

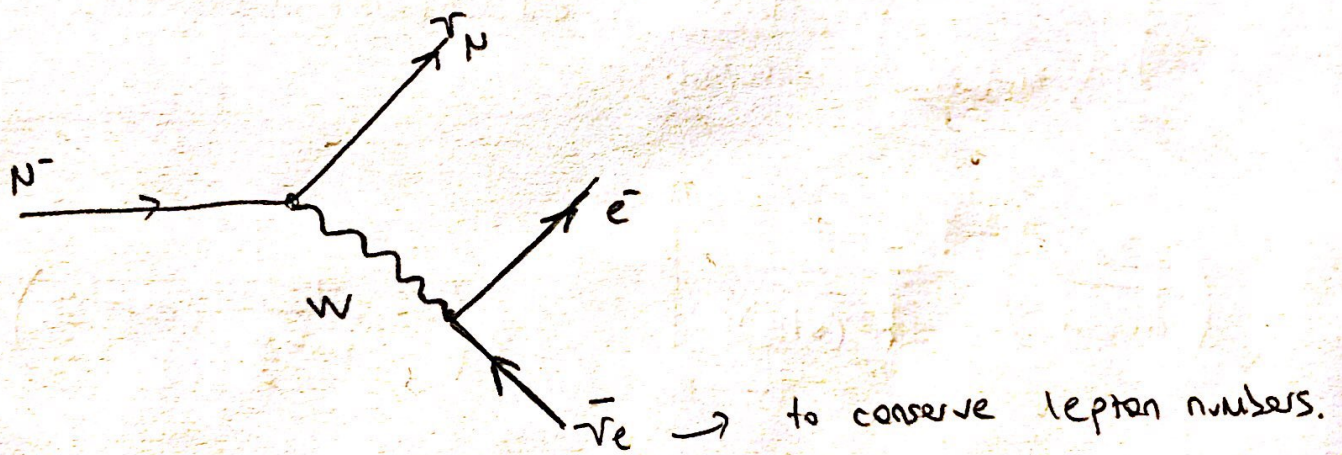
### Elementary Particle Decay

\* All decays of elementary matter particles go via weak charged current ( $W^\pm$ ): (changes flavour)

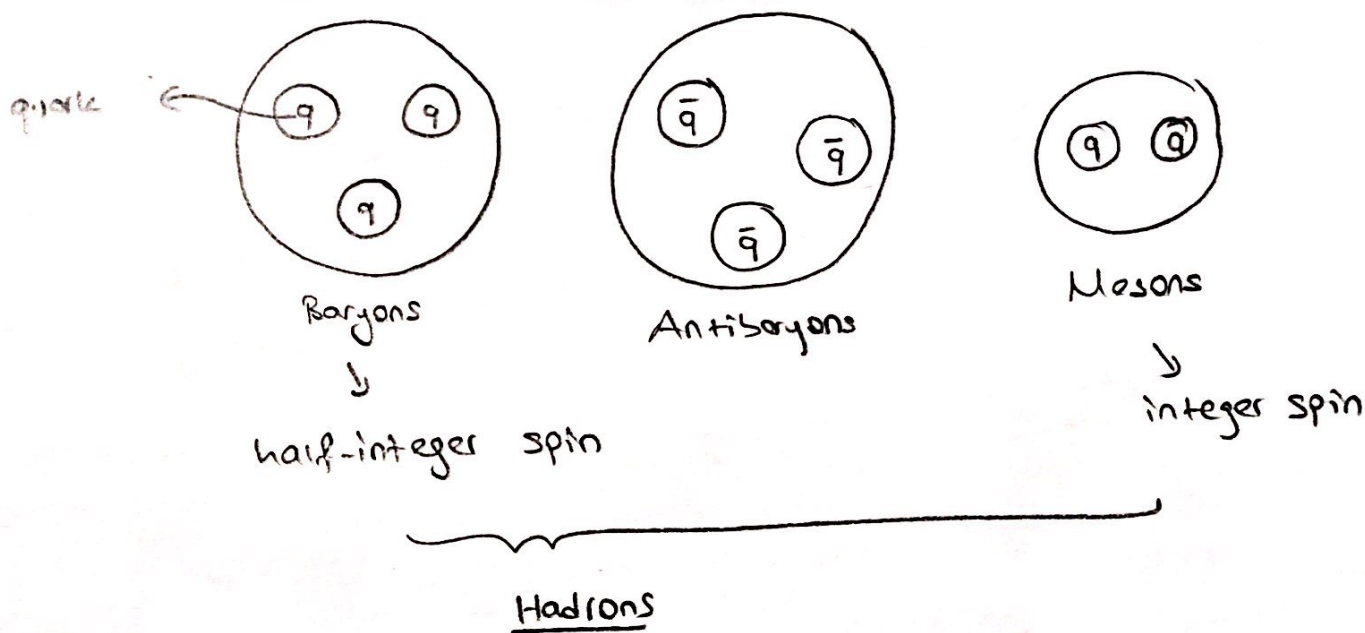
\* For a decay to occur, there need to be particles with lower mass.

