

Developing SPH Based Modeling of Volcanic Plumes using Modified SPH Schemes

Zhixuan Cao¹ Abani Patra^{1,2}

¹Department of Mechanical and Aerospace
University at Buffalo, Buffalo, New York, U.S.A.

²Center for Computational Research
University at Buffalo, Buffalo, New York, U.S.A.

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Outline

- 1 Motivation for Choosing SPH
- 2 Physics Model
- 3 SPH Discretization
- 4 Verification and Validation
- 5 Data Structure and Parallelism

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Volcanic plume development is essentially a multi-phase, turbulent mixing process coupled with heat transfer and other microphysics processes without pre-defined boundaries. SPH (Smoothed particle hydrodynamics), as a mesh-less method, is suitable for such problems for several reasons:

- SPH is able to automatically construct the interface.
- Multiphase modeling is easily accomplished using SPH
- Adding of new physics and new phases is easier in terms of programming in SPH
- With very limited global communication requirements, SPH solvers can scale better for parallel computing than grid based ones.

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The governing equations and EOS for closing the system of equations:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0 \quad (1)$$

$$\frac{\partial \rho \xi}{\partial t} + \nabla \cdot (\rho \xi \mathbf{v}) = 0 \quad (2)$$

$$\frac{\partial \rho \mathbf{v}}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} + p \mathbf{I}) = \rho \mathbf{g} \quad (3)$$

$$\frac{\partial \rho E}{\partial t} + \nabla \cdot [(\rho E + p) \mathbf{v}] = \rho \mathbf{g} \cdot \mathbf{v} \quad (4)$$

$$p = (\gamma_m - 1) \rho e \quad (5)$$

Specific heat ratio γ_m is updated based on massfraction of erupted material:

$$\gamma_m = R_m / C_{vm} + 1 \quad (6)$$

$$R_m = \xi_g R_g + \xi_a R_a \quad (7)$$

$$C_{vm} = \xi_s C_{vs} + \xi_g C_{vg} + \xi_a C_{va} \quad (8)$$

$$\xi_a = 1 - \xi \quad (9)$$

$$\xi_g = \xi \cdot \xi_{g0} \quad (10)$$

$$\xi_s = \xi - \xi_g \quad (11)$$

Boundary Conditions

Eruption BC

$$\rho = \text{const} = p / (R_m T) \quad (12)$$

$$\xi = \text{const} = 1 \quad (13)$$

$$\mathbf{v} = \text{const} = \{u, v, w\}^T \quad (14)$$

$$\frac{\partial e}{\partial n} = \dot{M}e / (\pi r^2) \quad (15)$$

No-slip wall BC

$$\frac{\partial \rho}{\partial n} = \text{const} = 0 \quad (16)$$

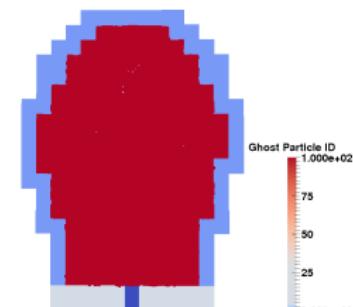
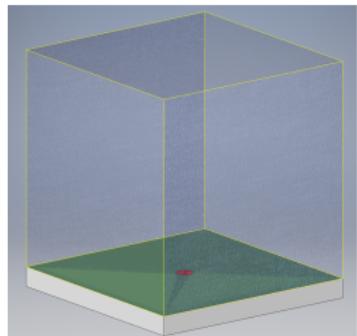
$$\frac{\partial \xi}{\partial n} = \text{const} = 0 \quad (17)$$

$$\mathbf{v} = \text{const} = \{0, 0, 0\}^T \quad (18)$$

$$\frac{\partial e}{\partial n} = 0 \quad (19)$$

Pressure outlet BC

$$p = p_a(z) \quad (20)$$



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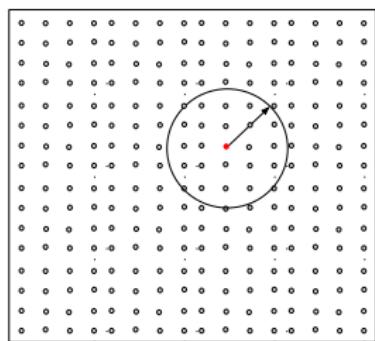
Any function A and its derivative can be approximated by:

$$\langle A(\mathbf{x}) \rangle \approx \sum_b m_b \frac{A_b}{\rho_b} w(\mathbf{x} - \mathbf{x}_b, h) \quad (21)$$

$$\langle \nabla A(\mathbf{x}) \rangle \approx \sum_b m_b \frac{A_b}{\rho_b} \nabla w(\mathbf{x} - \mathbf{x}_b, h) \quad (22)$$

Truncated Gaussian kernel functions

$$w(\mathbf{x} - \mathbf{x}_b) = \begin{cases} \frac{1}{(h\sqrt{\pi})^d} \exp \left[-\left(\frac{\mathbf{x} - \mathbf{x}_b}{h} \right)^2 \right] & |\mathbf{x} - \mathbf{x}_b| \leq 3h \\ 0 & \text{Otherwise} \end{cases} \quad (23)$$



We adopt the following artificial viscosity developed based on Riemann Solvers,:

$$\Pi_{ab}^{\beta} = \begin{cases} \frac{-\alpha \mu_{ab} \bar{c}_{ab} + \beta \mu_{ab}^2}{\bar{\rho}_{ab}} & \mathbf{v}_{ab} \cdot \mathbf{x}_{ab} < 0 \\ 0 & \mathbf{v}_{ab} \cdot \mathbf{x}_{ab} > 0 \end{cases} \quad (24)$$

Where

$$\mu_{ab} = \frac{h \mathbf{v}_{ab} \cdot \mathbf{x}_{ab}}{\mathbf{x}_{ab}^2 + (\eta h)^2} \quad (25)$$

Time step is constrained by CFL condition

$$\Delta t = \text{CFL} \min_a \left\{ \frac{\left[\frac{m_a}{\rho_a} \right]^{\frac{1}{d}}}{c_a} \right\} \quad (26)$$

The classical SPH method was known to suffer from tensile instability and boundary deficiency. To address these difficulties, we adopted a corrected SPH formulation.

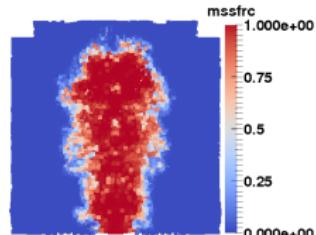
$$A_a = \frac{\sum_b m_b \frac{A_b}{\rho_b} w(x_a - x_b, h)}{\sum_b m_b \frac{1}{\rho_b} w(x_a - x_b, h)} \quad (27)$$

$$\nabla A_a = \frac{\sum_b m_b \frac{A_b - A_a}{\rho_b} \nabla_a w(x_a - x_b, h)}{\sum_b m_b \frac{x_b - x_a}{\rho_b} \nabla_a w(x_a - x_b, h)} \quad (28)$$

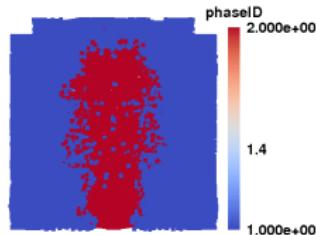
Multiphase SPH

In our model, only density needs to be updated respectively for each phase:

$$\langle \rho_{\alpha}^a \rangle = \frac{\sum m_b w_{\alpha b}(h)}{\sum \frac{m_b}{\rho_b} w_{\alpha b}(h) + \sum \frac{m_j}{\rho_j} w_{\alpha j}(h)} \quad (29)$$



$$\langle \rho_{\alpha}^{sg} \rangle = \frac{\sum_j m_j w_{\alpha j}(h)}{\sum \frac{m_b}{\rho_b} w_{\alpha b}(h) + \sum \frac{m_j}{\rho_j} w_{\alpha j}(h)} \quad (30)$$



- The mass discontinuity and density discontinuity issues in traditional formulations of SPH is fixed by this new formulation
- In areas far away from the interface, updating of density is exactly the same as that for single phase flow

To capture sub-particle scale momentum and energy exchange due to turbulence, we adopt a LANS type turbulence model, the $SPH - \varepsilon$ turbulence model and extend it for compressible flow.

