1. Estimate of Dirichlet Green Tensor

We need the following slight generalization of Van der Corput lemma for the oscillatory integral [2, P.152].

Lemma 1.1 Let $-\infty < a < b < \infty$, and u is a C^k function u in (a, b).

1. If $|u'(t)| \ge 1$ for $t \in (a,b)$ and u' is monotone in (a,b), then for any $\phi(t)$ in (a,b) with integrable derivatives

$$\left| \int_a^b e^{\mathbf{i}\lambda u(t)} \phi(t) dt \right| \le 3\lambda^{-1} \left[|\phi(b)| + \int_a^b |\phi'(t)| dt \right].$$

2. For all $k \geq 2$, if $|u^{(k)}(t)| \geq 1$ for $t \in (a,b)$, then for any $\phi(t)$ in (a,b) with integrable derivatives

$$\left| \int_{a}^{b} e^{\mathbf{i}\lambda u(t)} \phi(t) dt \right| \leq 12k\lambda^{-1/k} \left[|\phi(b)| + \int_{a}^{b} |\phi'(t)| dt \right].$$

Proof. The assertion can be proved by extending the Van der Corptut lemma in [2]. Here we omit the details.

We recall following lemma, see e.g. [3]:

Lemma 1.2 Let $F(\rho, a) = \int_0^a t^{\alpha - 1} f(t) e^{-i\rho t} dt$ where $0 < a \le +\infty$, $0 < \alpha < 1$, $\rho > 0$ and $t^{\alpha - 1} f \in L^1(0, a)$, then we have

$$|F(\rho, a)| \le C(\frac{1}{\rho^{\alpha}} f(0) + \frac{1}{\rho} (a^{\alpha - 1} f(a) + |t^{\alpha - 1} f|_{L^{1}(0, a)})$$
(1.1)

Proof. Put

$$g_0(t) = t^{\alpha - 1} e^{-\mathbf{i}\rho t} \tag{1.2}$$

and define

$$g_1(t) = -\int_t^{t-i\infty} x^{\alpha-1} e^{-i\rho x} dx \tag{1.3}$$

where the path of integration is the vertical line $x = t - \mathbf{i}y$, $y \ge 0$. It is easy to show that $g_1(t)' = g_0(t)$. Substituting $x = t - \mathbf{i}y$ into $g_1(t)$, we have

$$g_1(t) = \mathbf{i} \int_0^\infty (t - \mathbf{i}y)^{\alpha - 1} e^{-\mathbf{i}\rho t} e^{-\rho y} dy$$
 (1.4)

Upon integration by parts, we have

$$F(\rho, a) = \int_0^a f(t)dg_1(t)$$

$$= e^{-\mathbf{i}\frac{\alpha\pi}{2}}f(0)\Gamma(\alpha)\frac{1}{\rho^{\alpha}} + f(a)g_1(a) - \int_0^a f'(t)g_1(t)dt$$

$$= e^{-\mathbf{i}\frac{\alpha\pi}{2}}f(0)\Gamma(\alpha)\frac{1}{\rho^{\alpha}} - \mathbf{i}\int_0^\infty e^{-\rho y}dy\int_0^a f'(t)(t - \mathbf{i}y)^{\alpha - 1}e^{-\mathbf{i}\rho t}dt$$

Let

$$h(y) = \int_0^a f'(t)(t - \mathbf{i}y)^{\alpha - 1} e^{-\mathbf{i}\rho t} dt$$

and observe that

$$|h(y)| \le \int_0^a |f'(t)|(t^2 + y^2)^{\frac{\alpha - 1}{2}} dt$$

Lemma 1.3 Let $F(\rho, a) = \int_0^a t^{-1/2} f(t) e^{-i\rho t} dt$ where $0 < a \le +\infty$ and $\rho > 0$, then we have

$$|F(\rho, a) - e^{-i\frac{\pi}{4}} f(0)\Gamma(1/2) \frac{1}{\rho^{1/2}}|$$
 (1.5)

$$\leq C\left(\int_{0}^{\infty} e^{-\rho y} dy \int_{0}^{a} |f'(t)| (t^{2} + y^{2})^{-\frac{1}{4}} dt + \frac{1}{\rho} a^{-1/2} f(a)\right)$$
(1.6)

Proof. Put

$$g_0(t) = t^{-1/2}e^{-\mathbf{i}\rho t}$$
 (1.7)

and define

$$g_1(t) = -\int_t^{t-i\infty} x^{-1/2} e^{-i\rho x} dx \tag{1.8}$$

where the path of integration is the vertical line $x = t - \mathbf{i}y, y \ge 0$. It is easy to show that $g_1(t) = g_0(t)$. Substituting $x = t - \mathbf{i}y$ into $g_1(t)$, we have

$$g_1(t) = \mathbf{i} \int_0^\infty (t - \mathbf{i}y)^{-1/2} e^{-\mathbf{i}\rho t} e^{-\rho y} dy$$
(1.9)

Upon integration by parts, we have

$$F(\rho, a) = \int_0^a f(t)dg_1(t)$$

$$= e^{-\mathbf{i}\frac{\pi}{4}} f(0)\Gamma(1/2) \frac{1}{\rho^{1/2}} + f(a)g_1(a) - \int_0^a f'(t)g_1(t)dt$$

$$= e^{-\mathbf{i}\frac{\pi}{4}} f(0)\Gamma(1/2) \frac{1}{\rho^{1/2}} + \mathbf{i}f(a) \int_0^\infty (a - \mathbf{i}y)^{-1/2} e^{-\mathbf{i}\rho t} e^{-\rho y} dy$$

$$- \mathbf{i} \int_0^\infty e^{-\rho y} dy \int_0^a f'(t)(t - \mathbf{i}y)^{-1/2} e^{-\mathbf{i}\rho t} dt$$

Let

$$h(y) = \int_0^a f'(t)(t - \mathbf{i}y)^{-1/2} e^{-\mathbf{i}\rho t} dt$$

and observe that

$$|h(y)| \le \int_0^a |f'(t)|(t^2 + y^2)^{-\frac{1}{4}} dt$$

It is easy to see that

$$|g_1(a)| \le a^{-1/2} \int_0^\infty e^{-\rho y} dy \le C \frac{1}{\rho}$$

Lemma 1.4 Assume that $0 < \kappa := \sin \phi_{\kappa} < 1, 0 < \phi_{\kappa} < \pi/2, \ 0 \le \phi \le \pi/2$ and $-\pi/2 < t_1 < t_2 < \pi/2$ satisfy that $\kappa^2 = \sin^2(\phi + t_1) = \sin^2(\phi + t_2)$. Let $f(\theta)$:

$$f(t,\phi) := F(\sin(t+\phi), \cos(t+\phi), (\kappa^2 - \sin^2(t+\phi))^{1/2})$$
(1.10)

be a function in $C^{\infty}(([-\pi/2, \pi/2] \setminus \{t_1, t_2\}) \times [0, \pi/2])$. Moreover, there exits $\epsilon > 0$ such that $f(\theta)$ can be represented as

$$f(t,\phi) = g_1(t,\phi) + g_2(t,\phi)(\kappa^2 - \sin^2(t+\phi))^{1/2})^{N/2}$$
(1.11)

where $g_1, g_2 \in C^{\infty}((\bigcup_{i=1,2} (t_i - \epsilon, t_i + \epsilon)) \times [0, \pi/2]))$ and $N = \pm 1$. Then for any $\rho \geq 1$, we have

$$\left| I(\rho, \phi) := \int_{-\pi/2}^{\pi/2} f(\theta) e^{\mathbf{i}\rho\cos\theta} d\theta - \frac{N+1}{2} \left(\frac{2\pi}{\rho}\right)^{1/2} f(0) e^{\mathbf{i}\rho - \mathbf{i}\pi/4} \right|$$

$$\leq C \frac{1}{\rho^{(2+N)/4}} \tag{1.12}$$

Proof. The proof will be split into two parts about whether ϕ equal to ϕ_{κ} .

