

# DETERMINING OPTIMAL TEST FUNCTIONS FOR 2-LEVEL DENSITIES

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## Abstract

It is of great interest to study the order of vanishing at the central point of  $L$ -functions, such as the famous Riemann zeta function and its generalizations. For example, the Birch and Swinnerton-Dyer conjecture relates the order of vanishing for an elliptic curve  $L$ -function to the rank of its group of rational solutions. Katz and Sarnak conjectured a correspondence between the  $n$ -level density statistics of zeros from families of  $L$ -functions (which essentially involves  $n$ -tuples of zeros) with eigenvalues from random matrix ensembles, and in many cases that sums of smooth test functions, whose Fourier transforms are finitely supported over scaled zeros in a family, converge to an integral of the test function against a density  $W_{n,G}$  depending on the symmetry  $G$  of the family (unitary, symplectic or orthogonal). This integral bounds the average order of vanishing at the central point of the corresponding family of  $L$ -functions.

We can obtain better estimates on this vanishing in two ways. The first is to do more number theory, and prove results for larger  $n$  and greater support; the second is to do functional analysis and obtain better test functions to minimize the resulting integrals. We pursue the latter here when  $n = 2$ , minimizing

$$\frac{1}{\Phi(0,0)} \int_{\mathbb{R}^2} W_{2,G}(x,y) \Phi(x,y) dx dy$$

over test functions  $\Phi : \mathbb{R}^2 \rightarrow [0, \infty)$  with compactly supported Fourier transform. We study a restricted version of this optimization problem, imposing that our test functions take the form  $\phi(x)\psi(y)$  for some fixed admissible  $\psi(y)$  and  $\text{supp } \hat{\phi} \subseteq [-1, 1]$ , thereby reducing the problem to one analogous to the 1-level density case in optimizing over one-variable test functions  $\phi(x)$ . Continuing with the analogy, Devlin and Miller extended the functional analytic arguments of Iwaniec, Luo and Sarnak, converting the restricted optimization problem to finding the unique  $g \in L^2[-1/2, 1/2]$  satisfying a Fredholm integral equation of the second kind. From here we take two approaches. First, showing  $g$  satisfies a homogeneous linear system of differential equations, a method introduced by Freeman and Miller, from which we can derive a closed form expression for  $g$ . Second, iterating to obtain a series representation of  $g$  and truncating to compute explicit estimates on the minimum value. We conclude by discussing improvements to previous bounds.

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# 1 Preparatory material

Denote by  $G$  any one of the classical compact matrix groups, namely the orthogonal group  $O$ , the even special orthogonal group  $SO(\text{even})$ , the odd special orthogonal group  $SO(\text{odd})$ , and the symplectic group  $Sp$ . The 1-level densities of each of these groups are known to be

$$\begin{aligned} W_O(x) &= 1 + \frac{1}{2}\delta_0(x), \\ W_{SO(\text{even})}(x) &= 1 + \frac{\sin 2\pi x}{2\pi x}, \\ W_{SO(\text{odd})}(x) &= 1 - \frac{\sin 2\pi x}{2\pi x} + \delta_0(x), \\ W_{Sp}(x) &= 1 - \frac{\sin 2\pi x}{2\pi x}. \end{aligned}$$

Here  $\delta_0(x)$  is the Dirac delta centered at zero. The corresponding Fourier transforms take the form

$$\widehat{W}_G(\xi) = \delta_0(\xi) + m_G(\xi)$$

where

$$\begin{aligned} m_O(\xi) &= \frac{1}{2}, \\ m_{SO(\text{even})}(\xi) &= \frac{1}{2}\mathbb{1}_{[-1,1]}(\xi), \\ m_{SO(\text{odd})}(\xi) &= 1 - \frac{1}{2}\mathbb{1}_{[-1,1]}(\xi), \\ m_{Sp}(\xi) &= -\frac{1}{2}\mathbb{1}_{[-1,1]}(\xi). \end{aligned}$$

