

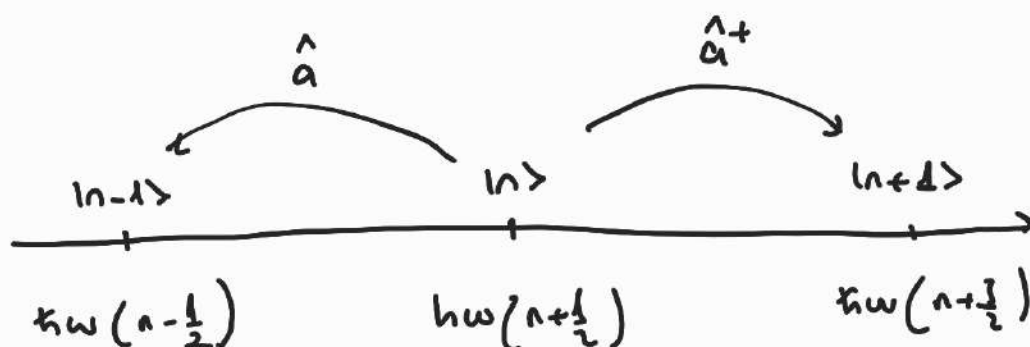
COHERENT STATE WAVE FUNCTION

Quantum Harmonic Oscillator:

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2 \hat{x}^2 = \hbar\omega \left(\hat{a}^\dagger \hat{a} + \frac{1}{2} \right) = \hbar\omega \left(\hat{N} + \frac{1}{2} \right)$$

$$\hat{H}|n\rangle = E_n |n\rangle$$

$$E_n = \hbar\omega \left(n + \frac{1}{2} \right) \rightarrow n = 0, 1, 2, \dots$$



Coherent States

$$\hat{a}|\alpha\rangle = \alpha|\alpha\rangle \quad \alpha \in \mathbb{C}$$

→ In the energy basis: $|\alpha\rangle = e^{-|\alpha|^2/2} \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} |n\rangle$

→ With displacement operator: $\hat{D}(\alpha) = e^{\alpha \hat{a}^\dagger - \alpha^* \hat{a}}$

$$|\alpha\rangle = \hat{D}(\alpha)|0\rangle$$

$$\hat{a}|\alpha\rangle = \alpha|\alpha\rangle$$

$$|\alpha\rangle = \hat{D}(\alpha)|0\rangle$$

Position Representation: $\hat{x}|x\rangle = x|x\rangle$

$$|\psi\rangle = \int dx \psi(x)|x\rangle$$

$$\psi(x) = \langle x|\psi\rangle$$

"wave function"

Coherent state wave function:

$$\psi_\alpha(x) = \langle x|\alpha\rangle = \langle x|\hat{D}(\alpha)|0\rangle$$

$$\hat{D}(\alpha) = e^{\alpha\hat{a}^\dagger - \alpha^*\hat{a}}$$

$$\hat{a} = \sqrt{\frac{m\omega}{2\hbar}} \hat{x} + i \frac{1}{\sqrt{2m\hbar\omega}} \hat{p}$$

$$\hat{a}^\dagger = \sqrt{\frac{m\omega}{2\hbar}} \hat{x} - i \frac{1}{\sqrt{2m\hbar\omega}} \hat{p}$$

$$\alpha\hat{a}^\dagger - \alpha^*\hat{a} = \alpha \left(\sqrt{\frac{m\omega}{2\hbar}} \hat{x} - i \frac{1}{\sqrt{2m\hbar\omega}} \hat{p} \right)$$

$$- \alpha^* \left(\sqrt{\frac{m\omega}{2\hbar}} \hat{x} + i \frac{1}{\sqrt{2m\hbar\omega}} \hat{p} \right)$$

$$\alpha \hat{a}^\dagger - \alpha^* \hat{a} = \sqrt{\frac{m\omega}{2\hbar}} (\alpha - \alpha^*) \hat{x} - \frac{i}{\sqrt{2m\hbar\omega}} (\alpha + \alpha^*) \hat{p}$$

