

Particle Physics 1: Exercise 7

Exercise 1

Verify the statement that the Einstein energy-momentum relationship is recovered if any of the four orthogonal Dirac spinors derived in the lecture are substituted into the Dirac equation written in terms of momentum, $(\gamma^\mu p_\mu - m)u = 0$.

Exercise 2

For a particle with four-momentum $p^\mu = (e, \vec{p})$, the general solution to the free-particle Dirac equation can be written

$$\psi(p) = [au_1(p) + bu_2(p)]e^{i(\vec{p}\cdot\vec{x} - Et)}$$

Using the explicit form for u_1 and u_2 , show that the four-vector current $j^\mu = (\rho, \vec{j})$ is given by

$$j^\mu = 2p^\mu.$$

Furthermore, show that the resulting probability density and probability current are consistent with a particle moving with velocity $\beta = p/E$.

Exercise 3

Show that up to an overall phase factor

$$\hat{P}u_\uparrow(\theta, \phi) = u_\downarrow(\pi - \theta, \pi + \phi)$$

where u_\uparrow and u_\downarrow are, respectively, the right-handed and left-handed helicity particle spinors. Comment on the result.

Exercise 4

Under the combined operation of parity and charge conjugation ($\hat{C}\hat{P}$) spinors transform as

$$\psi \rightarrow \psi^C = \hat{C}\hat{P}\psi = i\gamma^2\gamma^0\psi^*.$$

Show that up to an overall phase factor

$$\hat{C}\hat{P}u_\uparrow(\theta, \phi) = v_\downarrow(\pi - \theta, \pi + \phi).$$

Exercise 1

$$\vec{p}_N = (\bar{\epsilon}, -\vec{p}_x, -\vec{p}_y, -\vec{p}_z)$$

Verify the statement that the Einstein energy-momentum relationship is recovered if any of the four orthogonal Dirac spinors derived in the lecture are substituted into the Dirac equation written in terms of momentum, $(\gamma^\mu p_\mu - m)u = 0$.

$$\gamma^\mu p_\mu - m = \bar{\epsilon} \begin{bmatrix} 1 & & & \\ & 1 & & \\ & & -1 & \\ & & & -1 \end{bmatrix} - p_x \begin{bmatrix} & 1 & & \\ & & 1 & \\ & & & -1 \\ & -1 & & \end{bmatrix} - p_y \begin{bmatrix} & & 1 & \\ & & & i \\ & i & & \\ & & -i & \end{bmatrix}$$

$\underbrace{\hspace{1cm}}_{\gamma^0}$ $\underbrace{\hspace{1cm}}_{\gamma^1}$

$$-p_z \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} - m I \rightarrow \begin{bmatrix} 1 & & & \\ & 1 & & \\ & & 1 & \\ & & & 1 \end{bmatrix}$$

$$(\gamma^\mu p_\mu - m) u = \begin{bmatrix} \bar{\epsilon} - m & 0 & -p_z & -p_x + ip_y \\ 0 & \bar{\epsilon} - m & -p_x - ip_y & +p_z \\ p_z & p_x - ip_y & -(\bar{\epsilon} + m) & 0 \\ p_x + ip_y & -p_z & 0 & -(\bar{\epsilon} + m) \end{bmatrix} \begin{bmatrix} 1 \\ 0 \\ \frac{p_z}{\bar{\epsilon} + m} \\ \frac{p_x + ip_y}{\bar{\epsilon} + m} \end{bmatrix} = 0$$

$\tilde{\epsilon}^2 = p^2 + m^2$

$$\begin{bmatrix} (\bar{\epsilon} - m) + 0 - \frac{p_z^2}{\bar{\epsilon} + m} + \frac{-p_x^2 - p_y^2}{\bar{\epsilon} + m} \\ 0 + 0 - \frac{p_z(p_x + ip_y)}{\bar{\epsilon} + m} + \frac{p_z(p_x + ip_y)}{\bar{\epsilon} + m} \\ p_z + 0 - p_z + 0 \\ (p_x + ip_y) + 0 + 0 - (p_x + ip_y) \end{bmatrix} = \frac{1}{\bar{\epsilon} + m} \begin{bmatrix} \tilde{\epsilon}^2 - m^2 - p_x^2 - p_y^2 - p_z^2 \\ 0 \\ 0 \\ 0 \end{bmatrix} = 0$$