\* Electron is stable since it is the lightest charged particles. Neutrinos are also stable.

$$\mu^- \rightarrow e^- + \bar{\nu}_e + \nu_\mu$$



# Non elementary Particles



Proton =  $|p\rangle = |uud\rangle \rightarrow$  Stable

\* Baryon conservation in SM

Neutron =  $|n\rangle = |udd\rangle \rightarrow$  Decay down to proton

$$n \rightarrow p e^- \bar{\nu}_e$$

Pions  $\pi^\pm$  :

$$\begin{aligned} \rightarrow |\pi^+\rangle &= |u\bar{d}\rangle \\ \rightarrow |\pi^-\rangle &= |\bar{u}d\rangle \\ \rightarrow |\pi^0\rangle &= \frac{1}{\sqrt{2}} |u\bar{u} - d\bar{d}\rangle \end{aligned}$$

## Particle Lifetime

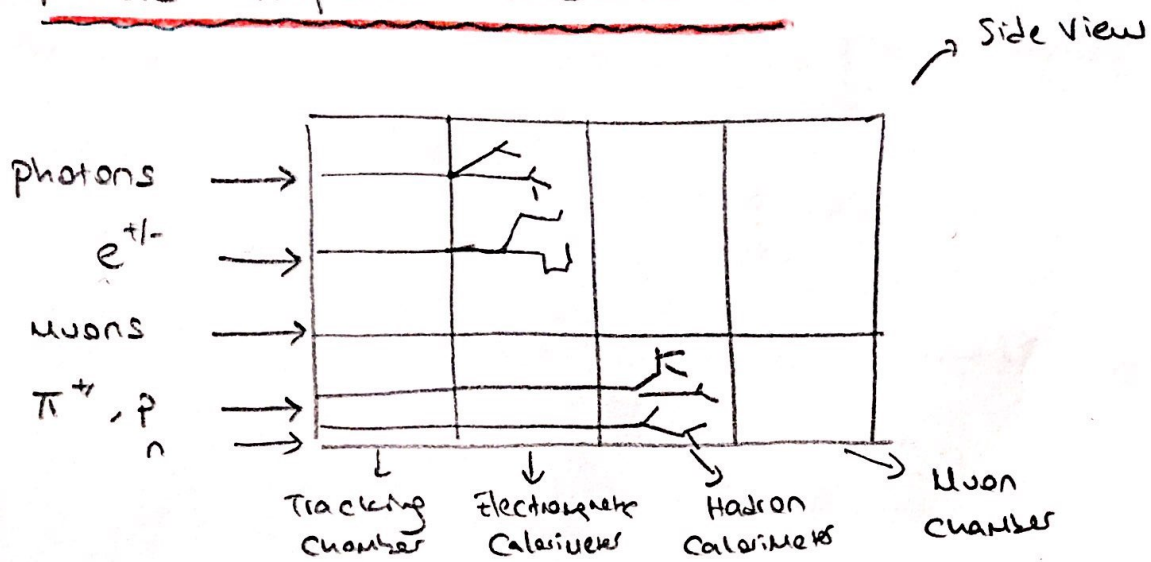
Weak  $\rightarrow$  large lifetimes, long-lived

