

Scattering Amplitudes as Differential Forms

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$\mathcal{N} = 4$ Supersymmetric Yang-Mills theory

- AdS/CFT Duality
- Integrability
- Twistor strings
- Amplituhderon/Grassmaniann
- Lots of Math

The content of $\mathcal{N} = 4$ SYM:

- 2 gauge bosons with $h = \pm 1$: $|a\rangle^{+1}, |a\rangle_{IJKL}^{-1}$
- 8 fermions with $h = \pm \frac{1}{2}$: $|a\rangle_I^{+1/2}, |a\rangle_{IJK}^{-1/2}$
- 6 spinless scalars: $|a\rangle_{IJ}^0$

They are all related by supersymmetric generators $Q_{I=1,2,3,4}^\alpha$ and

$\tilde{Q}_{I=1,2,3,4}^{\dot{\alpha}}$. For example, $|a\rangle_I^{+1/2} = \tilde{Q}_I |a\rangle^{+1}$

Thus, all on-shell states in $\mathcal{N} = 4$ SYM can be grouped into a *single* supermultiplet:

$$\begin{aligned} |a\rangle &\equiv \exp(\tilde{Q}_I \cdot \tilde{\lambda} \cdot \tilde{\eta}^I) |a\rangle^{+1} \\ &= |a\rangle^{+1} + \tilde{\eta}^I |a\rangle_I^{1/2} + \cdots + \frac{1}{4!} \tilde{\eta}^I \tilde{\eta}^J \tilde{\eta}^K \tilde{\eta}^L |a\rangle_{IJKL}^{-1} \end{aligned}$$

Superamplitudes, not amplitudes!

$$\mathcal{A}_{n,\text{tree}}^{k=2} = \frac{\delta^4(P)\delta^8(Q)}{\langle 12 \rangle \cdots \langle n1 \rangle}$$

$$\mathcal{A}_{n,\text{tree}}^{k=3} = \mathcal{A}_{n,\text{tree}}^{k=2} \sum_{2 < i < j < n-1} [1, i-1, i, j, j+1] \text{ (Momentum Twistor come in!)}$$

$\mathcal{A}_{n,\text{tree}}^k$ are known for any n and any k

$\mathcal{A}_{n,L=1}^k$ are known for any n and any k

$\mathcal{A}_{n,L=2}^k$ are known for any n and $k=2$

...

$\mathcal{A}_{6,5}^2$ are known.

[Bern, Caron-Huot, Dixon, Drummond, Duhr, Foster, Gürdogan, He, Henn, von Hippel, Golden, Kosower, Papathanasiou, Pennington, Roiban, Smirnov, Spradlin, Vergu, Volovich, ...]

Amplituhderon & Positive geometry

Positive geometry is defined by its boundary which also is a *positive geometry*. For each positive geometry, we define a canonical form such that it only has logarithmic residues on its boundary. For example:

$$\Omega[\{0 < x < 1\}] = \frac{dx}{x} - \frac{dx}{x-1} = d \log \frac{x}{1-x}$$

$$\Omega[\{0 < x < y < 1\}] = d \log \frac{x}{x-y} \wedge d \log \frac{x-y}{1-y}$$

$$\Omega[\{x_{ij}z_{ij} + y_{ij}w_{ij} < 0 | 1 \leq i < j \leq L\}] = \text{4pt L-loop planar integrand!}$$

Amplituhderon is a positive geometry defined in Momentum twistor space.

[Arkani-Hamed, Bai, Thomas, Trnka, ...]