Define

$$\mathfrak{I}_G(\sigma) = \inf_{\phi} \frac{1}{\phi(0)} \int_{\mathbb{R}} \phi(x) W_G(x) dx$$

where the infimum is taken over non-negative Schwarz functions  $\phi : \mathbb{R} \rightarrow [0, \infty)$  with compactly supported Fourier transform satisfying  $\text{supp}(\widehat{\phi}) \subseteq [-2\sigma, 2\sigma]$ . We refer to such  $\phi$  as TEST FUNCTIONS. We are interested in studying the following:

- a. Finding test functions witnessing the infimum  $\mathfrak{I}_G(\sigma)$  for fixed  $\sigma$  and group  $G$ , i.e. finding  $\phi$  satisfying  $\text{supp}(\widehat{\phi}) \subseteq [-2\sigma, 2\sigma]$  and

$$\mathfrak{I}_G(\sigma) = \frac{1}{\phi(0)} \int_{\mathbb{R}} \phi(x) W_G(x) dx$$

- b. Computing the value of  $\mathfrak{I}_G(\sigma)$ .
- c. Analyzing the behavior of  $\mathfrak{I}_G$  as a real function on  $(0, \infty)$ . In particular, Freeman conjectured that  $\mathfrak{I}_G$  is continuous and real analytic except at integers and half-integers.

Fix  $\sigma > 0$  and one of our aforementioned compact groups  $G$ ; for brevity we suppress the subscripts  $W := W_G$  and  $m := m_G$ . Moreover, whenever we define a new object in terms of  $W$  and  $m$ , there is an implicit subscript  $G$  to denote the dependence on the group.

In the literature, no one tries to directly find the optimal test function  $\phi$ , instead appealing to functional analytic arguments. For example, it follows from a theorem of Ahiezer and the Paley-Wiener

theorem that the optimal test function admits the form  $\phi(z) = |h(z)|^2$ , where  $h : \mathbb{C} \rightarrow \mathbb{C}$  is an entire function of exponential type 1 and  $h \in L^2(\mathbb{R})$ . That is, its Fourier transform admits the form

$$\widehat{\phi}(\xi) = (g * g^\smile)(\xi)$$

where  $\text{supp } g \subseteq [-\sigma, \sigma]$  and  $g \in L^2[-\sigma, \sigma]$  and  $g^\smile(\xi) = \overline{g(-\xi)}$ . Notice that the Fourier transforms of the distributions  $W$  are significantly easier to work with as they are step functions, so naturally we appeal to the Plancharel theorem to write

$$\frac{1}{\phi(0)} \int_{\mathbb{R}} \phi(x) W(x) dx = \frac{1}{\int_{\mathbb{R}} \widehat{\phi}(\xi) d\xi} \int_{\mathbb{R}} \widehat{\phi}(\xi) \widehat{W}(\xi) d\xi.$$

Rewriting  $\widehat{\phi}$  in terms of  $g \in L^2[-\sigma, \sigma]$  converts the problem to an equivalent optimization over the Hilbert space  $L^2$  (for details of derivation, cf. Freeman Proposition 3.2), namely

$$\mathfrak{I}_G(\sigma) = \inf_{g \in L^2[-\sigma, \sigma]} R(g),$$

where we define a quadratic form  $R : L^2[-\sigma, \sigma] \rightarrow \mathbb{R}$  and a self-adjoint operator  $K : L^2[-\sigma, \sigma] \rightarrow L^2[-\sigma, \sigma]$  by

$$R(g) = \frac{\langle (I + K)g, g \rangle}{|\langle g, 1 \rangle|^2}, \quad (Kg)(x) = \int_{-\sigma}^{\sigma} m(x - y)g(y)dy,$$

where  $I$  is the identity operator on  $L^2[-\sigma, \sigma]$ . By functional analysis hocus pocus, the constant function one is in the image of  $I + K$ . Moreover, there exists  $f \in (\ker(I + K))^\perp$  satisfying the equation

$$(I + K)f \equiv 1.$$

This implies  $A := \langle (I + K)f, f \rangle = \langle 1, f \rangle$  is a positive constant. ILS proved that in fact

$$\mathfrak{I}_G(\sigma) = \inf_{g \in L^2[-\sigma, \sigma]} R(g) = R(f) = \frac{1}{A}.$$

