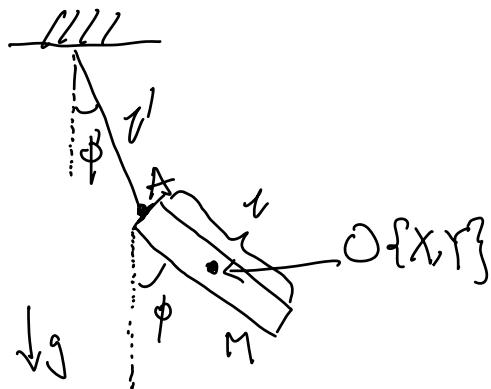


i) A bar on a thread



a) Find eom

$$X = l' \sin \phi' + \frac{l}{2} \sin \phi$$

$$Y = -\left(l' \cos \phi' + \frac{l}{2} \cos \phi\right)$$

$$\dot{X} = l' \cos \phi' \dot{\phi}' + \frac{l}{2} \cos \phi \dot{\phi}$$

$$\dot{Y} = l' \sin \phi' \dot{\phi}' + \frac{l}{2} \sin \phi \dot{\phi}$$

$$T = \frac{1}{2} M (\dot{X}^2 + \dot{Y}^2) + \frac{1}{2} I \dot{\phi}^2$$

$$V = MgY$$

$$L = \frac{1}{2} M \left(l'^2 \cos^2 \phi \dot{\phi}'^2 + \frac{l^2}{4} \cos^2 \phi \dot{\phi}^2 + l' l \cos \phi' \cos \phi \dot{\phi}' \dot{\phi} \right.$$

$$\left. + l'^2 \sin^2 \phi \dot{\phi}'^2 + \frac{l^2}{4} \sin^2 \phi \dot{\phi}^2 + l' l \sin \phi' \sin \phi \dot{\phi}' \dot{\phi} \right)$$

$$+ \frac{1}{2} I \dot{\phi}^2 + Mg \left(l' \cos \phi + \frac{l}{2} \cos \phi \right)$$

$$\underline{L} = \frac{1}{2} M \left(l'^2 \dot{\phi}'^2 + \frac{l^2}{4} \dot{\phi}^2 + l' l \dot{\phi}' \dot{\phi} (\cos \phi' \cos \phi + \sin \phi' \sin \phi) \right)$$

$$+\frac{1}{2}I_A\dot{\phi}^2 + Mg(l^2\cos\phi + \frac{l}{2}\cos\phi) \underbrace{\cos(\phi - \phi')}_1$$

$$\frac{d^2}{dt^2} = Ml^2\ddot{\phi}' + \frac{l^2M}{2}\cos(\phi - \phi')\dot{\phi}$$

$$\frac{d^2}{dt^2} = Ml^2\ddot{\phi}'\dot{\phi} \sin(\phi - \phi') - Mg l \sin\phi$$

$$\frac{d}{dt} \left(Ml^2\dot{\phi}' \frac{Ml^2}{2}\cos(\phi - \phi')\dot{\phi} \right) = \frac{Ml^2}{2}\dot{\phi}'\dot{\phi} \sin(\phi - \phi') - Mg l \sin\phi$$

$$Ml^2\ddot{\phi}' + \frac{Ml^2}{2}\cos(\phi - \phi')\ddot{\phi} - \frac{Ml^2}{2}\sin(\phi - \phi')\dot{\phi}(\dot{\phi} - \dot{\phi}')$$

$$= \frac{Ml^2}{2}\dot{\phi}'\dot{\phi} \sin(\phi - \phi') - Mg l \sin\phi$$

$$\boxed{Ml^2\ddot{\phi}' + \frac{1}{2}Ml^2\cos(\phi - \phi')\ddot{\phi} - \frac{1}{2}Ml^2\sin(\phi - \phi')\dot{\phi}^2 + Mg l \sin\phi = 0}$$

$$\frac{d^2}{dt^2} = \frac{Ml^2}{4}\dot{\phi}' + \frac{1}{2}Ml^2\cos(\phi - \phi')\dot{\phi}' + I_A\dot{\phi}$$

$$= I_A\dot{\phi} + l^2e\cos(\phi - \phi')\dot{\phi}' \frac{1}{2}M$$

$$\frac{d^2}{dt^2} = -\frac{1}{2}Ml^2\dot{\phi}'\dot{\phi} \sin(\phi - \phi') - \frac{Mgl}{2}\sin\phi$$