If $\phi \neq \phi_{\kappa}$, there exists $0 < \delta < \pi/4$ such that

$$|(\kappa^2 - \sin^2(t+\phi))^{1/2}| > \frac{1}{2}|(\kappa^2 - \sin^2\phi)^{1/2}| \tag{1.13}$$

for any $t \in (-\delta, \delta)$. Let $\chi_{\delta} \in C_0^{\infty}(-\pi/2, \pi/2)$ be the cut-off function with that $0 \le \chi_{\delta} \le 1$, $\chi_{\delta} = 1$ in $(-\delta/2, \delta/2)$ and $\chi_{\delta} = 0$ in $(-\pi/2, \pi/2) \setminus (-\delta, \delta)$. Then we can divide I into two parts such that

$$I = \int_{-\delta}^{\delta} f(t)\chi_{\delta}(t)e^{\mathbf{i}\rho\cos t}dt + \int_{-\pi/2}^{\pi/2} f(t)(1-\chi_{\delta}(t))e^{\mathbf{i}\rho\cos t}dt$$

=: $I_1 + I_2$

Subtitating $t(s) = 2 \arcsin s/2$ for t in I_1 , we can obtain

$$I_{1} = \int_{\mathbb{R}} f(t(s)) \chi_{\delta}(t(s)) \frac{1}{\sqrt{1 - s^{2}/4}} e^{\mathbf{i}\rho} e^{-\mathbf{i}\rho s^{2}/2} ds$$
 (1.14)

$$= \int_{\mathbb{R}} h_{\delta}(s)e^{\mathbf{i}\rho}e^{-\mathbf{i}\rho s^{2}/2}ds \tag{1.15}$$

It is easy to see that $h_{\delta}(s) \in C_0^4(\mathbb{R})$. By the lemma of the stationary phase for quadratic term in [1], we have

$$I_1 = e^{\mathbf{i}\rho} \int_{\mathbb{R}} h_{\delta}(s) e^{-\mathbf{i}\frac{\rho}{2}s^2} ds = e^{\mathbf{i}\rho} \int_{\mathbb{R}} \widehat{h_{\delta}}(y) \alpha(-y) dy$$
 (1.16)

where

$$\alpha(y) = \left(\frac{1}{2\pi\rho}\right)^{1/2} e^{-i\pi/4} e^{\frac{i}{2\rho}y^2} \tag{1.17}$$

$$= \left(\frac{1}{2\pi\rho}\right)^{1/2} e^{-i\pi/4} \left(1 + O\left(\frac{y^2}{\rho}\right)\right) \tag{1.18}$$

Consequently

$$I_{1} = \left(\frac{1}{2\pi\rho}\right)^{1/2} e^{\mathbf{i}\rho - \mathbf{i}\pi/4} \int_{\mathbb{R}} \widehat{h_{\delta}}(y) \left(1 + \frac{1}{\rho}O(y^{2})\right) dy \tag{1.19}$$

Moreover, $\int_{\mathbb{R}} \widehat{h_{\delta}}(y) dy = 2\pi h_{\delta}(0)$ and $|\int_{\mathbb{R}} \widehat{h_{\delta}}(y) y^2 dy| < C$ since $\widehat{h_{\delta}}(y) = O(1/y^4)$. Now, it turns to estimate I_2 .

When N = 1, using integration by parts, we have

$$|I_2| = \left| \int_{(-\frac{\pi}{2}, \frac{\pi}{2}) \setminus (-\frac{\delta}{2}, \frac{\delta}{2})} f(t)(1 - \chi_{\delta}(t)) / \sin t \ de^{\mathbf{i}\rho \cos t} \right|$$
 (1.20)

(1.21)

$$\leq C \frac{1}{\rho} + \left| \int_{\left(-\frac{\pi}{2}, \frac{\pi}{2}\right) \setminus \left(-\frac{\delta}{2}, \frac{\delta}{2}\right)} (f(t)(1 - \chi_{\delta}(t)) / \sin t)' e^{\mathbf{i}\rho \cos t} dt \right|$$
 (1.22)

$$\leq C \frac{1}{\rho} \tag{1.23}$$

From above analysis, we obtain

$$\left| I(\rho, \phi) - \left(\frac{2\pi}{\rho} \right)^{1/2} f(0) e^{\mathbf{i}\rho - \mathbf{i}\pi/4} \right| \le C(\phi) \frac{1}{\rho}$$
(1.24)

When N = -1, we can not use integration by parts again since $f'(\theta)$ is not integrable. However, for any $0 < \lambda_1 < 1$ and $1 < \lambda_2 < 1/\kappa$, there exists $0 < \sigma < \epsilon$, such that $\chi := ((t_1 - \sigma, t_1 + \sigma) \cup (t_2 - \sigma, t_2 + \sigma)) \cap (-\delta, \delta) = 0$, dependent on λ_1, λ_2 and

$$\lambda_1 \kappa < |\sin(t + \phi)| < \lambda_2 \kappa. \tag{1.25}$$

for any $t \in \chi$.

We only analysis the integral on $\chi_1 = (t_1 - \sigma, t_1 + \sigma) \cap [-\pi/2, \pi/2]$ here, which denoted by I_{χ_1} , the proof of I_{χ_2} is similar. It is easy to see that $\sin(t + \phi)$ is monotonic in χ_1 . Without loss of generality, we assume that $\sin(t_1 - \sigma + \phi) < \kappa < \sin(t_1 + \sigma + \phi)$. Let $\sin(t + \phi) = \kappa \sin \theta$ and the implicit mapping from θ to t is denoted by $t(\theta)$ while the inverse mapping by $\theta(t)$, taking the interval χ_1 onto $L_{\theta}: \theta_1 \to \pi/2 \to \pi/2 - \mathbf{i}\theta_2$ where $\sin(t_1 - \sigma + \phi) = \kappa \sin \theta_1, \sin(t_1 + \sigma + \phi) = \kappa \sin(\pi/2 - \mathbf{i}\theta_2)$. By substituting $t(\theta)$ into I_{χ_1} , we have

$$I_{\chi_1} = \int_{t_1 - \sigma}^{t_1 + \sigma} \frac{f(t)(\kappa^2 - \sin^2(t + \phi))^{1/2}}{(\kappa^2 - \sin^2(t + \phi))^{1/2}} e^{i\rho \cos t}$$
(1.26)

$$= \int_{L_{\theta}} \frac{\kappa f(t(\theta)) \cos \theta}{(1 - \kappa^2 \sin^2 \theta)^{1/2}} e^{\mathbf{i}\rho(\cos(t(\theta)))} d\theta$$
 (1.27)

$$= \int_{L_{\theta}} \frac{\kappa g_1(t(\theta)) \cos \theta + g_2(t(\theta))}{(1 - \kappa^2 \sin^2 \theta)^{1/2}} e^{\mathbf{i}\rho(\cos(t(\theta)))} d\theta$$
 (1.28)

$$:= \int_{L_{\theta}} \frac{h(\theta)}{(1 - \kappa^2 \sin^2 \theta)^{1/2}} e^{\mathbf{i}\rho(\cos(t(\theta)))} d\theta \tag{1.29}$$

Observe that $h(\theta)$ and $\partial h/\partial \theta$ are integrable on the path L_{θ} by (1.11). A simple computation show that

$$\frac{dt(\theta)}{d\theta} = \frac{\kappa \cos \theta}{\cos(t+\phi)} \qquad \frac{d^2t(\theta)}{dt^2} = \frac{\kappa^2 \cos^2 \theta \sin(t+\phi) - \kappa \sin \theta \cos^2(t+\phi)}{\cos^3(t+\phi)}$$