$$\hat{D}(\alpha) = \exp \left[\sqrt{\frac{m\omega}{2\hbar}} (\alpha - \alpha^*) \hat{x} - \frac{i}{\sqrt{2m\hbar\omega}} (\alpha + \alpha^*) \hat{p} \right]$$

$$[\hat{A}, [\hat{A}, \hat{B}]] = 0, [\hat{B}, [\hat{A}, \hat{B}]] = 0: e^{\hat{A} + \hat{B}} = e^{\hat{A}} e^{\hat{B}} e^{-\frac{1}{2}[\hat{A}, \hat{B}]}$$

✓
satisfied

$$[\hat{x}, \hat{p}] = i\hbar \rightarrow [\hat{x}, \underbrace{[\hat{x}, \hat{p}]}_{i\hbar}] = 0$$

$$[\hat{p}, \underbrace{[\hat{x}, \hat{p}]}_{i\hbar}] = 0$$

$$\hat{D}(\alpha) = \exp \left(\sqrt{\frac{m\omega}{2\hbar}} (\alpha - \alpha^*) \hat{x} \right) \exp \left(-\frac{i}{\sqrt{2m\hbar\omega}} (\alpha + \alpha^*) \hat{p} \right)$$

$$= \exp \left(-\frac{1}{2} \left[\sqrt{\frac{m\omega}{2\hbar}} (\alpha - \alpha^*) \hat{x}, -\frac{i}{\sqrt{2m\hbar\omega}} (\alpha + \alpha^*) \hat{p} \right] \right)$$

$$\sqrt{\frac{m\omega}{2\hbar}} \left(-\frac{i}{\sqrt{2m\hbar\omega}} \right) (\alpha - \alpha^*)(\alpha + \alpha^*) \underbrace{[\hat{x}, \hat{p}]}_{i\hbar}$$

$\alpha^2 - (\alpha^*)^2$

$$-\frac{i}{2\hbar} (\alpha^2 - (\alpha^*)^2) i\hbar = \frac{1}{2} (\alpha^2 - (\alpha^*)^2)$$

$$\hat{D}(\alpha) = \exp\left(\sqrt{\frac{m\omega}{2\hbar}} (\alpha - \alpha^\dagger) \hat{x}\right) \exp\left(-\frac{i}{\sqrt{2m\hbar\omega}} (\alpha + \alpha^\dagger) \hat{p}\right) \\ \times \exp\left(-\frac{1}{4} (\alpha^2 - (\alpha^\dagger)^2)\right)$$

$$\Psi_\alpha(x) = \langle x | \hat{D}(\alpha) | 0 \rangle$$

\uparrow
 scalar

$$\langle x | \exp\left(\sqrt{\frac{m\omega}{2\hbar}} (\alpha - \alpha^\dagger) x\right)$$

$$\Psi_\alpha(x) = \exp\left(-\frac{1}{4} (\alpha^2 - (\alpha^\dagger)^2)\right) \langle x | \exp\left(\sqrt{\frac{m\omega}{2\hbar}} (\alpha - \alpha^\dagger) x\right)$$

$$\exp\left(-\frac{i}{\sqrt{2m\hbar\omega}} (\alpha + \alpha^\dagger) \hat{p}\right) | 0 \rangle$$

$$\Psi_\alpha(x) = \exp\left(-\frac{1}{4} (\alpha^2 - (\alpha^\dagger)^2)\right) \exp\left(\sqrt{\frac{m\omega}{2\hbar}} (\alpha - \alpha^\dagger) x\right)$$

$$\langle x | \exp\left(-\frac{i}{\sqrt{2m\hbar\omega}} (\alpha + \alpha^\dagger) \hat{p}\right) | 0 \rangle$$

