Wright-Fisher stochastic heat equations with irregular drifts

Zhenyao Sun

Beijing Institute of Technology

July, 2025

Joint work with Clayton Barnes (AWS) and Leonid Mytnik (Technion)

Regularization by noise

• Consider the differential equation:

$$\begin{cases} \mathrm{d}X_t = b(X_t)\mathrm{d}t = |X_t|^\alpha \mathrm{d}t, \quad t>0, \\ X_0 = 0, \end{cases}$$

where $\alpha \in (0,1)$.

- The drift $b(x) = |x|^{\alpha}$ is not Lipschitz at 0 \implies non-uniqueness of the solutions.
- One solution $X_t \equiv 0$.
- The other solution $X_t = C_{\alpha} t^{\frac{1}{1-\alpha}}, t \geq 0$.

Regularization by non-degenerate noise

Zvonkin (1974, Mat. Sb. (N.S.)), Veretennikov (1979, Mat. Sb. (N.S.))

Suppose that

- b is a bounded measurable function, and
- B is a Brownian motion,

then there exists a unique strong solution to the SDE

$$\begin{cases} \mathrm{d}X_t = b(X_t)\mathrm{d}t + \mathrm{d}B_t, & t > 0, \\ X_0 = x \in \mathbb{R}. \end{cases}$$

Zvonkin's transform is not available for SPDEs.

Partial regularization effect by degenerate noise

- Uniqueness in law for one-dimensional SDE can be analyzed by Feller's test.
- For example, consider non-negative solutions to the SDE

$$\mathrm{d}X_t = b(X_t)\mathrm{d}t + \sqrt{2X_t}\mathrm{d}B_t; \quad X_0 = 0$$

where, with $\alpha > 0$ and $\beta > 0$,

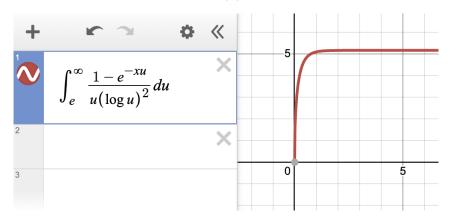
$$b(x) := \int_e^\infty \frac{1 - e^{-xu}}{\alpha u (\log u)^{1+\beta}} du, \quad x \ge 0.$$

Clement (2019, Electron. J. Probab.)

- If $\beta > 1$, the uniqueness in law holds;
- If $\beta = 1$ and $\alpha \ge 1$, the uniqueness in law holds;
- If $\beta = 1$ and $\alpha < 1$, the uniqueness in law fails;
- If β < 1, the uniqueness in law fails.

Regularization by multiplicative noise

• The shape of a "critical" drift b(x):



Wright-Fisher Stochastic Heat Equations (Wright-Fisher SHE)

 Quasi-linear heat equation perturbed by the Wright-Fisher space-time white noise

$$\begin{cases} \partial_t u = \frac{\Delta}{2} u + b(u) + \sqrt{|u(1-u)|} \dot{W}, & x \in \mathbb{R}, t \ge 0, \\ u_0 = f, & x \in \mathbb{R}. \end{cases}$$

- \dot{W} is the space-time white noise on $\mathbb{R}_+ \times \mathbb{R}$, i.e. a centered Gaussian process with $\mathbb{E}[\dot{W}_t(x)\dot{W}_s(y)] = \delta_0(t-s)\delta_0(x-y)$.
- The noise coefficient $\sqrt{|u(1-u)|}$
 - is non-Lipshitz at u = 0 and u = 1; and
 - is degenerate at u = 0 and u = 1.
- Challenging open problems:
 - the strong uniqueness?
 - the solution theory in higher dimensions?
- Question: How strong is the regularization effect of the Wright-Fisher white noise?

Motivation

- Shiga (1988, Math. Appl.): Wright-Fisher SHE = scaling limit of "genetic stepping stone model."
 - $b(u) = c_1(1-u) c_2u + c_3u(1-u)$.
 - $c_1 > 0$ and $c_2 > 0$ are called the mutation rates.
 - $c_3 \in \mathbb{R}$ is called the selection rate.
- Mueller-Tribe (1995, Probab. Theory Related Fields), Durrett-Fan (2016, Ann. Appl. Probab.): Wright-Fisher SHE = scaling limit of (biased) voter model.
 - $b(u) = c_3 u(1-u)$.
 - Unbiased $\implies c_3 = 0$.
- Brunet-Derrida (1997, Phys. Rev. E), Mueller-Mytnik-Quastel (2011, Invent. Math.): The Wright-Fisher SHE is the key to the proof of the Brunet-Derrida conjecture.