Electromagnetically  $\rightarrow$  Intermediate lifetime

Strong  $\rightarrow$  Short-lived



# Particle Properties Measurement



- Charge
- Energy & momentum
- Particle Identification (PID)
- Point where the particle was produced or decayed.
  - (production) vertex
  - (decay) vertex
- what was the mother particle

Having these properties, try to reconstruct an initial process in a collision.

## Reconstruction with kinematics

$$M^2 = \left( \sum E_i \right)^2 - \left| \sum \vec{p}_i \right|^2$$

→ Usually it is not possible to measure both energies and momenta.

→ Measure a combination: energy + mass or momentum + mass (3)

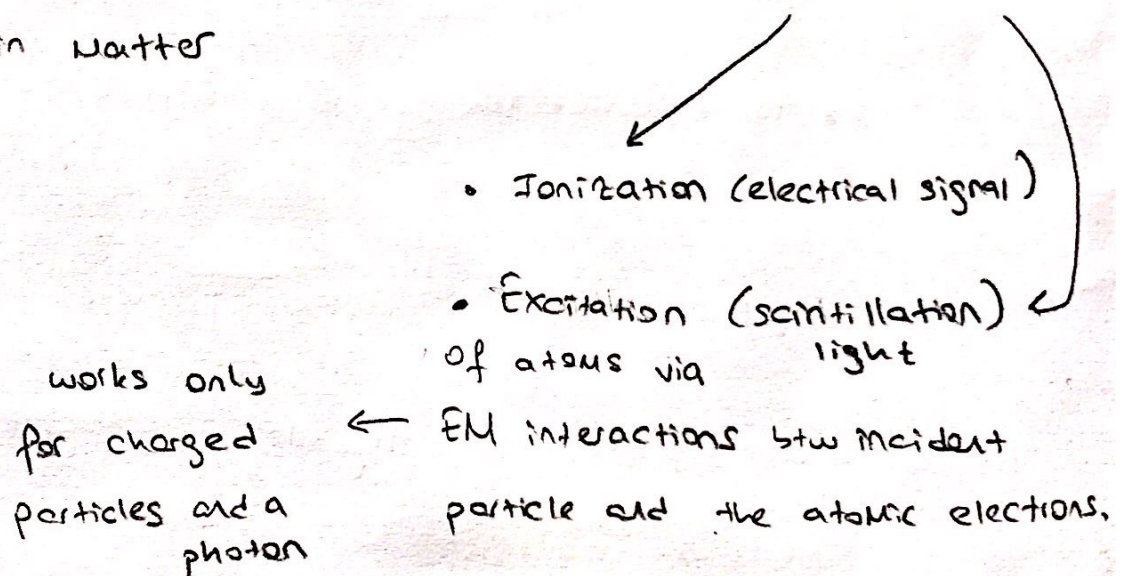
Decay Length  $L = \beta \gamma c T$  where  $T$  is particle lifetime.

$$\beta \gamma = \frac{p}{m}$$

Many different types of particle detectors:

- Scintillators
- Bubble or Cloud Chamber
- Wire chambers ...

