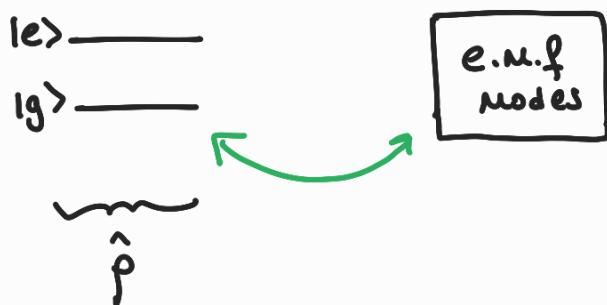


OPTICAL BLOCH EQUATIONS

Two level system + environment



I. Derivation

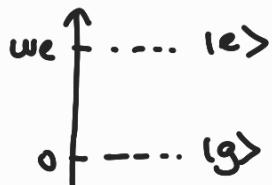
1) Model

$$\hat{H} = \Delta |e\rangle\langle e| + \frac{\Omega}{2} (|e\rangle\langle g| + |g\rangle\langle e|)$$

$\hat{\sigma}_x$ → can be anything

$$\hat{H} = \frac{\Delta}{2} \hat{\mathbb{I}} + \vec{h} \cdot \vec{\sigma}$$

$$\Delta = \omega_e - \omega_{\text{drive}}$$



- \hat{H} in the frame rotating at the drive frequency + RWA

- We have chosen the phase of the external drive

$$\vec{h} \in xy \text{ plane}$$

In general, $\frac{\Omega}{2} (|e\rangle\langle g| e^{i\phi} + |g\rangle\langle e| e^{-i\phi})$, ϕ arbitrary

Spontaneous emission

- $$\hat{L} = \sqrt{\Gamma} |g\rangle\langle e|$$
- jump operator
- jumps produce transitions $|e\rangle \rightarrow |g\rangle$
 - Note: $\text{Tr}(\hat{\rho} \hat{L}^+ \hat{L}) = \Gamma$. $\text{Tr}(\hat{\rho} |e\rangle\langle e|) = \Gamma p_e$
 \downarrow
prob. being
in $|e\rangle$ state

2. Justification

Extended Hilbert Space : $\mathcal{H} = \mathcal{H}_{\text{2-level sys}} \otimes \mathcal{H}_{\text{emf}}$ (fock space)

Light-Matter interaction : $\hat{H}_{\text{int}} = \hat{d} \otimes \vec{E}$

$(\text{dipole moment operator})$
 $(\text{electromagnetic field operator})$

$$\hat{d} = \vec{e}_x d_0 (|e\rangle\langle g| + |g\rangle\langle e|) \quad (\text{Dipolar approximation})$$

\downarrow
polarization vector
reduced dipole moment

$$\vec{E} = -i \sum_{\vec{k}, \lambda} \sqrt{\frac{\hbar \omega_{\vec{k}}}{2 \epsilon_0 V}} \vec{\epsilon}_{\lambda} (e^{-i(\vec{k} \cdot \vec{r} - \omega_{\vec{k}} t)} \hat{a}_{\vec{k}, \lambda} - \text{hermit conj})$$

$\omega_{\vec{k}} = c |\vec{k}|$

V quantization volume

Fermi Golden Rule $\Gamma_{E,\lambda} = \frac{2\pi}{\hbar} \left| \langle \{0\}_{E,\lambda}, e | \hat{j} \otimes \hat{E} | \{1\}_E, g \rangle \right|^2$

$$= \delta(\omega_e - \omega_E)$$

Lindblad Equation

$$\dot{\hat{\rho}} = -i [\hat{H}, \hat{\rho}] + \sum_{E,\lambda} \left(\hat{L}_{E,\lambda} \hat{\rho} \hat{L}_{E,\lambda}^+ - \frac{1}{2} \hat{L}_{E,\lambda}^+ \hat{L}_{E,\lambda} \hat{\rho} - \frac{1}{2} \hat{\rho} \hat{L}_{E,\lambda}^+ \hat{L}_{E,\lambda} \right)$$