$$\hat{T}\left(\sqrt{\frac{\hbar}{2m\omega}} (\alpha + \alpha^\dagger)\right)$$

Translation Operator: $\hat{T}(\lambda) = e^{-i\lambda \hat{p} / \hbar}$

$$\hat{T}(\lambda) | x \rangle = | x + \lambda \rangle \longleftrightarrow \langle x | \hat{T}(\lambda) = \langle x - \lambda |$$

$$\langle x | \hat{T} \left(\sqrt{\frac{\hbar}{2m\omega}} (\alpha + \alpha^\dagger) \right) = \langle x - \sqrt{\frac{\hbar}{2m\omega}} (\alpha + \alpha^\dagger) |$$

$$\Psi_\alpha(x) = \exp \left[-\frac{1}{4} (\alpha^2 - (\alpha^\dagger)^2) \right] \exp \left[\sqrt{\frac{m\omega}{2\hbar}} (\alpha - \alpha^\dagger) x \right]$$

$$\langle x - \sqrt{\frac{\hbar}{2m\omega}} (\alpha + \alpha^\dagger) | 0 \rangle$$

ground state wavefunction
the harmonic oscillator

$$\Psi_0 \left(x - \sqrt{\frac{\hbar}{2m\omega}} (\alpha + \alpha^\dagger) \right)$$

$$\Psi_\alpha(x) = \exp \left[-\frac{1}{4} (\alpha^2 - (\alpha^\dagger)^2) \right] \exp \left[\sqrt{\frac{m\omega}{2\hbar}} (\alpha - \alpha^\dagger) x \right]$$

$$\Psi_0 \left(x - \sqrt{\frac{\hbar}{2m\omega}} (\alpha + \alpha^\dagger) \right)$$

$$\langle \hat{x} \rangle_\alpha = \sqrt{\frac{\hbar}{2m\omega}} (\alpha + \alpha^\dagger)$$

$$\langle \hat{p} \rangle_\alpha = -i \sqrt{\frac{m\hbar\omega}{2}} (\alpha - \alpha^\dagger)$$

$$\Psi_\alpha(x) = \exp \left[-\frac{1}{4} (\alpha^2 - (\alpha^\dagger)^2) \right] \exp \left[\frac{i}{\hbar} \langle p \rangle_\alpha x \right] \Psi_0(x - \langle x \rangle_\alpha)$$

$$-\frac{1}{4} (\alpha^2 - (\alpha^*)^2) = -i \operatorname{Re}(\alpha) \operatorname{Im}(\alpha) = i\theta(\alpha)$$

$$(\alpha)^2 = (a+ib)^2 = \cancel{a^2} + 2iab - \cancel{b^2}$$

$$(\alpha^*)^2 = (a-ib)^2 = \cancel{a^2} - 2iab - \cancel{b^2}$$

$$\left(-\frac{1}{4}\right) (\alpha^2 - (\alpha^*)^2) = (4iab) \left(-\frac{1}{4}\right) = -iab$$

$$\theta(\alpha) = -\operatorname{Re}(\alpha) \operatorname{Im}(\alpha) \rightarrow \theta(\alpha) \in \mathbb{R}$$

$$\Psi_\alpha(x) = e^{i\theta_\alpha} e^{i \langle \hat{p} \rangle_\alpha x / \hbar} \Psi_0(x - \langle \hat{x} \rangle_\alpha)$$

$$\Psi_0(x) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-m\omega x^2 / 2\hbar}$$

$$|\Psi_\alpha(x)|^2 = |\Psi_0(x - \langle \hat{x} \rangle_\alpha)|^2$$

$$\Psi_0(x) = \left(\frac{m\omega}{\pi\hbar} \right)^{1/4} e^{-m\omega x^2 / 2\hbar}$$

$$|\Psi_\alpha(x)|^2 = \sqrt{\frac{m\omega}{\pi\hbar}} \exp\left(-\frac{m\omega}{\hbar} (x - \langle \hat{x} \rangle_\alpha)^2\right)$$

$$\Delta \hat{x}_\alpha = \sqrt{\frac{\hbar}{2m\omega}} \Rightarrow \frac{m\omega}{\hbar} = \frac{1}{2(\Delta \hat{x}_\alpha)^2}$$

QUASI-CLASSICAL STATES

DOES THE QUANTUM HARMONIC OSCILLATOR ACTUALLY
OSCILLATE?