\* They all rely on the detection of a perturbation induced in matter



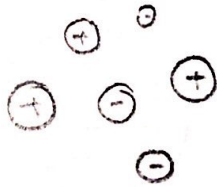
Sensitive volumes in which the perturbations occur

- gas
- liquid
- solid



# Ionization

Bethe-Bloch equation:



$$\frac{dE}{dx} \approx -4\pi k^2 c^2 \alpha^2 \frac{n z^2}{u^2 \beta^2} \left\{ \ln \left[ \frac{2\beta^2 \gamma^2 m_e}{I_e} \right] - \beta^2 \right\}$$

$v = \beta c$  is a speed of the travelling particle

$z$ : atomic number of the material

Material density

$$n = \frac{\rho}{A m_u}$$

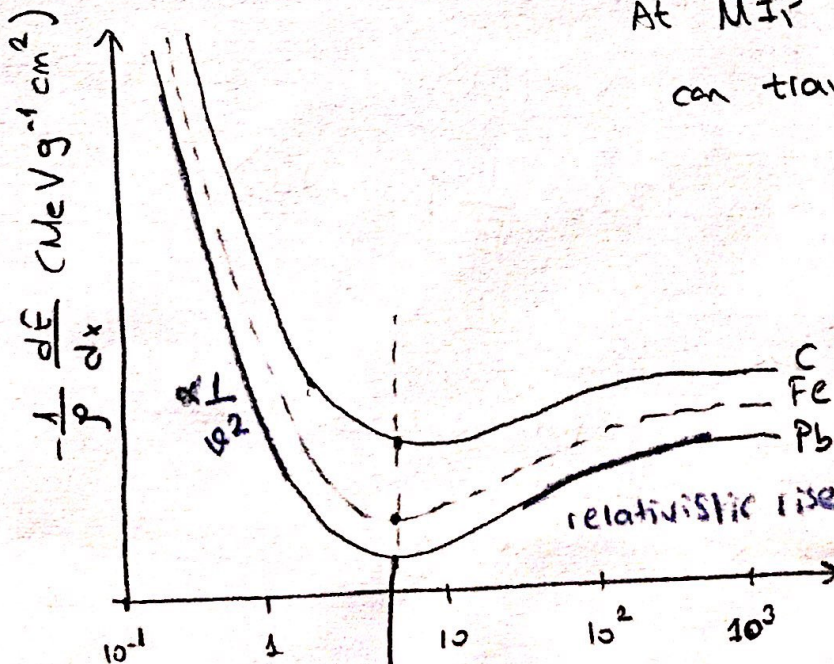
atomic mass number  $\rightarrow 1.66 \times 10^{-24}$   
unified atomic mass unit

$n$  its number density

$I_e \sim 102 \text{ eV} \rightarrow$  effective ionization potential of the material

$$\frac{1}{\rho} \frac{dE}{dx} \propto \frac{z^2}{A} \approx \text{const}$$

At MIP level, a particle can travel a long distance.



MIP (Minimum ionizing particles  $\beta\gamma \approx 3$ ) ⑤

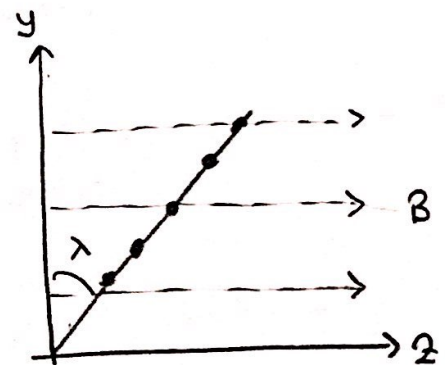
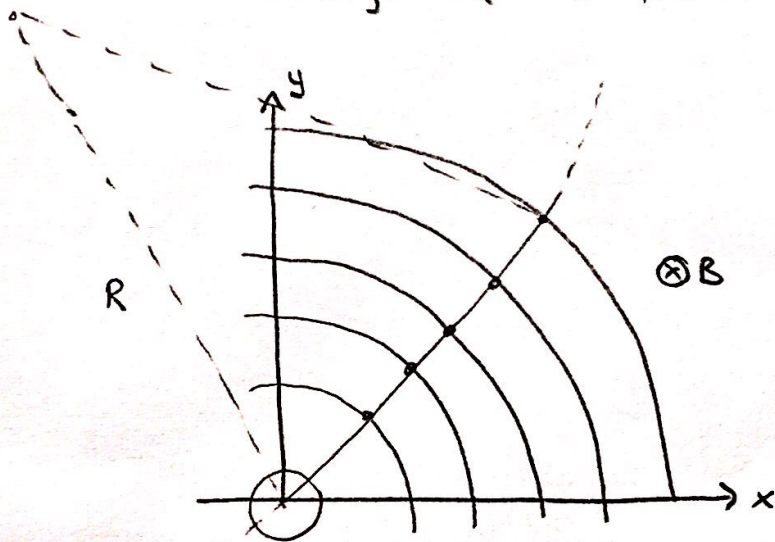
## Momentum Measurement