ILS showed that if  $\sigma = 1$ , then the optimal functions  $f \in L^2[-1, 1]$  are even, and are trigonometric polynomials when restricted to  $[0, 1]$ . Freeman extended this result to arbitrary  $\sigma > 0$ . The brute force approach to this problem is to find a trigonometric polynomial  $f : [0, \sigma] \rightarrow \mathbb{R}$ , i.e. taking the form

$$f(x) = \sum_{n \geq 0} a_n \cos(nx) + b_n \sin(nx)$$

where  $a_n, b_n = 0$  for all  $n \geq N$ , such that  $(I + K)f$  is a constant. Normalizing by this constant gives the desired equation  $(I + K)f \equiv 1$ . Since  $f$  is even, we can extend it to a function on  $[-\sigma, \sigma]$  to recover the desired optimal function.

## 2 One level

The key observation in deriving the optimal  $f \in L^2[-\sigma, \sigma]$  is, as remarked earlier, it suffices to find  $f$  such that  $(I + K)f$  is constant. Differentiating gives us a delay differential equation which can be solved by hand. Here we work out the example of SO(even) for  $\sigma = 1$ . We want to find  $f \in L^2[-1, 1]$  such that

$$f(x) + \frac{1}{2} \int_{-1}^1 f(y) \mathbb{1}_{[-1, 1]}(x - y) dy = 1 \tag{1}$$

whenever  $x \in [0, 1]$ . We know  $f$  is even, so symmetrizing allows us to reconstruct  $f$  on  $[-1, 0]$ . First we rewrite the equation above in a more workable form; split the integral

$$\int_{-1}^1 f(y) \mathbb{1}_{[-1,1]}(x-y) dy = \int_{-1}^0 f(y) \mathbb{1}_{[-1,1]}(x-y) dy + \int_0^1 f(y) \mathbb{1}_{[-1,1]}(x-y) dy.$$

For the second integral on the right,  $x, y \in [0, 1]$  so  $x-y \in [-1, 1]$ , so the characteristic function term is redundant. In the first integral on the left,  $y \in [-1, 0]$  and  $x \in [0, 1]$  implies  $x-y \in [-1, 1]$  if and only if  $y \in [x-1, 0]$ . By evenness, this is equivalent to integrating  $f(y)$  for  $y \in [0, 1-x]$ . Thus (1) is equivalent to

$$f(x) + \frac{1}{2} \int_0^1 f(y) dy + \frac{1}{2} \int_0^{1-x} f(y) dy = 1 \quad (2)$$

for  $x \in [0, 1]$ . We want to find  $f$  such that the left hand side is constant with respect to  $x$ . Afterwards, setting  $x = 1$  allows us to compute the desired normalizing constant. Differentiating (2) with respect to  $x$  gives the delay differential equation

$$f'(x) - \frac{1}{2} f(1-x) = 0 \quad (3)$$

for all  $x \in [0, 1]$ . Divine inspiration and a routine check using the identity  $\sin(x) = \cos(x - \pi/2)$  shows that

$$f(x) = \cos\left(\frac{x}{2} - \frac{\pi+1}{4}\right)$$

is a solution. Of course, for those who are not divinely inspired, we can solve by differentiating (3) to obtain

$$f''(x) = -\frac{1}{2} f'(1-x) = -\frac{1}{4} f(x). \quad (4)$$

the second equality follows directly from (3). This is a standard linear differential equation with a two-parameter family of solutions given by  $f(x) = a \cos(x/2) + b \sin(x/2)$ . Substituting these solutions into (3) yields the linear system

$$\begin{pmatrix} 2 \cos(1/2) & 2 \sin(1/2) - 1 \\ 2 \sin(1/2) + 1 & -2 \cos(1/2) \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

The matrix has determinant zero with non-trivial rank so the kernel of the matrix is a one-dimensional subspace, i.e. all solutions to (3) are scalar multiples of our divinely inspired answer.

