

2.4.9 Solve $\begin{cases} u_t = k u_{xx} \quad (*) \\ u(x, 0) = x^2 \end{cases}$ by the given method.

If u satisfies $(*)$, then

$$(u_{xxx})_t = (u_t)_{xxx} = (k u_{xx})_{xxx} = k (u_{xxx})_{xx},$$

\downarrow exchange order of differentiation \downarrow by $(*)$ \downarrow exchange order of differentiation

with $u_{xxx}(x, 0) = \frac{d^3}{dx^3} u(x, 0) = \frac{d^3}{dx^3} x^2 = 0$
i.e. $v \stackrel{\text{def}}{=} u_{xxx}$ satisfies $\begin{cases} v_t = k v_{xx} \\ v(x, 0) = 0 \end{cases}$

By uniqueness of heat equation with zero initial condition, we have $v \equiv 0$, i.e. $u_{xxx} = 0$.

Integrating $v \equiv 0$ thrice, $u(x, t) = A(t)x^2 + B(t)x + C(t)$.

Since $u(x, 0) = x^2$, $A(0) = 1$, $B(0) = C(0) = 0$.

Also by $u_t = k u_{xx}$, we have

$$A'(t)x^2 + B'(t)x + C'(t) = 2kA(t) \Rightarrow \begin{cases} A'(t) = B'(t) = 0 \\ C'(t) = 2kA(t) \end{cases}$$

Thus $A(t) \equiv 1$, $B(t) \equiv 0$, and

$C'(t) = 2k$, $C(0) = 0$ give $C(t) = 2kt$

Thus $u(x, t) = x^2 + 2kt$.

2.4.10. (a) By the formulae given by the heat kernel,

$$u(x,t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4kt}} \cdot y^2 dy$$

Fix x . Let $p = \frac{x-y}{\sqrt{4kt}}$, then $dp = -\frac{1}{\sqrt{4kt}} dy$

$$u(x,t) = \frac{1}{\sqrt{4\pi kt}} \int_{\infty}^{-\infty} e^{-p^2} \cdot (x - \sqrt{4kt} \cdot p)^2 \cdot (-\sqrt{4kt}) dp$$

$$= \frac{1}{\cancel{\sqrt{4\pi kt}}} \int_{-\infty}^{\infty} e^{-p^2} \cdot (x - \sqrt{4kt} \cdot p)^2 \cdot \cancel{\sqrt{4kt}} dp$$

$$= \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-p^2} (x^2 - 4x\cancel{p\sqrt{kt}} + 4kt \cdot p^2) dp$$

by symmetry

$$= \frac{1}{\sqrt{\pi}} \left[\underbrace{\left(\int_{-\infty}^{\infty} e^{-p^2} dp \right)}_{=\sqrt{\pi} \text{ (exercise)}} \cdot x^2 + 4kt \int_{-\infty}^{\infty} e^{-p^2} \cdot p^2 dp \right]$$

$$= x^2 + \frac{4kt}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-p^2} \cdot p^2 dp$$

(b) By uniqueness (which is not so rigorous but OKAY),

$$u(x,t) = x^2 + \frac{4kt}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-p^2} \cdot p^2 dp = x^2 + 2kt$$

by 2.4.9

$$\Rightarrow \int_{-\infty}^{\infty} e^{-p^2} \cdot p^2 dp = \frac{\sqrt{\pi}}{2}$$

2.4.14 $\phi(x)$ is continuous s.t. $|\phi(x)| \leq C e^{ax^2}$, then (*)

$$\left| \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4kt}} \cdot \phi(y) dy \right|$$

$$\leq \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4kt}} \cdot |\phi(y)| dy$$

by (*)

$$\leq \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4kt}} \cdot C e^{ay^2} dy$$

$$= \frac{C}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4kt} + ay^2} dy \quad (**)$$

If $0 < t < \frac{1}{4ak}$, then $-\infty < -\frac{1}{4kt} < -a$ (Note $a > 0$, and $k > 0$)

For simplicity I denote $\lambda \stackrel{\text{def}}{=} -\frac{1}{4kt}$, so the exponent is

$$\begin{aligned} \lambda(x-y)^2 + ay^2 &= (\lambda+a)y^2 - 2\lambda x \cdot y + \lambda x^2 \\ &= (\lambda+a) \left(y^2 - \frac{2\lambda x}{\lambda+a} y + \left(\frac{\lambda x}{\lambda+a} \right)^2 \right) - \frac{(\lambda x)^2}{\lambda+a} + \lambda x^2 \end{aligned}$$

$$= (\lambda+a) \left(y - \frac{\lambda x}{\lambda+a} \right)^2 + \frac{\lambda a}{\lambda+a}$$