$$\frac{d}{dt} \left(I_A\dot{\phi} + Ml^2\dot{\phi}'\cos(\phi - \phi')\dot{\phi}' \right) = -Ml^2\dot{\phi}'\dot{\phi} \sin(\phi - \phi') - Mg l \sin\phi$$

$$J_A \ddot{\phi} + \frac{1}{2} M I_C' \cos(\phi - \phi') \ddot{\phi}' - \frac{1}{2} M I_C' \sin(\phi - \phi') \dot{\phi} (\dot{\phi} - \dot{\phi}') = 0$$

$$J_A \ddot{\phi} + \frac{1}{2} M I_C' \cos(\phi - \phi') \ddot{\phi}' - \frac{1}{2} M I_C' \sin(\phi - \phi') \dot{\phi} (\dot{\phi} - \dot{\phi}') \\ = -\frac{1}{2} M I_C' \dot{\phi}' \dot{\phi} \sin(\phi - \phi') - \frac{M g l}{2} \sin \phi$$

$$\boxed{J_A \ddot{\phi} + \frac{1}{2} M I_C' \cos(\phi - \phi') \ddot{\phi}' + \frac{1}{2} M I_C' \sin(\phi - \phi') \dot{\phi}^2 + \frac{M g l}{2} \sin \phi = 0}$$

(i) Linearize around $\phi = \phi' = 0$ to first order.

$$\sin \phi' = \phi \\ \cos(\phi - \phi') = 1 - \frac{(\phi - \phi')^2}{2} =$$

$$M I_C' \ddot{\phi} + \frac{1}{2} M I_C' \cos(\phi - \phi') \ddot{\phi}' - \frac{1}{2} M I_C' \sin(\phi - \phi') \dot{\phi}^2 + M g l' \sin \phi = 0$$

$$\hookrightarrow \boxed{M I_C'^2 \ddot{\phi}' + \frac{1}{2} M I_C' \ddot{\phi} + M g l' \dot{\phi}' = 0}$$

$$J_A \ddot{\phi} + \frac{1}{2} M I_C' \cos(\phi - \phi') \ddot{\phi}' + \frac{1}{2} M I_C' \sin(\phi - \phi') \dot{\phi}^2 + \frac{M g l}{2} \sin \phi = 0$$

$$\hookrightarrow \boxed{J_A \ddot{\phi} + \frac{1}{2} M I_C' \ddot{\phi}' + \frac{M g l}{2} \phi = 0}$$

c) Find eigenfrequencies of small oscillations near equilibrium.

$$I_A \ddot{\phi} + \frac{Mg\ell}{2} \phi + \frac{1}{2} M\ell^2 \dot{\phi} = 0$$

$$M\ell'^2 \ddot{\phi}' + Mg\ell' \phi' + \frac{1}{2} M\ell' \dot{\phi}' = 0$$

$$\begin{pmatrix} I_A & \frac{1}{2} M\ell^2 \\ \frac{1}{2} M\ell^2 & M\ell'^2 \end{pmatrix} \begin{pmatrix} \ddot{\phi} \\ \ddot{\phi}' \end{pmatrix} + \begin{pmatrix} \frac{Mg\ell}{2} & 0 \\ 0 & Mg\ell' \end{pmatrix} \begin{pmatrix} \phi \\ \phi' \end{pmatrix} = 0$$

$$\text{Let } \ddot{\Phi} = A e^{j\omega t}$$

$$\left| -\omega^2 \begin{pmatrix} I_A & \frac{1}{2} M\ell^2 \\ \frac{1}{2} M\ell^2 & M\ell'^2 \end{pmatrix} + \begin{pmatrix} \frac{Mg\ell}{2} & 0 \\ 0 & Mg\ell' \end{pmatrix} \right| = 0$$

$$(-\omega^2 I_A + \frac{Mg\ell}{2})(-\omega^2 M\ell'^2 + Mg\ell') - \frac{1}{4} M^2 \ell^2 \ell'^2 \omega^4 = 0$$

$$\omega^4 \frac{1}{3} \ell^2 M\ell'^2 - \omega^2 Mg\ell - \omega^2 \frac{Mg\ell}{2} M\ell'^2 + \frac{M^2 g^2 \ell^2}{2} - \frac{1}{4} M^2 \ell^2 \ell'^2 \omega^4 = 0$$