Velocity is smoothed (averaged) by

$$\hat{\mathbf{v}}(\mathbf{x}) = \int \mathbf{v}(\mathbf{x}') G(|\mathbf{x}' - \mathbf{x}|, l) d\mathbf{x}' \quad (31)$$

Where G satisfies:

$$\int G(|\mathbf{x}' - \mathbf{x}|, l) d\mathbf{x}' = 1 \quad (32)$$

$$\lim_{h \rightarrow 0} G(\mathbf{x}' - \mathbf{x}, h) = \delta(\mathbf{x}' - \mathbf{x}) \quad (33)$$

Averaging introduces an extra turbulence stress term:

$$\Phi_{ab} = \frac{\varepsilon}{2} \frac{\mathbf{v}_{ab} \cdot \mathbf{v}_{ab}}{\rho_b} \quad (34)$$

Use Reynolds analogy to model the heat transfer due to turbulence.

$$\kappa = \frac{C_p \mu}{Pr} \quad (35)$$

Instead of an analytical expression, heat transfer coefficient is calculated for each pair of particles.

$$\kappa_{t,ab} = \begin{cases} 0 & \text{if } \mathbf{v}_{ab} \cdot \mathbf{x}_{ab} \leq 0 \\ \frac{\varepsilon \bar{C}_{p,ab} \bar{\rho}_{ab} x_{ab}^2 \mathbf{v}_{ab} \cdot \mathbf{v}_{ab}}{2SPr_t \bar{h}_{ab} \bar{c}_{ab} \mathbf{v}_{ab} \cdot \mathbf{x}_{ab}} & \text{otherwise} \end{cases} \quad (36)$$

$$\begin{aligned} \left\langle \frac{d\mathbf{v}_\alpha}{dt} \right\rangle = & - \sum_b \left[m_b \left(\frac{p_b}{\rho_b^2} + \frac{p_\alpha}{\rho_\alpha^2} + \Pi_{\alpha b}^\beta - \Phi_{\alpha b} \right) \nabla_\alpha w_{\alpha b}(h) \right] \\ & - \sum_j \left[m_j \left(\frac{p_j}{\rho_j^2} + \frac{p_\alpha}{\rho_\alpha^2} + \Pi_{\alpha j}^\beta - \Phi_{\alpha j} \right) \nabla_\alpha w_{\alpha j}(h) \right] + \mathbf{g} \end{aligned} \quad (37)$$

$$\begin{aligned} \left\langle \frac{de_\alpha}{dt} \right\rangle = & 0.5 \sum_b \left[m_b \widehat{\mathbf{v}_{\alpha b}} \left(\frac{p_b}{\rho_b^2} + \frac{p_\alpha}{\rho_\alpha^2} + \Pi_{\alpha b}^\beta - \Phi_{\alpha b} \right) \nabla_\alpha w_{\alpha b}(h) \right] \\ & + 2 \sum_b \frac{m_b}{\rho_\alpha \rho_b} \kappa_{t,\alpha b} (T_\alpha - T_b) F_{\alpha b}(h) \\ & + 0.5 \sum_j \left[m_j \widehat{\mathbf{v}_{\alpha b}} \left(\frac{p_j}{\rho_j^2} + \frac{p_\alpha}{\rho_\alpha^2} + \Pi_{\alpha j}^\beta - \Phi_{\alpha j} \right) \nabla_\alpha w_{\alpha j}(h) \right] \\ & + 2 \sum_j \frac{m_j}{\rho_\alpha \rho_j} \kappa_{t,\alpha j} (T_\alpha - T_j) F_{\alpha j}(h) \end{aligned} \quad (38)$$

