

1. Proof of Schmidt Theorem

$$\mathcal{H}_{\text{tot}} = \mathcal{H}_A \otimes \mathcal{H}_B \quad \dim \mathcal{H}_A = D_A, \quad \dim \mathcal{H}_B = D_B$$

Let the system be a pure vector state $|\Psi\rangle \in \mathcal{H}_{\text{tot}}$

Let $\begin{cases} g_A = \text{Tr}_B |\Psi\rangle \langle \Psi| \\ g_B = \text{Tr}_A |\Psi\rangle \langle \Psi| \end{cases}$

a) Non-zero eigenvalues of g_A and g_B are identical

possibly they have zero eigenvalues but in differing number

As a consequence :

$$S(g_A) = S(g_B)$$

b)
$$g_A = \sum_{i=1}^k \lambda_i |\psi_i\rangle_A \langle \psi_i|_A$$
 $\lambda_i \neq 0$

$$g_B = \sum_{i=1}^k \lambda_i |\psi_i\rangle_B \langle \psi_i|_B$$

Spectral Decomposition

with $\{|\psi_i\rangle_A\}$ orthonormal & $\{|\psi_i\rangle_B\}$ are orthonormal

c)
$$|\Psi\rangle = \sum_{i=1}^k \sqrt{\lambda_i} |\psi_i\rangle_A \otimes |\psi_i\rangle_B$$

↓
pure state

Proof:

First: we check that b implies a

($\lambda_i \neq 0$) are the same in \mathcal{g}_A & \mathcal{g}_B

$$S(\mathcal{g}_A) = S(\mathcal{g}_B) = - \sum_{i=1}^k \lambda_i \log \lambda_i$$

Second remark: c implies b

$$|\Psi\rangle\langle\Psi| = \sum_{i,j} \sqrt{\lambda_i} \sqrt{\lambda_j} |\varphi_i\rangle_A \otimes |\psi_j\rangle_B \langle\varphi_j|_A \otimes \langle\psi_j|_B$$

$$\mathcal{g}_A = \text{Tr}_B (|\Psi\rangle\langle\Psi|) = \sum_{i,j} \sqrt{\lambda_i} \sqrt{\lambda_j} |\varphi_i\rangle_A \langle\varphi_j|_A \underbrace{\text{Tr}_B (|\Psi\rangle\langle\Psi|_B)}_{\delta_{ji} = \begin{cases} 1 & i=j \\ 0 & i \neq j \end{cases}}$$

$$\boxed{\mathcal{g}_A = \sum_i \lambda_i |\varphi_i\rangle_A \langle\varphi_i|_A}$$

$$\boxed{\mathcal{g}_B = \sum_i \lambda_i |\psi_i\rangle_B \langle\psi_i|_B}$$

Prove c:

$$|\Psi\rangle = \sum_{i=1}^k \lambda_i |\varphi_i\rangle_A \otimes |\psi_i\rangle_B$$

$$|\Psi\rangle \in \underbrace{\mathcal{H}_A}_{\substack{\leftarrow D_A \\ \text{Dimension}}} \otimes \underbrace{\mathcal{H}_B}_{D_B}$$

Expand $|\Psi\rangle$ on tensor product basis:

$$|\Psi\rangle = \sum_{n,n'} c_{nn'} \underbrace{|n\rangle_A}_{\substack{\downarrow \\ \in \mathbb{C}}} \otimes \underbrace{|n'\rangle_B}_{\substack{\text{basis of} \\ \mathcal{H}_A}} \underbrace{\quad}_{\substack{\text{basis of} \\ \mathcal{H}_B}}$$

$$n = 1 \dots D_A$$

$$n' = 1 \dots D_B \quad c_{nn'} \text{ is a } D_A \times D_B \text{ matrix}$$

From linear algebra singular value decomposition:

$$C = U \sum V^+ \quad , \quad V^+ = V^T \star$$

\downarrow
 $D_A \times D_A$
 unitary matrix
 $D_B \times D_B$
 unitary matrix
 $D_A \times D_B$ matrix

contains the singular values of C

Structure of Matrix Σ :

$$\rightarrow D_A = D_B$$

$$\Sigma = \begin{pmatrix} \sigma_1 & & & 0 \\ & \ddots & & \\ 0 & & \ddots & \\ & & & \sigma_D \end{pmatrix}; \sigma_i \geq 0$$

$$\sigma_1, \dots, \sigma_k > 0$$

$$\sigma_{k+1}, \dots, \sigma_D = 0$$

it would be that $k=D \rightarrow$ no zero singular values

$$\rightarrow D_A > D_B$$

$$\Sigma = \begin{bmatrix} \sigma_1 & & & 0 \\ & \ddots & & \\ 0 & & \sigma_{D_B} & \\ & & & 0 \end{bmatrix}$$

$$\sigma_1, \dots, \sigma_k > 0 ; \sigma_{k+1}, \dots, \sigma_{D_B} = 0$$

(it would be that $k=D_B$)