$$\frac{1}{12} M^2 \ell^2 \ell'^2 \omega^4 - \omega^2 \left(\frac{M^2 g}{3} \left(\frac{1}{3} \ell^2 \ell'^2 + \frac{1}{2} \ell^2 \ell'^2 \right) + \frac{M^2 g^2 \ell^2}{2} \right) = 0$$

$$\omega^4 - \omega^2 \left(\frac{1}{3} \ell^2 \ell'^2 + \frac{1}{2} \ell^2 \ell'^2 \right) + \frac{M^2 g^2 \ell^2}{21} = 0$$

$$\omega^4 - \omega^2 \left(4\Omega'^2 + 4\Omega^2 \right) + 4\Omega^2 \Omega'^2 = 0$$

$$\begin{aligned} \frac{3}{2} \Omega' &= \Omega^2 \\ \Omega' &= \Omega^2 \end{aligned}$$

$$\omega^2 = \frac{4(\Omega'^2 + \Omega^2) \pm \sqrt{16(\Omega'^2 + \Omega^2)^2 - 16\Omega'^2 \Omega^2}}{2}$$

$$\boxed{\omega_{\pm} \doteq 2(\Omega^2 + \omega^2) \pm \sqrt{(\Omega^2 + \omega^2)^2 - \omega^2 \Omega'^2}}$$

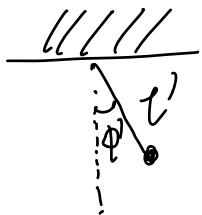
d) i) Discuss motion in limit $\ell \ll \ell'$

If $\ell \ll \ell'$: then there is no ϕ eq goes away.

$$\cancel{\frac{1}{3}M\ell^2\ddot{\phi} + \frac{1}{2}M\ell\ell'\ddot{\phi}' + \frac{Mgl}{2}\phi = 0}$$

$$M\ell'^2\ddot{\phi}' + \cancel{\frac{1}{2}M\ell\ell'\ddot{\phi}} + Mgl\phi' = 0$$

$$\ddot{\phi}' + \frac{g}{\ell'}\phi' = 0 \quad \leftarrow \text{a simple pendulum with } \omega^2 = \frac{g}{\ell'} = \Omega'^2$$

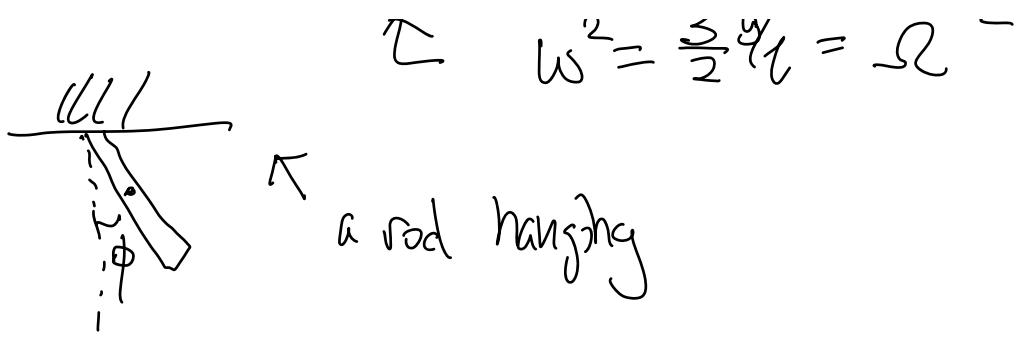


For $\ell' \ll \ell$:

we have

$$\cancel{\frac{1}{3}M\ell^2\ddot{\phi} + \frac{1}{2}M\ell\ell'\ddot{\phi}'} + \frac{Mgl}{2}\phi = 0$$

$$\ddot{\phi} + \frac{3}{2}\frac{g}{\ell}\phi = 0$$



e) find eigenmodes for $\ell = \ell'$

$$\omega_{\pm}^2 = 2(\Omega^2 + \Omega'^2) \pm 2\sqrt{(\Omega'^2 + \Omega^2)^2 - \Omega^2 \Omega'^2}$$

$$\hookrightarrow \omega_{\pm}^2 = 2\left(\frac{5}{2}\Omega'^2\right) \pm 2\sqrt{\left(\frac{5}{2}\Omega'^2\right)^2 - \frac{3}{2}\Omega'^4}$$

