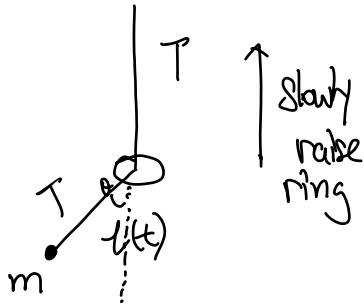


Adiabatic Invariance: (Slow is the key word).



→ There is upward force by string on ring.

→ Displacement Δz is up, so string does work, so string's energy goes down

Since $\omega = \sqrt{\frac{g}{l(t)}}$ and if ring increase height, then $l(t)$ increases.

so $\omega(t)$ goes down.

By Newton:

$$F_z = T - T \cos \theta = T \frac{\theta^2}{2} \propto E$$

$$\overline{F_z} = T \frac{\overline{\theta^2}}{2} = \frac{1}{2} m g \overline{\theta^2} = \frac{\overline{E}(t)}{2 l(t)}$$

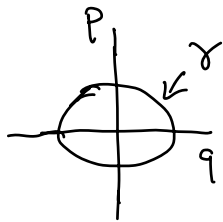
$$\overline{W} = -\frac{d\overline{E}}{dt} = \overline{F_z} \frac{dl}{dt} = \frac{\overline{E}}{2l} \frac{dl}{dt}$$

$$\text{then } \frac{d(\overline{E} l)}{dt} = 0 \quad \text{so} \quad \frac{d(\overline{E})}{dt} = 0$$

In this problem:

Hamiltonian $\Rightarrow H(q, p, \lambda(t))$ ↖ slow function of time.

If we have fixed energy, we have a closed orbit in phase space.



$$\text{Then } p_r = p_r(q, E, \lambda)$$

Determined by

$$E = H(q, p_r, \lambda)$$

Define a quantity? I , adiabatic invariant.

$$\star \left[\begin{array}{l} I(E, \lambda) = \int dq dp \quad \leftarrow \text{Area is constant due to Liouville Theorem.} \\ I = \oint_{\gamma} p_r dq \quad \leftarrow \text{Stokes Theorem, convert from area to line integral.} \end{array} \right]$$

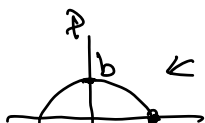
If $E(t)$ and $\lambda(t)$, both function of time

then

$$\frac{d}{dt} \left(\oint_{\gamma} p_r dq \right) = 0$$

Ex: For SHO

$$E = \frac{p^2}{2} + \frac{\omega^2}{2} q^2 \quad \leftarrow \text{Determines } p_r(E)$$



\leftarrow Ellipse

$$b = \sqrt{2E} \quad \text{When } q=0$$

$$\int_a^b \frac{1}{q}$$

$$a = \sqrt{\frac{2E}{\omega^2}} \text{ when } p=0$$

$$\text{then } I = \int dq dp = \pi ab = \pi \frac{2E(t)}{\omega(t)}$$

Proof:

$$\frac{dE}{dt} = \frac{\partial H}{\partial t} = \left(\frac{\partial H}{\partial \lambda} \right)_p \frac{d\lambda}{dt}$$

Need an average rate of change. with slow varying $\lambda(t)$

$$\overline{\frac{dE}{dt}} = \overline{\left(\frac{\partial H}{\partial \lambda} \right)_p} \frac{d\lambda}{dt}$$

← average over 1-cycle of orbit.