Here $\hat{L}_{E,\lambda} = \sqrt{\Gamma_{E,\lambda}} \lg X_{el}$

$$\dot{\hat{\rho}} = -i [\hat{H}, \hat{\rho}] + \left(\sum_{E,\lambda} \Gamma_{E,\lambda} \right) \lg X_{el} \hat{\rho} \lg X_{el} - \frac{1}{2} \lg X_{el} \hat{\rho} - \frac{1}{2} \hat{\rho} \lg X_{el}$$

Replace all $L_{E,\lambda}$ by $\hat{L} = \sqrt{\Gamma} \lg X_{el}$, $\Gamma = \sum_{E,\lambda} \Gamma_{E,\lambda}$

Remark Momentum is not conserved ($\hbar \vec{k}$ is taken by the field)

3. Bloch vector dynamics

$$\dot{\hat{\rho}} = -i [\hat{H}, \hat{\rho}] + \Gamma \left(\lg X_{el} \hat{\rho} \lg X_{el} - \frac{1}{2} \hat{\rho} \lg X_{el} - \frac{1}{2} \lg X_{el} \hat{\rho} \right)$$

Matrix element $\dot{\rho}_{ee} = \langle e | \dot{\hat{\rho}} | e \rangle$

$$= -i \langle e | \hat{H} \hat{\rho} - \hat{\rho} \hat{H} | e \rangle - \Gamma \rho_{ee}$$

$$= -i \frac{\Omega}{2} (\rho_{ge} - \rho_{eg}) - \Gamma \rho_{ee}$$

$$\dot{\rho}_{eg} = -i \langle e \vec{l} \hat{H} \vec{p} - \vec{p} \hat{H} l g \rangle - \frac{\Gamma}{2} \rho_{eg}$$

$$= -i \left(\Delta \rho_{eg} + \frac{\Omega}{2} (\rho_{gg} - \rho_{ee}) \right) - \frac{\Gamma}{2} \rho_{eg}$$

$$\rho_{ee} + \rho_{gg} = 1 \Rightarrow \dot{\rho}_{ee} = -\dot{\rho}_{gg}$$

$$\rho_{ge} = \rho_{eg}^*$$

?
(!)

lexel Igxg1 $(2^\circ)-(0^\circ)$

$$\alpha_z = \rho_{ee} - \rho_{gg} = 2\rho_{ee} - 1$$

$$\hat{p} = \frac{1}{2} (\mathbb{I} + \vec{a} \cdot \vec{\sigma})$$

$$\alpha_x = \rho_{eg} - \rho_{ge}$$

$$\alpha_y = i(\rho_{eg} - \rho_{ge})$$

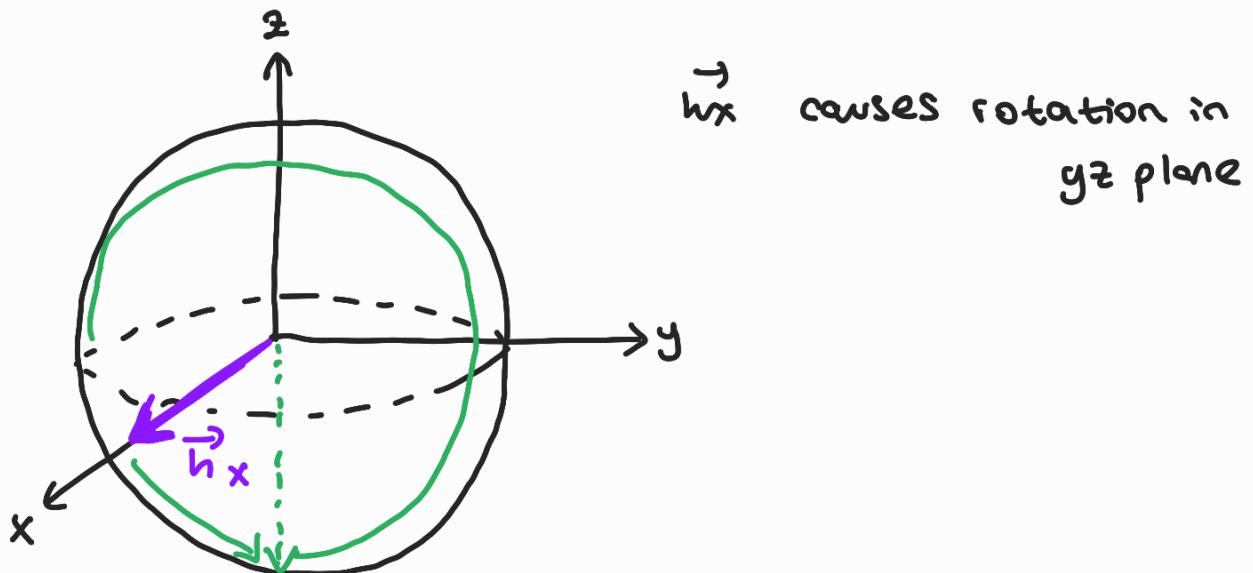
$$\dot{\alpha}_z = -\Omega \alpha_y - \Gamma (\alpha_z + 1)$$