Then we can obtain

$$\frac{d\cos t}{d\theta} = \frac{-\kappa\sin t\cos\theta}{\cos(t+\phi)}$$

$$\frac{d^2 \cos t}{d\theta^2} = \frac{d^2 \cos t}{dt^2} \left(\frac{dt}{d\theta}\right)^2 + \frac{d \cos t}{dt} \frac{d^2 t}{d\theta^2}$$

$$= \frac{-\kappa^2 \cos^2 \theta \cos t}{\cos^2(t+\phi)} + \frac{\kappa \sin \theta \cos^2(t+\phi) \sin t - \kappa^2 \cos^2 \theta \sin(t+\phi) \sin t}{\cos^3(t+\phi)}$$

$$= \frac{-\kappa^2 \cos^2 \theta \cos \phi + \kappa \sin \theta \cos^2(t+\phi) \sin t}{\cos^3(t+\phi)}$$

$$= \frac{(\sin^2(t+\phi) - \kappa^2) \cos \phi + \cos^2(t+\phi) \sin(t+\phi) \sin t}{\cos^3(t+\phi)}$$

Since $|\sin t| > |\sin \delta|$ and $1 - \lambda_2^2 \kappa^2 < \cos^2(t + \phi) < 1 - \lambda_1^2 \kappa^2$ for $t \in \chi_1$, it follows that $\theta = \pi/2$ is the only stationary point of $\cos(t(\theta))$ and

$$\left| \frac{d^2 \cos t}{d\theta^2} (\pi/2) \right| = \frac{(1 - \kappa^2)\kappa}{(1 - \kappa^2)^{3/2}} |\sin t| > \frac{(1 - \kappa^2)\kappa}{(1 - \kappa^2)^{3/2}} \sin \delta$$
 (1.30)

Therefore, we can choose appropriate λ_1, λ_2 such that

$$\left|\frac{d^2\cos t}{d\theta^2}\right| > \frac{(1-\kappa^2)\kappa}{(1-\kappa^2)^{3/2}}\sin\delta\tag{1.31}$$

for any $\theta \in \theta(\chi_1)$. According to lemma (6.1), we obtain $|I_{\chi_1}| \leq C \frac{1}{\rho^{1/2}}$, and also $|I_{\chi_2}| \leq C \frac{1}{\rho^{1/2}}$. Using integration by parts, we get

$$\left| \int_{[-\pi/2,\pi/2]\setminus((-\delta,\delta)\cup\chi)} f(t)(1-\chi_{\delta}(t))e^{\mathbf{i}\rho\cos t}dt \right| \le C\frac{1}{\rho}$$

Consequently, for N=-1 and $\phi\neq\phi_{\kappa}$, we get $|I(\rho,\phi)|\leq\frac{1}{\rho^{1/2}}$.

We now turn to the case of $\phi = \phi_{\kappa}$. By (1.11), we can define χ_{ϵ} similarly and also decompose I into I_1 and I_2 . Using the same agurement above, we can easily carry out that: for N=1, we have $|I_2| \leq C\frac{1}{\rho}$; for N=-1, we have $|I_2| \leq C\frac{1}{\rho^{1/2}}$. Finally, it remains to analysis I_1 . By (1.11), we have

$$I_{1} = \int_{-\epsilon}^{\epsilon} g_{1}\chi_{\epsilon} + g_{2}\chi_{\epsilon}(\sin^{2}\phi_{\kappa} - \sin^{2}(t + \phi_{\kappa}))^{N/2}e^{\mathbf{i}\rho\cos t}dt$$

$$= \int_{-\epsilon}^{\epsilon} g_{1}\chi_{\epsilon} + g_{2}\chi_{\epsilon}(-2(\sin\phi_{\kappa} + \sin(t + \phi_{\kappa}))\cos\frac{2\phi_{\kappa} + t}{2}\sin t/2)^{N/2}e^{\mathbf{i}\rho\cos t}dt$$

$$= \int_{\mathbb{R}} g_{1}\chi_{\epsilon} + g_{2}\chi_{\epsilon}((\sin\phi_{\kappa} + \sin(t + \phi_{\kappa}))\cos\frac{2\phi_{\kappa} + t}{2})^{N/2}(-2\sin t/2)^{N/2}e^{\mathbf{i}\rho\cos t}dt$$

Also, subtituting $t(s) = 2 \arcsin s/2$ for t in I_1 , it follows that

$$I_{1} = \int_{\mathbb{R}} h_{1}(s)e^{-i\rho\frac{s^{2}}{2}} + h_{2}(s)(-s)^{N/2}e^{-i\rho\frac{s^{2}}{2}}$$

$$= I_{11} + I_{12}$$
(1.32)

where

$$h_1(s) = g_1(t(s))\chi_{\epsilon}(t(s))\sqrt{1 - s^2/4} e^{i\rho}$$

$$h_2(s) = g_2\chi_{\epsilon}((\sin\phi_{\kappa} + \sin(t + \phi_{\kappa}))\cos\frac{2\phi_{\kappa} + t}{2})_{t=t(s)}^{N/2}\sqrt{1 - s^2/4} e^{i\rho}$$

and $h_1(s), h_2(s) \in C_c^{\infty}(\mathbb{R})$. Using stationary phase lemma similarly, if N = 1,

$$I_{11} = \left(\frac{2\pi}{\rho}\right)^{1/2} g_1(0)e^{\mathbf{i}\rho - \mathbf{i}\pi/4} + O(\frac{1}{\rho})$$
(1.34)

$$= \left(\frac{2\pi}{\rho}\right)^{1/2} f(0)e^{i\rho - i\pi/4} + O(\frac{1}{\rho})$$
 (1.35)

if N = -1, we get $|I_{11}| \leq C \frac{1}{\rho^{1/2}}$. For I_{12} , we have

$$I_{12} = \int_0^\infty (\mathbf{i}h_2(s) + h_2(-s))s^{N/2} e^{-\mathbf{i}\rho s^2/2} ds$$
 (1.36)

$$= \frac{1}{2} \int_0^\infty (\mathbf{i} h_2(\sqrt{s}) + h_2(-\sqrt{s})) s^{N/4 - 1/2} e^{-\mathbf{i}\rho s/2} ds$$
 (1.37)

By lemma (1.2), we get $|I_{12}| \leq C \frac{1}{\rho^{(N+2)/4}}$.

