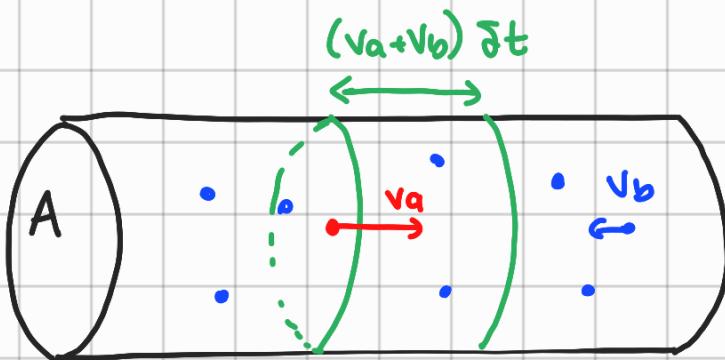
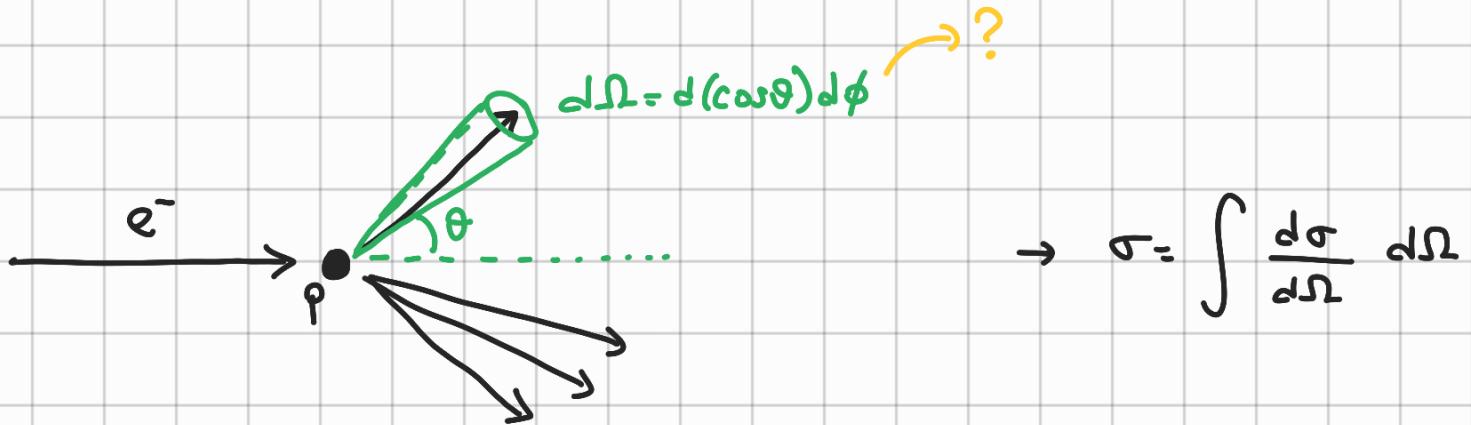


# Week 5 Particle Scattering

## Cross Sections

$$\sigma = \frac{\text{# of interactions per unit time per target}}{\text{incident flux}}$$

$$\frac{d\sigma}{d\Omega} = \frac{\text{# of interactions per unit time per target into } d\Omega}{\text{incident flux}}$$



in time  $\Delta t$  a single particle a traverses region containing  $n_b (v_a + v_b) A \Delta t$  particles of type b

$$\text{Interaction probability} = \frac{n_b (\overbrace{v_a + v_b}^{\downarrow} A) \delta t \sigma}{A}$$

$$= \frac{n_b v \delta t \sigma}{|V|}$$

$$\text{Rate for a single particle } a = n_b v \sigma$$

Consider volume  $V$ , total reaction rate:

$$(n_b v \sigma) \cdot (n_a V) = (n_b V)(n_a v) \sigma$$

$$= N_b \phi_a \sigma$$

Rate = (# of targets)<sub>b</sub> • (Flux)<sub>a</sub> • (Cross Section)

## Cross Section Calculations

"Scattering Process"



why only  $p_3$  &  $p_4$ ?