$$\dot{\alpha}_x = -\frac{\Gamma}{2} \alpha_x - \underbrace{\Delta \alpha_y}_{\text{detuning}}$$

$$\dot{\alpha}_y = -\frac{\Gamma}{2} \alpha_y + \underbrace{\Delta \alpha_x}_{\text{external drive}} - \Omega \alpha_z$$

"Optical Bloch"
Equation

* Choice of phase appears in the α_y equation



II. Solutions

Stationary solutions ($\dot{\vec{q}} = 0$)

Let $s = \frac{2\Omega^2}{\Gamma^2 + 4\Delta^2} = \text{saturation parameter}$
 (dimensionless)

$$a_x = -\frac{2\Delta}{\Omega} \frac{s}{1+s}$$

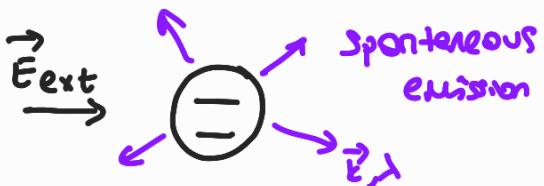
$$a_y = \frac{\Gamma}{\Omega} \frac{s}{(1+s)}$$

$$a_z = \frac{-1}{(1+s)}$$

Interpretation: $p_e = \langle e | \hat{p} | e \rangle = p_{ee} = \frac{1+a_z}{2} = \frac{1}{2} \frac{s}{(1+s)}$

\downarrow
prob. being in 1e>

↔ Rate of photon scattering

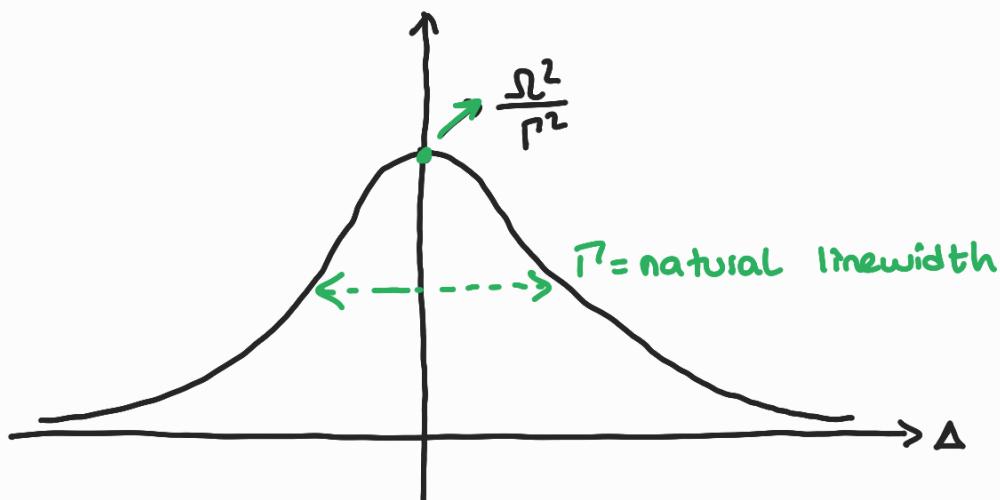


$$\Gamma_{\text{fee}} = \text{rate of spontaneous emission} = \text{photon scattering}$$