2. Some draft about Green Tensor Analysis

Let substitute $\xi = k \sin \theta$ into integral and shift the variable, we have

$$I(y) = \int_{\mathbb{R}} f(\xi) e^{\mathbf{i}\xi y_1 + \mu(\xi)y_2} d\xi = \int_{\mathbb{R}} f(\xi) e^{\mathbf{i}\xi(y_1 - z_1) + \mu(\xi)(y_2 - z_2)} e^{\mathbf{i}\xi z_1 + \mu(\xi)z_2} d\xi$$
(2.1)

$$= k \int_{L} f(k\sin\theta)\cos\theta e^{\mathbf{i}k|y-z|\cos(\theta-\eta)} e^{\mathbf{i}|z|\cos(\theta-\phi)} d\theta$$
 (2.2)

$$= k \int_{L_{\phi}} f(k \sin(\theta + \phi)) \cos(\theta + \phi) e^{ik|y - z|\cos(\theta + \phi - \eta)} e^{i|z|\cos\theta} d\theta$$
 (2.3)

$$= k \int_{L} f(k\sin(\theta + \phi))\cos(\theta + \phi)e^{\mathbf{i}k|y-z|\cos(\theta + \phi - \eta)}e^{\mathbf{i}|z|\cos\theta}d\theta$$
 (2.4)

where $y_1, y_2 > 0$, $\sin \phi = \frac{z_1}{|z|}$, $\cos \phi = \frac{z_2}{|z|}$, $0 < \phi < \pi/2$ and $\sin \eta = \frac{y_1 - z_1}{|y - z|}$, $\cos \eta = \frac{y_2 - z_2}{|y - z|}$, $0 < \eta < \pi$. It is easy to see that $\phi + \eta < \pi$. Roughly, using stationary phase lemma, we obtain:

$$I(y) = f(k\sin\phi)k\cos\phi e^{ik|y-z|\cos(\phi-\eta)} \left(\frac{2\pi}{|z|}\right)^{1/2} e^{i|z|-i\frac{\pi}{4}} \left(1 + O(\frac{1}{|z|})\right)$$
(2.5)

$$\cos(a + \mathbf{i}b) = \frac{e^b + e^{-b}}{2}\cos a + \mathbf{i}\frac{e^{-b} - e^b}{2}\sin a$$
 (2.6)

$$\sin(a + \mathbf{i}b) = \frac{e^b + e^{-b}}{2}\sin a + \mathbf{i}\frac{e^b - e^{-b}}{2}\cos a \tag{2.7}$$

When $\theta \in (-a - \pi/2, -a - \pi/2 + \mathbf{i}\infty)$, let $\theta = -a - \pi/2\mathbf{i}t$, where $t > 0, 0 \le a \le \phi$, then

$$-\mathrm{Im}\left(|z|\cos\theta + |y-z|\cos(\theta + \phi - \eta)\right)$$

$$= |z|\sin(a+\pi/2) + |y-z|\sin(a+\pi/2 - \phi + \eta)$$
(2.8)

$$= |z|\cos a + |y - z|\cos(a - \phi + \eta) \tag{2.9}$$

$$= |z|\cos a + \cos a|y - z|(\cos \phi \cos \eta + \sin \phi \sin \eta) \tag{2.10}$$

$$+\sin a|y-z|(\sin\phi\cos\eta-\cos\phi\sin\eta) \tag{2.11}$$

$$= |z|\cos a + \cos a((y_2 - z_2)\cos \phi + (y_1 - z_1)\sin \phi)$$
 (2.12)

$$+\sin a((y_2-z_2)\sin\phi-(y_1-z_1)\cos\phi)$$
 (2.13)

$$= y_1 \sin(\phi - a) + y_2 \cos(\phi - a) > 0 \tag{2.14}$$

Now, Using Cauchy Integral Theorem, we have

$$I(y) = k \int_{L} f(k\sin(\theta + \phi))\cos(\theta + \phi)e^{ik|y-z|\cos(\theta + \phi - \eta)}e^{i|z|\cos\theta}d\theta \qquad (2.15)$$

Let $L_1 = (-\pi/2, -\pi/2 + \mathbf{i}\infty)$ and $\theta = -\pi/2 + \mathbf{i}t, t > 0$, then

$$I_1(y) = k \int_{L_1} f(k \sin(\theta + \phi)) \cos(\theta + \phi) e^{\mathbf{i}k|y-z|\cos(\theta + \phi - \eta)} e^{\mathbf{i}|z|\cos\theta} d\theta \qquad (2.16)$$

$$= (2.17)$$

$$I(y) = f(k\sin\phi)k\cos\phi e^{ik|y-z|\cos(\phi-\eta)} (\frac{2\pi}{|z|})^{1/2} e^{i|z|-i\frac{\pi}{4}}$$
 (2.18)

$$+\frac{kz_2}{|z|}O(\left(\frac{1}{k|z|}\right)^{3/4} + \frac{1}{k|y|}) + \frac{kz_1}{|z|}O(\left(\frac{1}{k|z|}\right)^{5/4} + \left(\frac{1}{k|y|}\right)^2)$$
(2.19)

It is easy to see

$$\int_{-d}^{d} \frac{k}{(k|x-z|)^{\alpha}} \frac{1}{(k|x-y|)^{\beta}} dx_1 \le C\left(\frac{1}{(kz_2)^{\alpha+\beta-1}} + \frac{1}{(ky_2)^{\alpha+\beta-1}}\right)$$
(2.20)

where $z, y \in \mathbb{R}^2_+$, $x \in \Gamma_0$ and $\alpha + \beta > 0$.

$$e^{i\mu y_2 + i\xi(x_1 - y_1)} = e^{i\mu y_2 - iy_2/\tan\phi} = e^{iy_2(\mu - \xi/\tan\phi)}$$
 (2.21)

Another method

$$\int_{-\pi/2}^{\pi/2} f(k\sin(\theta + \psi))k\cos(\theta + \psi)e^{\mathbf{i}k|x-y|\cos\theta}$$
(2.22)

$$= \int_{-\pi/2}^{\pi/2} f(k\sin(\theta + \psi))k\cos(\theta + \psi)e^{\mathbf{i}k|x-y|\cos(\theta + \psi - \psi)}$$
(2.23)

$$= \int_{-\pi/2}^{\pi/2} f(k\sin(\theta + \psi))k\cos(\theta + \psi)e^{\mathbf{i}ky_2\cos(\theta + \psi) + \mathbf{i}k|x_1 - y_1|\sin(\theta + \psi)}$$
(2.24)

$$= \int_{-\pi/2}^{\pi/2} f(k\sin(\theta + \psi))k\cos(\theta + \psi)$$
 (2.25)

$$e^{ik(y_2-z_2)\cos(\theta+\psi)+ik(|x_1-y_1|-|x_1-z_1|)\sin(\theta+\psi)+ik|z|\cos(\theta+\psi-\phi)}$$
 (2.26)

3. Finite Aperture Point Spread Function

If $x \in \Gamma_0$ and $z, y \in \mathbb{R}^2_+$, by lemma (??) we have

$$G(x,y) = \frac{\mathbf{i}k_s}{\mu\sqrt{2\pi}} \frac{1}{\delta(\xi)} \begin{pmatrix} \mu_s \beta & \xi \beta \\ 2\xi \mu_s \mu_p & 2\xi^2 \mu_p \end{pmatrix}_{\xi = k_s \frac{x_1 - y_1}{|x - y|}} \frac{y_2}{|x - y|} \frac{1}{(k_s |x - y|)^{1/2}} e^{\mathbf{i}k_s |x - y| - \mathbf{i}\frac{\pi}{4}}$$

$$+ \frac{\mathbf{i}k_p}{\mu\sqrt{2\pi}} \frac{1}{\delta(\xi)} \begin{pmatrix} 2\xi^2 \mu_s & -2\xi \mu_s \mu_p \\ -\xi \beta & \mu_p \beta \end{pmatrix}_{\xi = k_p \frac{x_1 - y_1}{|x - y|}} \frac{y_2}{|x - y|} \frac{1}{(k_p |x - y|)^{1/2}} e^{\mathbf{i}k_p |x - y| - \mathbf{i}\frac{\pi}{4}} (3.1)$$

$$+O(\frac{y_2}{|x-y|}\frac{1}{(k_s|x-y|)^{3/4}} + \frac{|x_1-y_1|}{|x-y|}\frac{1}{(k_s|x-y|)^{5/4}})$$

$$:= \mathcal{G}_s(x,y) + \mathcal{G}_p(x,y) + O(\frac{y_2}{|x-y|}\frac{1}{(k_s|x-y|)^{3/4}} + \frac{|x_1-y_1|}{|x-y|}\frac{1}{(k_s|x-y|)^{5/4}})$$