$$= 5\Omega'^2 \pm \Omega'^2 \sqrt{\frac{25}{4} - \frac{3}{2}}$$

$$= 5\Omega'^2 \pm \Omega'^2 \sqrt{\frac{19}{4}}$$

$$\omega_{\pm}^2 = (5 \pm \sqrt{19})\Omega'^2$$

$$\text{For } \omega^2 = \omega_{\pm}^2 \quad \ell = \ell'$$

$$\left\{ -(5 \pm \sqrt{19}) \begin{pmatrix} I_x & \frac{1}{2} M \ell^2 \\ \frac{1}{2} M \ell^2 & M \ell^2 \end{pmatrix} + \begin{pmatrix} \frac{M \ell^4}{2} & 0 \\ 0 & M g \ell \end{pmatrix} \right\} \begin{pmatrix} E_1^+ \\ E_2^+ \end{pmatrix} = 0$$

$$\left[\begin{array}{cc} -\left(5+\sqrt{19}\right)\Omega^2 & \frac{1}{2} \\ \frac{1}{2} & 1 \end{array} \right] + \left[\begin{array}{cc} \Omega^2 \frac{1}{2} & 0 \\ 0 & \Omega^2 \end{array} \right] \left[\begin{array}{c} E_1^+ \\ E_2^+ \end{array} \right] = 0$$

$$\left[\begin{array}{cc} \left(-\frac{\left(5+\sqrt{19}\right)}{3} + \frac{1}{2} \right) \Omega^2 & -\frac{\left(5+\sqrt{19}\right)}{2} \Omega^2 \\ -\frac{\left(5+\sqrt{19}\right)}{2} \Omega^2 & \left(-\left(5+\sqrt{19}\right) + 1 \right) \Omega^2 \end{array} \right] \left[\begin{array}{c} E_1^+ \\ E_2^+ \end{array} \right] = 0$$

let $E_1^+ = 1$

$$-\frac{\left(5+\sqrt{19}\right)}{3} + \frac{1}{2} - \frac{\left(5+\sqrt{19}\right)}{2} E_2^+ = 0$$

$$-\left(\frac{7+2\sqrt{19}}{6}\right) = \left(\frac{5+\sqrt{19}}{2}\right) E_2^+$$

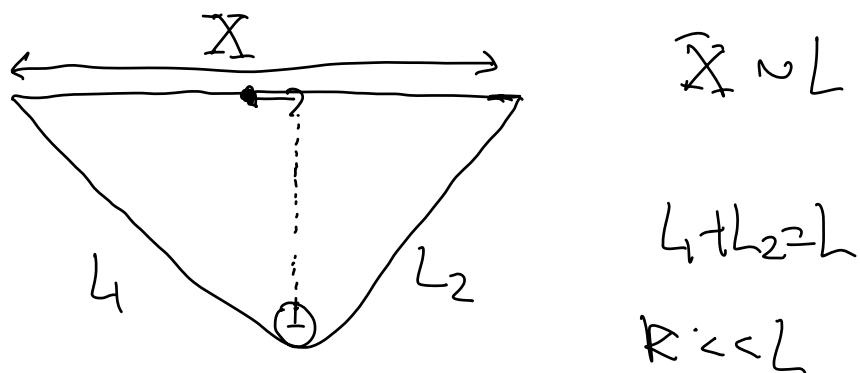
$$E_2^+ = -\left(\frac{7+2\sqrt{19}}{3(5+\sqrt{19})}\right)$$

$$= \frac{1}{7-2\sqrt{19}}$$

$$E_2 = \left[\frac{1}{3(\zeta - \sqrt{19})} \right] \quad E_1 = 1$$

$$\begin{pmatrix} \phi \\ \phi' \end{pmatrix} = A e^{-i\omega t} \begin{pmatrix} 1 \\ -\left(\frac{7+2\sqrt{19}}{3(\zeta + \sqrt{19})}\right) \end{pmatrix} + B e^{-i\omega t} \begin{pmatrix} 1 \\ -\left(\frac{7-2\sqrt{19}}{3(\zeta - \sqrt{19})}\right) \end{pmatrix}$$

2) Rolling of a spool:



a) know Ellipse:

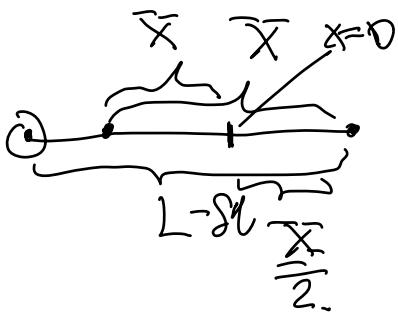
$$\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1$$

Determine a, b in terms of L and X

Show: $L_1 = \frac{L}{2} + \frac{X}{2} \times$

$$L_2 = \frac{L}{2} - \frac{1}{2}x$$

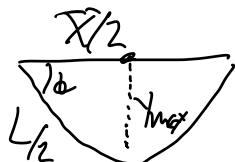
When $y=0$



$$x_{\max} = a = L - sL - \frac{x}{2}$$

Since $x \approx L$ $\frac{x}{2} \approx \frac{L}{2}$

When $x=0$ $y_{\max}^2 = b^2$



$$\left(\frac{x}{2}\right)^2 + y_{\max}^2 = (L_2)^2$$

$$y_{\max} = \sqrt{(L_2)^2 - \left(\frac{x}{2}\right)^2}$$

$$b = \frac{1}{2} \sqrt{L^2 - x^2}$$

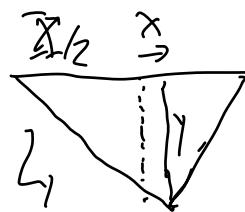
$$\frac{x^2}{\left(\frac{L}{2}\right)^2} + \frac{y^2}{\frac{1}{4}(L^2 - x^2)} = 1$$

$$y^2 = \left(1 - \frac{4x^2}{L^2}\right) \frac{1}{4}(L^2 - x^2)$$

$$= \left(\frac{1}{4} - \frac{x^2}{L^2}\right)(L^2 - x^2)$$

$$\boxed{\frac{1}{4} - \frac{x^2}{L^2} \quad | \quad L^2 - x^2 \quad | \quad L^2 - x^2 \quad |}$$

$$L = \sqrt{L^2} //$$



$$L_1^2 = (\cancel{x}/2 + x)^2 + y^2$$

$$\stackrel{!}{=} x^2 + \cancel{x}x + \frac{\cancel{x}^2}{4} + y^2$$

$$L_2^2 = (\cancel{x}/2 - x)^2 + y^2$$

$$\stackrel{!}{=} x^2 - \cancel{x}x + \frac{\cancel{x}^2}{4} + y^2$$

$$L_1^2 - L_2^2 = 2\cancel{x}x$$

$$L_1^2 - (L - L_1)^2 = L^2 + 2LL_1 = 2\cancel{x}x$$

$$L_1 = \frac{L}{2} + \frac{\cancel{x}}{L}x$$

$$L - L_1 = L_2$$

$$L_2 = \frac{L}{2} - \frac{\cancel{x}}{L}x$$

b) $T = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2) + \frac{1}{2}Iw^2$

$$I = \int dm \ r^2$$

$$\begin{aligned}
 &= \int \frac{M}{\pi R^2} r^2 r dr d\theta \\
 &= \frac{2\pi M}{\pi R^2} \frac{R^4}{4} \\
 I_{disk} &\stackrel{\perp}{=} \frac{MR^2}{2}
 \end{aligned}$$

$$V = MgY$$

$$\gamma = -\sqrt{\left(\frac{L^2}{4} - x^2\right)\left(\frac{L^2 - \bar{x}^2}{L^2}\right)}$$

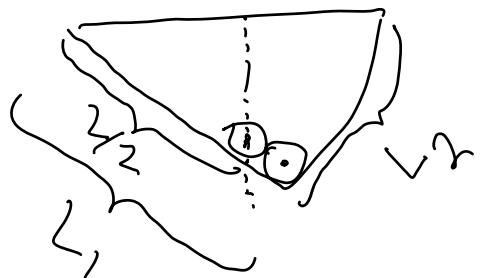
$$\dot{y} = \sqrt{\frac{L^2 - \bar{x}^2}{L^2}} \frac{\dot{x} \dot{\bar{x}}}{\sqrt{\frac{L^2}{4} - x^2}}$$

$$L = \frac{1}{2} M \left(\dot{x}^2 + \left(\frac{L^2 - \bar{x}^2}{L^2} \right) \frac{x^2 \dot{x}^2}{\left(\frac{L^2}{4} - x^2 \right)} \right)$$

$$+ \frac{1}{2} I \dot{\phi}^2 - MgY.$$

$$\begin{aligned}
 L &\stackrel{\perp}{=} \frac{1}{2} M \left(1 + \left(\frac{L^2 - \bar{x}^2}{L^2} \right) \frac{x^2}{\left(\frac{L^2}{4} - x^2 \right)} \right) \dot{x}^2 \\
 &+ \frac{1}{2} I \dot{\phi}^2 - Mg \sqrt{\frac{L^2 - \bar{x}^2}{L^2}} \sqrt{\frac{L^2}{4} - x^2}
 \end{aligned}$$

