Absense of Positive Eigenvalues for the Linearized Elasticity System in the Half Space

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Abstract. In this paper, we prove that the linearized elasticity system in the half-space with traction free boundry has no eigenvalues.

1. Introduction

section1

In this paper, we consider the linearized and isotropic elasticity system defined on an unbounded domain $\Omega = \mathbb{R}^2_+ \setminus \overline{D}$ with traction free surface $\Gamma_0 = \{(x_1, x_2)^T : x_1 \in \mathbb{R}, x_2 = 0\}$, where $D \subseteq \mathbb{R}^2_+$ is a bounded Lipschitz domain with the unit outer normal ν to its boundary Γ_D . We study the eigenvalues of the following elastic scattering problem in the isotropic homogeneous medium half space with $Lam\acute{e}$ constant λ and μ and constant density $\rho \equiv 1$:

$$\nabla \cdot \sigma(\mathbf{u}) + \rho \omega^2 \mathbf{u} = f \qquad \text{in } \mathbb{R}^2_+ \setminus \bar{D} \qquad (1.1) \quad \boxed{\text{elastic_eq}}$$

$$\mathbf{u} = 0 \text{ on } \Gamma_D \text{ and } \sigma(\mathbf{u}) \cdot e_2 = 0 \text{ on } \Gamma_0$$
 (1.2) elastic_bd

together with the constitutive relation (Hookes law)

$$\sigma(\mathbf{u}) = 2\mu\varepsilon(\mathbf{u}) + \lambda \operatorname{div}\mathbf{u}\mathbb{I}$$
$$\varepsilon(\mathbf{u}) = \frac{1}{2}(\nabla\mathbf{u} + (\nabla\mathbf{u})^T)$$

where ω is the circular frequency, $\mathbf{u}(x_1, x_2) = (u_1(x), u_2(x))^T \in \mathbb{C}^2$ denotes the displacement fields and $\sigma(u)$ is the stress tensor. We also need to define the surface traction $T_x^n(\cdot)$ on the normal direction n,

$$T_x^n \mathbf{u}(x) := \sigma \cdot n = 2\mu \frac{\partial \mathbf{u}}{\partial n} + \lambda n \operatorname{div} \mathbf{u} + \mu n \times \operatorname{curl} \mathbf{u}$$

For simplicity, let's introduce $Lam\acute{e}$ operator Δ_e as

$$\Delta_e \mathbf{u} = (\lambda + 2\mu)\nabla\nabla \cdot \mathbf{u} - \mu\nabla \times \nabla \times u = \nabla \cdot \sigma(\mathbf{u})$$

We remark that the results in this paper can be extended to other boundary conditions such as Neumann or mixed boundary conditions on Γ_D , or even to penetrable obstacle.

In order to complete the definition of the proble, we introduce the domain of the operator Δ_e

$$\mathcal{D}(\Delta_e, S) = \{ v \in H^1(S) : \Delta_e v \in L^2(S), \sigma(v)e_2 = 0 \text{ on } \Gamma_0 \}$$

where S is an unbounded domain in \mathbb{R}^2_+ . For the elasticity system, the study of eigenvalue is little. The layout of the paper is as follows. In section 2

2. Absence of Positive Eigenvalues

In this section. Throughout the paper, we will assume that for $z \in \mathbb{C}$, $z^{1/2}$ is the analytic branch of \sqrt{z} such that $\text{Im}(z^{1/2}) \geq 0$. This corresponds to the right half real axis as the branch cut in the complex plane. For $z = z_1 + \mathbf{i}z_2, z_1, z_2 \in \mathbb{R}$, we have

$$z^{1/2} = sgn(z_2)\sqrt{\frac{|z| + z_1}{2}} + i\sqrt{\frac{|z| - z_1}{2}}$$
 (2.1) [convention_1]

For z on the right half real axis, we take $z^{1/2}$ as the limit of $(z + i\varepsilon)^{1/2}$ as $\varepsilon \to 0^+$.

Theorem 2.1 Let $\omega \in \mathbb{R}$ and \mathbf{u} statisfy the equations (1.1-1.2) in $\mathcal{D}(\Delta_e, \Omega)$, then we assert $\mathbf{u} = 0$.

Lemma 2.1 The Rayleigh equation $\delta(\xi) = 0$ has only two zeros $\pm k_R$, $k_R > k_s$, in the complex plane.

Lemma 2.2 Suppose that $f \in L^2(\mathbb{R}^2_+)$ with compact support in $B \subsetneq \mathbb{R} \times (h, +\infty)$, h > 0. Let $\omega \in \mathbb{R}$ and $\mathbf{u} \in \mathcal{D}(\Delta_e, \mathbb{R}^2_+)$ such that:

$$\Delta_e \mathbf{u} + \omega^2 \mathbf{u} = f \tag{2.2}$$

then we assert $\mathbf{u} = 0$ in $(\mathbb{R} \times (h, +\infty))$.