$$T_{D}(x,z) = \frac{k_{s}}{\sqrt{2\pi}} \frac{1}{\gamma(\xi)} \begin{pmatrix} \mu_{s}\mu_{p} & \xi\mu_{p} \\ \xi\mu_{s} & \xi^{2} \end{pmatrix}_{\xi=k_{s}\frac{x_{1}-z_{1}}{|x-z|}} \frac{z_{2}}{|x-z|} \frac{1}{(k_{s}|x-z|)^{1/2}} e^{\mathbf{i}k_{s}|x-z|-\mathbf{i}\frac{\pi}{4}}$$

$$+ \frac{k_{p}}{\sqrt{2\pi}} \frac{1}{\gamma(\xi)} \begin{pmatrix} \xi^{2} & -\xi\mu_{p} \\ -\xi\mu_{s} & \mu_{p}\mu_{s} \end{pmatrix}_{\xi=k_{p}\frac{x_{1}-z_{1}}{|x-z|}} \frac{z_{2}}{|x-z|} \frac{1}{(k_{p}|x-z|)^{1/2}} e^{\mathbf{i}k_{p}|x-z|-\mathbf{i}\frac{\pi}{4}}$$

$$+ O(\frac{k_{s}z_{2}}{|x-z|} \frac{1}{(k_{s}|x-z|)^{3/4}} + \frac{k_{s}|x_{1}-z_{1}|}{|x-z|} \frac{1}{(k_{s}|x-z|)^{5/4}})$$

$$:= \mathcal{T}_{s}(x,z) + \mathcal{T}_{p}(x,z) + O(\frac{k_{s}z_{2}}{|x-z|} \frac{1}{(k_{s}|x-z|)^{3/4}} + \frac{k_{s}|x_{1}-z_{1}|}{|x-z|} \frac{1}{(k_{s}|x-z|)^{5/4}})$$

Now we consider the finite aperture point spread function $J_d(z, y)$:

$$\int_{-d}^{d} (T_D(x_1, 0; z_1, z_2))^T \overline{G(x_1, 0; y_1, y_2)} dx_1$$
(3.3)

Recall following standard asymptotic expansion:

$$|x - y| = |x - z| + \widehat{x - z} \cdot (z - y) + O(\frac{|y - z|^2}{|x - z|})$$
(3.4)

$$|y|^{-\alpha} = |z|^{-\alpha} \left(1 + \frac{|y| - |z|}{|z|}\right)^{-\alpha} = |z|^{-\alpha} \left(1 + O\left(\frac{|y - z|}{|z|}\right)\right)$$
(3.5)

$$e^{\mathbf{i}t} = 1 + O(t) \tag{3.6}$$

$$|a^{1/2} - b^{1/2}| \le C\sqrt{|a - b|} \tag{3.7}$$

where $x, y, z \in \mathbb{R}^2$, $t, a, b \in \mathbb{R}$ and $\alpha > 0$.

Lemma 3.1 For any $z, y \in \mathbb{R}^2_+$, $J_d(z, y) = F(z, y) + O((1 + \frac{|y-z|}{z_2})(\frac{1}{k_s z_2})^{1/4} + \frac{(k_s |y-z|)^2}{k_s z_2} + (\frac{|y-z|}{z_2})^{1/2})$, where

$$F(z,y) = -\frac{\mathbf{i}}{2\pi\mu} \int_{\theta_1^d}^{\theta_2^d} f_s(\theta) \begin{pmatrix} \sin^2\theta & \sin\theta\cos\theta \\ \sin\theta\cos\theta & \cos^2\theta \end{pmatrix} e^{\mathbf{i}k_s(z_1 - y_1)\cos\theta + \mathbf{i}k_s(z_2 - y_2)\sin\theta} d\theta$$
(3.8)
$$-\frac{\mathbf{i}}{2\pi\mu} \int_{\theta_1^d}^{\theta_2^d} f_p(\theta) \begin{pmatrix} \cos^2\theta & -\sin\theta\cos\theta \\ -\sin\theta\cos\theta & \sin^2\theta \end{pmatrix} e^{\mathbf{i}k_p(z_1 - y_1)\cos\theta + \mathbf{i}k_p(z_2 - y_2)\sin\theta} d\theta$$
(3.9)

and

$$f_s(\theta) = \frac{\sin\theta((\kappa^2 - \cos^2\theta)^{1/2}(1 - 2\cos^2\theta) + 2\overline{(\kappa^2 - \cos^2\theta)^{1/2}}\cos^2\theta)}{(\cos^2\theta + \sin\theta(\kappa^2 - \cos^2\theta)^{1/2})\overline{((1 - 2\cos^2\theta)^2 + 4\cos^2\theta\sin\theta(\kappa^2 - \cos\theta)^{1/2})}}$$

$$f_p(\theta) = \frac{\sin\theta(1/\kappa^2 - \cos^2\theta)^{1/2}}{(\cos^2\theta + \sin\theta(1/\kappa^2 - \cos^2\theta)^{1/2})((1/\kappa^2 - 2\cos^2\theta)^2 + 4\cos^2\theta\sin\theta(1/\kappa^2 - \cos\theta)^{1/2})}$$

$$where \ 0 < \theta_1^d < \pi/2 < \theta_2^d < \pi \ and \ z_2 = (d + z_1)\tan\theta_1^d = (z_1 - d)\tan\theta_2^d.$$

Proof.

$$\frac{y_2}{|x-y|} \frac{1}{(k_s|x-y|)^{3/4}} + \frac{|x_1-y_1|}{|x-y|} \frac{1}{(k_s|x-y|)^{5/4}}$$

$$= \left(\frac{z_2}{|x-z|} \frac{1}{(k_s|x-z|)^{3/4}} + \frac{|x_1-z_1|}{|x-z|} \frac{1}{(k_s|x-z|)^{5/4}}\right) \left(1 + O\left(\frac{|y-z|}{|x-z|}\right)\right)$$

$$|\mu_{\mathbf{i}}(k_j \frac{x_1-y_1}{|x-y|}) - \mu_{\mathbf{i}}(k_j \frac{x_1-z_1}{|x-z|})|$$

$$\leq Ck_j \sqrt{\left|\frac{x_1-y_1}{|x-y|} - \frac{x_1-z_1}{|x-z|}\right|} \leq Ck_j \left(\frac{|y-z|}{|x-z|}\right)^{1/2}$$

where i, j = s, p. By above, we can obtain

$$\mathcal{G}_s(x,y) = \mathcal{G}_s(x,z)e^{\mathbf{i}k_s\widehat{x-z}\cdot(z-y)} + O(\frac{(k_s|y-z|)^2}{(k_s|x-z|)^{3/2}}) + O(\frac{(k_s|y-z|)^{1/2}}{k_s|x-z|})$$
(3.10)

$$\mathcal{G}_p(x,y) = \mathcal{G}_p(x,z)e^{\mathbf{i}k_p\widehat{x-z}\cdot(z-y)} + O(\frac{(k_p|y-z|)^2}{(k_p|x-z|)^{3/2}}) + O(\frac{(k_p|y-z|)^{1/2}}{k_p|x-z|})$$
(3.11)

For l > 1, a simple computation show that

$$\int_{-d}^{d} \frac{k_s}{(k_s|x-z|)^l} dx_1 = \frac{1}{(k_s z_2)^{l-1}} \int_{\frac{-d-z_1}{z_2}}^{\frac{d-z_1}{z_2}} \frac{1}{(1+t^2)^{l/2}} dt \le C \frac{1}{(k_s z_2)^{l-1}}$$
(3.12)