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Verification and Validation

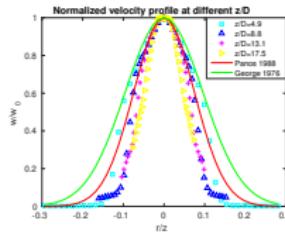
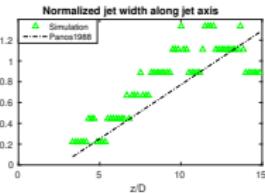
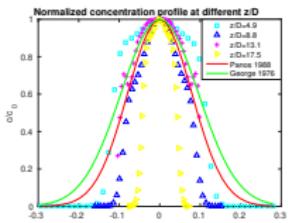
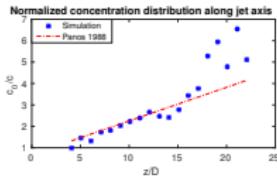


Figure: JPUE test verification: dimensionless velocity and concentration distribution along centralline and across the cross-section.

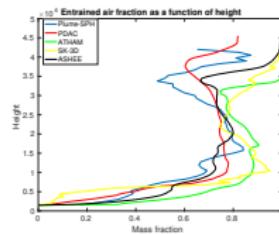
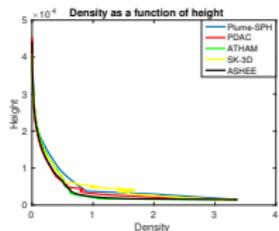
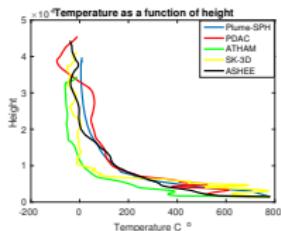
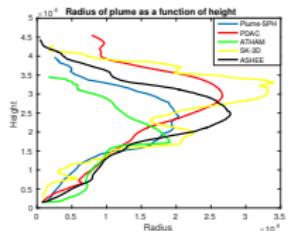


Figure: Simulation of Pinatubo (Philippines, 15 June 1991) : integrated are profiles compared with existing mesh-based models

Animation of Pinatubo Eruption

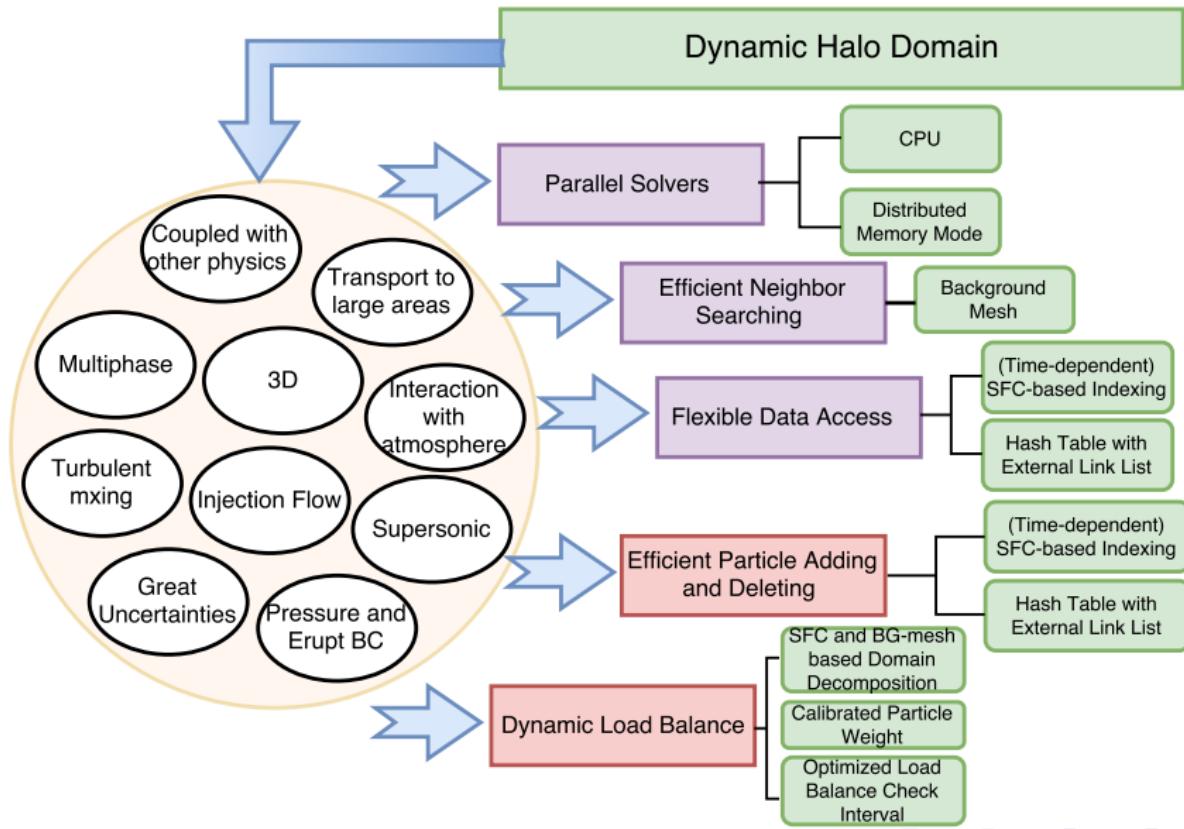
Eruption condition, atmosphere, material properties are mimic Pinatubo eruption (Philippines, 15 June 1991).

