

The force-free magnetosphere of a Kerr black hole: the role of boundary conditions

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The role of boundary conditions in the structure of a force-free black hole magnetosphere was rarely discussed, since previous studies have been focused on the the field lines entering the horizon which is causally disconnected and on which the boundary condition imposed usually makes no difference to the magnetosphere structure. However, recent high-resolution general relativistic (GR) force-free simulation shows that there are both field lines entering the horizon and field lines ending up on the equatorial current sheet within the ergosphere for asymptotic uniform field. For the latter, the equatorial boundary condition is well approximated being marginally force-free, i.e., $B^2 - E^2 = 0$, where B and E are the magnetic and electric field strength, respectively. In this paper, we revisit this problem by solving the GR Grad-Shafranov equation which governs the structure of the force-free magnetosphere in steady state and self-consistently imposing the marginally force-free equatorial condition. We also discuss the applicability of this boundary condition and the numerical algorithm proposed in this paper for general magnetic field configurations.

I. INTRODUCTION

The paper is organized as follows. In Section II, we outline the basic governing equations. In Section III, we clarify the radiation condition, boundary conditions and numerical algorithm for the uniform field solution. In Section IV, we discuss the solution uniqueness, near-horizon field line configuration and the application to general field line configurations. Summary is given in Section V.

II. BASIC EQUATIONS

In this paper, we adopt the Kerr-Schild coordinate with the line element

$$ds^2 = - \left(1 - \frac{2r}{\Sigma}\right) dt^2 + \left(\frac{4r}{\Sigma}\right) dr dt + \left(1 + \frac{2r}{\Sigma}\right) dr^2 \\ + \Sigma d\theta^2 - \frac{4ar \sin^2 \theta}{\Sigma} d\phi dt - 2a \left(1 + \frac{2r}{\Sigma}\right) \sin^2 \theta d\phi dr \\ + \frac{\beta}{\Sigma} \sin^2 \theta d\phi^2$$

where $\Sigma = r^2 + a^2 \mu^2$, $\Delta = r^2 - 2r + a^2$, $\beta = \Delta \Sigma + 2r(r^2 + a^2)$, and a is the dimensionless BH spin. In the force-free approximation, electromagnetic energy greatly exceeds that of matter. Consequently, the force-free magnetosphere is governed by energy conservation equation of electromagnetic field, or conventionally called as the GS equation. In the Kerr spacetime, the axisymmetric and

steady GS equation can be written in a compact form [1]

$$\begin{aligned} & \left[A_{\phi,rr} + \frac{\sin^2 \theta}{\Delta} A_{\phi,\mu\mu} \right] \mathcal{K}(r, \theta; \Omega) \\ & + \left[A_{\phi,r} \partial_r^\Omega + \frac{\sin^2 \theta}{\Delta} A_{\phi,\mu} \partial_\mu^\Omega \right] \mathcal{K}(r, \theta; \Omega) \\ & + \frac{1}{2} \left[A_{\phi,r}^2 + \frac{\sin^2 \theta}{\Delta} A_{\phi,\mu}^2 \right] \Omega' \partial_\Omega \mathcal{K}(r, \theta; \Omega) \\ & - \frac{\Sigma}{\Delta} II' = 0, \end{aligned} \quad (1)$$

with $\mathcal{K}(r, \theta; \Omega) = \left[\frac{\beta}{\Sigma} \Omega^2 \sin^2 \theta - \frac{4ra}{\Sigma} \Omega \sin^2 \theta - \left(1 - \frac{2r}{\Sigma}\right) \right]$ being the LS function, $\mu \equiv \cos \theta$, and the primes designate derivatives with respect to A_ϕ . $\partial_i^\Omega (i = r, \mu)$ denotes the partial derivative with respect to coordinate i with Ω fixed, and ∂_Ω is the derivative with respect to Ω .

III. UNIFORM FIELD SOLUTION

A. Solution uniqueness and radiation condition

For common field configurations, including split monopole and parabola, there exist two LSs where the LS function vanishes and the GS equation degrades from second order to first order. As proposed by Contopoulos et.al. [2], one can adjust the two free functions $\Omega(A_\phi)$ and $I(A_\phi)$ enabling field lines smoothly cross the two LSs, then the solution $\{\Omega(A_\phi), I(A_\phi), A_\phi(r, \theta)\}$ is uniquely found. But for the vertical field lines, their exist only one LS, which is insufficient to determine two free functions. Therefore people have been wondering the possibility of many solutions though the many-solutions scenario is not supported by force-free dynamics simulations. To tackle this problem, we proposed that the two functions are not independent; instead, they are related by the radiation condition [1, 3]

$$I = 2\Omega A_\phi, \quad (2)$$

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which has been confirmed by the recent high-accuracy simulations [4].