Let

$$\mathcal{G}_{\alpha}(x,y) = \frac{\mathbf{i}}{\sqrt{2\pi}\mu} g_{\alpha}(\frac{x_1 - y_1}{|x - y|}, \kappa) \frac{1}{(k_{\alpha}|x - y|)^{1/2}} e^{\mathbf{i}k_{\alpha}|x - y| - \mathbf{i}\frac{\pi}{4}}$$
(3.13)

$$\mathcal{T}_{\alpha}(x,y) = \frac{k_{\alpha}}{\sqrt{2\pi}} t_{\alpha}(\frac{x_1 - z_1}{|x - z|}, \kappa) \frac{1}{(k_s|x - z|)^{1/2}} e^{\mathbf{i}k_{\alpha}|x - z| - \mathbf{i}\frac{\pi}{4}}$$
(3.14)

where $\alpha = s, p$. Now, by substituting (3.10-3.11) into $J_d(z, y)$ and using inequality (3.12), we have

$$J_{d}(z,y) = \frac{-\mathbf{i}}{2\pi\mu} \int_{-d}^{d} t_{s} \left(\frac{x_{1}-z_{1}}{|x-z|},\kappa\right)^{T} \overline{g_{s}(\frac{x_{1}-z_{1}}{|x-z|},\kappa)} \frac{e^{\mathbf{i}k_{s}\widehat{x-z}\cdot(y-z)}}{|x-z|} + t_{p} \left(\frac{x_{1}-z_{1}}{|x-z|},\kappa\right)^{T} \overline{g_{p}(\frac{x_{1}-z_{1}}{|x-z|},\kappa)} \frac{e^{\mathbf{i}k_{p}\widehat{x-z}\cdot(y-z)}}{|x-z|} dx_{1}$$
(3.15)

$$-\frac{\mathbf{i}}{2\pi\mu} \int_{-d}^{d} t_{p} (\frac{x_{1}-z_{1}}{|x-z|}, \kappa)^{T} \overline{g_{s}(\frac{x_{1}-z_{1}}{|x-z|}, \kappa)} \frac{e^{\mathbf{i}k_{s}\widehat{x-z}\cdot(y-z)}}{|x-z|}$$
(3.16)

$$+t_s(\frac{x_1-z_1}{|x-z|},\kappa)^T \overline{g_p(\frac{x_1-z_1}{|x-z|},\kappa)} \frac{e^{ik_p \widehat{x}-\overline{z}\cdot(y-z)}}{|x-z|} dx_1$$
(3.17)

$$+O((1+\frac{|y-z|}{z_2})(\frac{1}{k_s z_2})^{1/4} + \frac{(k_s|y-z|)^2}{k_s z_2} + (\frac{|y-z|}{z_2})^{1/2})$$
(3.18)

$$:= F(z,y) + R(z,y)$$
 (3.19)

$$+O((1+\frac{|y-z|}{z_2})(\frac{1}{k_s z_2})^{1/4} + \frac{(k_s|y-z|)^2}{k_s z_2} + (\frac{|y-z|}{z_2})^{1/2})$$
(3.20)

We denote $\widehat{x-z} = x - z/|x-z| = (\cos(\phi+\pi), \sin(\phi+\pi))$, then taking the substitution $x_1 = z_1 - z_2 \cot \phi$, we obtain

$$F(z,y) = \frac{-\mathbf{i}}{2\pi\mu} \int_{\theta_1^d}^{\theta_2^d} A_s(\phi,\kappa) e^{\mathbf{i}k_s(z_1 - y_1)\cos\phi + \mathbf{i}k_s(z_2 - y_2)\sin\phi}$$
(3.21)

$$+\frac{-\mathbf{i}}{2\pi\mu} \int_{\theta_1^d}^{\theta_2^d} A_p(\phi,\kappa) e^{\mathbf{i}k_p(z_1-y_1)\cos\phi + \mathbf{i}k_p(z_2-y_2)\sin\phi}$$
(3.22)

$$R(z,y) = \frac{-\mathbf{i}}{2\pi\mu} \int_{\theta_2^d}^{\theta_2^d} B_s(\phi,\kappa) e^{\mathbf{i}k_s(z_1 - y_1)\cos\phi + \mathbf{i}k_s(z_2 - y_2)\sin\phi + (k_p - k_s)|x - z|}$$
(3.23)

$$+\frac{-\mathbf{i}}{2\pi\mu}\int_{\theta_2^d}^{\theta_2^d} B_p(\phi,\kappa)e^{\mathbf{i}k_p(z_1-y_1)\cos\phi+\mathbf{i}k_p(z_2-y_2)\sin\phi+(k_s-k_p)|x-z|}$$
(3.24)

It is easy to see that $|R(z,y)| \le C \frac{|z-y|}{z_2}$.

Let

$$g(x_1) = \frac{1}{((x_1 - z_1)^2 + z_2^2)^{3/4} ((x_1 - y_1)^2 + y_2^2)^{1/4}}$$

$$\phi(x_1) = ((x_1 - z_1)^2 + z_2^2)^{1/2} - ((x_1 - y_1)^2 + y_2^2)^{1/2}$$

Then, we have

$$g'(x_1) = -g(x_1) \left[\frac{3(x_1 - z_1)}{2((x_1 - z_1)^2 + z_2^2)} + \frac{(x_1 - y_1)}{2((x_1 - y_1)^2 + y_2^2)} \right]$$

$$\phi'(x_1) = \frac{x_1 - z_1}{((x_1 - z_1)^2 + z_2^2)^{1/2}} - \frac{x_1 - y_1}{((x_1 - y_1)^2 + y_2^2)^{1/2}}$$

$$= \frac{\frac{(x_1 - z_1)^2}{(x_1 - z_1)^2 + z_2^2} - \frac{(x_1 - y_1)^2}{(x_1 - y_1)^2 + y_2^2}}{\frac{x_1 - z_1}{((x_1 - z_1)^2 + z_2^2)^{1/2}} + \frac{x_1 - y_1}{((x_1 - y_1)^2 + y_2^2)^{1/2}}}$$

$$= \frac{(x_1 - z_1)^2 y_2^2 - (x_1 - y_1)^2 z_2^2}{(\frac{x_1 - z_1}{((x_1 - z_1)^2 + z_2^2)^{1/2}} + \frac{x_1 - y_1}{((x_1 - y_1)^2 + y_2^2)^{1/2}})((x_1 - z_1)^2 + z_2^2)((x_1 - y_1)^2 + y_2^2)}$$

$$\phi''(x_1) = \frac{z_2^2}{((x_1 - z_1)^2 + z_2^2)^{3/2}} - \frac{y_2^2}{((x_1 - y_1)^2 + y_2^2)^{3/2}}$$

Using integration by parts, we can obtain

$$\int_{-d}^{d} g(x_1)e^{\mathbf{i}\phi(x_1)}dx_1$$

$$= \left(\frac{g(d)}{\phi'(d)}e^{\mathbf{i}\phi(d)} - \frac{g(-d)}{\phi'(-d)}e^{\mathbf{i}\phi(-d)}\right) - \int_{-d}^{d} \frac{g'(x_1)}{\phi'(x_1)} - \frac{g(x_1)\phi''(x_1)}{(\phi'(x_1))^2}dx_1$$

Assume that

$$|y_1| \le c_0 d$$
 $|z_1| \le c_0 d$ $h \le y_2, z_2 \le c_1 h$ $d \le c_2 h$

where $0 < c_0 < 1$. Let define $0 < \theta_y, \theta_z < \pi$ such that

$$\cos \theta_y = \frac{x_1 - y_1}{((x_1 - y_1)^2 + y_2^2)^{1/2}}$$
$$\cos \theta_z = \frac{x_1 - z_1}{((x_1 - z_1)^2 + z_2^2)^{1/2}}$$

