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QM 2:

1) Suppose $g(t)$ is some matrix smoothly dependent on t .

a) prove $\partial_t(g^{-1}) = -g^{-1}(\partial_t g)g^{-1}$

$$\partial_t(g g^{-1}) = \partial_t(\mathbb{1}) = 0$$

$$0 = \partial_t(g) g^{-1} + g \partial_t(g^{-1})$$

$$-g \partial_t g^{-1} = \partial_t(g) g^{-1}$$

Multiply $-g^{-1}$
to both sides

$$\hookrightarrow \underbrace{g^{-1}g}_{\mathbb{1}} \partial_t g^{-1} = -g^{-1} \partial_t(g) g^{-1}$$

$$\hookrightarrow \boxed{\partial_t g^{-1} = -g^{-1} \partial_t(g) g^{-1}}$$

b) Assume $g(t)$ is orthogonal matrix $\rightarrow g^T = g^{-1}$

$$\text{prove: } (g^{-1} \partial_t g)^T = -g^{-1} \partial_t g$$

$$(g^{-1} \partial_t g)^T = (\partial_t g)^T (g^{-1})^T$$

$$= \partial_t(g^T) (g^T)^{-1}$$

For orthogonal
Matrix, $g^T = g^{-1}$

\hookrightarrow

$$= \partial_t(g^{-1}) \underbrace{(g^{-1})^{-1}}_g$$

$$\text{but } \partial_t(g^{-1}g) = \partial_t(\mathbb{1}) = 0$$

$$\hookrightarrow 0 = \partial_t(g^{-1})g + g^{-1}\partial_t g$$

$$\boxed{\Omega^T = (g^{-1} \partial_t g)^T = -g^{-1} \partial_t g = -\Omega}$$

c) If $g(t)$ is unitary matrix, $g^{-1} = g^\dagger$

$$\begin{aligned} (g^{-1} \partial_t g)^\dagger &= (\partial_t g)^\dagger (g^{-1})^\dagger \\ &= \underbrace{\partial_t (g^\dagger)}_{\partial_t (g^{-1})} \underbrace{(g^\dagger)^{-1}}_{(g^{-1})^{-1} = g} \\ &= \partial_t (g^{-1}) g \end{aligned}$$

again:

$$\begin{aligned} 0 &= \partial_t (g^{-1}) g + g^{-1} \partial_t g - g^{-1} \partial_t g \\ &\quad \left| \begin{array}{l} (g^{-1} \partial_t g)^\dagger = -g^{-1} \partial_t g \end{array} \right. \end{aligned}$$

2) a) Calculate $\text{Tr}\{(\vec{a} \cdot \vec{\sigma})(\vec{b} \cdot \vec{\sigma})(\vec{c} \cdot \vec{\sigma})\}$

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\vec{a} \cdot \vec{\sigma} = a_x \sigma_x + a_y \sigma_y + a_z \sigma_z$$

$$= \begin{pmatrix} 0 & a_x \\ a_x & 0 \end{pmatrix} + \begin{pmatrix} 0 & -ia_y \\ ia_y & 0 \end{pmatrix} + \begin{pmatrix} a_z & 0 \\ 0 & -a_z \end{pmatrix}$$

$$= \begin{pmatrix} a_z & a_x - ia_y \\ a_x + ia_y & -a_z \end{pmatrix}$$

$$\text{so } \text{Tr}\{(\vec{a} \cdot \vec{\sigma})(\vec{b} \cdot \vec{\sigma})(\vec{c} \cdot \vec{\sigma})\}$$

$$= \text{Tr} \left\{ \begin{pmatrix} a_z & a_x - ia_y \\ a_x + ia_y & -a_z \end{pmatrix} \begin{pmatrix} b_z & b_x - ib_y \\ b_x + ib_y & -b_z \end{pmatrix} \begin{pmatrix} c_z & c_x - ic_y \\ c_x + ic_y & -c_z \end{pmatrix} \right\}$$

$$\text{Tr} \left\{ \begin{pmatrix} a_z b_z + (a_x - ia_y)(b_x + ib_y) & a_z(b_x - ib_y) - b_z(a_x - ia_y) \\ b_z(a_x + ia_y) - a_z(b_x + ib_y) & (a_x + ia_y)(b_x - ib_y) + a_z b_z \end{pmatrix} \begin{pmatrix} c_z & c_x - ic_y \\ c_x + ic_y & -c_z \end{pmatrix} \right\}$$

$$= \left[\cancel{a_z b_z} + (a_x - ia_y)(b_x + ib_y) \right] c_z + \left[a_z(b_x - ib_y) - b_z(a_x - ia_y) \right] (c_x + ic_y) \\ + \left[b_z(a_x + ia_y) - a_z(b_x + ib_y) \right] (c_x - ic_y) - c_z \left[(a_x + ia_y)(b_x - ib_y) + \cancel{a_z b_z} \right]$$