### 3 Quadratic kernel

There are three key observations. First, the kernels take the form of quadratic polynomials in  $|x|$  on the interval  $[-1, 1]$ , i.e.  $m_{G,\psi}(x) = -a - b|x| - c|x|^2$  whenever  $x \in [-1, 1]$ . Second, the solutions found in [FM] and [ILS] to the 1-level density case for varying support took the form of piecewise trigonometric polynomials, and in particular continuously differentiable everywhere except finitely many points. Third, the right-hand side of (9) is constant with respect to  $x \in [-1/2, 1/2]$ . Therefore, assuming  $g$  is sufficiently smooth, we can differentiate the Fredholm integral equation (9) to obtain a corresponding system of linear homogeneous differential equations.

**Lemma 1.** *If  $g \in L^2[-1/2, 1/2]$  is smooth and solves*

$$1 = g(x) + \int_{-1/2}^{1/2} (a + b|x-y| + c|x-y|^2) g(y) dy, \quad (5)$$

then it satisfies for  $x \in [-1/2, 1/2]$  the following system of equations,

$$1 = g(0) + a \int_{-1/2}^{1/2} |y|g(y)dy, \quad (6)$$

$$0 = g''(x) + 2bg(x) + 2c \int_{-1/2}^{1/2} g(y)dy, \quad (7)$$

$$0 = g'''(x) + 2bg'(x). \quad (8)$$

*Proof.* Notice that

$$\int_{-1/2}^{1/2} |x-y|g(y)dy = \int_{-1/2}^x (x-y)g(y)dy + \int_x^{1/2} (y-x)g(y)dy.$$

Differentiating under the integral sign gives

$$\begin{aligned} \frac{d}{dx} \int_{-1/2}^{1/2} |x-y|g(y)dy &= \int_{-1/2}^x g(y)dy - \int_x^{1/2} g(y)dy, \\ \frac{d^2}{dx^2} \int_{-1/2}^{1/2} |x-y|g(y)dy &= 2g(x) \end{aligned}$$

and

$$\begin{aligned} \frac{d}{dx} \int_{-1/2}^{1/2} (x-y)^2 g(y)dy &= \int_{-1/2}^{1/2} (2x-2y)g(y)dy, \\ \frac{d^2}{dx^2} \int_{-1/2}^{1/2} (x-y)^2 g(y)dy &= 2 \int_{-1/2}^{1/2} g(y)dy. \end{aligned}$$

And on the First Day of Genesis, He the God of Math deemed the rest to be trivial. □

**Theorem 2.** *The solutions take the form*

$$g(x) = \frac{6b^{3/2}(b+c) \cos(\sqrt{2b}x) - 6\sqrt{2}bc \sin \sqrt{b/2}}{6\sqrt{b}(b+c)^2 \cos \sqrt{b/2} + \sqrt{2}(6ab^2 + 3b^3 + 3b^2c + bc(c-12) - 6c^2) \sin(\sqrt{b/2})}.$$

*Proof.* Assuming  $g$  is even, has solutions of (8) take the form  $A \cos(\sqrt{2b}x) + C$ . This has two degrees of freedom. Substituting into (7) reduces to one degree of freedom,

$$\begin{aligned} 0 &= g''(x) + 2bg(x) + 2c \int_{-1/2}^{1/2} g(y)dy \\ &= 2C(b+c) + \frac{4Ac}{\sqrt{2b}} \sin \left( \frac{\sqrt{2b}}{2} \right). \end{aligned}$$

This shows that  $g$  is a scalar multiple of

$$\cos(\sqrt{2b}x) - \frac{2c}{b+c} \frac{\sin(\sqrt{b/2})}{\sqrt{2b}}.$$

Obviously this implies that  $b \geq 0$  if we want to avoid nasty complexifications. Pumping into Wolfram Alpha, we obtain

$$g(x) = \frac{6b^{3/2}(b+c) \cos(\sqrt{2b}x) - 6\sqrt{2}bc \sin \sqrt{b/2}}{6\sqrt{b}(b+c)^2 \cos \sqrt{b/2} + \sqrt{2}(6ab^2 + 3b^3 + 3b^2c + bc(c-12) - 6c^2) \sin(\sqrt{b/2})}.$$