Proof. Let $\mathcal{F}_{x_1}(\cdot): L^2(R_+^2) \to L^2(R_+^2)$ be the partial Fourier transfor given by $\hat{g} := \mathcal{F}_{x_1}(g) := \int_{\mathbb{R}} g(x_1, x_2) e^{\mathbf{i}x_1\xi} dx_1$. By taking the Fourier transform of (2.2) and (1.2), we obtain ODEs for x_2 in R_+

$$\mu \frac{d^2 \hat{u}_1}{dx_2^2} + \mathbf{i}(\lambda + \mu)\xi \frac{d\hat{u}_2}{dx_2} + (\omega^2 - (\lambda + 2\mu)\xi^2)\hat{u}_1 = \hat{f}_1$$
 (2.3) [pp3]

$$(\lambda + 2\mu)\frac{d^2\hat{u}_2}{dx_2^2} + \mathbf{i}(\lambda + \mu)\xi\frac{d\hat{u}_1}{dx_2} + (\omega^2 - \mu\xi^2)\hat{u}_2 = \hat{f}_2$$
 (2.4) \[\text{pp4}\]

and the boundary coditions on $x_2 = 0$ are

$$\mu \frac{d\hat{u}_1}{dx_2} + \mathbf{i}\mu \xi \hat{u}_2 = 0 \tag{2.5}$$

$$(\lambda + 2\mu)\frac{d\hat{u}_2}{dx_2} + \mathbf{i}\lambda\xi\hat{u}_1 = 0 \tag{2.6}$$

In order to work with real coefficient, we use the change of variables:

$$v_1 = \mathbf{i}\hat{u}_1, \quad v_2 = \hat{u}_2, \quad \mathbf{v} = (v_1, v_2)^T$$

 $q_1 = \mathbf{i}\hat{f}_1, \quad q_2 = \hat{f}_2, \quad \mathbf{q} = (q_1, q_2)^T$

Then we have the following equations:

$$\left[\mathbb{A}_1 \frac{d^2}{dx_2^2} + (\mathbb{A}_2 - (\mathbb{A}_2)^T)\xi \frac{d}{dx_2} - \mathbb{A}_3\xi^2 + \omega^2\right] \mathbf{v} = \mathbf{q} \quad \text{in } \mathbb{R}_+$$
 (2.7) eq3

$$\left(\mathbb{A}_1 \frac{d}{dx_2} + \mathbb{A}_2 \xi\right) \quad \mathbf{v} = 0 \qquad \qquad \text{on } x_2 = 0 \tag{2.8}$$

where

$$\mathbb{A}_1 = \begin{pmatrix} \mu & 0 \\ 0 & \lambda + 2\mu \end{pmatrix}, \qquad \mathbb{A}_2 = \begin{pmatrix} 0 & -\mu \\ \lambda & 0 \end{pmatrix}, \qquad \mathbb{A}_3 = \begin{pmatrix} \lambda + 2\mu & 0 \\ 0 & \mu \end{pmatrix}$$

Let \mathbf{w} be the solution of the following equations:

$$[\mathbb{A}_1 \frac{d^2}{dx_2^2} + (\mathbb{A}_2 - (\mathbb{A}_2)^T) \xi \frac{d}{dx_2} - \mathbb{A}_3 \xi^2 + \omega^2] \mathbf{w} = \mathbf{q} \quad \text{in} \quad (0, h)$$

$$\mathbf{w} = 0, \quad (\mathbb{A}_1 \frac{d}{dx_2} + \mathbb{A}_2 \xi) \quad \mathbf{w} = 0 \quad \text{on} \quad x_2 = h$$

It is easy to transform above equations into a simpler form by variables substitution $\mathbf{W} = (\mathbf{w}, (\mathbb{A}_1 \frac{d}{dx_2} + \mathbb{A}_2 \xi) \mathbf{w})^T, \mathbf{Q} = (0, 0, \mathbf{q})^T$

$$\frac{d}{dx_2}\mathbf{W} = \mathbf{A}\mathbf{W} + \mathbf{Q} \quad \text{in } (0, h)$$

$$\mathbf{W} = 0 \quad \text{on } x_2 = h$$

where

$$\mathbb{A} = \begin{pmatrix} -\mathbb{A}_1^{-1} \mathbb{A}_2 \xi & \mathbb{A}_1^{-1} \\ -\mathbb{A}_2^T \mathbb{A}_1^{-1} \mathbb{A}_2 \xi^2 + \mathbb{A}_3 \xi^2 - \omega^2 & \mathbb{A}_2^T \mathbb{A}_1^{-1} \xi \end{pmatrix}$$