$$= (\cancel{a_x b_x} - ia_y b_x + ia_x b_y + \cancel{a_y b_y}) c_z - c_z (\cancel{a_x b_x} - ia_x b_y + ia_y b_x + \cancel{a_y b_y})$$

$$(a_z b_x - ia_z b_y - b_z a_x + ib_z a_y)(c_x + ic_y) + (b_z a_x + ib_z a_y - a_z b_x - ia_z b_y)(c_x - ic_y)$$

$$= c_z (2ia_x b_y - 2ia_y b_x)$$

$$+ c_x (\cancel{a_z b_x} - ia_z b_y - \cancel{b_z a_x} + ib_z a_y + \cancel{b_z a_x} + ib_z a_y - \cancel{a_z b_x} - ia_z b_y)$$

$$+ ic_y (a_z b_x - \cancel{ia_z b_y} - b_z a_x + \cancel{ib_z a_y} - b_z a_x - \cancel{ib_z a_y} + a_z b_x + \cancel{ia_z b_y})$$

$$= 2i(a_x b_y - a_y b_x) c_z + 2i(a_y b_z - a_z b_y) c_x + 2i(a_z b_x - a_x b_z) c_y$$

$$\boxed{\text{Tr}\{(\vec{a} \cdot \vec{\sigma})(\vec{b} \cdot \vec{\sigma})(\vec{c} \cdot \vec{\sigma})\} = 2i \epsilon^{ijk} a_i b_j c_k}$$

b) Calculate $\exp\begin{pmatrix} 3 & 4 \\ 2 & 1 \end{pmatrix}$ using Pauli - Matrices.

From part c, we know any 2×2 matrix can be decomposed into linear combinations of Pauli - Matrices

$$\vec{A} \rightarrow \vec{a} \cdot \vec{\sigma}$$

We know via Taylor:

$$e^{\vec{A}} = \sum_{n=0}^{\infty} \frac{A^n}{n!} \rightarrow \sum_{n=0}^{\infty} \frac{\vec{a} \cdot \vec{\sigma}}{n!}$$

split into even and odd terms.

$$= \sum_{p=0}^{\infty} \frac{(\vec{a} \cdot \vec{\sigma})^{2p}}{(2p)!} + \sum_{q=0}^{\infty} \frac{(\vec{a} \cdot \vec{\sigma})^{2q+1}}{(2q+1)!}$$

consider

$$(\vec{a} \cdot \vec{\sigma})(\vec{a} \cdot \vec{\sigma}) = (a_i \sigma_i)(a_j \sigma_j)$$

$$= \frac{a_i a_j}{2} \left(\underbrace{\sigma_i \sigma_j - \sigma_j \sigma_i}_{[\sigma_i, \sigma_j] = 2i\epsilon_{ijk}\sigma_k} + \underbrace{\sigma_i \sigma_j + \sigma_j \sigma_i}_{\{\sigma_i, \sigma_j\} = 2\delta_{ij}\mathbb{I}} \right)$$

$$= a_i a_j [i\epsilon_{ijk}\sigma_k + \delta_{ij}\mathbb{I}]$$

$$= a_j a_i + i \underbrace{a_i a_j \epsilon_{ijk}}_{\vec{a} \times \vec{a} = 0} \sigma_k$$

$$(\vec{a} \cdot \vec{\sigma})(\vec{a} \cdot \vec{\sigma}) = \vec{a} \cdot \vec{a}$$

Therefore:

$$\begin{aligned}
 &= \sum_{p=0}^{\infty} \frac{(\vec{a} \cdot \vec{\sigma})^{2p}}{(2p)!} + \sum_{q=0}^{\infty} \frac{(\vec{a} \cdot \vec{\sigma})^{2q+1}}{(2q+1)!} \\
 &= \sum_{p=0}^{\infty} \frac{(\vec{a} \cdot \vec{a})^p}{(2p)!} + (\vec{a} \cdot \vec{\sigma}) \sum_{q=0}^{\infty} \frac{(\vec{a} \cdot \vec{a})^q}{(2q+1)!} \\
 &= \underbrace{\sum_{p=0}^{\infty} \frac{(\sqrt{\vec{a} \cdot \vec{a}})^{2p}}{(2p)!}}_{\cosh(\sqrt{\vec{a} \cdot \vec{a}})} + (\vec{a} \cdot \vec{\sigma}) \underbrace{\frac{1}{\sqrt{\vec{a} \cdot \vec{a}}} \sum_{q=0}^{\infty} \frac{(\sqrt{\vec{a} \cdot \vec{a}})^{2q+1}}{(2q+1)!}}_{\sinh(\sqrt{\vec{a} \cdot \vec{a}})}
 \end{aligned}$$

$$e^{\vec{a} \cdot \vec{\sigma}} = \mathbb{I} \cosh(\sqrt{\vec{a} \cdot \vec{a}}) + \frac{(\vec{a} \cdot \vec{\sigma})}{\sqrt{\vec{a} \cdot \vec{a}}} \sinh(\sqrt{\vec{a} \cdot \vec{a}})$$

using results from part C: we can decompose $\begin{pmatrix} 3 & 4 \\ 2 & 1 \end{pmatrix}$ into

$$\hookrightarrow \begin{pmatrix} 3 & 4 \\ 2 & 1 \end{pmatrix} = 2 \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + 3 \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\hookrightarrow \exp \begin{pmatrix} 3 & 4 \\ 2 & 1 \end{pmatrix} = \exp \{ 2\mathbb{I} + 3\sigma_x + i\sigma_y + \sigma_z \}$$

since \mathbb{I} commutes with everything.