J V L - V 4



$$R\phi = l - l_0 \quad \text{length at origin}$$

$$\uparrow \quad \downarrow = L_1 - L_2$$

distance traveled

$$\uparrow = \frac{L}{2} + \frac{\bar{x}}{L} x - L_2$$

$$R\phi = \frac{\bar{x}}{L} x$$

$$\phi = \frac{\bar{x}}{RL} x$$

$$\hookrightarrow \dot{\phi} = \frac{\bar{x}}{RL} \dot{x}$$

$$L \doteq \frac{1}{2} M \left(1 + \left(\frac{L^2 - \bar{x}^2}{L} \right) \frac{x^2}{1 + 2 \cdot \frac{x^2}{L}} \right) \ddot{x}^2$$

$$\begin{aligned}
 & \text{Free Body Diagram: } \\
 & \text{Torque about center: } \\
 & +\frac{1}{2}I\left(\frac{x}{R}\right)^2 \dot{x}^2 + Mg \sqrt{\frac{L^2 - x^2}{L^2}} \sqrt{\frac{L^2}{4} - x^2} \\
 & \frac{M R^2}{2} \\
 & L = \frac{1}{2}M \left(1 + \frac{x^2}{2L^2} + \left(\frac{L^2 - x^2}{L^2} \right) \left(\frac{x^2}{\frac{L^2}{4} - x^2} \right) \right) \dot{x}^2 \\
 & + Mg \sqrt{\frac{L^2 - x^2}{L^2}} \sqrt{\frac{L^2}{4} - x^2}
 \end{aligned}$$

c) Determine Frequency around equilibrium.

Clearly equilibrium when $x_{eq} = 0$

$$x = 0 + \delta x \approx \delta x$$

$$L \approx \frac{1}{2}M(x=0) \dot{x}^2 + Mg \sqrt{\frac{L^2 - x^2}{L^2}} \sqrt{\frac{L^2}{4} - \delta x^2}$$

$$L = \frac{1}{2}M\left(1 + \frac{x^2}{2L^2}\right)\ddot{s}_x^2 + Mg\sqrt{\frac{L^2 - x^2}{L^2}} \frac{L}{2} \sqrt{1 - \frac{4s_x^2}{L^2}}$$

$$= \frac{1}{2}M\left(1 + \frac{x^2}{2L^2}\right)\ddot{s}_x^2 + Mg\sqrt{\frac{L^2 - x^2}{L^2}} \frac{L}{2} \left(1 - \frac{2s_x^2}{L^2}\right)$$

$$= \frac{1}{2}M\left(1 + \frac{x^2}{2L^2}\right)\ddot{s}_x^2 - \frac{Mg}{L}\sqrt{\frac{L^2 - x^2}{L^2}} s_x^2$$

$$\frac{dL}{ds_x} = M\left(1 + \frac{x^2}{2L^2}\right)\ddot{s}_x$$

$$\frac{dL}{ds_x} = -\frac{2Mg}{L}\sqrt{\frac{L^2 - x^2}{L^2}} \dot{s}_x$$

$$\ddot{s}_x = \underbrace{-\frac{2g}{L}\sqrt{\frac{L^2 - x^2}{L^2}}}_{W^2} \left(1 + \frac{x^2}{2L^2}\right)^{-1} \dot{s}_x$$

$$W^2 = \frac{2g}{L}\sqrt{\frac{L^2 - x^2}{L^2}} \left(\frac{2L^2}{2L^2 + x^2}\right)^{-1}$$

$$= \frac{2g}{L}\sqrt{\frac{L^2 - x^2}{L^2}} \frac{2L^2}{2L^2 + x^2}$$

$$\boxed{\omega^2 = \frac{4g\sqrt{L^2 - \dot{x}^2}}{2L^2 + \dot{x}^2}}$$

d) Now consider if \dot{x} is slowly increased.

$$\frac{1}{2\pi} \int dp dq = I$$

Area of phase space.