B. Boundary conditions

The boundary conditions at infinity (inner infinity $r = r_+$ and outer infinity $r \rightarrow \infty$) and on the polar axis can be simply set as

$$\begin{aligned} A_{\phi,r}|_{r=r_+,\infty} &= 0, \\ A_{\phi}|_{\mu=1} &= 0, \end{aligned} \quad (3)$$

where r_+ is the radius of the event horizon, while the equatorial boundary conditions are more unclear until recent simulations [4] come out which show that there exists an equator current sheet within the ergosphere where the magnetic dominance marginally loses, i.e., $B^2 - E^2 \approx 0$. Motivated by these simulations, we choose the following equatorial boundary condition in our numerical solutions,

$$A_{\phi,\mu}(\mu = 0, r > 2) = 0, \quad (4a)$$

$$B^2 - E^2(\mu = 0, r_+ \leq r \leq 2) = 0. \quad (4b)$$

In fact, Equation (4b) is neither a Dirichlet nor a Neumann boundary condition, since

$$\begin{aligned} &B^2 - E^2 \\ &= \frac{1}{\Sigma \sin^2 \theta} \left[-\mathcal{K} \left(A_{\phi,r}^2 + \frac{\sin^2 \theta}{\Delta} A_{\phi,\mu}^2 \right) + \frac{\Sigma}{\Delta} I^2 \right], \end{aligned} \quad (5)$$

which involves both derivatives $A_{\phi,\mu}$ and $A_{\phi,r}$ on the boundary. Therefore, it is numerically non-trivial to impose this boundary condition in computation.

In addition, there is a coordinate singularity $1/\Delta$ in the expression of $B^2 - E^2$. To avoid possible numerical difficulty, we use the prescription

$$\int_{A_{\phi}^{\text{HE}}}^{A_{\phi}^{\text{EE}}} \left(\frac{B^2 - E^2}{B^2 + E^2} \right)^2 dA_{\phi} \Big/ (A_{\phi}^{\text{EE}} - A_{\phi}^{\text{HE}}) < 10^{-3}, \quad (6)$$

in our computation, as a proxy of the marginally force-free equatorial boundary condition (4b), where A_{ϕ}^{HE} and A_{ϕ}^{EE} are the magnetic flux entering the horizon and the ergosphere, respectively; “HE” and “EE” are short for Horizon-Equator and Ergosphere-Equator, respectively. For definiteness, we choose $B^2 + E^2$ to be the energy density measured by the zero-angular-momentum-observers. Explicitly, we have

$$\begin{aligned} &B^2 + E^2 \\ &= \frac{1}{\Sigma \sin^2 \theta} \left[\left(\mathcal{K} + \frac{\Delta \Sigma}{\beta} \right) \left(A_{\phi,r}^2 + \frac{\sin^2 \theta}{\Delta} A_{\phi,\mu}^2 \right) + \frac{\Sigma}{\Delta} I^2 \right] \end{aligned} \quad (7)$$

C. Generic properties of the BH magnetosphere structure

Before delving into the details of numerically solving the GS equations, here we point out that from the radiation condition (2) and the marginally force-free boundary condition (4b) themselves contain rich information about the BH magnetosphere structure.

Let’s first find out where the LS intersect with the equator, $r_{\text{LS}}|_{\mu=0}$. From the expression of the LS function, it is straightforward to see $r_{\text{LS}}|_{\mu=0} \leq 2$ for $\Omega \geq 0$. On this point, the LS function \mathcal{K} vanishes, and I must also vanish for satisfying the boundary condition (4b), which in turns indicates a vanishing angular velocity Ω from the radiation condition (2). Due to the static limit, Ω must be non-zero inside the ergosphere. Therefore the only possibility is $r_{\text{LS}}|_{\mu=0} = 2$, i.e. to satisfy the boundary condition (4b), the LS must interact the equator at $r = 2$, which also justifies our choice of equatorial boundary conditions (4a, 4b).