Exercise 2

For a particle with four-momentum $p^\mu = (e, \vec{p})$, the general solution to the free-particle Dirac equation can be written

$$\psi(p) = [au_1(p) + bu_2(p)]e^{i(\vec{p} \cdot \vec{x} - Et)}$$

Using the explicit form for u_1 and u_2 , show that the four-vector current $j^\mu = (\rho, \vec{j})$ is given by

$$j^\mu = 2p^\mu.$$

Furthermore, show that the resulting probability density and probability current are consistent with a particle moving with velocity $\beta = p/E$.

for arbitrary spinors Ψ and ϕ

$$\bar{\Psi} \gamma^0 \phi = (\Psi_1^*, \Psi_2^*, -\Psi_3^*, -\Psi_4^*) \begin{pmatrix} 1 & & & \\ & 1 & & \\ \downarrow & & -1 & \\ \Psi_1^* \gamma^0 & & & -1 \end{pmatrix} \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \\ \phi_4 \end{pmatrix}$$

$$(\Psi_1^*, \Psi_2^*, \Psi_3^*, \Psi_4^*) \begin{pmatrix} 1 & & & \\ & 1 & & \\ & & -1 & \\ & & & -1 \end{pmatrix}$$

$$= (\Psi_1^*, \Psi_2^*, -\Psi_3^*, -\Psi_4^*) \begin{pmatrix} \phi_1 \\ \phi_2 \\ -\phi_3 \\ -\phi_4 \end{pmatrix}$$

$$\bar{\Psi} \gamma^0 \phi = \Psi_1^* \phi_1 - \Psi_2^* \phi_2 + \Psi_3^* \phi_3 + \Psi_4^* \phi_4$$

Repeat this procedure for $\mu = 1, 2, 3$

$$\bar{\Psi} \gamma^1 \phi = \Psi_1^* \phi_4 + \Psi_2^* \phi_3 + \Psi_3^* \phi_2 + \Psi_4^* \phi_1$$

$$\bar{\Psi} \gamma^2 \phi = -i (\Psi_1^* \phi_4 - \Psi_2^* \phi_3 + \Psi_3^* \phi_2 - \Psi_4^* \phi_1)$$

$$\bar{\Psi} \gamma^3 \phi = \Psi_1^* \phi_3 - \Psi_2^* \phi_4 + \Psi_3^* \phi_1 - \Psi_4^* \phi_2$$

For the free particle spinor $v_1(p) = \sqrt{E+m} \begin{pmatrix} 1 \\ 0 \\ p_z/E+m \\ (p_x+ip_y)/E+m \end{pmatrix}$

$$\begin{aligned}\bar{v}_1 \gamma^0 v_1 &= (E+m) \left[1 + \frac{p_z^2}{(E+m)^2} + \frac{(p_x^2 + p_y^2)}{(E+m)^2} \right] \\ &= (E+m) \left[1 + \frac{p^2}{(E+m)^2} \right] \\ &= \frac{(E+m)^2 + p^2}{E+m} = \frac{E^2 + 2Em + m^2 + p^2}{E+m} = \frac{2E^2 + 2Em}{E+m} \\ &= \frac{2E(E+m)}{(E+m)} = 2E\end{aligned}$$

Repeating this (n) $\bar{v}_1 \gamma^N v_1 = (2E, 2p_x, 2p_y, 2p_z) = 2p^N$

$$\bar{v}_1 \gamma^N v_1 = \bar{v}_2 \gamma^N v_2 = \bar{v}_1 \gamma^N v_1 = \bar{v}_2 \gamma^N v_2 = 2p^N$$