Fix x , make the change of variables $z = y - \frac{\lambda x}{\lambda+a}$.
then (**) becomes

$$\frac{C}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{(\lambda+a)z^2 + \frac{\lambda a}{\lambda+a}} dz$$

$$= \frac{C}{\sqrt{4\pi kt}} \cdot e^{\frac{\lambda a}{\lambda+a}} \cdot \int_{-\infty}^{\infty} e^{(\lambda+a)z^2} dz < \infty$$

if (and only if) $\lambda+a < 0$, i.e. $0 < t < \frac{1}{4ak}$.

$$2.4.16 \quad u_t - k u_{xx} + b u = 0 \quad -\infty < x < \infty$$

$$u(x, 0) = \phi(x),$$

where $b > 0$ is a constant.

Proof. By the hint, let $u(x, t) = e^{-bt} v(x, t)$,
then $u_t = -b e^{-bt} v + e^{-bt} v_t$

$$u_x = e^{-bt} v_x \text{ and } u_{xx} = e^{-bt} v_{xx}$$

Thus $u_t - k u_{xx} + b u = 0$ reads

$$-b e^{-bt} v + e^{-bt} v_t - k e^{-bt} v_{xx} + b e^{-bt} v = 0$$

$$\text{i.e., } v_t - k v_{xx} = 0 \quad (\text{Note } e^{-bt} > 0)$$

Also by $u(x, 0) = \phi(x)$, we have

$$v(x, 0) = \phi(x) \cdot e^{b \cdot 0} = \phi(x)$$

Thus v satisfies $\begin{cases} v_t - k v_{xx} = 0 \\ v(x, 0) = \phi(x). \end{cases}$

By the formula,

$$v(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4kt}} \cdot \phi(y) dy$$

$$\text{i.e., } u(x, t) = \frac{e^{-bt}}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4kt}} \cdot \phi(y) dy$$

2.5.4 Let $u(x,t)$ satisfy $u_t - c^2 u_{xx} = 0$
 $-\infty < x < \infty$.

with $|u''| \leq C$ for some $C > 0$.

Let

$$v(x,t) = \frac{c}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-s^2 c^2 / 4kt} \cdot u(x,s) ds$$

(a) Show that $v(x,t)$ solves $v_t - \frac{k}{c^2} v_{xx} = 0$.

~~Proof~~. By the hint, we may rewrite v
 as $v(x,t) = \int_{-\infty}^{\infty} H(s,t) u(x,s) ds$, where

$$H(s,t) \stackrel{\text{def}}{=} \frac{c}{\sqrt{4\pi kt}} e^{-s^2 c^2 / 4kt}$$

Then $\frac{\partial}{\partial t} H(s,t)$

$$= \frac{c}{\sqrt{4\pi k}} \cdot \left(-\frac{1}{2} t^{-3/2} \cdot e^{-s^2 c^2 / 4kt} + t^{-1/2} \cdot e^{-s^2 c^2 / 4kt} \cdot \frac{(s^2 c^2)}{4k} \left(-\frac{1}{t^2} \right) \right)$$

$$= \frac{c}{\sqrt{4\pi k}} \cdot e^{-s^2 c^2 / 4kt} \left(\frac{s^2 c^2}{4k} \cdot \frac{1}{t^{5/2}} - \frac{1}{2} \cdot \frac{1}{t^{3/2}} \right)$$

$$\frac{\partial}{\partial s} H(s,t) = \frac{c}{\sqrt{4\pi kt}} \cdot e^{-s^2 c^2 / 4kt} \cdot \left(-\frac{2sc^2}{4kt} \right)$$

$$\frac{\partial^2}{\partial s^2} H(s,t) = \frac{c}{\sqrt{4\pi kt}} \left[e^{-s^2 c^2 / 4kt} \cdot \left(-\frac{2sc^2}{4kt} \right)^2 + e^{-s^2 c^2 / 4kt} \cdot \left(-\frac{2c^2}{4kt} \right) \right]$$