From above analysis, we expect several generic features in the magnetospheres that are independent of the GS equation: (1) the LS runs from $r = r_+$ to $r = 2$ as θ varies from 0 to $\pi/2$, (2) since I vanishes at $r_{\text{LS}}|_{\mu=0}$, we expect no current sheet within the magnetosphere except the equatorial current sheet extending from r_+ to 2, which gives rise to a cusp ($A_{\phi,\mu} \neq 0$) to the equatorial magnetic field lines, (3) magnetic field lines entering the ergosphere end up either on the horizon or on the equatorial current sheet, both of which carry electric current and therefore Poynting flux.

With the guidance of the qualitative features above, we are to numerically solve the GS equation and quantify these features.

D. Numerical method

We aim to find a pair of $\Omega(A_{\phi})$ and $I(A_{\phi})$ satisfying the radiation condition (2) and enabling field lines smoothly crossing the LS, and suitable normal derivative $A_{\phi,\mu}(\mu = 0, r_+ \leq r \leq 2)$ on the equator guaranteeing the boundary condition (4b).

In our computation, we define a new radial coordinate $R = r/(1+r)$, confine our computation domain $R \times \mu$ in the region $[R(r_+), 1] \times [0, 1]$, and implement a uniform 512×64 grid. The numerical algorithm for searching desired $\{\Omega(A_{\phi}), I(A_{\phi}), A_{\phi,\mu}(\mu = 0, r_+ \leq r \leq 2)\}$ is detailed in the following steps.

1. We choose an initial guess for the field configuration, functions $\{\Omega(A_{\phi}), I(A_{\phi})\}$ and equatorial boundary

condition as follows

$$A_\phi = \frac{B_0}{2} r^2 \sin^2 \theta, \quad (8a)$$

$$\Omega = 0.5\Omega_H (1 - A_\phi/A_\phi^{\text{HE}}), \quad (8b)$$

$$I = \Omega_H A_\phi (1 - A_\phi/A_\phi^{\text{HE}}), \quad (8c)$$

$$A_{\phi,\mu}(\mu = 0, r_+ \leq r \leq 2) = -(r/r_+)^3, \quad (8d)$$

where $\Omega_H = a/2r_+$ is the angular velocity of the BH.

2. We evolve the GS equation (1) using the well-known relaxation method [5] and adjust $II'(A_\phi)$ until field lines smoothly cross the LS [see e.g. 2, 3, 6, 7, for more details].

3. Usually the current I found in Step 2 neither satisfies the radiation condition (2) nor guarantees the boundary condition (4b). We adjust $A_{\phi,\mu}(\mu = 0, r_+ \leq r \leq 2)$ as follows

$$A_{\phi,\mu}|_{\text{new}} = A_{\phi,\mu}|_{\text{old}} + \zeta_1 \times [2\Omega A_\phi (2\Omega A_\phi)' - II'], \quad (9)$$

where ζ_1 is an empirical step size. For each new $A_{\phi,\mu}$, we repeat Step 2 and iterative correction (9) until $A_{\phi,\mu}(\mu = 0, r_+ \leq r \leq 2)$ converges, i.e. the condition $2\Omega A_\phi (2\Omega A_\phi)' = II'$ is achieved for $A_\phi \in (A_\phi^{\text{HE}}, A_\phi^{\text{EE}})$.

4. The remaining task is to adjust $\Omega(0 < A_\phi < A_\phi^{\text{HE}})$ enabling the radiation condition (2) for $A_\phi \in (0, A_\phi^{\text{HE}})$ and to adjust $\Omega(A_\phi^{\text{HE}} \leq A_\phi \leq A_\phi^{\text{EE}})$ enabling the boundary condition (4b) for $A_\phi \in (A_\phi^{\text{HE}}, A_\phi^{\text{EE}})$. The first part is straightforward, i.e.,

$$2A_\phi\Omega_{\text{new}} = I|_{0 < A_\phi < A_\phi^{\text{HE}}}, \quad (10)$$

and the second part can be realized by iterative correction

$$2A_\phi(\Omega_{\text{new}} - \Omega_{\text{old}}) = -\zeta_2 \times \Delta(B^2 - E^2)|_{\mu=0, r_+ \leq r \leq 2}, \quad (11)$$

where ζ_2 is again an empirical step size, and we have multiplied factor Δ in the correction term to avoid numerical difficulty in the vicinity of the event horizon. To eliminate unphysical discontinuity in the angular velocity at A_ϕ^{HE} , we fit $\Omega_{\text{new}}(A_\phi)$ on the whole range $(0, A_\phi^{\text{EE}})$ via a fifth-order polynomial.