Weak existence

• Fix intitial value $f \in \mathcal{C}(\mathbb{R},[0,1])$.

Shiga (1994, Can. J. Math.)

If $b(\cdot)$ is continuous and $b(0) \ge 0 \ge b(1)$, then the weak existence holds for the SPDE

(*)
$$\begin{cases} \partial_t u = \frac{\Delta}{2} u + b(u) + \sqrt{u(1-u)} \dot{W}, & x \in \mathbb{R}, t \ge 0, \\ u_0 = f, & x \in \mathbb{R}. \end{cases}$$

That is, there exists a filtered probability space $(\Omega, \mathcal{F}, (\mathcal{F}_t), \mathbb{P})$, a space-times white noise \dot{W} , and an adapted continuous $\mathcal{C}(\mathbb{R}, [0,1])$ -valued process $(u_t)_{t\geq 0}$, such that (*) holds in an analytically weak sense.

• Question: What if $b(\cdot)$ is discontinuous? Even the existence is not clear.

Uniqueness in law: Duality Method

• uniqueness in law = the probability law induced by $(u_t)_{t\geq 0}$ on the path sapce $\mathcal{C}(\mathbb{R}_+,\mathcal{C}(\mathbb{R},[0,1]))$ is unique.

Shiga (1988, Math. Appl.)

The uniqueness in law of (*) holds provided

$$b(u) = c_1(1-u) - c_2u + c_3u(1-u)$$
 where $c_1 \ge 0, c_2 \ge 0$ and $c_3 \in \mathbb{R}$.

Athreya-Tribe (2000, Ann. Probab.)

The uniqueness in law of (*) holds provided

$$b(u) = \sum_{k=0}^{\infty} b_k u^k$$
, and $b_1 < -\sum_{k=0, k \neq 1}^{\infty} |b_k| R^{k-1}$ for some $R > 1$.

- Both Shiga (1988) and Athreya-Tribe (2000) used the duality method.
- The drifts are Lipshitz functions.

Uniqueness in law: Girsanov transformation

Mueller-Mytnik-Ryzhik (2021, Comm. Math. Phys.)

The uniqueness in law holds provided b is continuous,

$$\sup_{u\in(0,1)}\frac{|b(u)|}{\sqrt{u(1-u)}}<\infty, \text{ and } f(x)=1-f(-x)=0 \text{ for large enough } x.$$

- When the red part holds, we say the initial value f has a compact interface.
- The main tool is Girsanov transformation.
- The drift can be a non-Lipshitz Hölder continuous function.

Main Result

• Recall the condition in Athreya-Tribe (2000):

$$b(u) = \sum_{k=0}^{\infty} b_k u^k$$
, and $b_1 < -\sum_{k \in \{0\} \cup \mathbb{N} \setminus \{1\}} |b_k| R^{k-1}$ for some $R > 1$.

Barnes-Mytnik-S. (2025+, to appear in Probab. Theory Related Fields)

The weak existence and uniqueness in law holds for (*) provided

$$b(u) = \sum_{k \in \{0,\infty\} \cup \mathbb{N}} b_k u^k = \sum_{k=0}^{\infty} b_k u^k + b_{\infty} \mathbf{1}_{\{1\}}(u)$$

with
$$b_1 \leq -\sum_{k \in \{0,\infty\} \cup \mathbb{N} \setminus \{1\}}^{\infty} |b_k| R^{k-1}$$
 for some $R \geq 1$.

Examples of Hölder drift

Consider the Wright-Fisher SHE with Hölder continuous drift:

(2)
$$\begin{cases} \partial_t u = \frac{\Delta}{2} u + u^{\alpha} (1 - u) + \sqrt{u(1 - u)} \dot{W}, & x \in \mathbb{R}, t \ge 0, \\ u_0 = f, & x \in \mathbb{R}. \end{cases}$$

According to	uniqueness in law holds for (2) provided		
Shiga (1988) or	. 1		
Athreya-Tribe (2000)	$\alpha = 1$		
Mueller-Mytnik-Ryzhik (2021)	$lpha \in [rac{1}{2},1]$ and		
	f has compact interface		
Barnes-Mytnik-S. (2025+)	$\alpha \in (0,1]$		

This is expected, since the uniqueness in law holds for the SDE

$$\mathrm{d} X_t = X_t^{\alpha}(1-X_t)\mathrm{d} t + \sqrt{X_t(1-X_t)}\mathrm{d} B_t; \quad X_0 = x \in [0,1].$$

What if " $\alpha=0$ "? Pay attention that $u^{\alpha}(1-u)$ converges to the discontinuous drift $(1-u)-\mathbf{1}_{\{0\}}(u)$ when $\alpha\downarrow 0$.