$$\hookrightarrow \frac{\Delta E}{\Delta t} = \overline{\left(\frac{\partial H}{\partial \lambda} \right)_p} \frac{\Delta \lambda}{\Delta t}$$

since $\frac{d\lambda}{dt}$ is small, only work with first order

$$\text{so } \frac{d\lambda}{dt} \approx \frac{\Delta \lambda}{\Delta t}$$

$$\overline{\left(\frac{\partial H}{\partial \lambda} \right)_p} = \frac{1}{T} \oint \left(\frac{\partial H}{\partial \lambda} \right)_p dt$$

$$\text{For } dt = \frac{dq}{\dot{q}} = \frac{dq}{\frac{\partial H}{\partial p}}$$

$$\frac{\Delta E}{\Delta t} = \frac{\oint \left(\frac{\partial H}{\partial \lambda} \right)_p \left(\frac{\partial p}{\partial \lambda} \right)_x dq \frac{\Delta \lambda}{\Delta t}}{\oint \left(\frac{\partial p}{\partial \lambda} \right)_x dq}$$

$$\text{Since } \left(\frac{\partial E}{\partial \lambda} \right) = \frac{\partial H}{\partial \lambda} + \frac{\partial H}{\partial p} \frac{\partial p}{\partial \lambda} \quad \dots$$

$$\left(\frac{\partial H}{\partial E} \right) \frac{\partial E}{\partial \lambda} + \left(\frac{\partial H}{\partial \lambda} \right) \frac{\partial \lambda}{\partial E} = 0$$

then
$$\frac{\partial P}{\partial \lambda} = - \frac{\partial H / \partial \lambda}{\partial H / \partial P}$$

then
$$\frac{\Delta E}{\Delta t} = \frac{- \oint \frac{\partial P}{\partial \lambda} dq}{\oint \left(\frac{\partial P}{\partial E} \right) dq} \frac{\Delta \lambda}{\Delta t}$$

$$\oint \underbrace{\left(\left(\frac{\partial P}{\partial E} \right) \frac{\Delta E}{\Delta t} + \left(\frac{\partial P}{\partial \lambda} \right) \frac{\Delta \lambda}{\Delta t} \right)}_{\oint \frac{\Delta P}{\Delta t} dq} dq = 0.$$

So $I = \oint P dq$ is invariant.

Ex 2:

A slowly change magnetic field:

Consider the circular orbits in xy plane with $x > 0$ of a particle mass m and charge q , in a constant field B in \hat{z} ,

a) use Hamilton formulation to determine the radius

and angular frequency of circular orbit.

use gauge: $\vec{A} = B(0, x, 0)$

$$L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) + \frac{q}{c} \vec{r} \cdot \vec{A}$$

$$= \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) + \frac{q}{c} B x \dot{y}$$

$$P_x = \frac{\partial L}{\partial \dot{x}} = m \dot{x}$$

$$P_y = \frac{\partial L}{\partial \dot{y}} = m \dot{y} + m \omega_B x \Rightarrow \dot{y} = \frac{P_y}{m} - \omega_B x$$

$$H = P_x \dot{x} + P_y \dot{y} - L$$

$$= \frac{P_x^2}{m} + P_y \left(\frac{P_y}{m} - \omega_B x \right) - \left(\frac{P_x^2}{2m} + \frac{m}{2} \left(\frac{P_y}{m} - \omega_B x \right)^2 \right) - m \omega_B x \left(\frac{P_y}{m} - \omega_B x \right)$$

$$H = \frac{1}{2m} P_x^2 + \frac{P_y^2}{2m} - \omega_B x P_y + \cancel{P_y \omega_B x} - \frac{m}{2} \omega_B^2 x^2 - \omega_B x P_y + m \omega_B^2 x^2$$