$$\begin{aligned}
 &= \exp \{ 2\mathbb{I} \} \exp \{ 3\sigma_x + i\sigma_y + \sigma_z \} \\
 &= \begin{pmatrix} e^2 & 0 \\ 0 & e^2 \end{pmatrix} \exp \{ 3\sigma_x + i\sigma_y + \sigma_z \}
 \end{aligned}$$

$$\text{let } \vec{a} = (3, i, 1) \rightarrow \sqrt{\vec{a} \cdot \vec{a}} = \sqrt{3^2 - 1 + 1} = 3$$

$$e^{\vec{a} \cdot \vec{\sigma}} = \cosh(\sqrt{\vec{a} \cdot \vec{a}}) \mathbb{I} + \frac{\vec{a} \cdot \vec{\sigma}}{\sqrt{\vec{a} \cdot \vec{a}}} \sinh(\sqrt{\vec{a} \cdot \vec{a}})$$

$$\hookrightarrow = \begin{pmatrix} e^2 & 0 \\ 0 & e^2 \end{pmatrix} \left\{ \begin{pmatrix} \cosh 3 & 0 \\ 0 & \cosh 3 \end{pmatrix} + \sinh 3 \left[\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + \begin{pmatrix} 0 & \frac{1}{3} \\ -\frac{1}{3} & 0 \end{pmatrix} + \begin{pmatrix} \frac{1}{3} & 0 \\ 0 & -\frac{1}{3} \end{pmatrix} \right] \right\}$$

$$\boxed{\exp \begin{pmatrix} 3 & 4 \\ 2 & 1 \end{pmatrix} = \begin{pmatrix} e^2 \left(\cosh 3 + \frac{\sinh 3}{3} \right) & e^2 \frac{4}{3} \sinh 3 \\ e^2 \frac{2}{3} \sinh 3 & e^2 \left(\cosh 3 - \frac{\sinh 3}{3} \right) \end{pmatrix}}$$

c) express $\begin{pmatrix} a & b \\ c & d \end{pmatrix}$ using pauli - Matrices:

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} = \alpha \underbrace{\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}}_{\mathbb{I}} + \beta \underbrace{\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}} + \gamma \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + \delta \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} = \begin{pmatrix} \alpha + \beta & \beta - i\gamma \\ \beta + i\gamma & \alpha - \beta \end{pmatrix}$$

so

$$\left. \begin{array}{l} \alpha + \beta = a \\ d = \alpha - \beta \end{array} \right\} \quad \beta + d + \beta = a \rightarrow \boxed{\beta = \frac{1}{2}(a-d)}$$

$$\text{and } \alpha = \beta + d = \frac{1}{2}(a-d) + d = \boxed{\frac{1}{2}(a+d) = \alpha}$$

$$\text{also } \left. \begin{array}{l} \beta - i\gamma = b \\ \beta + i\gamma = c \end{array} \right\} \quad b + i\gamma + i\gamma = c \rightarrow \boxed{\gamma = \frac{-i}{2}(c-b)}$$

$$Q = b + i\gamma = b + \frac{1}{2}(c-b) = \boxed{\frac{1}{2}(c+b) = Q}$$

then

$$\boxed{\begin{pmatrix} a & b \\ c & d \end{pmatrix} = \frac{1}{2}(a+d) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \frac{1}{2}(a-d) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + \frac{1}{2}(c+b) \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} - \frac{i}{2}(c-b) \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}}$$

$$3) \quad \psi = N e^{-\alpha r^2} (x+y) z, \quad r^2 = x^2 + y^2 + z^2$$

$$a) \quad \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dx dy dz \quad \psi^* \psi$$

$$\hookrightarrow = \iiint dx dy dz \quad |N|^2 e^{-2\alpha(x^2+y^2+z^2)} (x+y)^2 z^2$$

|

$$= |N|^2 \iiint e^{-2\alpha x^2} e^{-2\alpha y^2} e^{-2\alpha z^2} (x^2 + 2xy + y^2) z^2$$

|

$$= |N|^2 \int dy dz e^{-2\alpha y^2} e^{-2\alpha z^2} z^2 \left(\frac{\sqrt{\pi}}{2^{5/2} \alpha^{3/2}} + \sqrt{\frac{\pi}{2\alpha}} y^2 \right) z^2$$