Weak existence is non-trivial for discontinous drifts

- Denote by $u^{(\alpha)}$ the solution to the SPDE (2) with parameter α .
- It is standard to verify that the family of random elements $\{u^{(\alpha)}: \alpha \in (0,1]\}$ is tight in the path space $\mathcal{C}(\mathbb{R}_+, \mathcal{C}(\mathbb{R},[0,1]))$.
- By Skorohod's embedding, we can assume WLOG that there exists a sequence $\alpha_n \downarrow 0$ such that almost surely $u := \lim_{n \to \infty} u^{(\alpha_n)}$ exists in $\mathcal{C}(\mathbb{R}_+, \mathcal{C}(\mathbb{R}, [0, 1]))$.
- However, $z_n \to z$ in [0,1] does not necessarily imply that $z_n^{\alpha_n} \to z^0 = \mathbf{1}_{(0,1]}(z)$. For a counter example, consider $z_n := \exp(-\frac{\log 2}{\alpha_n}) \to 0 =: z$, but $z_n^{\alpha_n} = \frac{1}{2} \not\to \mathbf{1}_{(0,1]}(z) = 0$.
- So, from the standard "martingale problem argument", it is not clear if u solve the SPDE

$$\begin{cases} \partial_t u = \frac{\Delta}{2}u + (1-u) - \mathbf{1}_{\{0\}}(u) + \sqrt{u(1-u)}\dot{W}, & x \in \mathbb{R}, t \ge 0, \\ u_0 = f, & x \in \mathbb{R}. \end{cases}$$

Examples for discontinuous drifts

Nevertheless, our result implies the following:

Barnes-Mytnik-S. (2025+, to appear in Probab. Theory Related Fields)

For each $\delta \in [-1,1]$, the weak existence and uniqueness in law hold for the SPDE

(3)
$$\begin{cases} \partial_t u = \frac{\Delta}{2} u + (1-u) + \delta \mathbf{1}_{\{0\}}(u) + \sqrt{u(1-u)} \dot{W}, & x \in \mathbb{R}, t \ge 0, \\ u_0 = f, & x \in \mathbb{R}. \end{cases}$$

We can also show that δ is a relevant parameter!

Barnes-Mytnik-S. (2025+, to appear in Probab. Theory Related Fields)

Suppose that $f \not\equiv 1$. The distributions of the solution to the SPDE (3) are different for each $\delta \in [-1,1]$.

Examples for discontinuous drifts

This is drastically different from the SDE

(4)
$$dX_t = [(1 - X_t) + \delta \mathbf{1}_{\{0\}}(X_t)]dt + \sqrt{X_t(1 - X_t)}dB_t.$$

where δ is basically irrelevant.

Simple fact

When $\delta=-1$, the uniquess in law does not hold for the SDE (4). When $\delta\in(-1.1]$, the uniquess in law does hold, but the distributions of the solution to the SDE (4) are the same for different $\delta\in(-1,1]$.

Insight: The Wright-Fisher noise has a very different regularizing effect in the SPDE setting compared to the SDE setting!

Overview

Here is an overview of what we know about the 1-D SHE

$$\partial_t u_t = \frac{\Delta}{2} u_t + b(u) + \sigma(u) \dot{W}.$$

$\sigma(u)$ $b(u)$	0	Lipschitz	Hölder	Discontinuous	Measurable
deterministic: 0	well-posed	well-posed	non-uniqueness for some drift	non-uniqueness for some drift	non-uniqueness for some drift
additive: 1	well-posed	well-posed	well-posed	well-posed	well-posed
Lipschitz & non-degenerate	well-posed	well-posed	well-posed	well-posed	well-posed
3/4-Hölder & non-degenerate	well-posed	weakly well-posed	weakly well-posed	?	?
Feller noise: \sqrt{u}	weakly well-posed	weakly well-posed for some drift	?	?	?
Wright-Fisher noise: $\sqrt{u(1-u)}$	weakly well-posed	weakly well-posed	weakly well-posed for some drift	weakly well-posed for some drift	?