Fermi's Golden Rule:

$$T_{fi} = (2\pi)^4 \int |T_{fi}|^2 \delta(E_1 + E_2 - E_3 - E_4) \delta^3(\vec{p}_1 + \vec{p}_2 - \vec{p}_3 - \vec{p}_4) \frac{d^3 p_3}{(2\pi)^3} \frac{d^3 p_4}{(2\pi)^3}$$

$$\frac{\text{Rate}}{\text{Volume}} = (\# \text{ Density of 2}) \cdot (\text{Flux of 1}) \times \sigma = n_2 \cdot n_1 (v_1 + v_2) \cdot \sigma$$

For 1 target particle per unit volume



$$\text{Rate} = (v_1 + v_2) \sigma \Rightarrow$$

$$\sigma = \frac{\Gamma f_i}{v_1 + v_2}$$

$$\sigma = \frac{(2\pi)^4}{v_1 + v_2} \int |T_{fi}|^2 \delta(E_1 + E_2 - E_3 - E_4) \delta^3(p_1 + p_2 - p_3 - p_4) \frac{d^3 p_3}{(2\pi)^3} \frac{d^3 p_4}{(2\pi)^3}$$

Lorentz Invariant Parts

To obtain a Lorentz invariant form use wave-functions

normalised to  $2\bar{E}$  particles per unit volume  $\Psi' = (2\bar{E})^{1/2} \Psi$

Again, define  $M_{fi} = (2\bar{E}_1 2\bar{E}_2 2\bar{E}_3 2\bar{E}_4)^{1/2} T_{fi}$

$$\sigma = \frac{(2\pi)^{-2}}{2\bar{E}_1 2\bar{E}_2 (v_1 + v_2)} \int |M_{fi}|^2 \delta(E_1 + E_2 - E_3 - E_4) \delta^3(p_1 + p_2 - p_3 - p_4) \frac{d^3 p_3}{2\bar{E}_3} \frac{d^3 p_4}{2\bar{E}_4}$$

Lorentz-Invariant Form

$F = 2\bar{E}_1 2\bar{E}_2 (v_1 + v_2)$  in the form  
 of 4-vector scalar product  $\rightsquigarrow F = 4 \left[ (p_1^\mu p_{2\mu})^2 - m_1^2 m_2^2 \right]^{1/2}$

## Two Special Cases of Lorentz Invariant Flux

1. CoM frame:  $F = 4 E_1 E_2 (v_1 + v_2)$  How?

$$= 4 E_1 E_2 \left( \frac{|\vec{p}_1^+|}{\epsilon_1} + \frac{|\vec{p}_2^+|}{\epsilon_2} \right)$$

$$= 4 |\vec{p}^+| (\epsilon_1 + \epsilon_2) = 4 |\vec{p}^+| \sqrt{s}$$

Nambu-Goto Variable "s"  
 $\sqrt{s} = (p_1 + p_2)^2$

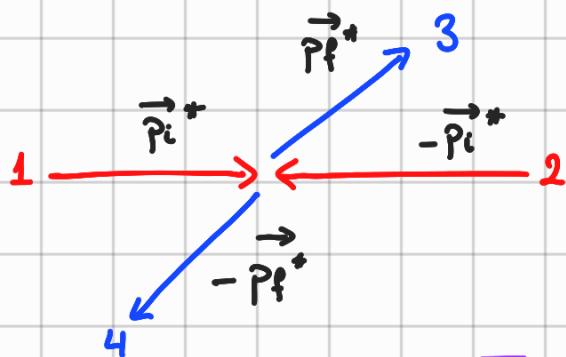
2. Target particle (particle 2) at rest  $\rightarrow \epsilon_2 = M_2, v_2 = 0$

$$F = 4 E_1 \bar{\epsilon}_2 (v_1 + v_2)$$

$$= 4 \bar{\epsilon}_2 M_2 v_1$$

$$= 4 \bar{\epsilon}_2 M_2 \frac{|\vec{p}_1|}{\bar{\epsilon}_1} = 4 M_2 |\vec{p}_1|$$