Classical Harmonic Oscillator:

$$x(t) = x_0 \cos(\omega t - \varphi)$$

$$p(t) = m \dot{x}(t) = -m\omega x_0 \sin(\omega t - \varphi)$$

$$E(t) = \frac{1}{2m} [p(t)]^2 + \frac{1}{2} m\omega^2 [x(t)]^2$$

$$E(t) = \frac{1}{2\mu} \mu^2 \omega^2 x_0^2 \sin^2(\omega t - \varphi) + \frac{1}{2} \mu \omega^2 x_0^2 \cos^2(\omega t - \varphi)$$

$$= \frac{1}{2} \mu \omega^2 x_0^2 \left[\underbrace{\sin^2(\omega t - \varphi) + \cos^2(\omega t - \varphi)}_1 \right] = \frac{1}{2} \mu \omega^2 x_0^2 \quad \leftrightarrow$$

Quantum harmonic oscillator:

$$\hat{H} = \frac{\hat{p}^2}{2\mu} + \frac{1}{2} \mu \omega^2 \hat{x}^2 = \hbar \omega \left(\hat{a}^\dagger \hat{a} + \frac{1}{2} \right) = \hbar \omega \left(\hat{N} + \frac{1}{2} \right)$$

$$\hat{H}|n\rangle = E_n|n\rangle \quad E_n = \hbar \omega \left(n + \frac{1}{2} \right), \quad n=0, 1, 2, \dots$$

How to connect classical and quantum oscillator?

position operator

$$\langle \hat{x} \rangle_n = \langle n | \sqrt{\frac{\hbar}{2\mu\omega}} (\hat{a} + \hat{a}^\dagger) | n \rangle = 0$$

$$\langle n | \underbrace{\hat{a}}_{\sqrt{n}|n-1\rangle} | n \rangle = \sqrt{n} \langle n | n-1 \rangle = 0$$

$$\langle n | \underbrace{\hat{a}^\dagger}_{\sqrt{n+1}|n+1\rangle} | n \rangle = \sqrt{n+1} \langle n | n+1 \rangle = 0$$

$$\langle \hat{p} \rangle_n = 0$$

$$\langle \hat{H} \rangle_n = \hbar \omega \left(n + \frac{1}{2} \right)$$

Classical

$|n\rangle$

$$x(t) = x_0 \cos(\omega t - \varphi)$$

$$p(t) = -m\omega x_0 \sin(\omega t - \varphi)$$

$$E(t) = \frac{1}{2} m \omega^2 x_0^2$$

no resemble
of
oscillation

$$\langle \hat{x} \rangle_n(t) = 0$$

$$\langle \hat{p} \rangle_n(t) = 0$$

$$\langle \hat{H} \rangle_n(t) = \hbar\omega\left(n + \frac{1}{2}\right)$$

THEREFORE WE NEED COHERENT STATE

Coherent States: (Quasi-classical States)

$$\hat{a}|\alpha\rangle = \alpha|\alpha\rangle, \quad \alpha \in \mathbb{C}$$

$$|\alpha\rangle = e^{-|\alpha|^2/2} \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} |n\rangle$$

$$\hat{a}|\alpha\rangle = \alpha|\alpha\rangle \quad \xleftrightarrow{\text{Dual space}} \quad \langle\alpha|\hat{a}^\dagger = \langle\alpha|\alpha^*$$

$$\langle \hat{a} \rangle_\alpha = \langle \alpha | \underbrace{\hat{a}}_{\alpha|\alpha} | \alpha \rangle = \alpha \langle \alpha | \alpha \rangle = \alpha$$

