

We have:

$$m \frac{d^2 q}{dt^2} + m\eta \frac{dq}{dt} + m\omega^2 q = F = f_0 \cos \omega t$$

When near resonance, $\omega \sim \omega_0$, nonlinearity becomes important.

We have eq: $\ddot{x} + \eta \dot{x} + \omega^2 x + \beta x^3 = \frac{f(t)}{m}$

First assume $\beta=0$, work with linear case

then $\ddot{x} + \eta \dot{x} + \omega^2 x = \frac{f(t)}{m} = \frac{f_0}{m} \cos \omega t$

let $x^{(0)} = a(t) \cos(-\omega t + \phi(t))$

$$\dot{x}^{(0)} = \dot{a} \cos(-\omega t + \phi) - a \sin(-\omega t + \phi) \{-\omega + \dot{\phi}\}$$

$$\begin{aligned} \ddot{x}^{(0)} = & \ddot{a} \cos(-\omega t + \phi) - \dot{a} \sin(-\omega t + \phi) \{-\omega + \dot{\phi}\} \\ & - \dot{a} \sin(-\omega t + \phi) \{-\omega + \dot{\phi}\} - a \{\cos(-\omega t + \phi) (-\omega + \dot{\phi})^2 \\ & + \ddot{\phi} \sin(-\omega t + \phi)\} \end{aligned}$$

$$\begin{aligned} \downarrow \quad & = -2\dot{a} \sin(-\omega t + \phi) - \omega^2 a \cos(-\omega t + \phi) \\ & + 2a\omega \dot{\phi} \cos(-\omega t + \phi) + \mathcal{O} \end{aligned}$$

Assume $\dot{a} = \dot{\phi} = 0$

$$\ddot{x} + \gamma \dot{x} + \omega_0^2 x = \frac{f_0}{m} \cos(\omega t)$$

$$(-\omega^2 + \omega_0^2) a \cos(-\omega t + \phi) + \gamma \omega a \sin(-\omega t + \phi) = \frac{f_0}{m} \cos(\omega t)$$

$$\begin{aligned} \text{let } \omega &= \omega_0 + (\omega - \omega_0) \\ \hookrightarrow \omega^2 &= \omega_0^2 + (\omega - \omega_0)^2 + 2\omega_0(\omega - \omega_0) \end{aligned}$$

$$\left[\underbrace{(\omega - \omega_0)^2}_{\downarrow} - 2\omega_0(\omega - \omega_0) \right] a \cos(-\omega t + \phi) + \gamma \omega a \sin(-\omega t + \phi) = \frac{f_0}{m} \cos(\omega t)$$

↑
let $\omega \approx \omega_0$

Since $\omega - \omega_0 = \mathcal{O}(\epsilon)$
ignore $\mathcal{O}(\epsilon^2)$

$$\hookrightarrow -2\omega_0(\omega - \omega_0) a \cos(-\omega t + \phi) + \gamma \omega_0 a \sin(-\omega t + \phi) = \frac{f_0}{m} \cos(\omega t)$$

$$\hookrightarrow -(\omega - \omega_0) a \cos(-\omega t + \phi) + \frac{\gamma}{2} a \sin(-\omega t + \phi) = \frac{f_0}{2m\omega_0} \cos(\omega t)$$

$$\text{let } r = \sqrt{(\omega - \omega_0)^2 + (\gamma/2)^2}$$

$$\begin{aligned} -(\omega - \omega_0) &= r \cos \phi_0 \\ \gamma/2 &= r \sin \phi_0 \end{aligned}$$

$$r \cos \phi_0 \cos(-\omega t + \phi_0) + r \sin \phi_0 \sin(-\omega t + \phi_0) = \frac{f_0}{2m\omega_0} \cos \omega t$$

$$\text{gel } a = \frac{f_0}{2m\omega_0} \frac{1}{[(\omega - \omega_0)^2 + (\frac{\gamma}{2})^2]^{1/2}} \Rightarrow a_{\max} = \frac{f_0}{m\omega_0 \gamma}$$

$$\tan \phi_0 = \frac{\gamma/2}{\omega - \omega_0}$$

$$\omega_0 - \omega$$

For nonlinear oscillator: $\ddot{x} + \gamma \dot{x} + \omega_0^2 x + \beta x^3 = \frac{f(t)}{m}$

$$\Delta\omega = \frac{3}{8} \frac{\beta a^2}{\omega_0} = K a^2$$

↗ frequency shift depends on amplitude.

Therefore, nonlinearity becomes important when

$$\begin{aligned} \Delta\omega &\sim \eta/2 \\ \hookrightarrow K a_{\max}^2 &\sim \eta/2 \\ \left(\frac{K}{\eta/2}\right) \left(\frac{f_0}{m\omega\eta}\right)^2 &\sim 1 \end{aligned}$$

Now for $\ddot{x} + \gamma \dot{x} + \omega_0^2 x + \beta x^3 = \frac{f(t)}{m}$

use $x = a \cos(\omega t + \phi)$

To zeroth order, find

$$\begin{aligned} \beta x^3 &= 2\omega_0 \Delta\omega a \cos(\omega t + \phi) \\ &\stackrel{!}{=} 2\omega_0 K a^2 \cos(\omega t + \phi) \end{aligned}$$

Then

$$\begin{aligned} -2\omega_0(\omega - \omega_0)a \cos(\omega t + \phi) + \gamma\omega_0 a \sin(\omega t + \phi) + 2\omega_0 \Delta\omega a \cos(\omega t + \phi) \\ = \frac{f_0}{m} \cos(\omega t + \phi) \end{aligned}$$

$$\hookrightarrow = [-(\omega - \omega_0 - k a^2) a \cos(-\omega t + \phi) + \frac{\gamma}{2} \sin(\omega t + \phi)] = \frac{f_0}{2m\omega_0} \cos(-\omega t + \phi)$$

$$\hookrightarrow a \gamma \underbrace{[\cos \phi \cos(-\omega t + \phi) + \sin \phi \sin(-\omega t + \phi)]}_{\cos(\omega t)} = \frac{f_0 \cos \omega t}{2m\omega_0}$$

$$\gamma = \sqrt{(\omega - \omega_0 - k a^2)^2 + (\gamma/2)^2}$$

$$\text{then } \tan \phi = \frac{\gamma/2}{\omega + k a^2 - \omega}$$

$$a^2 \gamma^2 = \left(\frac{f_0}{2m\omega_0} \right)^2$$

$$a^2 (\omega - \omega_0 - k a^2)^2 + a^2 (\gamma/2)^2 = \left(\frac{f_0}{2m\omega_0} \right)^2$$

→
Solve to get amplitude.

$$\text{define: } \bar{f}^2 = \left(\frac{k}{\gamma/2} \right) \left(\frac{f_0}{m\omega_0 \gamma} \right)^2$$

$$\bar{a}^2 = \frac{k}{\gamma/2} a^2$$

$$\bar{\gamma} = \frac{\omega - \omega_0}{\gamma/2}$$

$$\bar{a}^2 (\bar{\gamma} - \bar{a}^2)^2 + \bar{a}^2 = \bar{f}^2$$

$$\text{so } \bar{f} \ll 1, \bar{a} \ll 1$$

$$\text{then } \bar{a}^2 \bar{\gamma}^2 + \bar{a}^2 = \bar{f}^2$$

$$\bar{a} = \frac{0}{\bar{\gamma}^2 + 1} \quad \leftarrow \text{simple harmonic oscillator result.}$$

For small $\bar{\gamma} \Rightarrow$ one real root.
For large $\bar{\gamma} \Rightarrow$ three real roots