Slowly increase \rightarrow adiabatic invariant.

$$\oint \frac{pdq}{2\pi} = I$$

Adiabatic Invariant.

$I = \pi ab$ for ellipse.

If E is constant, then phase space is closed.

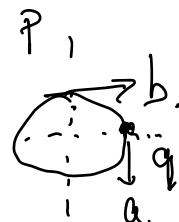
$$\ddot{x} + \omega^2 x = 0 \Rightarrow x = A_0 \cos(-\omega t + \phi)$$

$$E = \frac{1}{2} m(x) \dot{x}^2 + \frac{1}{2} m(x) \omega^2 x^2$$

$$= \frac{1}{2} M \left(1 + \frac{\dot{x}^2}{2L^2} \right) \dot{x}^2 + \frac{1}{2} M \left(1 + \frac{\dot{x}^2}{2L^2} \right) \omega^2 x^2$$

$$\text{We know } x = A_0 \cos(-\omega t + \phi)$$

$$\therefore \max(x) = A_0 = a$$



$$P = \frac{dL}{dx} = m(x) \dot{x}$$

$$= M\left(1 + \frac{\dot{x}^2}{2L^2}\right) W A_0 \sin(\omega t + \phi).$$

$$P_{max} = M\left(1 + \frac{\dot{x}^2}{2L^2}\right) W A_0 = b$$

$$I_f = \pi a b = M A_0^2 \left(1 + \frac{\dot{x}_0^2}{2L^2}\right) \sqrt{\frac{4g \sqrt{L^2 - \dot{x}_0^2}}{2L^2 + \dot{x}_0^2}}$$

$$I_f = I$$

$$M A_f^2 \left(1 + \frac{\dot{x}_f^2}{2L^2}\right) \sqrt{\frac{4g \sqrt{L^2 - \dot{x}_f^2}}{2L^2 - \dot{x}_f^2}} =$$

$$M A_0^2 \left(1 + \frac{\dot{x}_0^2}{2L^2}\right) \sqrt{\frac{4g \sqrt{L^2 - \dot{x}_0^2}}{2L^2 + \dot{x}_0^2}}$$

$$A_f = A_0 \sqrt{\frac{\left(1 + \frac{\dot{x}_0^2}{2L^2}\right) \sqrt{\frac{4g \sqrt{L^2 - \dot{x}_0^2}}{2L^2 + \dot{x}_0^2}}}{\left(1 + \frac{\dot{x}_f^2}{2L^2}\right) \sqrt{\frac{4g \sqrt{L^2 - \dot{x}_f^2}}{2L^2 - \dot{x}_f^2}}}}$$

3) Dissipation from an external field.

mass moving in 1D with $V(q)$ and

$$F_{\text{friction}} = -2m\gamma\dot{q}$$

a) $\frac{d\vec{L}}{dq} = f(t) \frac{d}{dq} \vec{L}_0$

$$\frac{d}{dt} \left(\frac{d\vec{L}}{dq} \right) = \frac{d}{dt} \left(f(t) \frac{d}{dq} \vec{L}_0 \right) = \dot{f} \frac{d}{dq} \vec{L}_0 + f \frac{d}{dt} \left(\frac{d}{dq} \vec{L}_0 \right)$$

$$\dot{f} \frac{d}{dq} \vec{L}_0 + f \frac{d}{dt} \left(\frac{d}{dq} \vec{L}_0 \right) = f(t) \frac{d}{dq} \vec{L}_0$$

$$\vec{L} = \frac{1}{2} m \dot{q}^2 - V(q)$$

$$\dot{f}(t) m \dot{q} + f m \ddot{q} = f - \frac{d}{dq} V(q)$$

$$\frac{\dot{f}}{f} m \dot{q} + m \ddot{q} + \frac{d}{dq} V(q) = 0$$

$$\frac{\dot{f}}{f} m \dot{q} = 2m\gamma\dot{q}$$

∴

$$\frac{v}{f} = 2\gamma$$

$$\ln f = 2\gamma t$$

$$e^{2\gamma t} = f(t)$$

$$L(H) = e^{2\gamma t} \left[\frac{1}{2} m \dot{q}^2 - V(q) \right]$$

b) $H = P\dot{q} - L$

$$P = \frac{\partial L}{\partial \dot{q}} = e^{2\gamma t} m \dot{q} \Rightarrow \dot{q} = \frac{P}{m} e^{-2\gamma t}$$