By mean value theorem and the law of sines, we get

$$|\phi'(x_1)| = |\cos \theta_z - \cos \theta_y| = |\sin \theta'| |\theta_z - \theta_y|$$

$$\ge \frac{h}{(1+c_0)d} |\sin(\theta_z - \theta_y)|$$

$$= \frac{h}{(1+c_0)d} \frac{|z-y|}{|x-y|} \sin \theta_{|x-y|}$$

$$= \frac{h}{(1+c_0)d} \frac{|z-y|}{|x-z|} \sin \theta_{|x-z|}$$

$$\ge \frac{h^2}{(1+c_0)^2 d^2} \frac{|z-y|}{|x-y|}$$
or
$$\ge \frac{h^2}{(1+c_0)^2 d^2} \frac{|z-y|}{|x-z|}$$

Then we have

$$\left| \frac{g(x_1)}{\phi'(x_1)} \right| \le \frac{(1+c_0)^2 d^2}{h^2} \frac{1}{|z-y||x-y|^{1/2}|x-z|^{1/2}}$$

$$\le C \frac{d^2}{h^3} \frac{1}{|z-y|}$$

Moreover, by mean value theorem agian, we have

$$|\phi''(x_1)| = \left| \frac{\sin^2 \theta_z}{|x - z|} - \frac{\sin^2 \theta_y}{|x - y|} \right|$$

$$= \left| \frac{2 \sin \theta' \cos \theta'}{|x - y'|} (\theta_z - \theta_y) - \frac{\sin^2 \theta'}{|x - y'|^2} (|x - z| - |x - y|) \right|$$

$$\leq \pi \frac{|\sin(\theta_z - \theta_y)|}{h} + \frac{|z - y|}{h^2}$$

$$\leq \pi \frac{|\sin \theta_{|x - z|}||z - y|}{h|x - z|} + \frac{|z - y|}{h^2}$$

$$\leq C \frac{|z - y|}{h^2}$$

Now, it is easy to see that

$$\left| \int_{-d}^{d} \frac{g'(x_1)}{\phi'(x_1)} - \frac{g(x_1)\phi''(x_1)}{(\phi'(x_1))^2} dx_1 \right|$$

$$\leq C \frac{d^3}{h^4} \frac{1}{|z-y|} + C \frac{d^3}{h^3} \frac{1}{|z-y|} \frac{d^2}{h^3}$$

Based on the above analysis, we can obtain

$$\left| \int_{-d}^{d} z_{2} g(x_{1}) e^{\mathbf{i}\phi(x_{1})} \right| \leq C\left(\left(\frac{d}{h}\right)^{2} + \left(\frac{d}{h}\right)^{3} + \left(\frac{d}{h}\right)^{5}\right) \frac{1}{|z - y|}$$

4. 2017.11.08

$$\sin \phi_{\kappa} - \sin(t + \phi) = -2\cos(\frac{\phi_{\kappa} + \phi + t}{2})\sin(\frac{t + \phi - \phi_{\kappa}}{2})$$
$$\sin(\frac{t + \phi - \phi_{\kappa}}{2}) = \sin\frac{t}{2}\cos(\frac{\phi - \phi_{\kappa}}{2}) + \cos\frac{t}{2}\sin(\frac{\phi - \phi_{\kappa}}{2})$$

Some think, substituting $t = 2 \arcsin s/2$ into following integral

$$\int_{0}^{\infty} \chi(t)(\sin \phi_{\kappa} - \sin(t + \phi))^{1/2} e^{-i\rho \cos t}$$

$$= \int_{0}^{\infty} \chi(t(s))(-s\cos(\frac{\phi - \phi_{\kappa}}{2}) - \sqrt{4 - s^{2}}\sin(\frac{\phi - \phi_{\kappa}}{2})^{1/2} e^{-i\rho s^{2}/2}$$

$$= \int_{0}^{\infty} \chi(t)(-\sqrt{t}\cos(\frac{\phi - \phi_{\kappa}}{2}) - \sqrt{4 - t}\sin(\frac{\phi - \phi_{\kappa}}{2})^{1/2} t^{-1/2} e^{-i\rho t/2}$$

Let

$$f(t) = t^{-1/2}e^{-\mathbf{i}\rho t/2}$$

$$g(t) = -\int_{t}^{t-\mathbf{i}\infty} x^{-1/2}e^{-\mathbf{i}\rho x/2}dx$$

$$= \mathbf{i}\int_{0}^{\infty} (t-\mathbf{i}x)^{-1/2}e^{-\mathbf{i}\rho t-\rho x}dx$$

It is easy to see that g'(t) = f(t). Then we have

$$= \int_{0}^{\infty} \chi(t)(-\sqrt{t}\cos(\frac{\phi - \phi_{\kappa}}{2}) - \sqrt{4 - t}\sin(\frac{\phi - \phi_{\kappa}}{2})^{1/2}t^{-1/2}e^{-i\rho t/2}$$

$$= \chi(0)(-2\sin(\frac{\phi - \phi_{\kappa}}{2}))^{1/2}g(0)$$

$$- \int_{0}^{\infty} (\chi(t)(-\sqrt{t}\cos(\frac{\phi - \phi_{\kappa}}{2}) - \sqrt{4 - t}\sin(\frac{\phi - \phi_{\kappa}}{2})^{1/2})'g(t)dt$$

We get

$$g(x) = \int_0^\infty \chi(t) (-\sqrt{t} \cos(\frac{\phi - \phi_{\kappa}}{2}) - \sqrt{4 - t} \sin(\frac{\phi - \phi_{\kappa}}{2})^{-1/2} t^{-1/2} (t - \mathbf{i}x)^{-1/2} e^{-\mathbf{i}\rho t} dt$$

$$R(\rho) = \int_0^\infty g(x) e^{-\rho x} dx$$

Because $\chi(t)$ has compact support $(-\delta, \delta)$, we obtain

$$gg(x) = \int_0^{\delta} (\sqrt{t}\cos(\theta) - \sqrt{4-t}\sin\theta)^{-1/2}t^{-1/2}(t^2 + x^2)^{-1/4}dt$$

where $\theta = \frac{\phi - \phi_{\kappa}}{2}$. For x > 0, Put L(x):

$$\int_0^a \frac{1}{t^{3/4}} \frac{1}{(t^2 + x^2)^{1/4}} dt$$

$$= 4 \int_0^a \frac{1}{(t^2 + x^2)^{1/4}} dt^{1/4}$$

$$= 4 \int_0^{a^{1/4}} \frac{1}{(t^8 + x^2)^{1/4}} dt$$

$$= 4x^{-1/4} \int_0^{(\frac{a}{x})^{1/4}} \frac{1}{(t^8 + 1)^{1/4}} dt$$

$$= 4x^{-1/4} \int_0^{(\frac{a}{x})^{1/4}} \frac{1}{(t^8 + 1)^{1/4}} dt$$

$$\leq 4x^{-1/4} \int_0^\infty \frac{1}{(t^8 + 1)^{1/4}} dt$$

Back to analysis qq(x), we have

$$\begin{split} gg(x) &\leq \int_0^\delta \left| \frac{\sqrt{t} + 2|\sin\theta}{t - 4\sin^2\theta} \right|^{1/2} t^{-1/2} (t^2 + x^2)^{-1/4} dt \\ &= \int_0^\delta \left| \frac{1}{\sqrt{t} - 2|\sin\theta} \right|^{1/2} t^{-1/2} (t^2 + x^2)^{-1/4} dt \\ &= 2 \int_0^{\sqrt{\delta}} \left| \frac{1}{t - 2|\sin\theta} \right|^{1/2} (t^4 + x^2)^{-1/4} dt \\ &= 2 \int_{-2|\sin\theta|}^{\sqrt{\delta} - 2|\sin\theta|} |t|^{-1/2} ((t + 2|\sin\theta|)^4 + x^2)^{-1/4} dt \\ &\leq 4 \int_0^{\delta^{1/4}} (t^8 + x^2)^{-1/4} dt + 4 \int_0^{\sqrt{2|\sin\theta|}} ((t^2 - 2|\sin\theta|)^4 + x^2)^{-1/4} dt \\ &\leq C x^{-1/4} (1 + \int_0^{\sqrt{2|\sin\theta|}} ((t^2 - 2|\sin\theta|)^4 / x + x)^{-1/4} dt) \\ &\leq C x^{-1/4} (1 + \int_0^{\sqrt{2|\sin\theta|}} (t^2 - 2|\sin\theta|)^{-1/2} dt) \\ &= C x^{-1/4} (1 + \int_0^1 (1 - t^2)^{-1/2} dt) \leq C x^{-1/4} \end{split}$$