$$\langle \hat{a}^n \rangle_\alpha = \langle \alpha | \hat{a}^n | \alpha \rangle = \alpha^n$$

$$\langle \hat{a}^\dagger \rangle_\alpha = \underbrace{\langle \alpha | \hat{a}^\dagger | \alpha \rangle}_{\langle \alpha | \alpha^* \rangle} = \alpha^*$$

$$\langle \hat{a}^{\dagger n} \rangle_\alpha = (\alpha^*)^n$$

$$\underbrace{\langle \alpha | \hat{a}^\dagger \hat{a} | \alpha \rangle}_{\langle \alpha | \alpha^* \alpha | \alpha \rangle} = |\alpha|^2$$

$$\langle \alpha | \hat{a} \hat{a}^\dagger | \alpha \rangle = \underbrace{\langle \alpha | \hat{a}^\dagger \hat{a} | \alpha \rangle}_{|\alpha|^2} + 1 = |\alpha|^2 + 1$$

$$[\hat{a}, \hat{a}^\dagger] = 1 \Rightarrow \hat{a} \hat{a}^\dagger = \hat{a}^\dagger \hat{a} + 1$$

$$\langle \hat{x} \rangle_\alpha = \langle \alpha | \sqrt{\frac{\hbar}{2m\omega}} (\hat{a} + \hat{a}^\dagger) | \alpha \rangle$$

$$= \sqrt{\frac{\hbar}{2m\omega}} \left(\underbrace{\langle \alpha | \hat{a} | \alpha \rangle}_\alpha + \underbrace{\langle \alpha | \hat{a}^\dagger | \alpha \rangle}_{\alpha^*} \right)$$

$$\langle \hat{x} \rangle_\alpha = \sqrt{\frac{\hbar}{2m\omega}} (\alpha + \alpha^*) = \sqrt{\frac{2\hbar}{m\omega}} \operatorname{Re}(\alpha)$$

$$\langle \hat{x}^2 \rangle_\alpha = \langle \alpha | \frac{\hbar}{2m\omega} (\hat{a} + \hat{a}^\dagger)^2 | \alpha \rangle$$

$$= \frac{\hbar}{2m\omega} \left(\underbrace{\langle \alpha | \hat{a}^2 | \alpha \rangle}_{\alpha^2} + \underbrace{\langle \alpha | \hat{a}^{\dagger 2} | \alpha \rangle}_{(\alpha^*)^2} + \underbrace{\langle \alpha | \hat{a} \hat{a}^\dagger | \alpha \rangle}_{1 + \alpha \alpha^*} + \underbrace{\langle \alpha | \hat{a}^\dagger \hat{a} | \alpha \rangle}_{\alpha \alpha^*} \right)$$

$$\langle \hat{x}^2 \rangle_\alpha = \frac{\hbar}{2m\omega} \left(\alpha^2 + (\alpha^*)^2 + 2\alpha\alpha^* + 1 \right) = \frac{\hbar}{2m\omega} \left((\alpha + \alpha^*)^2 + 1 \right)$$

$$\Delta \hat{x}_\alpha = \sqrt{\langle \hat{x}^2 \rangle_\alpha - \langle \hat{x} \rangle_\alpha^2}$$

$$\Delta \hat{x}_\alpha = \sqrt{\frac{\hbar}{2m\omega} \left((\alpha + \alpha^*)^2 + 1 \right) - \frac{2\hbar}{m\omega} (\alpha + \alpha^*)^2}$$

$$\Delta \hat{x}_\alpha = \sqrt{\frac{\hbar}{2m\omega}}$$

does not depend on α

Assume:

$$|\psi(0)\rangle = |\alpha_0\rangle \quad \alpha_0 = |\alpha_0| e^{i\varphi}$$

$$|\psi(t)\rangle = e^{-i\omega t/2} |\alpha\rangle, \quad \alpha = \alpha_0 e^{-i\omega t} = |\alpha_0| e^{-i(\omega t - \varphi)}$$

$$\langle x \rangle_\alpha(t) = \sqrt{\frac{\hbar}{2m\omega}} \left(|\alpha_0| e^{-i(\omega t - \varphi)} + |\alpha_0| e^{i(\omega t - \varphi)} \right)$$