|

$$= |N|^2 \int dz e^{-2\alpha z^2} z^2 \underbrace{\left(\frac{\sqrt{\pi}}{2^{5/2} \alpha^{3/2}} \sqrt{\frac{\pi}{2\alpha}} 2 \right)}_{\frac{\pi}{(2\alpha)^2}}$$

|

$$= |N|^2 \frac{\sqrt{\pi}}{2^{5/2} \alpha^{3/2}} \frac{\pi}{(2\alpha)^2}$$

$$1 = |N|^2 \pi^{3/2} \frac{1}{2^{9/2} \alpha^{7/2}}$$

$$\hookrightarrow \boxed{|N| = \sqrt{\frac{2^{9/2} \alpha^{7/2}}{\pi^{3/2}}} = \frac{2^{9/4} \alpha^{7/4}}{\pi^{3/4}}}$$

$$b) \quad \vec{L} = \vec{r} \times \vec{p} = -i\hbar \epsilon^{ijk} r_i \partial_j \hat{e}_k$$

$$L_x = -i\hbar (y\partial_z - z\partial_y)$$

$$L_y = -i\hbar (z\partial_x - x\partial_z)$$

$$L_z = -i\hbar (x\partial_y - y\partial_x)$$

Calculate derivatives: $\psi = N e^{-\alpha(x^2+y^2+z^2)} (x+y)z$

$$\partial_z \psi = N(x+y) e^{-\alpha(x^2+y^2)} \left[(-2\alpha z) e^{-\alpha z^2} z + e^{-\alpha z^2} \right]$$

$$= N e^{-\alpha(x^2+y^2+z^2)} \left[-2\alpha z^2 + 1 \right] (x+y)$$

$$\partial_x \psi = N z e^{-\alpha(y^2+z^2)} \left[-2\alpha x e^{-\alpha x^2} (x+y) + e^{-\alpha x^2} \right]$$

$$= N e^{-\alpha(x^2+y^2+z^2)} \left[-2\alpha x(x+y) + 1 \right] z$$

$$\partial_y \psi = N z e^{-\alpha(x^2+z^2)} \left[-2\alpha y e^{-\alpha y^2} (x+y) + e^{-\alpha y^2} \right]$$

$$= N e^{-\alpha(x^2+y^2+z^2)} \left[-2\alpha y(x+y) + 1 \right] z$$

$$\langle L_x \rangle = \iiint dx dy dz \psi^* (-i\hbar)(y\partial_z - z\partial_y) \psi$$

$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dx dy dz (-i\hbar) N^2 e^{-2\alpha(x^2+y^2+z^2)} (x+y)z$$

$$\left[y(x+y)(-2\alpha z^2 + 1) - z^2(-2\alpha y(x+y) + 1) \right]$$

$$= \iiint dx dy dz (-i\hbar) N^2 e^{-2\alpha(x^2+y^2+z^2)} (xz + yz)$$

$$\left[-2\alpha(\cancel{xy z^2} + y^2 z^2) + xy + y^2 + 2\alpha(\cancel{xy z^2} + y^2 z^2) - z^2 \right]$$

$$= \iiint dx dy dz (-i\hbar) N^2 \underbrace{e^{-2\alpha(x^2+y^2+z^2)}}_{\text{even term}} \underbrace{(x^2 y z + x y^2 z - x z^3 + x y^2 z + y^3 z - y z^3)}_{\text{Each term has at one least one odd term in either } x, y, \text{ or } z}$$

$$= 0$$

Since entire integrand is odd and we integrate over $(-\infty, +\infty)$.

$$\langle L_y \rangle = \iiint dx dy dz \psi^* (-i\hbar)(z\partial_x - x\partial_z) \psi$$

$$= \iiint dx dy dz (-i\hbar) N^2 e^{-2\alpha(x^2+y^2+z^2)} (x+y)z$$

$$\left[z^2(-2\alpha x(x+y) + 1) - x(x+y)(-2\alpha z^2 + 1) \right]$$

$$= \iiint dx dy dz (-i\hbar) N^2 e^{-2\alpha(x^2+y^2+z^2)} (xz + yz)$$

$$(-2\alpha(\cancel{x^2 z^2} + xy z^2) + z^2 + 2\alpha(\cancel{x^2 z^2} + xy z^2) - x^2 - xy)$$

$$= 0 \quad \text{odd integrand just like } L_x$$

$$\begin{aligned}\langle L_z \rangle &= \iiint dx dy dz \psi^* (-i\hbar) (x\partial_y - y\partial_x) \psi \\ &= \iiint dx dy dz (-i\hbar) N^2 e^{-2\alpha(x^2+y^2+z^2)} (x+y) z\end{aligned}$$

$$[xz(-2\alpha y(x+y)+1) - yz(-2\alpha x(x+y)+1)]$$

$$= \iiint dx dy dz (-i\hbar) N^2 e^{-2\alpha(x^2+y^2+z^2)} (xz + yz)$$

$$(-2\alpha(\cancel{x^2 y z} + x y^2 z) + xz + 2\alpha(\cancel{x^2 y z} + x y^2 z) - yz)$$

$$= \iiint dx dy dz (-i\hbar) N^2 e^{-2\alpha(x^2+y^2+z^2)} \{ \underbrace{(xz)^2 - (yz)^2}_{= z^2(x^2 - y^2)} \}$$

$$= z^2(x^2 - y^2)$$

Since there is a symmetry in x and y , $x^2 - y^2$ leads to 0 expectation value.