$$H = \frac{P^2}{m} e^{-2\gamma t} - e^{2\gamma t} \left[\frac{1}{2} \frac{P^2}{m} e^{-4\gamma t} - V(q) \right]$$

$$= \frac{P^2}{2m} e^{-2\gamma t} + e^{2\gamma t} V(q)$$

c) let $V(q) = \frac{1}{2} m \omega^2 q^2 = \frac{1}{2} q^2$ by $m = \omega = 1$

Find canonical transform $Q = Q(q, p, t)$

$$P = P(q, p, t)$$

such that $H(q, p, t) \rightarrow K(Q, P)$

$$H = \frac{P^2}{2} e^{-2\gamma t} + e^{2\gamma t} \frac{1}{2} q^2$$

$$\frac{\partial H}{\partial P} = \dot{q} \quad -\frac{\partial H}{\partial q} = \dot{P}$$

$$\stackrel{!}{=} pe^{-2\gamma t} \quad \stackrel{!}{=} -qe^{2\gamma t}$$

Find $\tilde{H}(Q, P)$ such that

$$\frac{\partial \tilde{H}}{\partial P} = \dot{Q} \quad -\frac{\partial \tilde{H}}{\partial Q} = \dot{P}$$

The variation in action should remain unchanged

$$S_1 = \int P dq - H dt$$

$$S_2 = \int P dQ - \tilde{H} dt$$

The difference must be a total derivative, else the action is different.

$$\int P dq - P dQ - (H - \tilde{H}) dt = \frac{dF}{dt}$$

So:

$$\dot{P}q - H(P, q, t) = \dot{P}\dot{Q} - \tilde{H}(P, Q, t) + \frac{dF}{dt}$$

$$\text{if we let } P = pe^{-2\gamma t} \quad Q = qe^{2\gamma t}$$

$$\hookrightarrow P = Pe^{rt} \quad \hookrightarrow q = Qe^{-rt}$$

$$\dot{q} = \dot{Q}e^{-rt} + -rQe^{-rt}$$

$$P\dot{q} - \mathcal{H} = Pe^{rt}(Qe^{-rt} - rQe^{-rt}) - \frac{P^2 e^{2rt}}{2} - \frac{Q^2 e^{-2rt}}{2}$$

$$\stackrel{!}{=} \dot{P}Q - P\gamma Q - \left(\frac{P^2}{2} + \frac{Q^2}{2} \right)$$

know

$$P\dot{q} - \mathcal{H} = \dot{P}Q - \mathcal{H} + \frac{df}{dt}$$

$$\dot{P}Q - P\gamma Q - \left(\frac{P^2}{2} + \frac{Q^2}{2} \right) = \dot{P}Q - \mathcal{H} + \frac{df}{dt}$$

$$\mathcal{H} = \left(\frac{P^2}{2} + \frac{Q^2}{2} \right) + \gamma PQ + \cancel{\frac{df}{dt}} \quad \text{let } f=0$$

Procedure:

let

$$\textcircled{1} \quad P\dot{q} - \mathcal{H} = \dot{P}Q - \mathcal{H} + \frac{df}{dt}$$

\textcircled{2} Have some ansatz on P and Q on q and p .

\textcircled{3} Revert to $q(Q, P)$, $p(Q, P)$. Plug into \textcircled{1}

Then determine \mathcal{H}

$$d) \dot{P} = -\frac{\partial H}{\partial Q} = -Q - \gamma P$$

$$\dot{Q} = \frac{\partial H}{\partial P} = P + \gamma Q$$

$$\ddot{Q} = \dot{P} + \gamma \dot{Q}$$

$$\ddot{Q} = -Q - \gamma P + \gamma \dot{Q}$$

$$\ddot{Q} = -Q - \gamma P + \gamma(P + \gamma Q)$$

$$\ddot{Q} = -Q + \gamma^2 Q$$

$$\ddot{Q} = -(1 - \gamma^2)Q$$

$$Q(t) = A \cos(\sqrt{1-\gamma^2}t - \phi)$$

$$\text{for } Q(t) = q e^{\lambda t}$$

$$q(t) = e^{-\lambda t} A \cos(\sqrt{1-\gamma^2}t - \phi)$$