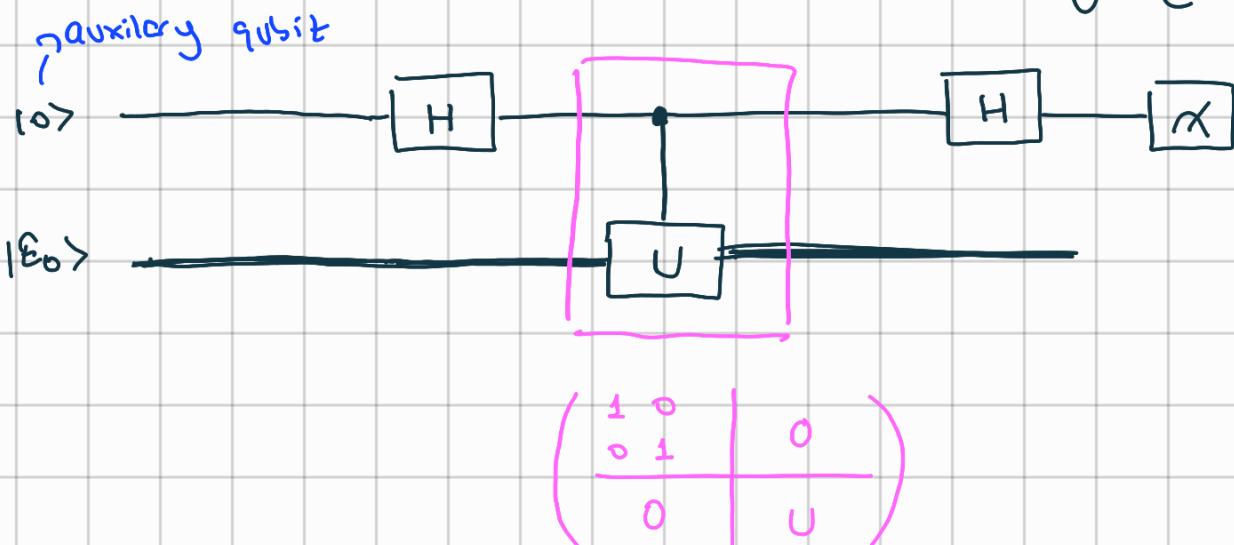
$$e^{-i\delta_t J z_1 z_3} e^{-i\delta_t z_2 z_3} e^{-i\delta_t J z_2 z_3} e^{i\delta_t \Gamma x_1} e^{i\delta_t \Gamma x_2} e^{i\delta_t \Gamma x_3} |000\rangle$$

QUANTUM PHASE ESTIMATION



$|E_0\rangle$, Problem: Find $\langle E_0 |$, $|\Psi_{in}\rangle = |0\rangle|E_0\rangle$

$$\hat{U} = e^{-iH\hat{A}}$$



$$|\Psi_{in}\rangle = |0\rangle|E_0\rangle$$

$$|\Psi_1\rangle = \frac{1}{\sqrt{2}} (|0\rangle|E_0\rangle + |1\rangle|E_0\rangle)$$

$$|\Psi_2\rangle = \frac{1}{\sqrt{2}} \left(|0\rangle|E_0\rangle + |1\rangle|E_0\rangle \right)$$

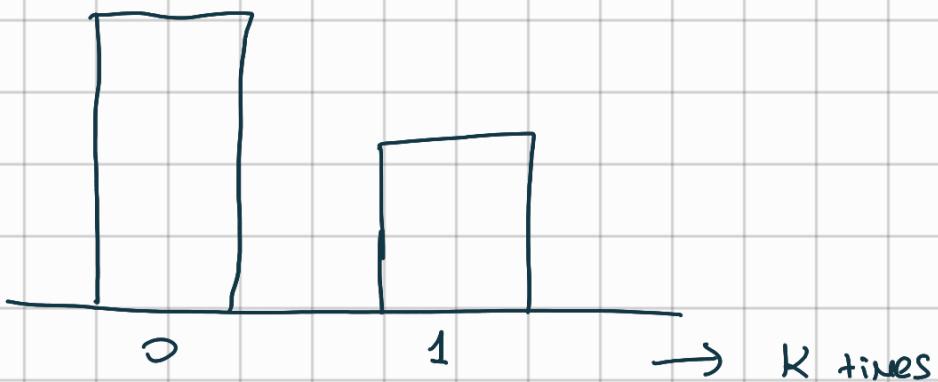
$$|\Psi_2\rangle = \frac{1}{\sqrt{2}} \left(|0\rangle|E_0\rangle + e^{-itE_0} |1\rangle|E_0\rangle \right)$$