Eruption conditions: vent velocity is $275 \text{ m} \cdot \text{s}^{-1}$, vent gas mass fraction is 0.05, vent temperature is 1053 K , vent height is 1500 km , mass discharge rate is $1.5 \times 10^9 \text{ kg} \cdot \text{s}^{-1}$.

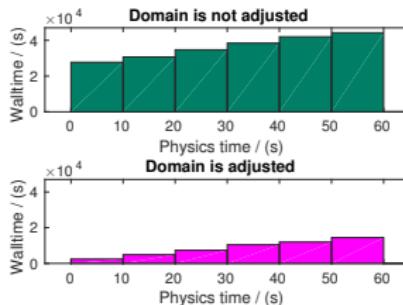
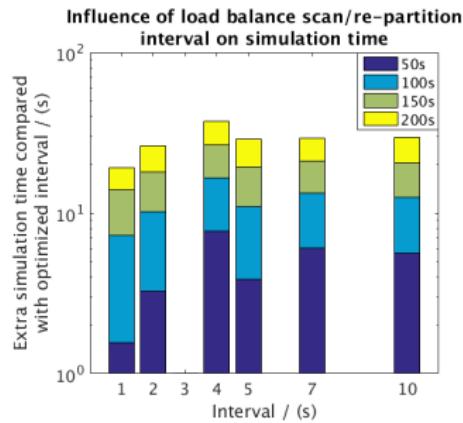
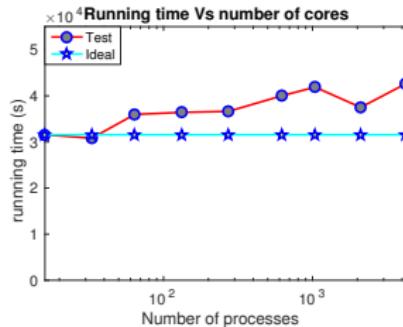
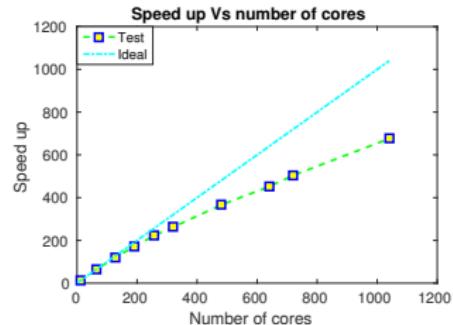
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Requirement and Our Strategies



Scalability and Computational Performance



Summary

- We present an initial effort and results towards developing a first principle based plume model with comprehensive physics, adopting proper numerical tools and high performance computing.
- More advanced numerical techniques, such as adaptive particle size, Godunov-SPH, semi-explicit time advancing scheme and better data management strategies and algorithms are on our list to exploit in the future.

Thank you!
Questions are welcome.

zhixuanc@buffalo.edu