By the standard arguments in ODEs, we can obtain

$$\mathbf{W}(\xi, x_2) = -\Phi(\xi, x_2) \int_h^{x_2} \Phi^{-1}(\xi, t) \mathbf{Q}(\xi, t) dt$$

where

$$\Phi(\xi,t) = \begin{pmatrix} -\mu_s(\xi)e^{\mathbf{i}\mu_s t} & -\xi e^{\mathbf{i}\mu_p(\xi)t} & -\mu_s(\xi)e^{-\mathbf{i}\mu_s(\xi)t} & \xi e^{-\mathbf{i}\mu_p t} \\ -\mathbf{i}\xi e^{\mathbf{i}\mu_s(\xi)t} & \mathbf{i}\mu_p(\xi)e^{\mathbf{i}\mu_p(\xi)t} & \mathbf{i}\xi e^{-\mathbf{i}\mu_s(\xi)t} & \mathbf{i}\mu_p(\xi)e^{-\mathbf{i}\mu_p(\xi)t} \\ -\mathbf{i}\mu\beta(\xi)e^{\mathbf{i}\mu_s t} & -2\mathbf{i}\mu\xi\mu_p(\xi)e^{\mathbf{i}\mu_p(\xi)t} & \mathbf{i}\mu\beta(\xi)e^{-\mathbf{i}\mu_s(\xi)t} & -2\mathbf{i}\mu\xi\mu_p(\xi)e^{-\mathbf{i}\mu_p t} \\ 2\mu\xi\mu_s(\xi)e^{\mathbf{i}\mu_s t} & -\mu\beta(\xi)e^{\mathbf{i}\mu_p(\xi)t} & 2\mu\xi\mu_s(\xi)e^{-\mathbf{i}\mu_s t} & \mu\beta(\xi)e^{-\mathbf{i}\mu_p(\xi)t} \end{pmatrix}$$

Here $k_p = \omega/\sqrt{\lambda + 2\mu}$, $k_s = \omega/\sqrt{\mu}$ are wave number of p-wave and s-wave, and $\mu_{\alpha} = (k_{\alpha}^2 - \xi^2)^{1/2}$ for $\alpha = s, p$.

We extend $\mathbf{w}(\xi, x_2)$ by zero in (h, ∞) . Therefore, $\mathbf{w}(\xi, x_2)$ satisfy equation 2.7 in \mathbb{R}_+ . Since $\Phi(\xi, t)$ are analytic w.r.t ξ in $\mathbb{R}\setminus\{k_p, k_s\}$ and $f(\mathbf{x})$ have compact support, we deduce that for almost every $\xi \in \mathbb{R}$, $\mathbf{w}(\xi, x_2)$ are analytic and so $(\mathbb{A}_1 \frac{d}{dx_2} + \mathbb{A}_2 \xi)\mathbf{w}$ are.

We set $\mathbf{U} = \mathbf{v} - \mathbf{w}$ and $\mathbf{U} = (U_1, U_2)^T$. Then \mathbf{U} satisfy the following Cauchy problem:

$$\left[\mathbb{A}_{1} \frac{d^{2}}{dx_{2}^{2}} + (\mathbb{A}_{2} - (\mathbb{A}_{2})^{T})\xi \frac{d}{dx_{2}} - \mathbb{A}_{3}\xi^{2} + \omega^{2}\right]\mathbf{U} = 0 \quad \text{in } \mathbb{R}_{+}$$

$$(2.9) \quad \boxed{\text{eq4}}$$

$$\left(\mathbb{A}_1 \frac{d}{dx_2} + \mathbb{A}_2 \xi\right) \mathbf{U} = \left(\mathbb{A}_1 \frac{d}{dx_2} + \mathbb{A}_2 \xi\right) \mathbf{w} \qquad \text{on } x_2 = 0 \qquad (2.10)$$

Since the coefficients of above equations are constants, we can represent $\mathbf{U}(\xi, x_2)$ in the following form:

$$\mathbf{U}(\xi, x_2) = c_1(\xi) \begin{pmatrix} -\mu_s \\ -\mathbf{i}\xi \end{pmatrix} e^{\mathbf{i}\mu_s x_2} + c_2(\xi) \begin{pmatrix} -\xi \\ \mathbf{i}\mu_p \end{pmatrix} e^{\mathbf{i}\mu_p x_2} + c_3(\xi) \begin{pmatrix} -\mu_s \\ \mathbf{i}\xi \end{pmatrix} e^{-\mathbf{i}\mu_s x_2} + c_4(\xi) \begin{pmatrix} \xi \\ \mathbf{i}\mu_p \end{pmatrix} e^{-\mathbf{i}\mu_p x_2}$$