$\alpha(t) + \alpha^*(t)$

$$e^{i\delta} + e^{-i\delta} = 2\cos(\delta)$$

$$\langle x \rangle_\alpha(t) = \sqrt{\frac{2\hbar}{m\omega}} |\alpha_0| \cos(\omega t - \varphi)$$

$$\langle \hat{p} \rangle_{\alpha} = -i \sqrt{\frac{m \hbar \omega}{2}} (\alpha - \alpha^*) = \sqrt{2m \hbar \omega} \operatorname{Im}(\alpha)$$

$$\langle \hat{p}^2 \rangle_{\alpha} = \frac{m \hbar \omega}{2} [1 - (\alpha - \alpha^*)^2]$$

$$\Delta \hat{p}_{\alpha} = \sqrt{\frac{m \hbar \omega}{2}}$$

$$|\psi(0)\rangle = |\alpha_0\rangle \longrightarrow |\psi(t)\rangle = e^{-i\omega t/2} |\alpha\rangle$$

$\alpha = |\alpha_0| e^{-i(\omega t - \varphi)}$

$$\langle \hat{p} \rangle_{\alpha}(t) = -\sqrt{2m \hbar \omega} |\alpha_0| \sin(\omega t - \varphi)$$

$$\langle \hat{H} \rangle_{\alpha} = \hbar \omega \left(|\alpha|^2 + \frac{1}{2} \right)$$

$$\langle \hat{H}^2 \rangle_{\alpha} = \hbar^2 \omega^2 \left(|\alpha|^4 + 2|\alpha|^2 + \frac{1}{4} \right)$$

$$\Delta \hat{H}_{\alpha} = \hbar \omega |\alpha|$$

$$|\psi(0)\rangle = |\alpha_0\rangle \longrightarrow |\psi(t)\rangle = e^{-i\omega t/2} |\alpha\rangle \quad \text{for } \alpha = \alpha_0 e^{-i\omega t}$$

$$\langle \hat{H} \rangle_{\alpha}(t) = \hbar \omega \left(|\alpha_0|^2 + \frac{1}{2} \right) \rightsquigarrow \text{TIME INDEPENDENT}$$

Classical

$$x(t) = x_0 \cos(\omega t - \phi)$$

$$p(t) = -m\omega x_0 \sin(\omega t - \phi)$$

$$E(t) = \frac{1}{2} m\omega^2 x_0^2$$

$$x_0 = \sqrt{\frac{2\hbar}{m\omega}} |\alpha_0| \Rightarrow x(t) = \langle \hat{x} \rangle_\alpha(t)$$

$$p(t) = \langle \hat{p} \rangle_\alpha(t)$$

$$\frac{1}{2} m\omega^2 \left(\frac{2\hbar}{m\omega} \right) |\alpha_0|^2 = \hbar\omega |\alpha_0|^2$$

$$E(t) = \langle \hat{H}_\alpha(t) \rangle - \frac{1}{2} \hbar\omega$$

→ zero-point energy
of quantum
harmonic
oscillator
"PURELY QUANTUM"
PHENOMENON

for classical $\rightarrow |\alpha_0| \gg \frac{1}{2} \hbar\omega$

RETURN TO $|\psi_\alpha(x)|^2$



QUASI-CLASSICAL STATES

$|\alpha\rangle$

$$\langle \hat{x} \rangle_\alpha(t) = \sqrt{\frac{2\hbar}{m\omega}} |\alpha_0| \cos(\omega t - \phi)$$

$$\langle \hat{p} \rangle_\alpha(t) = -\sqrt{2m\hbar\omega} |\alpha_0| \sin(\omega t - \phi)$$