5. For the new angular velocity $\Omega_{\text{new}}(A_\phi)$ obtained in Step 4, we repeat Step 2 to Step 4, until both the radiation condition (2) and the numerical prescription (6) for the boundary condition (4b) is satisfied.

E. Numerical results

In Figure 1, we plot the magnetic field lines enclosing a BH with spin $a = 0.99$ as an example, which explicitly displays the features we anticipated in Section III C.

In Figure 2, we show the angular velocity function $\Omega(A_\phi)$ for different BH spins and compare it with the counterpart obtained from recent high-resolution simulations done by East and Yang [4]. For reference, we

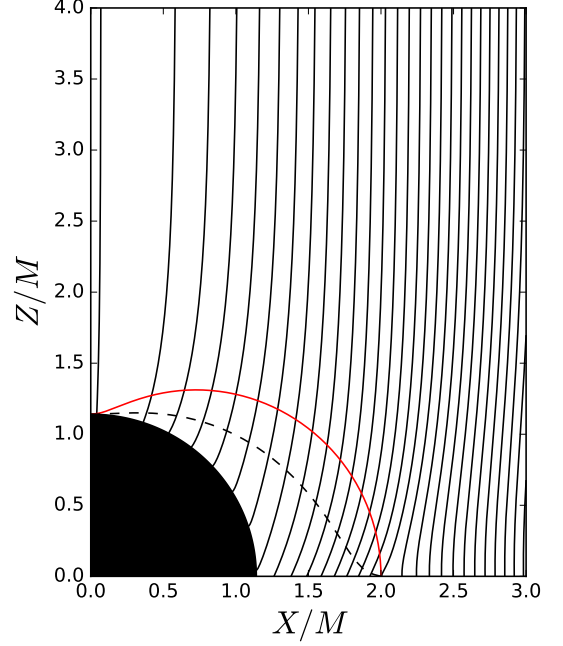


FIG. 1. The configuration of field lines for the magnetosphere of a Kerr BH with spin $a = 0.99$, where the solid/red line is the ergosphere and the dashed/black line is the LS, both of which intersect with the equator at $r = 2M$.

also plot the leading-order analytic solution in the slow-rotation limit [4, 8–10],

$$\Omega = \Omega_H \frac{\sqrt{1-\psi}}{1+\sqrt{1-\psi}}, \quad (12)$$

where $\psi = A_\phi/(2B_0M^2)$. From our numerical solutions, we find the magnetic flux entering the ergosphere A_ϕ^{EE} increases with BH spin and approach $2.75B_0M^2$ for extremal spins (upper panel of Figure 2), which is about $\approx 5\%$ lower than the simulation result (Figure 3 of [4]), while the angular velocity Ω as a function of normalized magnetic flux $A_\phi/A_\phi^{\text{EE}}$ shows a good agreement with the simulation results.

With the angular velocity $\Omega(A_\phi)$ obtained, the energy extraction rate from the BH is given by

$$\dot{E} = 4\pi \int_0^{A_\phi^{\text{EE}}} \Omega \times I dA_\phi. \quad (13)$$

It is straightforward to obtain the energy extraction rate in the slow-rotation limit

$$\dot{E} = 128\pi \left(\frac{17}{24} - \ln 2 \right) B_0^2 M^4 \Omega_H^2. \quad (14)$$

In Figure 3, we compare the energy extraction rate $\dot{E}(\Omega_H)$ derived from our numerical solutions with East and Yang's simulation results [4], where the data points are taken from either the simulations or our numerical

