

## Inverted Pendulum:



$$m\ddot{x} = -\frac{\partial U}{\partial x} + f(x,t) \quad , \quad f(x,t) = f_1(x) \cos(-\Omega t + \phi) \\ \stackrel{!}{=} f_1(x) \cos \Omega t + f_2(x) \sin \Omega t$$

Here  $\Omega$  is rapid,  $\Omega \gg \omega_0 \sim \frac{1}{T} \sim \sqrt{\frac{g}{L}}$   
 so the jitter is rapid.

Then let:  $x(t) = \bar{x}(t) + \zeta(t)$

$\nearrow$  large scale slow motion       $\nwarrow$  small scale fast motion.

$$\hookrightarrow m(\ddot{\bar{x}} + \ddot{\zeta}) = \frac{\partial U}{\partial x}(\bar{x} + \zeta) = \frac{\partial U}{\partial x} \Big|_{\bar{x}} + \frac{\partial^2 U}{\partial x^2} \Big|_{\bar{x}} \zeta \\ \stackrel{!}{=} \frac{\partial U}{\partial \bar{x}} + \frac{\partial^2 U}{\partial \bar{x}^2} \zeta$$

$$\hookrightarrow f(\bar{x} + \zeta, t) = f(\bar{x}, t) + \frac{\partial f(\bar{x}, t)}{\partial \bar{x}} \zeta$$

$$m\ddot{\bar{x}} + m\ddot{\zeta} = -\frac{\partial U}{\partial \bar{x}} - \frac{\partial^2 U}{\partial \bar{x}^2} \zeta + f(\bar{x}, t) + \zeta \frac{\partial f}{\partial \bar{x}}(\bar{x}, t)$$

$\uparrow$   
 fast

but larger than  $\zeta \frac{\partial^2 U}{\partial \bar{x}^2}$  or  $\zeta \frac{\partial f}{\partial \bar{x}}$ , since  $\ddot{\zeta} \sim \omega^2 \zeta$

## Fast Terms:

$$m\ddot{\zeta} = \underbrace{f(\bar{x}, t)}_{\text{N.A.}} = f_1(\bar{x}) \cos \Omega t + f_2(\bar{x}) \sin \Omega t$$

Find  $\left[ \zeta = \frac{-f(x,t)}{m\Omega^2} \right]$

Slow Terms: Average over time  $\Delta t$ :  $\Omega^{-1} \ll \Delta t \ll T$

longer than jitter, but shorter compared to overall motion

$$m\ddot{x} = -\frac{\partial U}{\partial x}(x) - \overline{\zeta \frac{\partial^2 U}{\partial x^2}} + \overline{\zeta \frac{\partial f}{\partial x}}$$

Since  $\zeta = \frac{-1}{m\Omega^2} (f_1(x)\cos\Omega t + f_2(x)\sin\Omega t)$  take average over  $\zeta$  terms. since they are fast

then  $\overline{\zeta \frac{\partial^2 U}{\partial x^2}} = 0$

For  $\overline{\zeta \frac{\partial f}{\partial x}} = \frac{-f}{m\Omega^2} \frac{\partial f}{\partial x} = -\frac{\partial}{\partial x} \left( \frac{f^2(x,t)}{2m\Omega^2} \right)$

since  $f(x,t) = f_1(x)\cos\Omega t + f_2(x)\sin\Omega t$

then  $\overline{f^2} = \frac{f_1^2 + f_2^2}{2}$

$$\hookrightarrow -\frac{\partial}{\partial x} \left( \frac{f^2(x,t)}{2m\Omega^2} \right) = -\frac{\partial}{\partial x} \left( \frac{f_1^2 + f_2^2}{4m\Omega^2} \right)$$

All together:

$$m\ddot{x} = -\frac{\partial}{\partial x} \left( U + \frac{f_1^2 + f_2^2}{4} \right) = -\frac{\partial}{\partial x} (V_{\text{eff}})$$

$$\text{or } 4m\Omega^2 / \text{or } m_{\text{eff}} /$$

$$\text{For } |\dot{\zeta}| = \Omega |\zeta| = \left| \frac{f}{m\Omega} \right|$$

$$\text{then } \frac{\overline{f^2}}{2m\Omega^2} = \frac{1}{2} m |\dot{\zeta}|^2$$

In general:

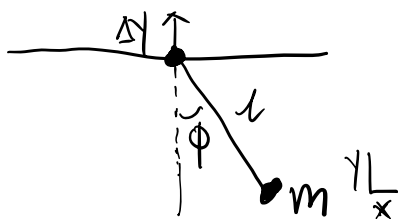
$$L = \frac{1}{2} m_{\text{eff}} \dot{q}^2 - U(q) + L_{\text{fast}}(q, t)$$

$$L_{\text{fast}}(q, t) = L_1(q) \cos \omega t + L_2(q) \sin \omega t$$

Then

$$\begin{aligned} m_{\text{eff}} \ddot{q} &= -\frac{\partial U}{\partial q} + \frac{\partial L_1}{\partial q} \cos \omega t + \frac{\partial L_2}{\partial q} \sin \omega t \\ &= -\frac{\partial}{\partial q} \left( U + \left( \frac{\partial L}{\partial q} \right)^2 / 4m_{\text{eff}} \Omega^2 \right) \end{aligned}$$

Inverted Pendulum



$$y = \Delta y - l \cos \phi \quad x = l \sin \phi$$

$$\Delta y = a \cos \Omega t$$

$$L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) - mgy$$

$$\begin{aligned} &= \frac{1}{2} m l^2 \dot{\phi}^2 + m l a \Omega \sin \Omega t \frac{d}{dt} \cos \phi + \frac{1}{2} m a^2 \Omega^2 \sin^2 \Omega t \\ &\quad - m g (a \cos \Omega t - l \cos \phi) \end{aligned}$$

total derivative

total derivative.

$$= \frac{1}{2} m l^2 \dot{\phi}^2 + \frac{d}{dt} (m l a \Omega \sin \Omega t \cos \phi) - m l a \Omega^2 \cos \Omega t \cos \phi + m g l \cos \phi$$

$$L = \frac{1}{2} m l^2 \dot{\phi}^2 - m l a \Omega^2 \cos \Omega t \cos \phi + m g l \cos \phi$$

$$\underbrace{m l^2}_{m_{\text{eff}}} \ddot{\phi} = - \underbrace{m g l \sin \phi}_{\frac{\partial U}{\partial \phi}} + \underbrace{(m a \Omega^2 \sin \phi \cos \Omega t)}_{f_{\phi}(\phi)}$$

then  $\rightarrow U_{\text{eff}} = -m g l \cos \phi + \frac{\overline{f_{\phi}^2(\phi)}}{4 m_{\text{eff}} \Omega^2}$

Effective  
Rosenbluth  
Potential.

$$= -m g l \cos \phi + \frac{(m a \Omega^2)^2}{4 m l^2 \Omega^2} \sin^2 \phi$$

$$= -m g l \cos \phi + \frac{m a^2 \Omega^2}{4} \sin^2 \phi$$

Notice:  $\gamma_{\text{eff}} = (m a \Omega^2 \sin \phi \cos \Omega t)$

$$\perp F_{\text{eff}} \cos \phi$$

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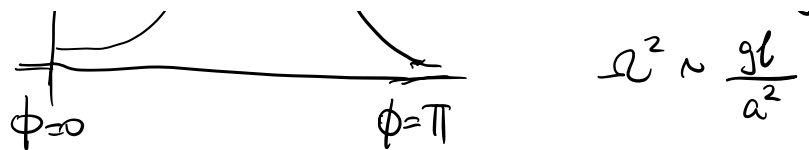
$(m a \Omega^2 \cos \Omega t)$

Suppose:  $m a^2 \Omega^2 \gg m g l$



loose stability:

$$m a^2 \Omega^2 \sim m g l$$



How fast for  $\Omega$  to see stable point.

Expand  $U_{\text{eff}}$  around  $\phi = \pi$ ,  $\cos\phi|_{\phi=\pi} \approx -1 + \frac{\phi^2}{2}$   
 $\rightarrow$  Perturbation near  $U_{\text{eff}}$ .

$$U = -mgl\cos\phi + \frac{ma^2\Omega^4}{4}\sin^2\phi$$

$$U = -mgl\left(-1 + \frac{\phi^2}{2}\right) + \frac{ma^2\Omega^2}{4}\phi^2$$

$$\stackrel{!}{=} mgl\left(1 - \frac{\phi^2}{2} + \frac{a^2\Omega^2}{4gl}\phi^2\right)$$

$$U(\phi) \stackrel{!}{=} mgl\left(1 + \left(-\frac{1}{2} + \frac{a^2\Omega^2}{4gl}\right)\phi^2\right)$$

In order for  $U(\phi)$  to be stable  $\frac{\partial^2}{\partial\phi^2}U > 0$ , so we have concave up shape.

$$\text{Thus: } -\frac{1}{2} + \frac{a^2\Omega^2}{4gl} > 0$$

$$\Omega^2 > \frac{2gl}{a^2} \quad \text{for stability}$$

### Electron Ion Trap:

With only electrostatic force, there is no equilibrium.

$$\text{let } \phi = \frac{1}{2} A \cos \Omega t (x^2 + y^2 - 2z^2)$$