$$F = \frac{mv^2}{R} = q(\vec{v} \times \vec{B})$$

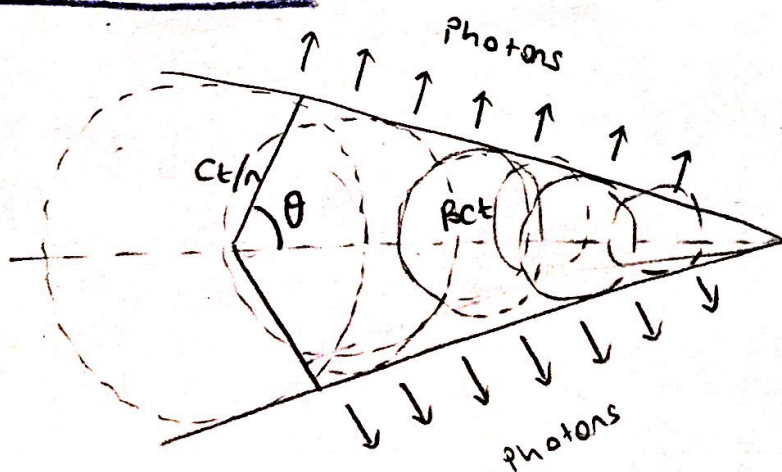
$$p \cos \lambda = qBR$$

$$p \cos \lambda [\text{GeV}/c] = 0.3 BR [\text{Tm}]$$

Obtain  $p$  and  $\cos \lambda$  from constructed particle trajectory and knowledge of B-field.



## Cerenkov Radiation



Occurs when  $v > c/n$

$$\cos \theta = \frac{1}{n\beta}$$

Light is emitted if:  
 $\beta > 1/n$

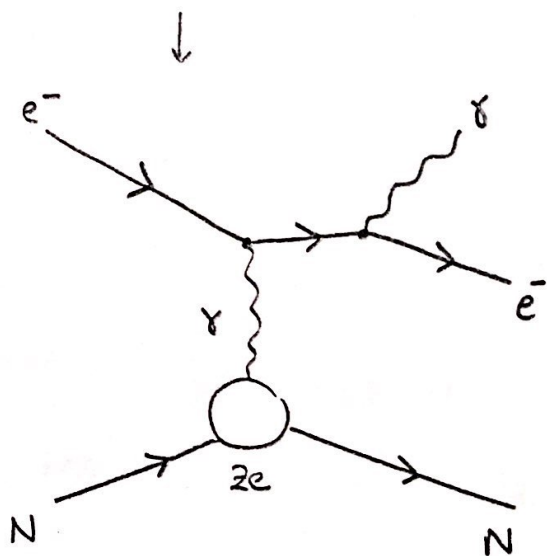
$$\beta = \frac{pc}{E} = \frac{p}{\sqrt{p^2 + m^2 c^2}}$$