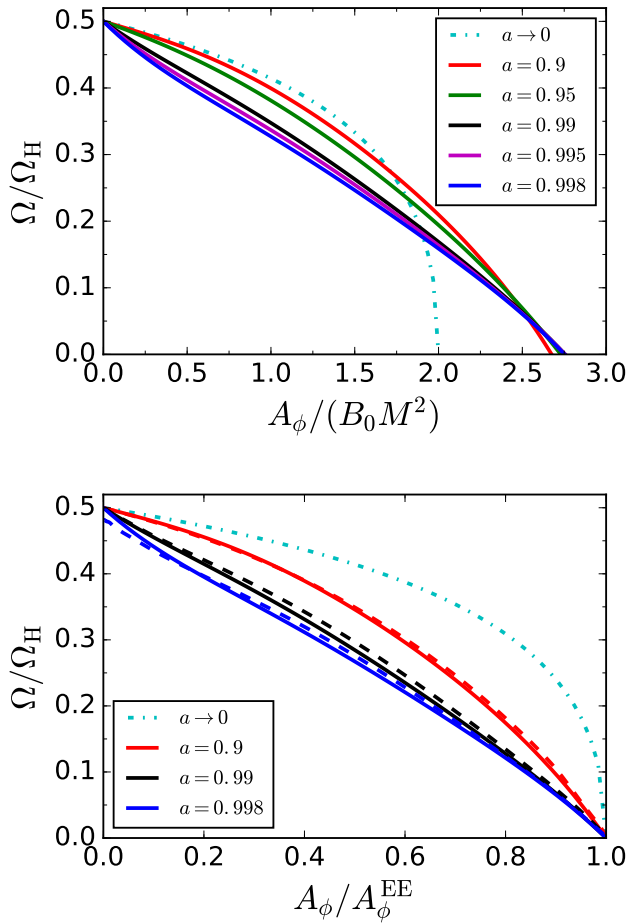


FIG. 2. Upper panel: the angular velocity $\Omega(A_\phi)$ for different BH spins. Lower Panel: comparison of our numerical results (solid lines) with the simulation results of Ref. [4] (dashed lines). For reference, we also plot the leading order analytic solution in dash-dotted lines.

solutions, while the solid lines are corresponding polynomial fitting curves which we require to approach Eq.(14) for small spins and to be flat for extremal spins. As expected, our energy extraction rate $\dot{E}(\Omega_H)$ is $\approx 10\%$ lower than the corresponding simulation results, due to the $\approx 5\%$ smaller magnetic flux A_ϕ^{EE} .

IV. DISCUSSION

A. Application to general field configurations

In real astrophysical environment, we expect the field lines are more close to parabolas instead of being strictly vertical. In many previous studies of such field configurations [e.g. 6, 7, 11], due to lacking knowledge of the equatorial boundary condition, the equator within the ergosphere was intentionally excluded out of the com-

putation domain by manually introducing a “wall” ex-

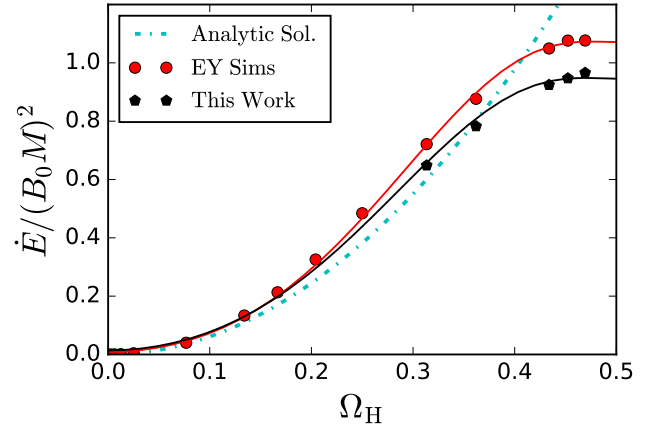


FIG. 3. Comparison of the energy extraction rates $\dot{E}(\Omega_H)$ obtained from three different approaches: the leading-order analytic solution (14), our numerical solutions and the high-resolution force-free simulations [4].

tending from the horizon-equator intersection to infinity. Such simplification obviously misses magnetic field lines rooting on the equatorial current sheet, which contribute about half of the total BZ flux for extremal spins in the case of vertical field lines.

Due to the resemblance of the vertical field configuration and general parabolic field configurations, it is natural to expect an equatorial current develops within the ergosphere, where the magnetic dominance loses, therefore the marginally force-free boundary condition (4b) should also be a good work approximation here. It is straightforward to solve the GS equation and self-consistently impose the marginally force-free boundary condition following the algorithm detailed in Section IIID.

The analysis of Section IIIC

B. Near-horizon field configuration

In one of our previous papers [3], we made a claim that “in the steady axisymmetric force-free magnetosphere around a Kerr BH, all magnetic field lines that cross the infinite-redshift surface must intersect the event horizon”. However, simulations shows that an equatorial current sheet inevitably develops within the ergosphere, where the force-free condition breaks down.

V. SUMMARY

In this paper, we revisit the problem of the uniform field solution to the BZ mechanism.

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