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- The result of this simple replacement is
every amplitude with n -point and L -loop turn out to be $(2n - 4 + 4L)$ - $d \log$ form

MHV case

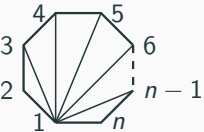
some simple examples:

$$\Omega_{3,\text{MHV}} = d \log \frac{\langle 12 \rangle}{\langle 23 \rangle} \wedge d \log \frac{\langle 23 \rangle}{\langle 13 \rangle}$$

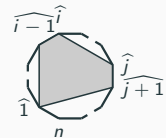
$$\Omega_{4,\text{MHV}} = d \log \frac{\langle 12 \rangle}{\langle 23 \rangle} \wedge d \log \frac{\langle 23 \rangle}{\langle 13 \rangle} \wedge d \log \frac{\langle 13 \rangle}{\langle 34 \rangle} \wedge d \log \frac{\langle 34 \rangle}{\langle 14 \rangle}$$

The result for MHV n -pt corresponds to a triangulation of n -gon :

$$\Omega_{n,\text{MHV}} = \bigwedge_{i=1}^{n-2} d \log \frac{\langle a_i b_i \rangle}{\langle b_i c_i \rangle} \wedge d \log \frac{\langle b_i c_i \rangle}{\langle a_i b_i \rangle}$$



$$= \bigwedge_{i=2}^{n-1} d \log \frac{\langle 1 i \rangle}{\langle i+1 1 \rangle} \wedge d \log \frac{\langle i i+1 \rangle}{\langle i+1 i+1 \rangle} .$$

$$\sum_{1 < i < j < n} \text{Diagram} = \sum_{1 < i < j < n} \Omega_{\text{MHV}}(1, \dots, i-1) \wedge \Omega_{\text{MHV}}(i, \dots, j) \\ \wedge \Omega_{\text{MHV}}(j+1, \dots, n, 1) \wedge \Omega_{\overline{\text{MHV}}}(\widehat{1}, \widehat{i-1}, \widehat{i}, \widehat{j}, \widehat{j+1}),$$


where

$$\begin{aligned} \tilde{\lambda}_{\widehat{1}} &= \tilde{\lambda}_1 + \sum_{a=2}^{i-1} \frac{\langle i a \rangle}{\langle i 1 \rangle} \tilde{\lambda}_a + \sum_{a=j+2}^n \frac{\langle a j+1 \rangle}{\langle 1 j+1 \rangle} \tilde{\lambda}_a, \\ \tilde{\lambda}_{\widehat{i}} &= \tilde{\lambda}_i + \sum_{a=2}^{i-1} \frac{\langle a 1 \rangle}{\langle i 1 \rangle} \tilde{\lambda}_a, \quad \tilde{\lambda}_{\widehat{i+1}} = \tilde{\lambda}_{i+1} + \sum_{a=i+2}^{j-1} \frac{\langle j a \rangle}{\langle j i+1 \rangle} \tilde{\lambda}_a, \\ \tilde{\lambda}_{\widehat{j}} &= \tilde{\lambda}_j + \sum_{a=m+1}^{n-1} \frac{\langle a i+1 \rangle}{\langle j i+1 \rangle} \tilde{\lambda}_a, \quad \tilde{\lambda}_{\widehat{j+1}} = \tilde{\lambda}_{j+1} + \sum_{a=m+1}^{n-1} \frac{\langle 1 a \rangle}{\langle 1 j+1 \rangle} \tilde{\lambda}_a. \end{aligned}$$

There is an algorithm to obtain $\Omega_{n,k}^{\text{tree}}$

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- Recently, Cachazo *et al.* find the canonical form for MHV case can be written as

$$\Omega_{n,\text{MHV}} = \frac{1}{(n-2)!} \left(\sum_{i=2}^{n-1} d \log \frac{\langle 1 i \rangle}{\langle i+1 1 \rangle} \wedge d \log \frac{\langle i i+1 \rangle}{\langle i+1 1 \rangle} \right)^{n-2}$$

$$\sim \frac{1}{(n-2)!} (\text{canonical form for } n\text{-gon})^{n-2}$$

- The connection with amplituhedron $4k$ -form.
- How to intrinsically define amplituhedron in momentum space.
- Generalization to loop level and more general theories.