All corresponding cross terms give zero

$$\bar{v}_1 \gamma^N v_2 = \bar{v}_2 \gamma^N v_1 = \dots$$

For a particle $\Psi = v(p) e^{ip \cdot x}$

$$\bar{\Psi} = \Psi^\dagger \gamma^0 = v(p)^\dagger \gamma^0 e^{-ip \cdot x} = \bar{v}(p) e^{-ip \cdot x}$$

$$j^N = \bar{\Psi} \gamma^N \Psi = \bar{v} \gamma^N v = \bar{v} \gamma^N v$$

↓
for anti-particle
spin

Given $v = \alpha_1 v_1 + \alpha_2 v_2$ with $|\alpha_1|^2 + |\alpha_2|^2 = 1$

$$\bar{v} \gamma^N v = |\alpha_1|^2 \underbrace{\bar{v}_1 \gamma^N v_1}_{2\rho^N} + |\alpha_2|^2 \underbrace{\bar{v}_2 \gamma^N v_2}_{2\rho^N} = 2\rho^N$$

$$\dot{j}^N = \bar{v} \gamma^N v = 2\rho^N$$

$$\dot{j}^N = (\rho, \vec{j})$$

$$\rho = 2\bar{v} \quad \vec{j} = 2\vec{p}$$

* $E = \gamma_{NC} = \gamma_N$ $\vec{p} = \gamma_N \vec{v} = \bar{v} \vec{v}$ $\rightsquigarrow \vec{j} = \frac{2\bar{v}}{\rho} \vec{v}$

$\vec{j} = \rho \vec{v}$

Exercise 3

Show that up to an overall phase factor

$$\hat{P} u_\uparrow(\theta, \phi) = u_\downarrow(\pi - \theta, \pi + \phi)$$

where u_\uparrow and u_\downarrow are, respectively, the right-handed and left-handed helicity particle spinors.

Comment on the result.

$$\hat{P} = \gamma^0 = \begin{pmatrix} 1 & & & \\ & 1 & & \\ & & -1 & \\ & & & -1 \end{pmatrix} \quad u_\uparrow(\theta, \phi) = \sqrt{E+m} \begin{pmatrix} \cos(\frac{\theta}{2}) \\ \sin(\frac{\theta}{2}) e^{i\phi} \\ \frac{p}{E+m} \cos(\frac{\theta}{2}) \\ \frac{p}{E+m} e^{i\phi} \sin(\frac{\theta}{2}) \end{pmatrix}$$

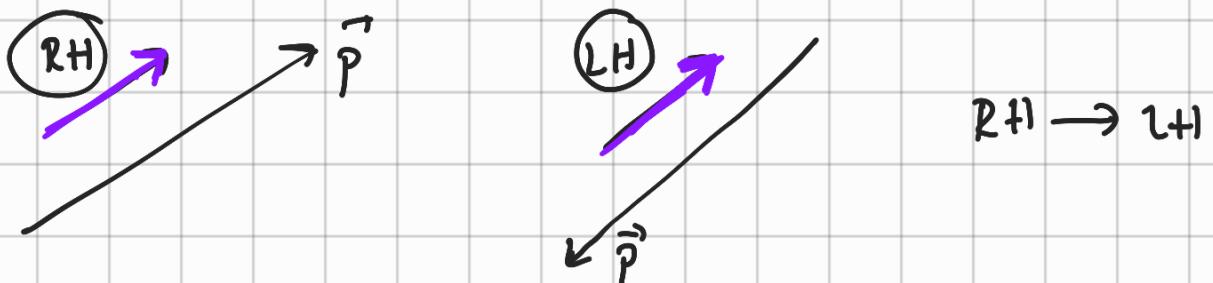
$$\hat{P} u_\downarrow(\theta, \phi) = \sqrt{E+m} \begin{pmatrix} c \\ s e^{i\phi} \\ -\frac{p}{E+m} c \\ -\frac{p}{E+m} s e^{i\phi} \end{pmatrix}$$