$2 \rightarrow 2$  Body Scattering in CoM frame



$$\sigma = \frac{(2\pi)^{-2}}{(2\epsilon_1)(2\epsilon_2)(v_1 + v_2)} \frac{\sqrt{s}}{4 |\vec{p}_1^+| \sqrt{s}} \int |M_{fi}|^2 \delta(\epsilon_1 + \epsilon_2 - \epsilon_3 - \epsilon_4) \delta^3(\vec{p}_1 + \vec{p}_2 - \vec{p}_3 - \vec{p}_4) \frac{d^3 p_3}{2\epsilon_3} \frac{d^3 p_4}{2\epsilon_4}$$

$$\vec{p}_1 + \vec{p}_2 = 0 \rightarrow \epsilon_1 + \epsilon_2 = \sqrt{s}$$

$$\sigma = \frac{(2\pi)^{-2}}{4 |\vec{p}_i^*| \sqrt{s}} \int |M_{fi}|^2 \delta(\sqrt{s} - \vec{E}_3 - \vec{E}_4) \delta^3(\vec{p}_3 + \vec{p}_4) \frac{d^3 p_3}{2\epsilon_3} \frac{d^3 p_4}{2\epsilon_4}$$

↓  
similar to case in particle decay calculation

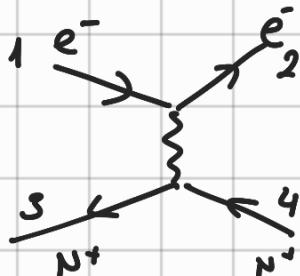
$$\sigma = \frac{(2\pi)^{-2}}{4 |\vec{p}_i^*| \sqrt{s}} \frac{|\vec{p}_f^*|}{4 m_i} \int |M_{fi}|^2 d\Omega^*$$

$$\sigma = \frac{1}{64 \pi^2 s} \frac{|\vec{p}_f^*|}{|\vec{p}_i^*|} \int |M_{fi}|^2 d\Omega^*$$

For an elastic scattering  $|\vec{p}_i^*| = |\vec{p}_f^*|$

total cross section

$$\sigma_{\text{elastic}} = \frac{1}{64\pi^2 s} \int |M_{fi}|^2 d\Omega^*$$



$$d\sigma = \frac{1}{64\pi^2 s} \frac{|\vec{p}_f^*|}{|\vec{p}_i^*|} |M_{fi}|^2 d\Omega^*$$

! Let's find a LI expression  
for  $d\sigma$

$$d(\cos\theta^*) d\phi^*$$

in the C.o.M frame

Mandelstam "t"

$$t = q^2 = (\vec{p}_1 - \vec{p}_3)^2 = \vec{p}_1^2 + \vec{p}_3^2 - 2\vec{p}_1 \cdot \vec{p}_3 = m_1^2 + m_3^2 - 2\vec{p}_1 \cdot \vec{p}_3$$

$$\vec{p}_1^N = (\vec{\epsilon}_1^*, 0, 0, |\vec{p}_1^*|)$$

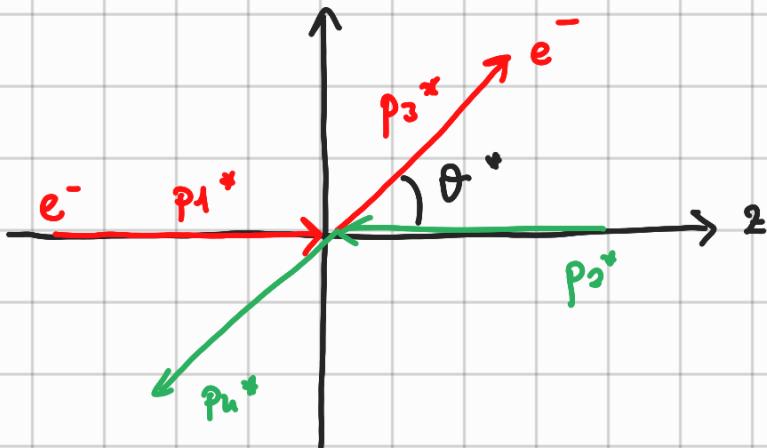
$$\vec{p}_3^N = (\vec{\epsilon}_3^*, |\vec{p}_3^*| \sin\theta^*, 0, |\vec{p}_3^*| \cos\theta^*)$$