$$\langle \hat{H} \rangle_\alpha(t) = \hbar\omega \left(|\alpha_0|^2 + \frac{1}{2} \right)$$

$$|\Psi_\alpha(x)|^2 = |\Psi_0(x - \langle \hat{x} \rangle_\alpha)|^2$$

$$\Psi_0(x) = \left(\frac{m\omega}{\pi\hbar} \right)^{1/4} e^{-m\omega x^2 / 2\hbar}$$

$$|\Psi_\alpha(x)|^2 = \sqrt{\frac{m\omega}{\pi\hbar}} \exp\left(-\frac{m\omega}{\hbar} (x - \langle \hat{x} \rangle_\alpha)^2\right)$$

$$\Delta \hat{x}_\alpha = \sqrt{\frac{\hbar}{2m\omega}} \Rightarrow \frac{m\omega}{\hbar} = \frac{1}{2(\Delta \hat{x}_\alpha)^2}$$

$$|\Psi_\alpha(x)|^2 = \frac{1}{\sqrt{2\pi} \Delta \hat{x}_\alpha} \exp\left(-\frac{1}{2} \left(\frac{x - \langle \hat{x} \rangle_\alpha}{\Delta \hat{x}_\alpha}\right)^2\right)$$

"Gaussian"

$$\Delta \hat{x}_\alpha = \sqrt{\frac{\hbar}{2m\omega}}$$

$$\Delta \hat{p}_\alpha = \sqrt{\frac{m\hbar\omega}{2}}$$

$$\Delta \hat{x}_\alpha \Delta \hat{p}_\alpha = \sqrt{\frac{\hbar}{2m\omega}} \sqrt{\frac{m\hbar\omega}{2}} = \frac{\hbar}{2}$$

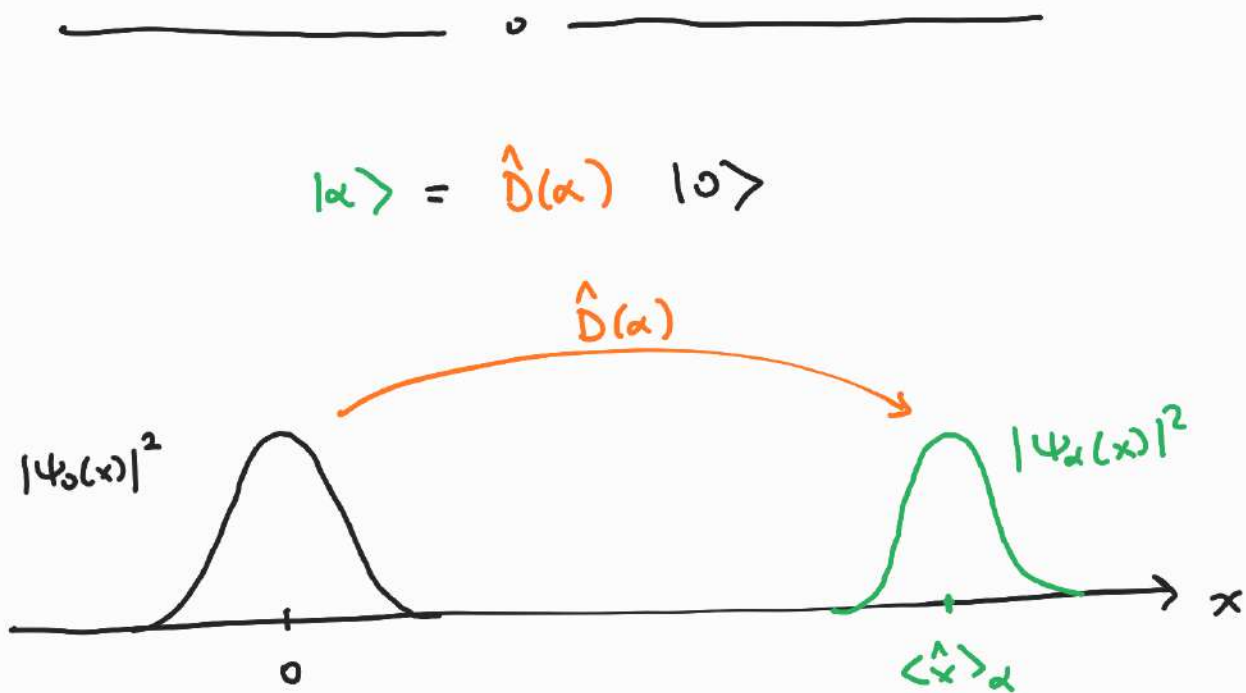
Heisenberg Uncertainty Principle:

$$\Delta \hat{x} \Delta \hat{p} \geq \frac{\hbar}{2}$$

For coherent states $\Delta \hat{x}_\alpha \Delta \hat{p}_\alpha = \frac{\hbar}{2}$

★ MINIMUM UNCERTAINTY STATES "

GAUSSIAN?



Time Evolution:

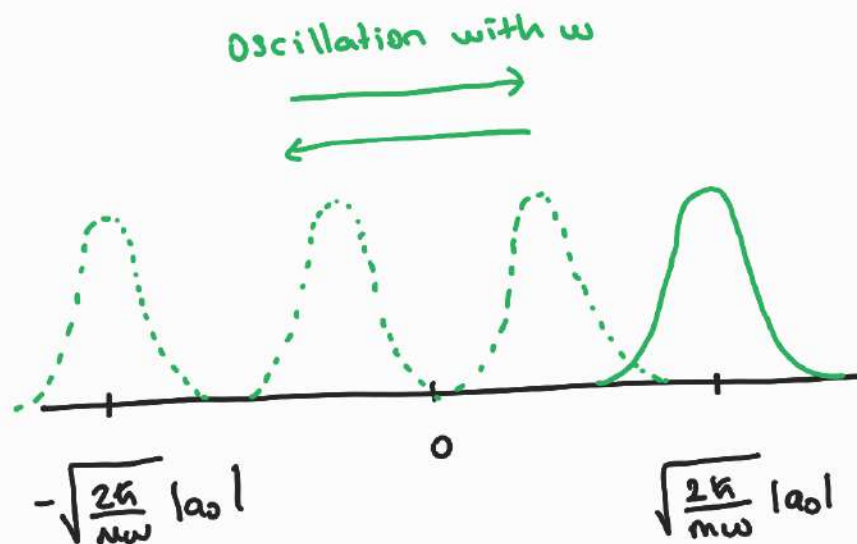
$$|\Psi(0)\rangle = |\alpha_0\rangle \Rightarrow |\Psi(t)\rangle = e^{-i\omega t/2} |\alpha\rangle, \quad \alpha = \alpha_0 e^{-i\omega t}$$

$$|\Psi_\alpha(x, t)|^2 = |\Psi_0(x - \langle \hat{x} \rangle_\alpha(t))|^2$$

$$|\Psi_\alpha(x, t)|^2 = \frac{1}{\sqrt{2\pi} \Delta \hat{x}_\alpha} \exp\left(-\frac{1}{2} \left(\frac{x - \langle \hat{x} \rangle_\alpha(t)}{\Delta \hat{x}_\alpha}\right)^2\right)$$

$$\langle \hat{x} \rangle_\alpha(t) = \sqrt{\frac{2\hbar}{m\omega}} |\alpha_0| \cos(\omega t - \varphi)$$

$$\Delta \hat{x}_\alpha = \sqrt{\frac{\hbar}{2m\omega}} \rightarrow \text{Time Independent, Thus Gaussian does not spread over time.}$$



Where do we use coherent states?

How entanglement
is achieved?

→ In Quantum Optics: Laser beam

. well defined phase and amplitude

→ In QFT

→ In Quantum-Information: creating entangled states
(SPDC)