$$v_{\downarrow}(\pi-\theta, \pi+\phi) = \sqrt{E+m}$$

$$\left(\begin{array}{c} -\sin\left(\frac{\pi-\theta}{2}\right) \\ \cos\left(\frac{\pi-\theta}{2}\right) e^{i(\pi+\phi)} \\ \left(\frac{p}{E+m}\right) \sin\left(\frac{\pi-\theta}{2}\right) \cos\frac{\theta}{2} \\ -\frac{p}{E+m} \cos\left(\frac{\pi-\theta}{2}\right) e^{i(\phi+\pi)} \\ \sin\left(\frac{\theta}{2}\right) \end{array} \right)$$

$$= \sqrt{E+m} \left(\begin{array}{c} -c \\ -se^{i\phi} \\ \frac{p}{E+m} c \\ \frac{p}{E+m} se^{i\phi} \end{array} \right)$$

$\vec{p} \rightarrow -\vec{p}$ → The parity operator reverses \vec{p}
 but leaves the spin orientation
 UNCHANGED



Exercise 4

Under the combined operation of parity and charge conjugation ($\hat{C}\hat{P}$) spinors transform as

$$\psi \rightarrow \psi^C = \hat{C}\hat{P}\psi = i\gamma^2\gamma^0\psi^*.$$

Show that up to an overall phase factor

$$\hat{C}\hat{P}u_{\uparrow}(\theta, \phi) = v_{\downarrow}(\pi - \theta, \pi + \phi).$$

$$\gamma^0 = \begin{pmatrix} 1 & & & \\ & 1 & & \\ & & -1 & \\ & & & -1 \end{pmatrix} \quad \gamma^2 = \begin{pmatrix} & & -i \\ & i & \\ -i & & \end{pmatrix} \quad u_{\uparrow}(\theta, \phi) = \sqrt{E+m} \begin{pmatrix} c \\ se^{i\phi} \\ \frac{p}{E+m} c \\ \frac{p}{E+m} se^{i\phi} \end{pmatrix}$$

$$\hat{C} \hat{P} v_{\uparrow}(\theta, \phi) = i \gamma^2 \gamma^0 v_{\uparrow}(\theta, \phi)^* = i \underbrace{\begin{pmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \\ 0 & i & 0 & 0 \\ -i & 0 & 0 & 0 \end{pmatrix}}_{\text{Matrix}} \underbrace{\begin{pmatrix} 1 & 1 \\ 1 & -1 \\ -1 & -1 \end{pmatrix}}_{v_{\uparrow}^*}$$

\downarrow

$i \gamma^2 \psi^*$ $\gamma^0 \psi$

$$= \sqrt{E+m} \begin{pmatrix} 1 & & & \\ & -1 & 1 & \\ & & -1 & \\ 1 & & & \end{pmatrix} \begin{pmatrix} c \\ s e^{-i\phi} \\ -p/E+m c \\ -p/E+m s e^{-i\phi} \end{pmatrix}$$

$$v_{\downarrow}(\pi, \theta, \pi + \phi) = \sqrt{E+m} \begin{pmatrix} \frac{-p}{E+m} s e^{-i\phi} \\ \frac{p}{E+m} c \\ -s e^{i\phi} \\ c \end{pmatrix}$$

$$v_{\downarrow}(\pi - \theta, \phi + \pi) = \sqrt{E+m} \begin{pmatrix} \frac{p}{E+m} s \\ -\frac{p}{E+m} c e^{i\phi} \\ s \\ -c e^{i\phi} \end{pmatrix} = -e^{i\phi} \sqrt{E+m} \begin{pmatrix} \frac{-p}{E+m} s e^{-i\phi} \\ \frac{p}{E+m} c \\ -s e^{-i\phi} \\ c \end{pmatrix}$$

$$\hat{C} \hat{P} v_{\uparrow}(\theta, \phi) = -e^{i\phi} v_{\downarrow}(\pi - \theta, \pi + \phi)$$

\downarrow

Overall phase factor