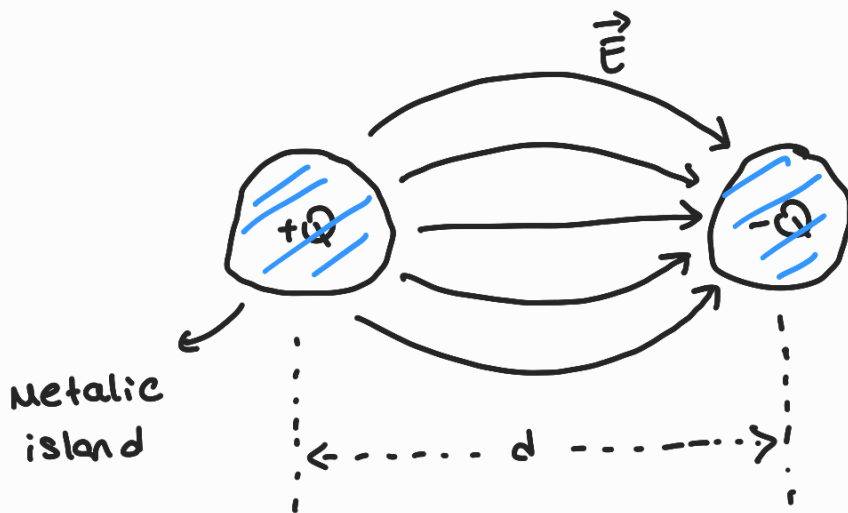


Quantization of electrical circuits

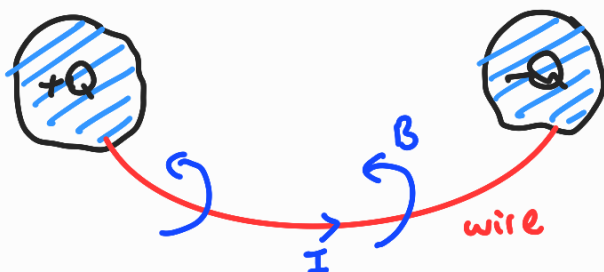


The energy required to move one unit charge from one to the other island is path-independent and given by the voltage V .

$$V = \frac{Q}{C}$$

Total energy stored in the electric field:

$$E_{el} = \int_0^Q dQ' V(Q') = \int_0^Q dQ' \frac{Q'}{C} = \frac{Q^2}{2C}$$



$$I = - \frac{dQ}{dt} = -\dot{Q}$$

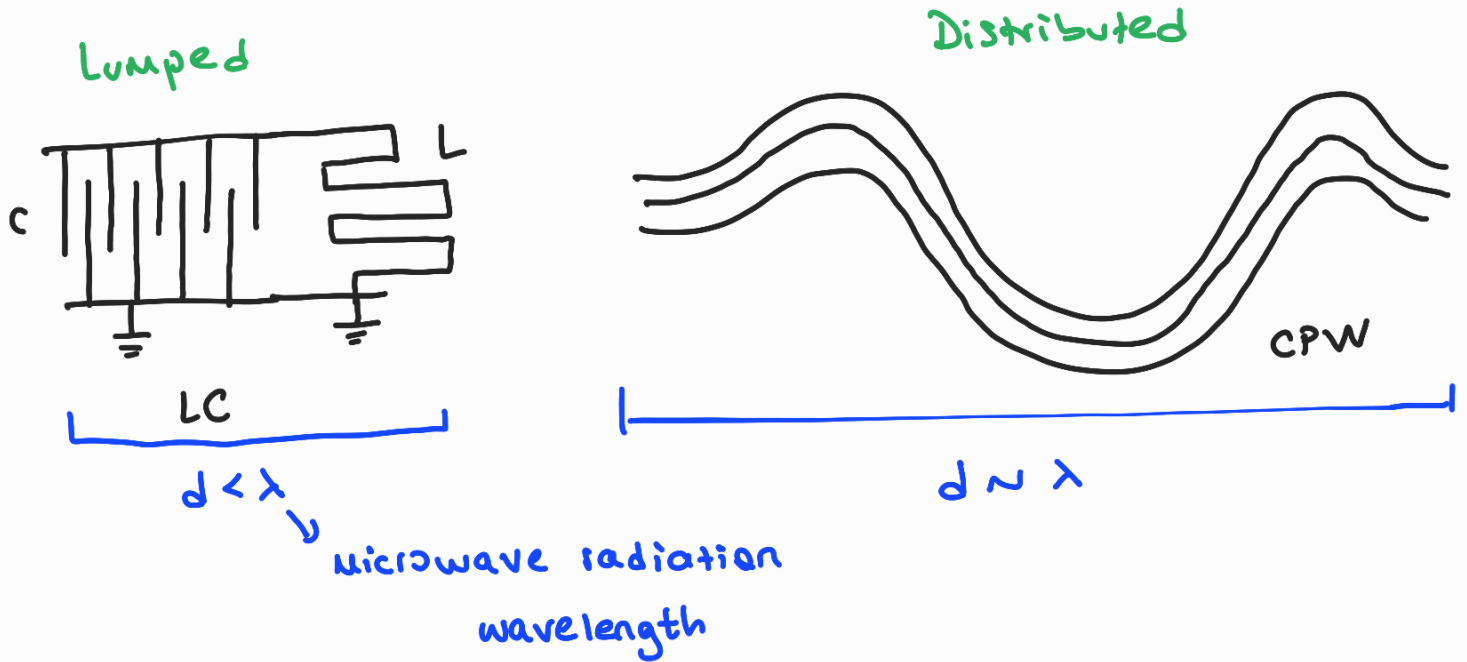
$$\Phi_B = L I$$

Φ_B ← Magnetic flux
 L ← inductance (geometry dependent)
 I ← induced voltage

$$\Phi_B(t) = \int_0^t V(t') dt' \quad (?)$$

Energy stored in magnetic field \vec{B} :

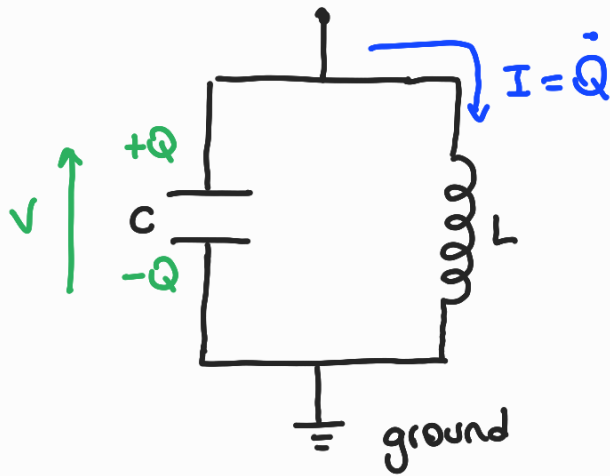
$$\mathcal{E}_{\text{mag}} = \int_0^I dI' \Phi(I') = \frac{1}{2} LI^2 = \frac{\Phi^2}{2L}$$



For GHz photons $\rightarrow \lambda \sim \text{cm}$

For the low frequency dynamics of this system are fully captured by two effective parameters L, C .

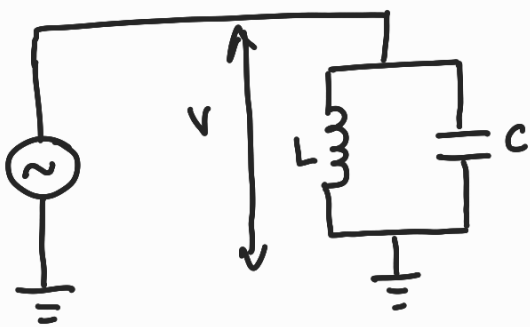
Quantization of the LC Resonator



Faraday's law of induction $\dot{\Phi} = V$

$$\ddot{\Phi} = \dot{V} = \frac{\dot{Q}}{C} = -\frac{I}{C} = -\frac{\Phi}{LC} = -\omega_0^2 \Phi, \quad \omega_0 = \frac{1}{\sqrt{LC}}$$

Electronic Harmonic Oscillator



$$Q = CV$$

$$\Phi = LI$$

$$V = \frac{Q}{C} = -L\dot{I} = -\dot{\Phi}$$

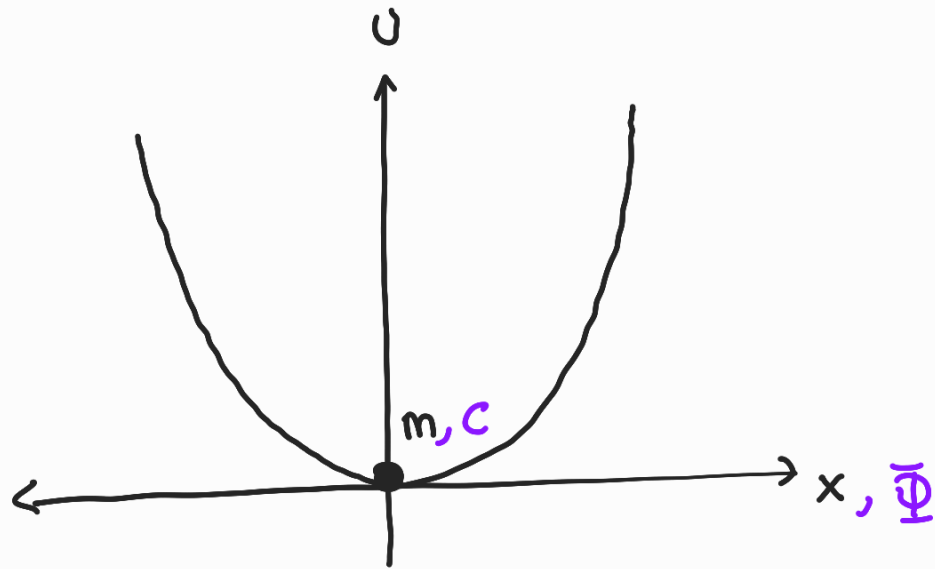
Hamiltonian:

$$\frac{CV^2}{2} + \frac{LI^2}{2} = \frac{Q^2}{2C} + \frac{\Phi^2}{2L}$$

Mechanical
Oscillator

$$H = \frac{p^2}{2m} + \frac{kx^2}{2}$$

\swarrow kinetic, K \searrow potential, U



Characteristic Quantities

Mechanical

Position x

Momentum p

Mass m

Spring constant k

resonance freq $\omega = \sqrt{\frac{k}{m}}$

$$\hat{x} = x$$

$$\hat{p} = -i\hbar \frac{\partial}{\partial x}$$

Electronic

Flux Φ

charge Q

capacitance C

inverse inductance $1/L$

$$\omega = \sqrt{\frac{1}{LC}}$$

$$\hat{\Phi} = \Phi$$

$$\hat{Q} = -i\hbar \frac{\partial}{\partial \Phi}$$

$$[\hat{x}, \hat{p}] = (\hat{x}\hat{p} - \hat{p}\hat{x})\psi$$

$$= x(-i\hbar)\frac{\partial}{\partial x}\psi - (-i\hbar)\frac{\partial}{\partial x}(x\psi)$$

$$= \cancel{-i\hbar x \frac{\partial \psi}{\partial x}} + i\hbar\psi + i\hbar \cancel{\frac{\partial \psi}{\partial x} x} = i\hbar$$

$$[\hat{\Phi}, \hat{Q}] = i\hbar$$

HAMILTONIAN OPERATOR

Using conjugate variables Q, Φ

$$\hat{H} = \frac{\hat{Q}^2}{2C} + \frac{\Phi^2}{2L} = -\frac{\hbar^2}{2C} \frac{\partial^2}{\partial \Phi^2} + \frac{1}{2L} \Phi^2$$

$m = \text{number operators}$

$$\hat{H} = \hbar\omega \left[\hat{a}^\dagger \hat{a} + \frac{1}{2} \right]$$

$$\hat{a}^\dagger = \frac{1}{\sqrt{2\hbar Z_C}} \left(Z_C \hat{Q}^\dagger - i \hat{\Phi}^\dagger \right) \rightarrow \text{creation}$$

$$\hat{a} = \frac{1}{\sqrt{2\hbar Z_C}} \left(Z_C \hat{Q} + i \hat{\Phi} \right) \rightarrow \text{annihilation}$$

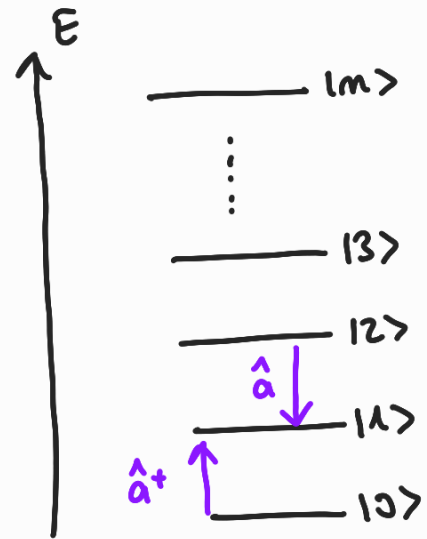
with $Z_c = \sqrt{\frac{L}{C}}$ impedance of oscillator

Properties of \hat{a}^+ and \hat{a}

$$\hat{a}^+ |m\rangle = \sqrt{m+1} |m+1\rangle$$

$$\hat{a} |m\rangle = \sqrt{m} |m-1\rangle$$

$$\hat{a}^+ \hat{a} |m\rangle = m |m\rangle$$



with $|m\rangle$ number (Fock) state of harmonic oscillator

$$\hat{Q} = \sqrt{\frac{\hbar}{2Z_c}} (\hat{a}^+ + \hat{a})$$

$$\hat{\Phi} = i \sqrt{\frac{\hbar Z_c}{2}} (\hat{a}^+ - \hat{a})$$

$$\hat{V} = \sqrt{\frac{\hbar \omega}{2C}} (\hat{a}^+ + \hat{a})$$

$$\hat{I} = i \sqrt{\frac{\hbar \omega}{2L}} (\hat{a}^+ - \hat{a})$$

$$\left. \begin{array}{l} \hat{V} \\ \hat{I} \end{array} \right\} \text{ where } \omega = \frac{1}{\sqrt{LC}} \quad V = \frac{Q}{C} \quad I = \frac{\Phi}{L}$$

Lagrangian Formalism

$$\mathcal{L}(x, \dot{x})$$

$$\mathcal{L} = E_k - E_p$$

Hamiltonian $\leftarrow H = \mathcal{H}(x, p)$

Conjugate variable

$$p = \frac{\partial \mathcal{L}}{\partial \dot{x}}$$

(generalized momentum)

Connection between Lagrangian and Hamiltonian via Legendre form:

$$\mathcal{H}(x, p) = \dot{x}p - \mathcal{L}(x, \dot{x})$$

↓
for complicated circuits (?)