$$= 0$$

Therefore

$$\boxed{\langle \vec{L} \rangle = \langle L_x \rangle + \langle L_y \rangle + \langle L_z \rangle = 0}$$

$$\text{Find } \langle L^2 \rangle = \langle L_x^2 \rangle + \langle L_y^2 \rangle + \langle L_z^2 \rangle$$

$$L^2 = \underbrace{(-i\hbar)^2}_{-\hbar^2} \left\{ (y\partial_z - z\partial_y)^2 + (z\partial_x - x\partial_z)^2 + (x\partial_y - y\partial_x)^2 \right\}$$

$$L_x^2 \psi = L_x L_x \psi$$

$$= (-i\hbar)^2 (y\partial_z - z\partial_y) N e^{-\alpha(x^2+y^2+z^2)} (y^2 + xy - z^2)$$

$$= (-i\hbar)^2 N e^{-\alpha(x^2+y^2+z^2)} \left\{ y \left[(-2\alpha z)(y^2 + xy - z^2) - 2z \right] \right. \\ \left. - z \left[(-2\alpha y)(y^2 + xy - z^2) + 2y + x \right] \right\}$$

$$L_x^2 \psi = \hbar^2 N e^{-\alpha(x^2+y^2+z^2)} (x + 4y)z$$

$$L_y^2 \psi = (-i\hbar)^2 L_y L_y \psi$$

$$= -\hbar^2 (z\partial_x - x\partial_z) (N e^{-\alpha(x^2+y^2+z^2)} (-x^2 - xy + z^2))$$

$$= -\hbar^2 N e^{-\alpha(x^2+y^2+z^2)} \left\{ z \left[-2\alpha x (-x^2 - xy + z^2) - 2x - y \right] \right. \\ \left. - x \left[-2\alpha z (-x^2 - xy + z^2) + 2z \right] \right\}$$

$$L_y^2 \psi = \hbar^2 N e^{-\alpha(x^2+y^2+z^2)} (4x + y)z$$

$$\begin{aligned}
 L_z^2 \psi &= -\hbar^2 (x \partial_y - y \partial_x) \text{Ne}^{-\alpha(x^2+y^2+z^2)} (x-y)z \\
 &\stackrel{!}{=} -\hbar^2 \text{Ne}^{-\alpha(x^2+y^2+z^2)} z (-x-y) \\
 &\stackrel{!}{=} \hbar^2 \text{Ne}^{-\alpha(x^2+y^2+z^2)} (x+y)z
 \end{aligned}$$

$$\begin{aligned}
 L^2 \psi &= L_x^2 \psi + L_y^2 \psi + L_z^2 \psi \\
 &\stackrel{!}{=} \hbar^2 \text{Ne}^{-\alpha(x^2+y^2+z^2)} z (4x+y+4y+x+x+y) \\
 &\stackrel{!}{=} 6\hbar^2 \underbrace{\text{Ne}^{-\alpha(x^2+y^2+z^2)} (x+y)z}_{=\psi} \\
 &\stackrel{!}{=} 6\hbar^2 \psi
 \end{aligned}$$

$$\begin{aligned}
 \langle L^2 \rangle &= \iiint dx dy dz \psi^* \underbrace{L^2 \psi}_{6\hbar^2 \psi} \\
 &\stackrel{!}{=} 6\hbar^2 \int \psi^* \psi
 \end{aligned}$$

$$\boxed{\langle L^2 \rangle = 6\hbar^2}$$

c) variance of L , $\Delta^2(L) = \langle L^2 \rangle - \langle L \rangle^2$

$$\boxed{\Delta^2(L) \stackrel{!}{=} 6\hbar^2}$$

$$\boxed{\Delta^2(L^2) = \langle L^4 \rangle - \langle L^2 \rangle^2 = 0}$$

since L^2 commutes with L^4
 $[L^4, L^2] = [L^2 L^2, L^2] = 0$

4) work out J-Matrices for $\bar{j} = \frac{1}{2}$, $\bar{j} = 1$, $\bar{j} = \frac{3}{2}$