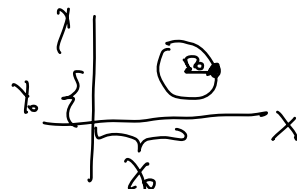
$$H = \frac{P_x^2}{2m} + \frac{P_y^2}{2m} - \omega_B x P_y + \frac{1}{2} m \omega_B^2 x^2$$

$$-\frac{\partial H}{\partial x} = \dot{P}_x = m \ddot{x} = \omega_B P_y - m \omega_B^2 x$$

$$-\frac{\partial H}{\partial y} = \dot{P}_y = 0 \Rightarrow P_y \text{ is constant.}$$

$$\frac{\partial H}{\partial P_x} = \dot{x} = \frac{P_x}{m}$$

$$\frac{\partial H}{\partial P_y} = \dot{y} = \frac{P_y}{m}$$



$$H = \frac{P_x^2}{2m} + \frac{P_y^2}{2m} - W_B x P_y + \frac{1}{2} m W_B^2 x^2$$

$$= \frac{1}{2m} P_x^2 + \frac{1}{2} m W_B^2 \left(x^2 - \frac{2}{m W_B} x P_y + \frac{P_y^2}{m^2 W_B^2} \right)$$

$$H = \frac{1}{2m} P_x^2 + \frac{1}{2} m W_B^2 \left(x - \frac{P_y}{m W_B} \right)^2$$

a shift of origin
in potential.

$$P_y = \frac{\partial L}{\partial \dot{y}} = m \dot{y} + m W_B x = \text{constant}$$

when $x = x_0, \dot{y} = 0$

So $P_y = m W_B x_0 \Rightarrow \frac{P_y}{m W_B} = x_0$

We know at $x = x_0 + r_0, P_x = 0$

So $E_{\text{tot}} = H(x = x_0 + r_0) = \frac{1}{2} m W_B^2 r_0^2$

So $r_0 = \sqrt{\frac{2 E_{\text{tot}}}{m W_B^2}}$

$$m \ddot{x} = W_B P_y - m W_B^2 x$$

$$\ddot{x} + W_B^2 x - \frac{W_B}{m} P_y = 0$$

↓
constant.

So $x = A \cos(W_B t) + \frac{W_B}{m} P_y$

↑

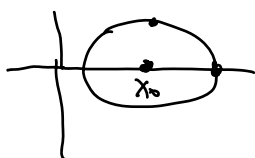
Clearly angular frequency is ω_B

b) How does radius and the center of the circular orbit changes as $B(t)$ slowly increased

$$I = \oint p_y dq.$$

know $p_y = m\dot{y} + m\omega_B x$ is conserved.

$$\text{know } H = \frac{1}{2m} p_x^2 + \frac{1}{2} m\omega_B^2 \left(x - \frac{p_y}{m\omega_B} \right)^2 = \frac{1}{2} m\omega_B^2 r_0^2 = E$$



$$\begin{array}{l} \text{when } x = x_0 \\ E = \frac{1}{2} m\omega_B^2 r_0^2 = \frac{p_{x_0}^2}{2m} \\ p_{x_0}^2 = m^2 \omega_B^2 r_0^2 \\ p_{x_0} = m\omega_B r_0 = \sqrt{2mE} \end{array} \quad \left| \begin{array}{l} \text{when } x = x_0 + r_0 \\ E = \frac{1}{2} m\omega_B^2 r_0^2 \\ \sqrt{\frac{2E}{m\omega_B^2}} = r_0 \end{array} \right.$$

Then

$$\begin{aligned} I &= \int dq dp \frac{1}{2\pi} \\ &= \pi \sqrt{2mE} \sqrt{\frac{2E}{m\omega_B^2}} \frac{1}{2} \\ I &= \frac{2\pi E}{\omega_B} \end{aligned}$$

$$I = \frac{\frac{1}{2} m\omega_B^2 r^2}{\omega_B} \propto \omega_B r^2 \propto B r^2$$

$$\sqrt{\frac{2I}{m\omega_B}} = r(t)$$

$$I_0 = B r_0^2 = B r^2 \Rightarrow r = \sqrt{\frac{B_0}{B}} r_0$$

When B increase r^2 should decrease proportionally
to let I remain constant.

Also we $P_y = m \omega_b X_0$ to be constant

$$\text{or } X = \frac{P_y}{m \omega_B}$$

So if B increase X_0 also decrease proportionally.
to remain constant.