Duality method

• We say two Markov processes $(X_t)_{t\geq 0}$ and $(Y_t)_{t\geq 0}$ are dual to each other if there exists a large class of functions H(x,y) such that

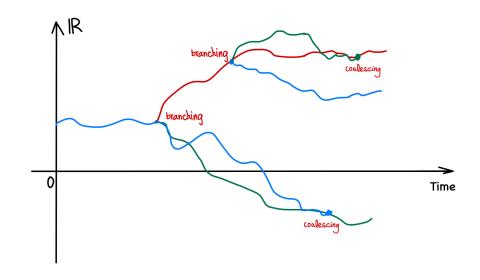
$$\mathbb{E}[H(X_t, Y_0)] = \mathbb{E}[H(X_0, Y_t)].$$

- Bachelier (1900, Ann. Sci. École Norm. Sup.): Brownian motion and the heat equation $\partial_t h = \frac{\Delta}{2}h$.
- McKean (1975, Comm. Pure Appl. Math.): Branching Brownian motion and the FKPP equation $\partial_t v = \frac{\Delta}{2}v + v(1-v)$.
- Harris (1978, Ann. Probab.):
 Coalescing random walk and the voter model.
- Shiga (1986, Math. Appl.): LCBM and the stochastic FKPP equation $\partial_t v = \frac{\Delta}{2} v + \sqrt{v(1-v)} \dot{W}$.
- Tóth-Werner (1998, Probab. Theory Relat. Fields): (Hard) Coalescing Brownian motions and itself.
- **Folklore**: Stochastic heat equation $\partial_t u = \frac{\Delta}{2} u + u \dot{W}$ and itself.
- . . .

The Dual of the Wright-Fisher SHEs

- The dual of Wright-Fisher SHEs are coalescing-branching Brownian motions (CBBMs).
- Two parameters:
 - Branching rate $\mu > 0$.
 - Offspring distribution $(p_k)_{k \in \{0,\infty\} \cup \mathbb{N}}$.
- Three dynamics:
 - Spatial movement: Particle move as independent Brownian motions.
 - Branching: Each particle branches into a random number of particles with the rate μ . The offspring number is sampled according to the distribution $(p_k)_{k \in \{0,\infty\} \cup \mathbb{N}}$.
 - *Coalescing:* Each pair of particles coalesces as one particle with rate 1/2 according to their intersection local time.

An illustration of the dual particle system



Explosion in CBBM

 To build a duality relation between CBBMs and the Wright-Fisher SHEs, we take

$$\mu:=\sum_{k\in\{0,\infty\}\cup\mathbb{N}\setminus\{1\}}^\infty |b_k|$$

and $p_1 := 0$, $p_k := |b_k|/\mu$ for $k \in \{0, \infty\} \cup \mathbb{N} \setminus \{1\}$.

• The dynamic is well-defined up to the explosion time

$$au_{\infty} := \lim_{n \to \infty} \inf\{t \ge 0 : \# \mathsf{particles} \ge n\}.$$

- (b_k) satisfies AT's condition $\implies p_{\infty} = 0$ and (p_k) has exponential moment $\implies \tau_{\infty} = \infty$ a.s.
- If AT's condition does not hold (especially when $p_{\infty}=|b_{\infty}|/\mu>0$) the explosion might happen in finite time.
- The definition of the particle system needs more justification!

Coming down from infinity

- A coalescing Brownian motion (CBM) is CBBM with $p_1 = 1$.
- Define CBM with infinitely many initial particles as the limit of a sequence of CBMs with finite initial particles.
- Denote by $Z_t(A)$ the number of particles in a domain A at time t of a CBM with infinitely many initial particles, i.e. $Z_0(\mathbb{R}) = \infty$.

Barnes-Mytnik-S. (2024, Ann. Probab.)

The total population $Z_t(\mathbb{R}) < \infty$ for every t > 0 $\iff Z_0(\cdot)$ is compactly supported.

Moreover, in this case

$$\left(\int v_t(x)\mathrm{d}x\right)^{-1}Z_t(\mathbb{R})\xrightarrow{L^1}1,\quad t\downarrow 0$$

where $(v_t(x))_{t\geq 0,x\in\mathbb{R}}$ is the unique non-negative solution to the 1d PDE $\partial_t v_t = \frac{\Delta}{2} v_t - v_t^2/2$ with initial value $v_0 = Z_0$.

Reflecting from infinity

- Similarly, we can justify the definition of the CBBM for arbitrary offspring distribution (allowing $p_{\infty} > 0$).
- It is defined as the limit of a sequence of CBBMs with truncated offspring distributions.
- Denote by $X_t(\mathbb{R})$ the total population of a CBBM with arbitrary branching rate and arbitrary offspring distribution.

```
Barnes-Mytnik-S. (2025+, to appear in Probab. Theory Related Fields)
```

If $X_0(\mathbb{R}) < \infty$, then $X_t(\mathbb{R})$ is "reflecting" from ∞ .

• This "reflecting from infinity" property of the dual particle system is the key to the well-posedness of the corersponding SPDE.

Thanks!