□

## 4 Worked example: two-level SO(even) for $\sigma = 1/2$

For the two level density, we are interested in finding an optimal test function  $\Phi : \mathbb{R}^2 \rightarrow \mathbb{R}$  with Fourier transform supported in  $[-2\sigma, -2\sigma] \times [-2\sigma, -2\sigma]$  minimizing

$$\frac{1}{\Phi(0,0)} \int_{\mathbb{R}^2} \Phi(x,y) W_G(x,y) dx dy.$$

Putting restrictions on  $\Phi$ , such as assuming it is of the form  $\Phi(x,y) = \phi(x)\psi(y)$ , makes the problem much easier while hopefully giving us a good upper bound on the true value of the minimum. Charlie was interested in  $\sigma = 1/2$  with

$$\psi(x) = \left( \frac{\sin(\pi x)}{\pi x} \right)^2, \quad \hat{\psi}(x) = (1 - |x|) \mathbb{1}_{[-1,1]}(x).$$

Analogous to the set-up in Appendix A of ILS and Jesse Freeman's thesis, Charlie found that finding the optimal  $\phi : \mathbb{R} \rightarrow \mathbb{R}$  was equivalent to finding the unique solution  $g \in L^2[-1/2, 1/2]$  satisfying equation

$$g(x) - \int_{-1/2}^{1/2} m_G(x-y) g(y) dy = 1, \quad (9)$$

for all  $x \in [-1/2, 1/2]$ . Charlie proved that  $g$  is even, increasing on  $[-1/2, 0]$  and decreasing on  $[0, 1/2]$ . This seems to suggest that solutions are of the form  $g(x) = a \cos(px) + c$  for some period  $p > 0$  and constants  $a, c \in \mathbb{R}$ . We show that this is indeed the case for SO(even), which has corresponding kernel

$$m_{\text{SO(even)}}(x) = -\frac{1}{2} + \frac{4}{3}(1 - |x|)^2 + \frac{2}{3}(1 - x^2). \quad (10)$$

Substituting (10) into (9) gives

$$-\frac{3}{2} \int_{-1/2}^{1/2} g(y) dy + \frac{8}{3} \int_{-1/2}^{1/2} |x-y| g(y) dy - \frac{2}{3} \int_{-1/2}^{1/2} (x-y)^2 g(y) dy = 1$$

for all  $x \in [-1/2, 1/2]$ . Notice that

$$\int_{-1/2}^{1/2} |x-y| g(y) dy = \int_{-1/2}^x (x-y) g(y) dy + \int_x^{1/2} (y-x) g(y) dy.$$

Differentiating under the integral sign gives

$$\begin{aligned} \frac{d}{dx} \int_{-1/2}^{1/2} |x-y| g(y) dy &= \int_{-1/2}^x g(y) dy - \int_x^{1/2} g(y) dy, \\ \frac{d^2}{dx^2} \int_{-1/2}^{1/2} |x-y| g(y) dy &= 2g(x) \end{aligned}$$

and

$$\begin{aligned} \frac{d}{dx} \int_{-1/2}^{1/2} (x-y)^2 g(y) dy &= \int_{-1/2}^{1/2} (2x-2y) g(y) dy, \\ \frac{d^2}{dx^2} \int_{-1/2}^{1/2} (x-y)^2 g(y) dy &= 2 \int_{-1/2}^{1/2} g(y) dy. \end{aligned}$$

Thus, (9) becomes, after differentiating three times with respect to  $x$ ,

$$g'''(x) + \frac{16}{3} g'(x) = 0$$

for all  $x \in [-1/2, 1/2]$ . The solution set consists of the trigonometric polynomials  $a \cos(4x/\sqrt{3}) + b \sin(4x/\sqrt{3}) + c$  for some constants  $a, b, c \in \mathbb{R}$ . We know  $g$  is even, so this forces  $b = 0$ . To obtain the remaining coefficients, we differentiate (9) twice to obtain

$$\begin{aligned} 0 &= g''(x) + \frac{16}{3}g(x) - \frac{4}{3} \int_{-1/2}^{1/2} g(y) dy \\ &= 4c - \frac{4}{3}a \int_{-1/2}^{1/2} \cos(4y/\sqrt{3}) dy. \end{aligned}$$

Rearranging,

$$c = \frac{1}{3}a \int_{-1/2}^{1/2} \cos(4y/\sqrt{3}) dy = \frac{2}{3}a \int_0^{1/2} \cos(4y/\sqrt{3}) dy = \frac{\sin(2/\sqrt{3})}{2\sqrt{3}}a.$$

This shows that  $g$  is a scalar multiple of  $\cos(4x/\sqrt{3}) + \frac{\sin(2/\sqrt{3})}{2\sqrt{3}}$ . Plugging into (??) for  $x = 0$  should give the desired multiplicative constant.