. $S \ll 1$: low saturation

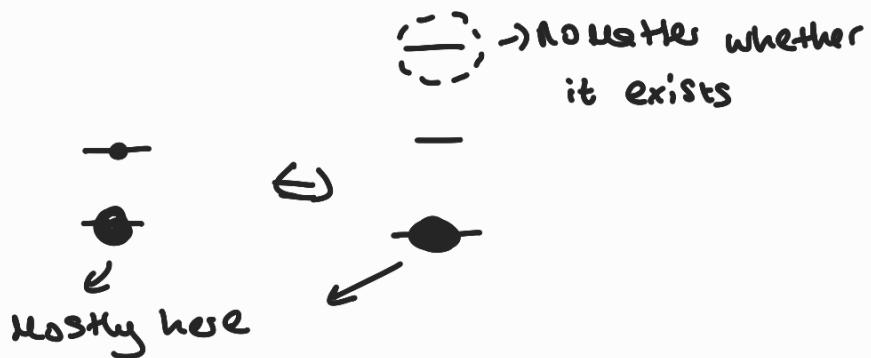
either $\Omega \ll \Gamma, \Delta \dots$
→ large Δ

Then $\rho_{\text{ee}} = \frac{S}{2} = \frac{\Omega^2}{\Gamma^2 + 4\Delta^2} \ll 1$ "Lorentzian Profile"



Remark: Classical harmonic oscillator with damping rate Γ

Low $S \Rightarrow \rho_{\text{ee}}$ low



. $S \gg 1$ Strong saturation $P_{ee} = \frac{1}{2}$, $\Omega \gg \Gamma, \Delta$

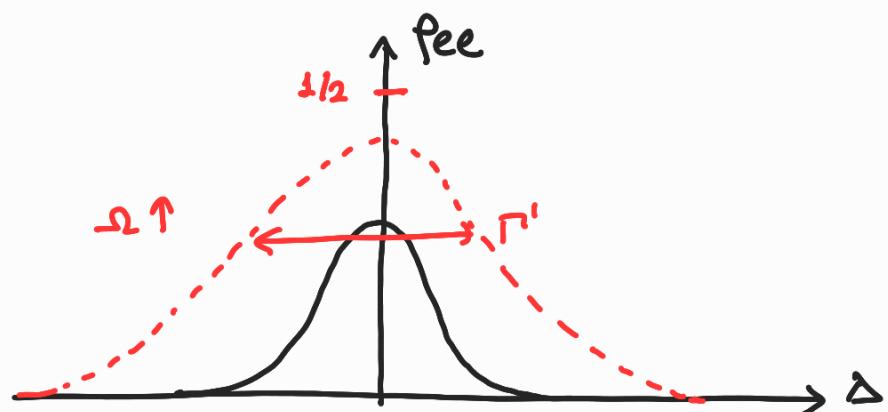
forget Γ : at the scale of $T_{\text{Rabi}} = \frac{2\pi}{\Omega}$,

there is no spontaneous emission.