$$\rightarrow D_B > D_A$$

$$\Sigma = \begin{bmatrix} \sigma_1 & & & 0 \\ & \ddots & & \\ 0 & & \sigma_k & \\ & & & 0 \end{bmatrix}$$

$$\sigma_1, \dots, \sigma_k > 0$$

$$\sigma_{k+1}, \dots, \sigma_{D_A} = 0$$

$$|\Psi\rangle = \sum_{n,n'} c_{n,n'} |n\rangle_A \otimes |n'\rangle_B^?$$

$$= \sum_{n,n'=1}^{D_A, D_B} \sum_{i=1}^k U_{n_i} \sigma_i (V^+)_{i,n'} |n\rangle_A \otimes |n'\rangle_B$$

$$= \sum_{i=1}^k \sigma_i \sum_{n,n'=1}^{D_A, D_B} U_{n_i} (V^+)_{i,n'} |n\rangle_A \otimes |n'\rangle_B$$

$$= \sum_{i=1}^k \sigma_i \left\{ \sum_{n=1}^{D_A} U_{n_i} |n\rangle_A \right\} \otimes \left\{ \sum_{n'=1}^{D_B} (V^+)_{i,n'} |n'\rangle_B \right\}$$

$|v_i\rangle_A$
 $\sqrt{v_i}$
 $|v_i\rangle_B$

$$|\Psi\rangle = \sum_{i=1}^k \sigma_i |\varphi_i\rangle_A \otimes |\psi_i\rangle_B$$

\downarrow
 $\sigma_i = \sqrt{\lambda_i}$

This explains many examples of bipartite systems

where we noticed $S(\rho_A) = S(\rho_B)$ for

$$|\Psi\rangle = |B_{00}\rangle = \frac{1}{\sqrt{2}} (|0\rangle_A \otimes |0\rangle_B + |1\rangle_A \otimes |1\rangle_B)$$

$$\begin{aligned} \rightarrow \rho_A &= \frac{1}{2} (|0\rangle_A \langle 0|_A + \\ &\quad |1\rangle_A \langle 1|_A) \\ \hookrightarrow \rho_A &= \rho_B \end{aligned}$$

$$|W\rangle = \frac{1}{\sqrt{3}} (|001\rangle + |010\rangle + |100\rangle)$$

$$\rho_A = \text{Tr}_{BC} |W\rangle \langle W|$$

$$\rho_{BC} = \text{Tr}_A |W\rangle \langle W|$$

We computed eigenvalues of ρ_A & ρ_{BC} $\lambda_i = \frac{1}{3}, \frac{2}{3}$

$$S(\rho_A) = S(\rho_{BC}) = +\frac{1}{3} \log 3 + \frac{2}{3} \log \frac{3}{2}$$

$$|W\rangle = \sqrt{\frac{1}{3}} |\Psi_0\rangle_A \otimes |\Psi_0\rangle_{BC} + \sqrt{\frac{2}{3}} |\Psi_1\rangle_A \otimes |\Psi_1\rangle_{BC}$$

↓
Schmidt Decomposition

2. Properties of von Neumann Entropy

Recall that $\rho \rightarrow S(\rho) = -\text{Tr}(\rho \log \rho)$

Spectral Decomposition

$$\rho^+ = \rho ; \rho \geq 0 ; \text{Tr}(\rho) = 1$$

. Eigenvalue $\lambda_i \in \mathbb{R}, \lambda_i \geq 0, \sum_i \lambda_i = 1$

$$. S(\rho) = - \sum_i \lambda_i \log \lambda_i$$

= Shannon Entropy for prob. distribution defined by the set of $\{\lambda_i\}$.

Following Properties:

a) $0 \leq S(\rho) \leq \log (\dim \mathcal{H})$

b) Concavity $S(t\rho + (1-t)\sigma) \leq tS(\rho) + (1-t)S(\sigma)$
 $0 \leq t \leq 1$