If $\xi^2 \leq k_p^2$, then it's simple to see that $\mathbf{U} = 0$ in $L_{x_2}^2(\mathbb{R}_+)$. So, for $\xi^2 < k_p^2$, $(\mathbb{A}_1 \frac{d}{dx_2} + \mathbb{A}_2 \xi) \mathbf{U} = 0$ which implies $(\mathbb{A}_1 \frac{d}{dx_2} + \mathbb{A}_2 \xi) \mathbf{w} = 0$. Since $(\mathbb{A}_1 \frac{d}{dx_2} + \mathbb{A}_2 \xi) \mathbf{w}$ are analytic for almost every $\xi \in \mathbb{R}$, we deduce that

$$\left(\mathbb{A}_{1} \frac{d}{dx_{2}} + \mathbb{A}_{2} \xi\right) \mathbf{U} = 0 \quad \text{on} \quad x_{2} = 0$$

$$(2.11) \quad \boxed{\text{bd_1}}$$

for almost every $\xi \in \mathbb{R}$. Therefore, we can obtain

$$\mathbf{U}(\xi, x_2) = \begin{cases} c(\xi) \begin{pmatrix} -\xi \\ \mathbf{i}\mu_p \end{pmatrix} e^{\mathbf{i}\mu_p x_2}, & k_p^2 < \xi^2 \le k_s^2 \\ c_1(\xi) \begin{pmatrix} -\mu_s \\ -\mathbf{i}\xi \end{pmatrix} e^{\mathbf{i}\mu_s x_2} + c_2(\xi) \begin{pmatrix} -\xi \\ \mathbf{i}\mu_p \end{pmatrix} e^{\mathbf{i}\mu_p x_2}, & \xi^2 > k_s^2 \end{cases}$$

$$(\mathbb{A}_1 \frac{d}{dx_2} + \mathbb{A}_2 \xi) \mathbf{U} = \begin{cases} c(\xi) \begin{pmatrix} -2\mathbf{i}\mu\xi\mu_p \\ -\mu\beta \end{pmatrix} e^{\mathbf{i}\mu_p x_2}, & k_p^2 < \xi^2 \le k_s^2 \\ c_1(\xi) \begin{pmatrix} -\mathbf{i}\mu\beta \\ 2\mu\xi\mu_s \end{pmatrix} e^{\mathbf{i}\mu_s x_2} + c_2(\xi) \begin{pmatrix} -2\mathbf{i}\mu\xi\mu_p \\ -\mu\beta \end{pmatrix} e^{\mathbf{i}\mu_p x_2}, & \xi^2 > k_s^2 \end{cases}$$

By boundary condition $\overset{\text{bd}}{2.11}$, we have $c(\xi)=0$ for $k_p^2<\xi^2\leq k_s^2$ and

$$\det \begin{pmatrix} -\mathbf{i}\mu\beta & -2\mathbf{i}\mu\xi\mu_p \\ 2\mu\xi\mu_s & -\mu\beta \end{pmatrix} = -\mathbf{i}\mu(\beta^2 + 4\xi^2\mu_s\mu_p) = 0 \quad \text{for } \xi^2 > k_s^2$$
 (2.12)

Therefore, by lemma $\frac{\text{llem 2.1}}{2.1 \text{ we}}$ have $\mathbf{U}(\xi, x_2) = 0$ for almost every $\xi \in \mathbb{R}$ which implies $\mathbf{v}(\xi, x_2) = 0$ for almost every $\xi \in \mathbb{R}$ and $x_2 \in (h, +\infty)$. This completes the proof by taking the inverse Fourier tranformation of $\hat{\mathbf{u}}(\xi, x_2)$.

proof of Theorem 2.1: Since $D \subseteq \mathbb{R}^2_+$, we can find two concentric circles B_{R_1} , B_{R_2} such that $D \subseteq B_{R_1} \subseteq B_{R_2} \subseteq \mathbb{R}^2_+$. Let $\chi \in C_0^{\infty}(\mathbb{R}^2_+)$ be the cut-off function such that $0 \le \chi \le 1$, $\chi = 0$ in B_{R_1} , and $\chi = 1$ outside of B_{R_2} . Let $v = \chi u$. Then v satisfies (2.2) with $f = \sigma(u)\nabla\chi + (\lambda + \mu)(\nabla^2\chi u + \nabla u\nabla\chi) + \mu\Delta\chi u + \mu \text{div} u\nabla\chi$, where $\nabla^2\chi$ is the Hessian matrix of χ . Clearly q has compact support. By lemme 2.2, we have u = v = 0 in $\mathbb{R} \times (h, +\infty)$. Finally, the unique continuation principle implies that u = 0 in \mathbb{R}^2_+ . This completes the proof.

References