$$J_z |j, m\rangle = \hbar m |j, m\rangle$$

$$\left. \begin{aligned} J_+ |j, m\rangle &= \hbar \sqrt{(j-m)(j+m+1)} |j, m+1\rangle \\ J_- |j, m\rangle &= \hbar \sqrt{(j+m)(j-m+1)} |j, m-1\rangle \end{aligned} \right\} \begin{aligned} J_x &= \frac{J_+ + J_-}{2} \\ J_y &= \frac{J_+ - J_-}{2i} \end{aligned}$$

$$\underline{\bar{j} = \frac{1}{2}}, \quad m = \frac{1}{2}, -\frac{1}{2}$$

$$J_z^{\frac{1}{2}} = \hbar \begin{pmatrix} \langle \frac{1}{2}, \frac{1}{2} | J_z | \frac{1}{2}, \frac{1}{2} \rangle & \langle \frac{1}{2}, \frac{1}{2} | J_z | \frac{1}{2}, -\frac{1}{2} \rangle \\ \langle \frac{1}{2}, -\frac{1}{2} | J_z | \frac{1}{2}, \frac{1}{2} \rangle & \langle \frac{1}{2}, -\frac{1}{2} | J_z | \frac{1}{2}, -\frac{1}{2} \rangle \end{pmatrix}$$

$$J_z^{\frac{1}{2}} = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$J_+^{\frac{1}{2}} = \begin{pmatrix} \langle \frac{1}{2}, \frac{1}{2} | J_+ | \frac{1}{2}, \frac{1}{2} \rangle & \langle \frac{1}{2}, \frac{1}{2} | J_+ | \frac{1}{2}, -\frac{1}{2} \rangle \\ \langle \frac{1}{2}, -\frac{1}{2} | J_+ | \frac{1}{2}, \frac{1}{2} \rangle & \langle \frac{1}{2}, -\frac{1}{2} | J_+ | \frac{1}{2}, -\frac{1}{2} \rangle \end{pmatrix}$$

$$= \hbar \begin{pmatrix} 0 & \sqrt{\left(\frac{1}{2} - (-\frac{1}{2})\right)\left(\frac{1}{2} - \frac{1}{2} + 1\right)} = 1 \\ 0 & 0 \end{pmatrix}$$

$$J_-^{\frac{1}{2}} = \begin{pmatrix} \langle \frac{1}{2}, \frac{1}{2} | J_- | \frac{1}{2}, \frac{1}{2} \rangle & \langle \frac{1}{2}, \frac{1}{2} | J_- | \frac{1}{2}, -\frac{1}{2} \rangle \\ \langle \frac{1}{2}, -\frac{1}{2} | J_- | \frac{1}{2}, \frac{1}{2} \rangle & \langle \frac{1}{2}, -\frac{1}{2} | J_- | \frac{1}{2}, -\frac{1}{2} \rangle \end{pmatrix}$$

$$= \hbar \begin{pmatrix} 0 & 0 \\ \sqrt{\left(\frac{1}{2} + \frac{1}{2}\right)\left(\frac{1}{2} - \frac{1}{2} + 1\right)} = 1 & 0 \end{pmatrix}$$

$$J_x^{\frac{1}{2}} = \frac{J_+ + J_-}{2} = \frac{\hbar}{2} \left(\begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \right) = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$J_y^{\frac{1}{2}} = \frac{J_+ - J_-}{2i} = \frac{\hbar}{2i} \left(\begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} - \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \right) = \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

Summary: $\hat{j} = \frac{1}{2}$:

$$J_x = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$J_y = \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$J_z = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\hat{j} = 1: \quad m = -1, 0, 1$$

$$J_z^1 = \hbar \begin{pmatrix} \langle 1 | J_z | 1 \rangle & \langle 1 | J_z | 0 \rangle & \langle 1 | J_z | -1 \rangle \\ \langle 0 | J_z | 1 \rangle & \langle 0 | J_z | 0 \rangle & \langle 0 | J_z | -1 \rangle \\ \langle -1 | J_z | 1 \rangle & \langle -1 | J_z | 0 \rangle & \langle -1 | J_z | -1 \rangle \end{pmatrix}$$

$$= \hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

$$J_+ = \hbar \begin{pmatrix} \langle 1 | J_+ | 1 \rangle & \langle 1 | J_+ | 0 \rangle & \langle 1 | J_+ | -1 \rangle \\ \langle 0 | J_+ | 1 \rangle & \langle 0 | J_+ | 0 \rangle & \langle 0 | J_+ | -1 \rangle \\ \langle -1 | J_+ | 1 \rangle & \langle -1 | J_+ | 0 \rangle & \langle -1 | J_+ | -1 \rangle \end{pmatrix}$$

$$= \hbar \begin{pmatrix} 0 & \sqrt{(1-0)(1+0+1)} = \sqrt{2} & 0 \\ 0 & 0 & \sqrt{(1-(-1))(1+(-1)+1)} = \sqrt{2} \\ 0 & 0 & 0 \end{pmatrix}$$