**Theorem 3.** *The solutions to the integral equation*

$$g_G(x) - \int_{-1/2}^{1/2} m_G(x-y)g_G(y)dy = 1$$

for  $x \in [-1/2, 1/2]$  are

$$\begin{aligned} g_{\text{SO(even)}}(x) &= \frac{216 \cos(4x/\sqrt{3}) + 36\sqrt{3} \sin(2/\sqrt{3})}{-162 \cos(2/\sqrt{3}) + 5\sqrt{3} \sin(2/\sqrt{3})}, \\ g_{\text{SO(odd)}}(x) &= \frac{8 \cos(4x) + 12 \sin(2)}{2 \cos(2) + 3 \sin(2)}, \\ g_U(x) &= \frac{6 \cos(2x) + 6 \sin(1)}{3 \cos(1) + 4 \sin(1)}, \\ g_O(x) &= \frac{6 \cos(2x\sqrt{2}) + 3\sqrt{2} \sin(\sqrt{2})}{3 \cos(\sqrt{2}) + \sqrt{2} \sin(\sqrt{2})} \end{aligned}$$

## 5 Iteration

We have solved the minimization problem for fixed  $\psi$ ; our goal now is to generalize by optimizing over a larger class of  $\psi$ . Supposing that we can write  $\hat{\psi} = g * g$  for some even real valued function  $g \in L^2[-1/2, 1/2]$ , we consider  $g$  which are solutions to the earlier minimization problem, namely functions of the form

$$g(x) = a \cos(bx) + c.$$

Since the problem is scaling invariant, we can normalize the constant term  $c = 1$ . For brevity, denote  $h : \mathbb{R} \rightarrow \mathbb{R}$  the Fourier inverse of the map  $x \mapsto \cos(bx)\mathbb{1}_{[-1/2, 1/2]}(x)$ . WolframAlpha gives us explicitly

$$h(x) = \int_{-1/2}^{1/2} e^{2\pi ixy} \cos(by) dy = \frac{2b \sin(b/2) \cos(\pi x) - 4\pi x \cos(b/2) \sin(\pi x)}{b^2 - 4\pi^2 x^2}. \quad (11)$$

For notational convenience, denote  $\mathcal{F}$  the Fourier transform operator, i.e.  $\mathcal{F}\phi = \hat{\phi}$ . Viewing  $g$  as a function on  $\mathbb{R}$  supported on  $[-1/2, 1/2]$ , its Fourier inverse takes the form

$$\mathcal{F}^{-1}(g)(x) = a \int_{-1/2}^{1/2} e^{2\pi ixy} \cos(by) dy + \int_{-1/2}^{1/2} e^{2\pi ixy} dy = ah(x) + \frac{\sin(\pi x)}{\pi x}.$$

As the Fourier transform and its inverse take convolutions to products and vice versa, we have

$$\psi(x) = \mathcal{F}^{-1}g * g(x) = \left[ \mathcal{F}^{-1}(g)(x) \right]^2 = [ah(x)]^2 + 2ah(x) \frac{\sin(\pi x)}{\pi x} + \left( \frac{\sin(\pi x)}{\pi x} \right)^2. \quad (12)$$

Notice the last term is exactly our original  $\psi$  that was used earlier. This is unsurprising since we are pushing around convolutions and Fourier transforms of trigonometric polynomials, so by linearity the constant term should pop out.

$$\widehat{\psi}(x) = a\mathcal{F}([h(x)]^2) + a\mathcal{F}\left(h(x) \frac{\sin(\pi x)}{\pi x}\right) + (1 - |x|)\mathbb{1}_{[-1,1]}(x).$$

## References

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- [FM] A. Y. Khinchin, *Continued Fractions*, Third Edition, The University of Chicago Press, Chicago 1964.