$\frac{1}{2}$ = time average of the Rabi oscillation

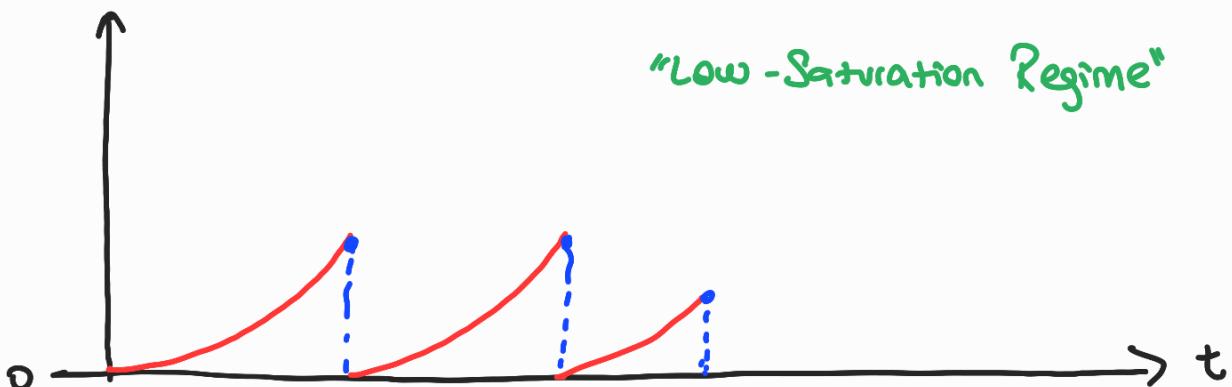
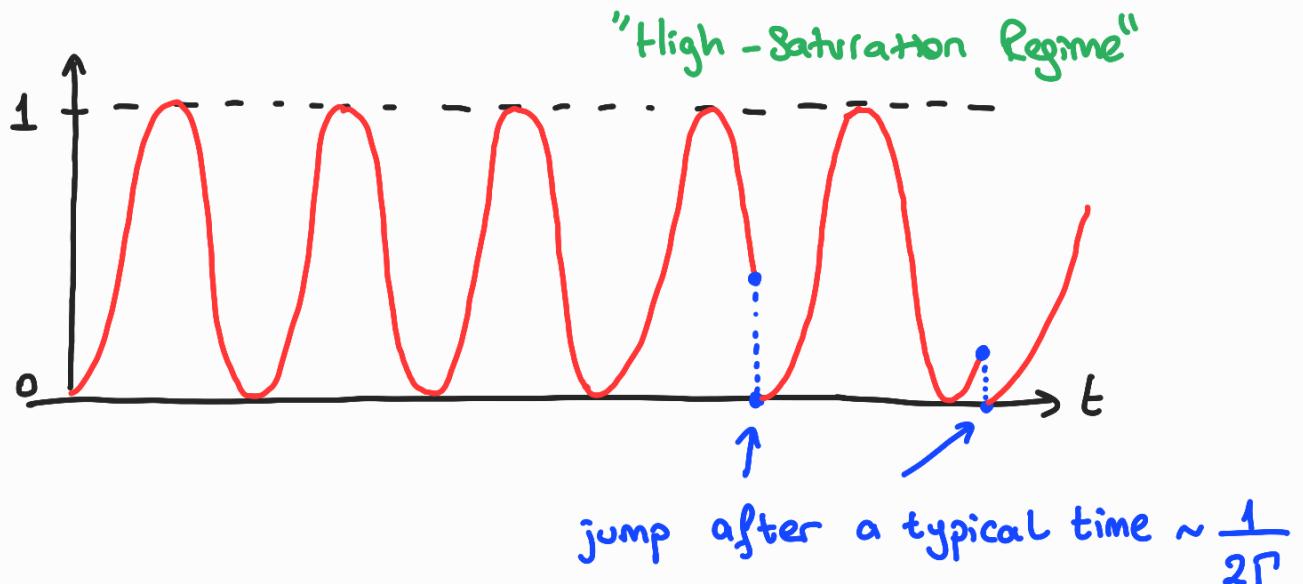
. Intermediate region: $P_{ee} = \frac{\frac{1}{2}\Omega^2}{4\Delta^2 + \Gamma'^2}$, $\Gamma' = \sqrt{\Gamma^2 + 2\Omega^2}$

power broadening



Quantum Trajectory Interpretation

$$|\langle e | \Psi(t) \rangle|^2$$



Remarks: We can include other processes than sp. emission

$$\hat{L}_1 = \sqrt{\Gamma_1} \lg \chi_{el}$$

$$\hat{L}_2 = \sqrt{\Gamma_2} \lg \chi_{el} \quad \text{Dephasing}$$

$$\text{then: } \dot{\rho}_{ee} = -i \langle e | \hat{H}^{\dagger} \hat{p} - \hat{p} \hat{H} | e \rangle - \Gamma_1 \rho_{ee}$$

$$\dot{\rho}_{eg} = -i \langle e | \hat{H}^{\dagger} \hat{p} - \hat{p} \hat{H} | g \rangle - \frac{\Gamma_1}{2} \rho_{eg} - \frac{\Gamma_2}{2} \rho_{eg}$$

In general, T_1 time \rightarrow decay time of PEE
 T_2 time \rightarrow " " " feg

Without dephasing : $T_2 = 2T_1 \rightarrow$ performance Criteria