$$\vec{p}_1^N \cdot \vec{p}_3^N = \vec{\epsilon}_1^* \vec{\epsilon}_3^* - |\vec{p}_1^*| |\vec{p}_3^*| \cos\theta^*$$

only free parameter

$$t = m_1^2 + m_3^2 - 2\vec{\epsilon}_1^* \vec{\epsilon}_3^* + 2\vec{p}_1^* \vec{p}_3^* \cos\theta^*$$

$$dt = 2\vec{p}_1^* \vec{p}_3^* d(\cos\theta^*)$$



$$d\Omega^* = d(\cos\theta^*) d\phi^*$$

$$d\Omega^* = \frac{dt d\phi^*}{2|\vec{p}_1^*||\vec{p}_3^*|}$$

Mandelstam "t"

$$d\sigma = \frac{1}{64\pi^2 s} \frac{|\vec{p}_f^*|}{|\vec{p}_i^*|} |M_{fi}|^2 d\Omega^* = \frac{1}{2 \cdot 64\pi^2 s |\vec{p}_i^*|^2} |M_{fi}|^2 d\phi^* dt$$

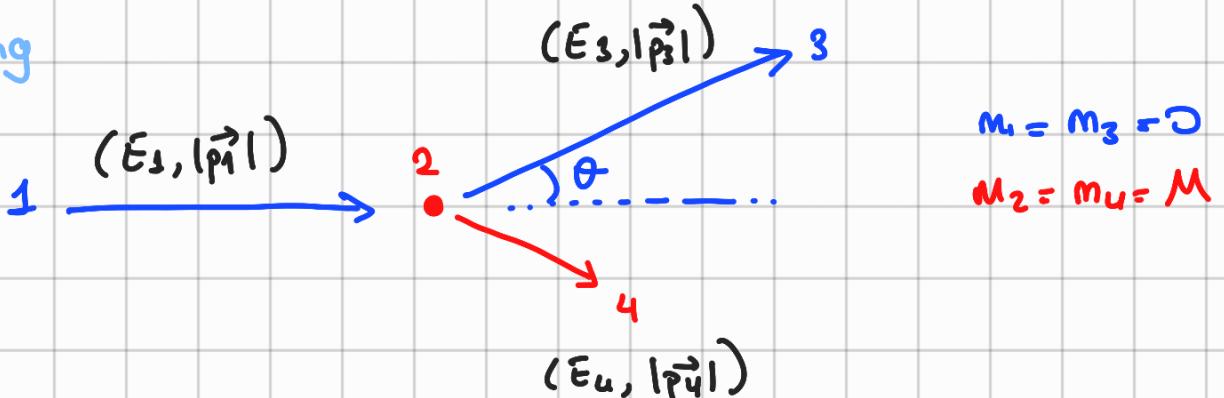
Assume no  $\phi^*$  dependence  
of  $|M_{fi}|^2$

↓ Integrate over  $d\phi^*$ ,  
introduces a factor of  $2\pi$

$$\frac{d\sigma}{dt} = \frac{1}{64\pi s} \frac{1}{p_i^{*2}} |M_{fi}|^2 \rightsquigarrow \text{Lorentz Invariant}$$

$$p_i^{*2} = \frac{1}{4s} [s - (m_1 + m_2)^2] [s - (m_1 - m_2)^2]$$

Scattering



$$d\Omega = 2\pi d(\cos\theta)$$

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma}{dt} \frac{dt}{d\Omega} = \frac{1}{2\pi} \frac{t}{d(\cos\theta)} \frac{d\sigma}{dt}$$

(?)