$$\begin{aligned}
&= \frac{c}{\sqrt{4\pi kt}} e^{-s^2 c^2 / 4kt} \left[\frac{s^2 c^4}{4k^2 t^2} + \frac{c^2}{2kt} \right] \\
\Rightarrow \frac{\partial}{\partial t} H(s, t) - \frac{k}{c^2} \frac{\partial^2}{\partial s^2} H(s, t) \\
&= \frac{c}{\sqrt{4\pi kt}} e^{-s^2 c^2 / 4kt} \left(\frac{s^2 c^2}{4kt^2} - \frac{1}{2t} \right) \\
&\quad - \frac{k}{c^2} \cdot \frac{c}{\sqrt{4\pi kt}} e^{-s^2 c^2 / 4kt} \left[\frac{s^2 c^4}{4k^2 t^2} + \frac{c^2}{2kt} \right] = 0 \\
&\text{i.e. } H(s, t) \text{ satisfies } \underline{H_t - \frac{k}{c^2} H_{ss} = 0}
\end{aligned}$$

$$\begin{aligned}
\text{Then } v_t &= \int_{-\infty}^{\infty} \frac{\partial}{\partial t} H(s, t) \cdot u(x, s) ds \\
&= \frac{k}{c^2} \int_{-\infty}^{\infty} \left(\frac{\partial^2}{\partial s^2} H(s, t) \right) \cdot u(x, s) ds
\end{aligned}$$

$$\begin{aligned}
&\stackrel{\text{IBP}}{=} \frac{k}{c^2} \left[\underbrace{\left(\frac{\partial}{\partial s} H(s, t) \right) \cdot u(x, s)}_{=0 \text{ as } t \rightarrow \pm \infty \text{ by L'Hopital}} \Big|_{s=-\infty}^{\infty} - \int_{-\infty}^{\infty} \frac{\partial}{\partial s} H(s, t) \cdot \frac{\partial}{\partial s} u(x, s) ds \right] \\
&= -\frac{k}{c^2} \int_{-\infty}^{\infty} \frac{\partial}{\partial s} H(s, t) \cdot \frac{\partial}{\partial s} u(x, s) ds \\
&\stackrel{\text{IBP}}{=} -\frac{k}{c^2} \left[\underbrace{H(s, t) \cdot \frac{\partial}{\partial s} u(x, s)}_{=0 \text{ as } t \rightarrow \pm \infty \text{ by L'Hopital}} \Big|_{s=-\infty}^{\infty} - \int_{-\infty}^{\infty} H(s, t) \cdot \frac{\partial^2}{\partial s^2} u(x, s) ds \right]
\end{aligned}$$

$$= \frac{k}{c^2} \int_{-\infty}^{\infty} H(s, t) \cdot \frac{\partial^2}{\partial s^2} u(x, s) ds.$$

$$\text{While } v_{xx} = \int_{-\infty}^{\infty} H(s, t) \cdot \frac{\partial^2}{\partial x^2} u(x, s) ds$$

$$\text{Thus } v_t - k v_{xx}$$

$$= \int_{-\infty}^{\infty} H(s, t) \cdot \left(\frac{k}{c^2} \frac{\partial^2}{\partial s^2} u(x, s) - k \frac{\partial^2}{\partial x^2} u(x, s) \right) ds$$

$$= \int_{-\infty}^{\infty} H(s, t) \cdot \frac{k}{c^2} (u_{ss} - c^2 u_{xx}) ds$$

$$= \int_{-\infty}^{\infty} H(s, t) \cdot \frac{k}{c^2} \cdot 0 ds = 0, \quad \text{by assumption of } u$$

$$\Rightarrow v_t - k v_{xx} = 0.$$

(b) Full marks will be given if you state (by the hint),

$$\text{for any } s, \lim_{t \rightarrow 0^+} H(s, t) = \delta(s),$$

i.e. the 'dirac delta' function at s .

Then performing some formal computation based on $\lim_{t \rightarrow 0^+} H(s, t) = \delta(s)$. Namely,

$$\begin{aligned} \lim_{t \rightarrow 0} v(x, t) &= \lim_{t \rightarrow 0} \int_{-\infty}^{\infty} H(s, t) u(x, s) ds \\ &= \int_{-\infty}^{\infty} \lim_{t \rightarrow 0} H(s, t) u(x, s) ds \end{aligned}$$

$$= \int_{-\infty}^{\infty} \delta(s) u(x, s) ds$$

$$= u(x, 0).$$

This is, essentially, the definition of $\delta(s)$.

For a rigorous treatment, see Strauss Chp 3.5, Theorem 1.