$$J_- = \hbar \begin{pmatrix} \langle 1 | J_- | 1 \rangle & \langle 1 | J_- | 0 \rangle & \langle 1 | J_- | -1 \rangle \\ \langle 0 | J_- | 1 \rangle & \langle 0 | J_- | 0 \rangle & \langle 0 | J_- | -1 \rangle \\ \langle -1 | J_- | 1 \rangle & \langle -1 | J_- | 0 \rangle & \langle -1 | J_- | -1 \rangle \end{pmatrix}$$

$$= \hbar \begin{pmatrix} 0 & 0 & 0 \\ \sqrt{(1+1)(1-1+1)} = \sqrt{2} & 0 & 0 \\ 0 & \sqrt{(1+0)(1-0+1)} = \sqrt{2} & 0 \end{pmatrix}$$

$$J_x = \frac{1}{2} (J_+ + J_-) = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

$$J_y = \frac{1}{2i} (J_+ - J_-) = \frac{\hbar}{2i} \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 1 \\ 0 & -1 & 0 \end{pmatrix}$$

$$\text{with } J_z = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

$$\bar{j} = \frac{3}{2}, \quad m = -\frac{3}{2}, -\frac{1}{2}, \frac{1}{2}, \frac{3}{2}$$

$$\bar{J}_z^{3/2} = \begin{pmatrix} \langle \frac{3}{2} | \bar{J}_z | \frac{3}{2} \rangle & \langle \frac{3}{2} | \bar{J}_z | \frac{1}{2} \rangle & \langle \frac{3}{2} | \bar{J}_z | -\frac{1}{2} \rangle & \langle \frac{3}{2} | \bar{J}_z | -\frac{3}{2} \rangle \\ \langle \frac{1}{2} | \bar{J}_z | \frac{3}{2} \rangle & \langle \frac{1}{2} | \bar{J}_z | \frac{1}{2} \rangle & \langle \frac{1}{2} | \bar{J}_z | -\frac{1}{2} \rangle & \langle \frac{1}{2} | \bar{J}_z | -\frac{3}{2} \rangle \\ \langle -\frac{1}{2} | \bar{J}_z | \frac{3}{2} \rangle & \langle -\frac{1}{2} | \bar{J}_z | \frac{1}{2} \rangle & \langle -\frac{1}{2} | \bar{J}_z | -\frac{1}{2} \rangle & \langle -\frac{1}{2} | \bar{J}_z | -\frac{3}{2} \rangle \\ \langle -\frac{3}{2} | \bar{J}_z | \frac{3}{2} \rangle & \langle -\frac{3}{2} | \bar{J}_z | \frac{1}{2} \rangle & \langle -\frac{3}{2} | \bar{J}_z | -\frac{1}{2} \rangle & \langle -\frac{3}{2} | \bar{J}_z | -\frac{3}{2} \rangle \end{pmatrix}$$

$$\bar{J}_z = \hbar \begin{pmatrix} 3/2 & 0 & 0 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & -1/2 & 0 \\ 0 & 0 & 0 & -3/2 \end{pmatrix}$$

$$\bar{J}_+^{3/2} = \hbar \begin{pmatrix} \langle \frac{3}{2} | \bar{J}_+ | \frac{3}{2} \rangle & \langle \frac{3}{2} | \bar{J}_+ | \frac{1}{2} \rangle & \langle \frac{3}{2} | \bar{J}_+ | -\frac{1}{2} \rangle & \langle \frac{3}{2} | \bar{J}_+ | -\frac{3}{2} \rangle \\ \langle \frac{1}{2} | \bar{J}_+ | \frac{3}{2} \rangle & \langle \frac{1}{2} | \bar{J}_+ | \frac{1}{2} \rangle & \langle \frac{1}{2} | \bar{J}_+ | -\frac{1}{2} \rangle & \langle \frac{1}{2} | \bar{J}_+ | -\frac{3}{2} \rangle \\ \langle -\frac{1}{2} | \bar{J}_+ | \frac{3}{2} \rangle & \langle -\frac{1}{2} | \bar{J}_+ | \frac{1}{2} \rangle & \langle -\frac{1}{2} | \bar{J}_+ | -\frac{1}{2} \rangle & \langle -\frac{1}{2} | \bar{J}_+ | -\frac{3}{2} \rangle \\ \langle -\frac{3}{2} | \bar{J}_+ | \frac{3}{2} \rangle & \langle -\frac{3}{2} | \bar{J}_+ | \frac{1}{2} \rangle & \langle -\frac{3}{2} | \bar{J}_+ | -\frac{1}{2} \rangle & \langle -\frac{3}{2} | \bar{J}_+ | -\frac{3}{2} \rangle \end{pmatrix}$$