c) Subadditivity $S(\rho_{AB}) \leq S(\rho_A) + S(\rho_B)$

d) Araki-Lieb inequality $|S(\rho_A) - S(\rho_B)| \leq S(\rho_{AB})$

e) Strong subadditivity

a) $0 \leq S(g) \leq \log(\dim J_t)$

↓
obvious

$S(g) = -\sum_{i=1}^D \lambda_i \log \lambda_i$

$D = \dim J_t$

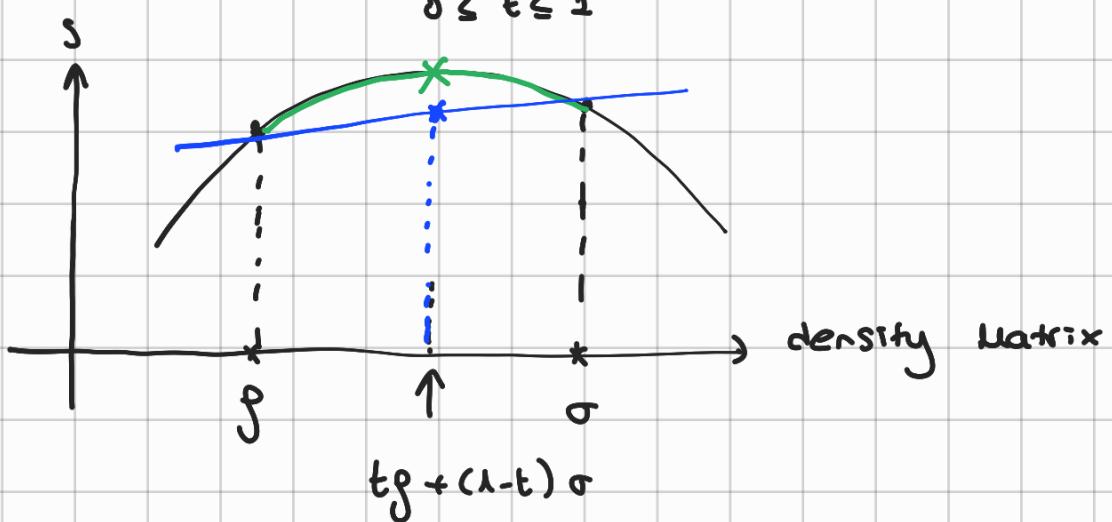
This is maximized for $\lambda_i = \frac{1}{D}$

$S_{\max} = -\sum_{i=1}^D \frac{1}{D} \log \frac{1}{D} = \log D$

b) Concavity:

$$S(tg + (1-t)\sigma) \geq tS(g) + (1-t)S(\sigma)$$

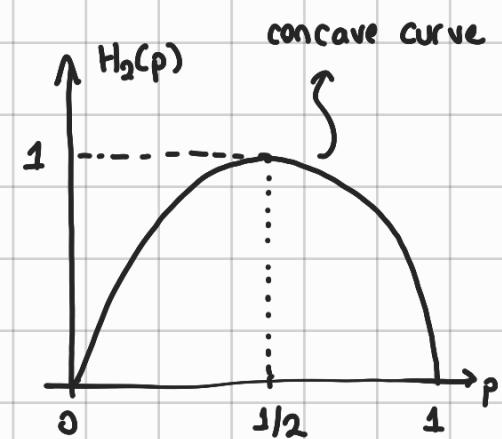
$0 \leq t \leq 1$



The same kind of concavity for the binary

Shannon entropy of $\text{Ber}(p)$ variable

$$H_2(p) = -p \log_2 p - (1-p) \log_2 (1-p)$$



c) Subadditivity $S(\rho_{AB}) \leq S(\rho_A) + S(\rho_B)$

(follows from concavity)

d) Araki - Lieb Inequality

$$S(\rho_{AB}) \geq |S(\rho_A) - S(\rho_B)|$$

Example:

$$\rho_{AB} = \underbrace{|B_00\rangle\langle B_00|}_{\text{pure State Vector}} ;$$

$$\rho_A = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \quad S(A) = \log 2$$

$$\rho_B = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \quad S(B) = \log 2$$

$$S(\rho_{AB}) = 0$$

$$0 \geq |\log 2 - \log 2|$$

Also Remark: $S(\rho_{AB}) \leq S(A)$ or $S(B)$ (not always true for all quantum system)

In classical info. theory, it is NEVER TRUE

Hall mark of entanglement in AB system

Classically $\rightarrow H(X,Y) \geq H(X) \text{ or } H(Y)$

Proof of Araki-Lieb Inequality:

Technique of purification:

→ Take ρ_{AB} density matrix, I construct a bigger state on bigger Hilbert Space $\mathcal{H}_{AB} \otimes \mathcal{H}_R$

$|\Psi\rangle_{ABR}$ which is pure s.t :

$$\rho_{AB} = \text{Tr}_R |\Psi\rangle_{ABR} \langle \Psi|_{ABR}$$

→ Why does $|\Psi\rangle_{ABR}$ exist? (Not unique)

take spectral decomposition of $\rho_{AB} = \sum_i \mu_i |x_i\rangle \langle x_i|$

$$|\Psi\rangle_{ABR} = \sum_i \sqrt{\mu_i} |x_i\rangle_{AB} \otimes |x_i\rangle_R$$

$$\text{easy to check } \rho_{AB} = \text{Tr}_R |\Psi\rangle_{ABR} \langle \Psi|_{ABR}$$

→ Use Schmidt Theorem on the purified state & Subadditivity to prove Araki-Lieb inequality

$$S(ABR) = 0 \quad \text{since} \quad |\psi\rangle_{ABR}\langle\psi|_{ABR}$$

is rank-one Density Matrix
(pure state)

$$\begin{array}{ccc} S(A) = S(BR) & S(B) = S(AR) & S(R) = S(AB) \\ \downarrow & \downarrow & \downarrow \\ \text{Schmidt} & \text{Schmidt} & \text{Schmidt} \end{array}$$

Suppose without loss of generality $S(A) \geq S(B)$

We want to prove $S(AB) \geq S(A) - S(B)$

$$\Leftrightarrow S(A) \leq S(AB) + S(B)$$

$$\Leftrightarrow \underline{S(BR) = S(R) + S(B)}$$

\downarrow
sub additivity