$$mc < p \sqrt{n^2 - 1}$$

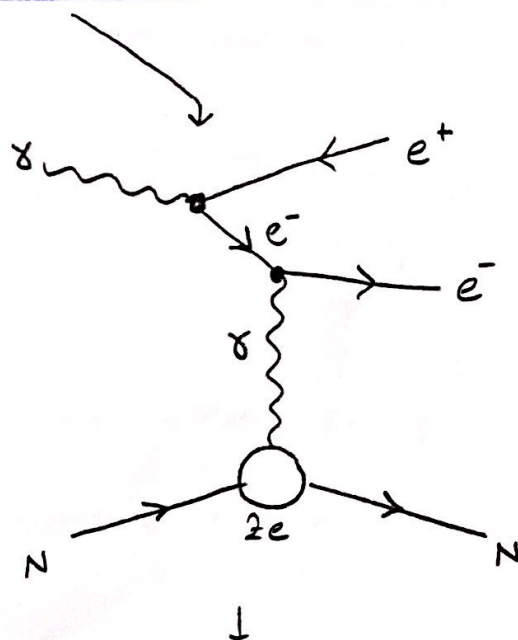
⑥



# Bremsstrahlung and $e^+e^-$ production



↳ Main energy loss mechanism for  $e^\pm$  with  $>E_c$  (800/2 MeV)



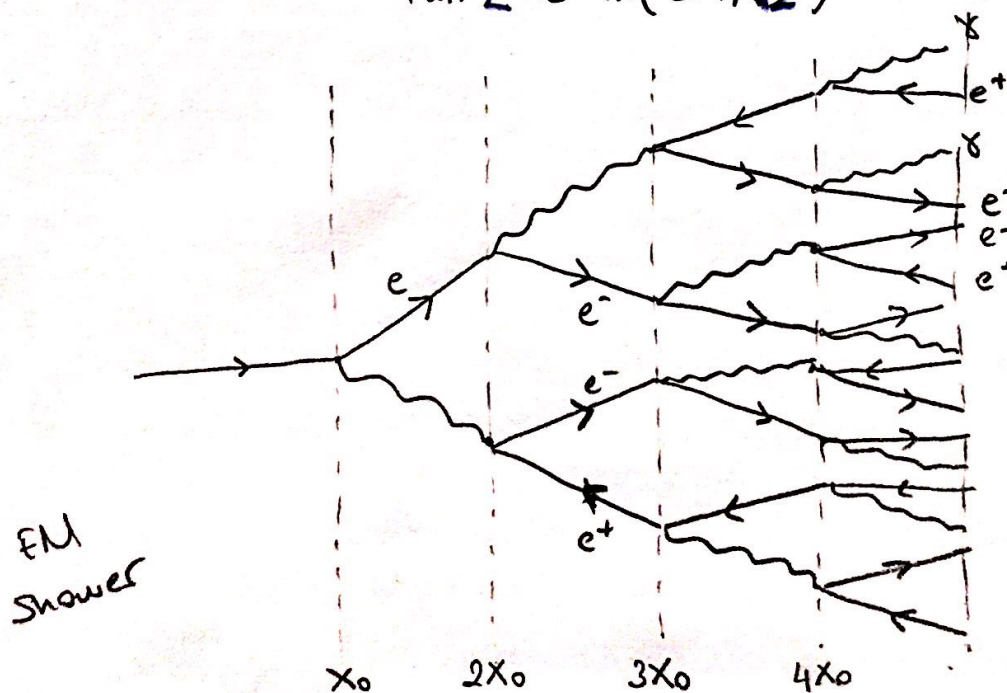
↳ Dominant energy loss mech for  $\gamma$  with  $E_\gamma > 10$  MeV

## Radiation Length ( $X_0$ )

Avg. distance over which the energy of an electron is reduced by bremsstrahlung by a factor of  $1/e$ .

$$X_0 \approx \frac{1}{4\alpha n Z^2 r_e^2 \ln(187/\sqrt{Z})}$$

$$r_e = \frac{e^2}{4\pi\epsilon_0 m c^2} = 2.8 \times 10^{-15} \text{ m}$$



EM Shower

\* Number of particles doubles after every  $X_0$ .

After  $x X_0$  avg. energy

$$\langle E \rangle \approx \frac{\bar{E}}{2^x}$$

Shower development stops after  $\langle E \rangle < \bar{E}_c$

$$X_{\text{max}} = \frac{\ln(E/\bar{E}_c)}{\ln 2}$$

(7)