$$\bar{J}_+^{3/2} = \hbar \begin{pmatrix} 0 & \sqrt{(\frac{3}{2} - \frac{1}{2})(\frac{3}{2} + \frac{1}{2} + 1)} = \sqrt{3} & 0 & 0 \\ 0 & 0 & \sqrt{(\frac{3}{2} - (-\frac{1}{2}))(\frac{3}{2} - \frac{1}{2} + 1)} = 2 & 0 \\ 0 & 0 & 0 & \sqrt{(\frac{3}{2} - (-\frac{3}{2}))(\frac{3}{2} - \frac{3}{2} + 1)} = \sqrt{3} \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$J_-^{3/2} = \hbar \begin{pmatrix} \langle \frac{3}{2} | J_- | \frac{3}{2} \rangle & \langle \frac{3}{2} | J_- | \frac{1}{2} \rangle & \langle \frac{3}{2} | J_- | -\frac{1}{2} \rangle & \langle \frac{3}{2} | J_- | -\frac{3}{2} \rangle \\ \langle \frac{1}{2} | J_- | \frac{3}{2} \rangle & \langle \frac{1}{2} | J_- | \frac{1}{2} \rangle & \langle \frac{1}{2} | J_- | -\frac{1}{2} \rangle & \langle \frac{1}{2} | J_- | -\frac{3}{2} \rangle \\ \langle -\frac{1}{2} | J_- | \frac{3}{2} \rangle & \langle -\frac{1}{2} | J_- | \frac{1}{2} \rangle & \langle -\frac{1}{2} | J_- | -\frac{1}{2} \rangle & \langle -\frac{1}{2} | J_- | -\frac{3}{2} \rangle \\ \langle -\frac{3}{2} | J_- | \frac{3}{2} \rangle & \langle -\frac{3}{2} | J_- | \frac{1}{2} \rangle & \langle -\frac{3}{2} | J_- | -\frac{1}{2} \rangle & \langle -\frac{3}{2} | J_- | -\frac{3}{2} \rangle \end{pmatrix}$$

$$J_-^{3/2} = \hbar \begin{pmatrix} 0 & 0 & 0 & 0 \\ \sqrt{(\frac{3}{2} + \frac{3}{2})(\frac{3}{2} - \frac{3}{2} + 1)} = \sqrt{3} & 0 & 0 & 0 \\ 0 & \sqrt{(\frac{3}{2} + \frac{1}{2})(\frac{3}{2} - \frac{1}{2} + 1)} = 2 & 0 & 0 \\ 0 & 0 & \sqrt{(\frac{3}{2} - \frac{1}{2})(\frac{3}{2} + \frac{1}{2} + 1)} = \sqrt{3} & 0 \end{pmatrix}$$

$$J_x^{3/2} = \frac{1}{2} (J_+^{3/2} + J_-^{3/2}) = \frac{\hbar}{2} \begin{pmatrix} 0 & \sqrt{3} & 0 & 0 \\ \sqrt{3} & 0 & 2 & 0 \\ 0 & 2 & 0 & \sqrt{3} \\ 0 & 0 & \sqrt{3} & 0 \end{pmatrix}$$

$$J_y^{3/2} = \frac{1}{2i} (J_+^{3/2} - J_-^{3/2}) = \frac{\hbar}{2i} \begin{pmatrix} 0 & \sqrt{3} & 0 & 0 \\ -\sqrt{3} & 0 & 2 & 0 \\ 0 & -2 & 0 & \sqrt{3} \\ 0 & 0 & -\sqrt{3} & 0 \end{pmatrix}$$

with

$$J_z^{3/2} = \frac{\hbar}{2} \begin{pmatrix} 3 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -3 \end{pmatrix}$$

5) show that in a state with definite value of L_z , $\langle L_x, y \rangle = 0$
 $\langle L_x \rangle \langle L_y \rangle = 0$

If in eigenstate with L_z , then $L_z |m\rangle = m |m\rangle$
or $\langle m | L_z = \langle m | m$

know

$$[L_z, L_x] = i\hbar L_y$$

$$\hookrightarrow \frac{1}{i\hbar} (L_z L_x - L_x L_z) = L_y$$

$$\frac{1}{i\hbar} \langle m | L_z L_x | m \rangle - \langle m | L_x L_z | m \rangle = \langle m | L_y | m \rangle$$

$$\hookrightarrow \frac{1}{i\hbar} \langle m | m L_x | m \rangle - \langle m | L_x m | m \rangle$$

$$\hookrightarrow \frac{1}{i\hbar} m \left(\underbrace{\langle m | L_x | m \rangle - \langle m | L_x | m \rangle}_{=0} \right) = \langle L_y \rangle$$

$$\hookrightarrow \boxed{\langle L_y \rangle = 0}$$

similarly, we know $[L_y, L_z] = i\hbar L_x$

$$\text{then } \frac{1}{i\hbar} \langle m | L_y L_z - L_z L_y | m \rangle = \langle L_x \rangle$$

$$\hookrightarrow \frac{1}{i\hbar} \underbrace{\langle m | L_y m - m L_y | m \rangle}_{=0} = \langle L_x \rangle$$

$$\boxed{\langle L_x \rangle = 0}$$