$$p_1 = (\bar{E}_1, 0, 0, \bar{E}_1)$$

$$p_2 = (M, 0, 0, 0)$$

$$p_3 = (\bar{E}_3, \bar{E}_3 \sin\theta, 0, \bar{E}_3 \cos\theta)$$

$$p_4 = (\bar{E}_4, \vec{p}_4)$$

$$t = (p_1 - p_3)^2 = -2p_1 \cdot p_3 = -2\bar{E}_1 \bar{E}_3 (1 - \cos\theta)$$

look at them later!

from  $(\vec{E}, \vec{p})$  conversion  $p_1 + p_2 = p_3 + p_4$

$$t = (p_2 - p_4)^2 = 2M^2 - 2p_2 p_4 = 2M^2 - 2M\bar{E}_4$$

$$t = 2M^2 - 2M(\bar{E}_1 + M - \bar{E}_3) = -2M(\bar{E}_1 - \bar{E}_3)$$

$\bar{E}_1$  is a constant (the energy of the incoming particle)

$$\frac{dt}{d(\cos\theta)} = 2M \frac{d\bar{E}_3}{d(\cos\theta)}$$

$$\bar{E}_3 = \frac{\bar{E}_1 M}{M + \bar{E}_1 - \bar{E}_1 \cos\theta}$$

$$\frac{d\bar{E}_3}{d\cos\theta} = \frac{\bar{E}_1^2 M}{(M + \bar{E}_1 - \bar{E}_1 \cos\theta)^2} = \bar{E}_1^2 M \left( \frac{\bar{E}_3}{\bar{E}_1 M} \right)^2 = \frac{\bar{E}_3^2}{M}$$

$$\frac{d\sigma}{d\Omega} = \frac{1}{2\pi} \frac{dt}{d(\cos\theta)} \frac{d\sigma}{dt} = \frac{1}{2\pi} 2M \frac{\bar{E}_3^2}{M} \frac{d\sigma}{dt}$$

$$= \frac{\bar{E}_3^2}{\pi} \frac{d\sigma}{dt} = \frac{\bar{E}_3^2}{\pi} \frac{1}{16\pi (s-M^2)^2} |N_{fi}|^2$$

)

equate  
these

$$S = (p_1 + p_2)^2 = M^2 + 2 \cdot p_1 \cdot p_2 = M^2 + 2M E_1$$

since  $p_2^2 = 0$

$$(S - M^2) = 2M E_1$$

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2} \left( \frac{\bar{E}_3}{M\bar{E}_1} \right)^2 |M_{fi}|^2$$

in the  
limit of  $M_1 \rightarrow 0$

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2} \left( \frac{1}{M + \bar{E}_1 - \bar{E}_1 \cos\theta} \right)^2 |M_{fi}|^2$$

# Summary

Particle decay:  $\Gamma = \frac{|\vec{p}^*|}{32\pi^2 m_i^2} \int |M_{fi}|^2 d\Omega$

where  $p^* = \frac{1}{2m_i} \sqrt{[m_i^2 - (m_1 + m_2)^2][m_i^2 - (m_1 - m_2)^2]}$  ?

$$\sigma = \frac{1}{64\pi^2 s} \frac{|\vec{p}_f^*|}{|\vec{p}_i^*|} \int |M_{fi}|^2 d\Omega^*$$

Lab frame ( $m_1=0$ )

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2} \left( \frac{\bar{E}_3}{M\bar{E}_1} \right)^2 |M_{fi}|^2 \Leftrightarrow \frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2} \left( \frac{1}{M + \bar{E}_1 - \bar{E}_3 \cos\theta} \right)^2 |M_{fi}|^2$$

Lab frame ( $m_1 \neq 0$ )

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2} \cdot \frac{1}{|\vec{p}_1|m_1} \cdot \frac{|\vec{p}_3|^2}{|\vec{p}_3|(\bar{E}_1 + m_2) - \bar{E}_3 |\vec{p}_1| \cos\theta} \cdot |M_{fi}|^2$$

$$\bar{E}_1 + m_2 = \sqrt{|\vec{p}_3|^2 + m_3^2} + \sqrt{|\vec{p}_1|^2 + |\vec{p}_3|^2 - 2|\vec{p}_1||\vec{p}_3| \cos\theta + m_4^2}$$
 ?