Applying
the time evolution operator

$$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} |1\rangle \\ \frac{1}{\sqrt{2}}e^{-itE_0} \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 + e^{-itE_0} \\ 1 - e^{-itE_0} \end{pmatrix}$$

$$|\Psi_3\rangle = \frac{1}{2} \left[(1 + e^{-itE_0}) |0\rangle|E_0\rangle + (1 - e^{-itE_0}) |1\rangle|E_0\rangle \right]$$

$$\text{Prob}(0) = \frac{1}{4} \left| 1 + e^{-itE_0} \right|^2 = \cos^2 \left(\frac{tE_0}{2} \right)$$



$$\left(\frac{N_0}{K} \right)$$

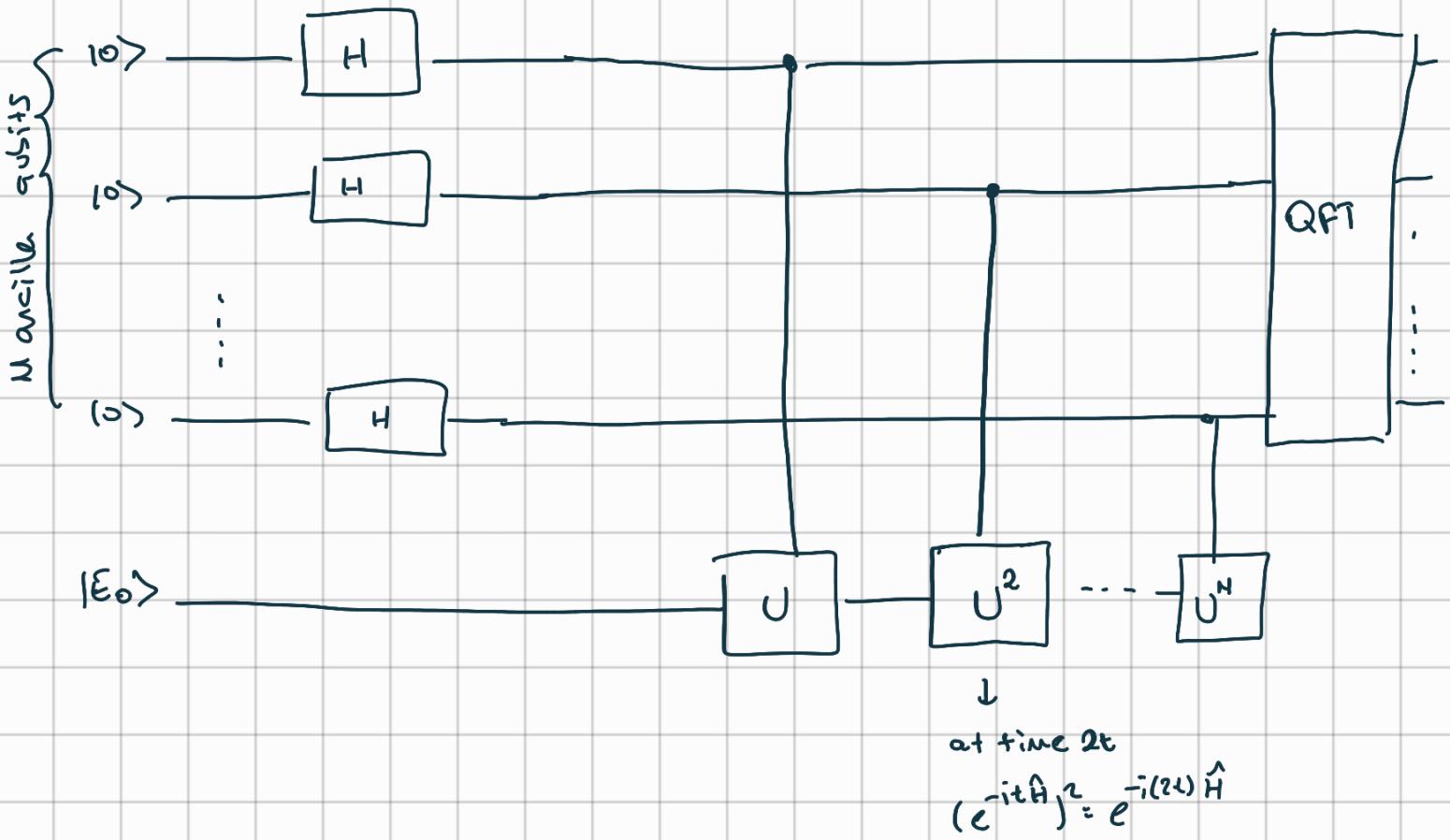
$$\left(\frac{N_1}{K} \right)$$



Proportional to the Prob(0)

FULL QUANTUM PHASE ESTIMATION

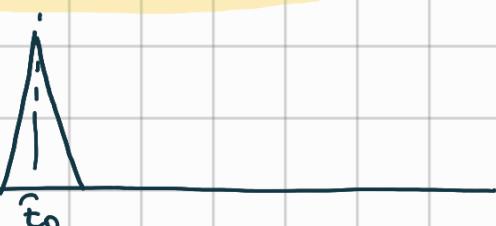
↑ several peaks at eigenstates



$$|\Psi_{out}\rangle = \frac{|E_0\rangle}{\sqrt{2^m}} \left[|0\rangle + e^{-itE_0} |1\rangle \right] \left[|0\rangle + e^{-i(2t)E_0} |1\rangle \right] \dots \left[|0\rangle + e^{-i((N-1)t)E_0} |1\rangle \right]$$

$|\Psi_{out}\rangle = \frac{1}{\sqrt{2^m}} \sum_{k=0}^{2^m-1} e^{-i(kt)E_0} |k\rangle |E_0\rangle$

Fourier transform



$$\text{PROB}(E_0) \sim |C_0|^2$$

$$C_0 \sim \frac{1}{\text{Poly}(N)}$$

$$\langle E_0 | \Psi_{(0)} \rangle = C_0$$

\downarrow polynomial
N qubits