Immediately, we can obtain

$$|g(x)| \le Cx^{-1/4}$$

It follows that

$$R(\rho) \le \int_0^\infty x^{-1/4} e^{-\rho x} \le C \rho^{-3/4}$$

5. stationary of phase lemma

Lemma 5.1 Assume that $0 < \kappa := \sin \phi_{\kappa} < 1, 0 < \phi_{\kappa} < \pi/2, 0 \le \phi \le \pi/2$. Let

$$f(t,\phi) := F(\sin(t+\phi), \cos(t+\phi), (\kappa^2 - \sin^2(t+\phi))^{1/2})$$
(5.1)

be a complexed function in $C([-\pi/2, \pi/2] \times [0, \pi/2])$. Moreover, its partial derivative with respect to t can be represented as

$$\frac{\partial f(t,\phi)}{\partial t} = g(t,\phi)(\kappa^2 - \sin^2(t+\phi))^{-1/2}$$
(5.2)

where $g(t, \phi)$ is uniformly bounded. Then for any $\rho \geq 1$, we have

$$\left| I(\rho, \phi) := \int_{-\pi/2}^{\pi/2} f(t)e^{\mathbf{i}\rho\cos t}dt - \left(\frac{2\pi}{\rho}\right)^{1/2} f(0)e^{\mathbf{i}\rho - \mathbf{i}\pi/4} \right|
\leq C \frac{1}{\rho^{3/4}}$$
(5.3)

Proof. Solving the following equation:

$$\kappa^2 - \sin^2(t + \phi) = 0$$

we have, if $0 < \phi < \pi/2 - \phi_{\kappa}$,

$$t_1(\phi) = \phi_{\kappa} - \phi$$
 $t_2(\phi) = -\phi_{\kappa} - \phi$

and if $\pi/2 - \phi_{\kappa} \leq \phi < \pi/2$,

$$t_1(\phi) = \phi_{\kappa} - \phi$$
 $t_2(\phi) = \pi - \phi_{\kappa} - \phi$

Since $|t_2(\phi)| < \phi_{\kappa}$ or $|t_2(\phi)| < \pi/2 - \phi_{\kappa}$, we now define $\delta := \min(\frac{\phi_{\kappa}}{2}, \frac{\pi/2 - \phi_{\kappa}}{2})$ and it is easy to see that

$$\kappa + \sin(t + \phi) \neq 0 \tag{5.4}$$

$$\cos(\frac{t+\phi+\phi_{\kappa}}{2}) \neq 0 \tag{5.5}$$

for any $(t, \phi) \in [-\delta, \delta] \times [0, \pi/2]$. Let $\chi_{\delta} \in C_0^{\infty}(-\pi/2, \pi/2)$ be the cut-off function with that $0 \le \chi_{\delta} \le 1$, $\chi_{\delta} = 1$ in $(-\delta/2, \delta/2)$ and $\chi_{\delta} = 0$ in $(-\pi/2, \pi/2) \setminus (-\delta, \delta)$. Then we can divide I into two parts such that

$$I = \int_{-\delta}^{\delta} f(t)\chi_{\delta}(t)e^{\mathbf{i}\rho\cos t}dt + \int_{-\pi/2}^{\pi/2} f(t)(1-\chi_{\delta}(t))e^{\mathbf{i}\rho\cos t}dt$$

=: $I_1 + I_2$

Subtitating $t(s) = 2 \arcsin s/2$ for t in I_1 , we can obtain

$$I_{1} = \int_{-2\sin\frac{\delta}{2}}^{2\sin\frac{\delta}{2}} f(t(s))\chi_{\delta}(t(s)) \frac{1}{\sqrt{1 - s^{2}/4}} e^{i\rho} e^{-i\rho s^{2}/2} ds$$

$$\int_{-2\sin\frac{\delta}{2}}^{2\sin\frac{\delta}{2}} f(t(s))\chi_{\delta}(t(s)) \frac{1}{\sqrt{1 - s^{2}/4}} e^{i\rho} e^{-i\rho s^{2}/2} ds$$
(5.6)

$$= \int_0^{2\sin\frac{\delta}{2}} (f(t(s))\chi_{\delta}(t(s)) + f(-t(s))\chi_{\delta}(-t(s))) \frac{1}{\sqrt{1-s^2/4}} e^{i\rho} e^{-i\rho s^2/2} ds (5.7)$$

$$:= I_{11} + I_{12} (5.8)$$

Taking substitution $s = \sqrt{x}$, we get

$$I_{11} = \frac{1}{2} \int_0^{(2\sin\frac{\delta}{2})^2} f(t(\sqrt{x})) \chi_{\delta}(t(\sqrt{x})) \frac{1}{\sqrt{1 - x/4}} x^{-1/2} e^{i\rho} e^{-i\rho x/2} dx$$

Observe that

$$\sin \phi_{\kappa} - \sin(t + \phi) = -2\cos(\frac{\phi_{\kappa} + \phi + t}{2})\sin(\frac{t + \phi - \phi_{\kappa}}{2})$$

$$\sin(\frac{t + \phi - \phi_{\kappa}}{2}) = \sin\frac{t}{2}\cos(\frac{\phi - \phi_{\kappa}}{2}) + \cos\frac{t}{2}\sin(\frac{\phi - \phi_{\kappa}}{2})$$

$$:= \sin\frac{t}{2}\cos\theta + \cos\frac{t}{2}\sin\theta$$

where $\theta = \frac{\phi - \phi_{\kappa}}{2}$. By lemma (1.3) and using representation (5.2), inequality (5.4-5.5), it follows that

$$|I_{11} - \frac{1}{2} \sqrt{\frac{2\pi}{\rho}} f(0) e^{\mathbf{i}\rho - \mathbf{i}\frac{\pi}{4}}|$$

$$\leq \int_{0}^{\infty} e^{-\rho y} dy \int_{0}^{(2\sin\frac{\delta}{2})^{2}} \left| \frac{\partial (f(t(\sqrt{x}))\chi_{\delta}(t(\sqrt{x})) \frac{1}{\sqrt{1-x/4}})}{\partial x} \right| (x^{2} + y^{2})^{-\frac{1}{4}} dx$$

$$\leq C \int_{0}^{\infty} e^{-\rho y} dy \int_{0}^{(2\sin\frac{\delta}{2})^{2}} |\sqrt{x}\cos\theta + \sqrt{4-x}\sin\theta|^{-1/2}x^{-1/2}(x^{2}+y^{2})^{-\frac{1}{4}} dx$$

$$\leq C \int_{0}^{\infty} e^{-\rho y} dy \int_{0}^{(2\sin\frac{\delta}{2})^{2}} \frac{(\sqrt{x}|\cos\theta| + \sqrt{4-x}|\sin\theta|)^{1/2}}{|x-4\sin^{2}\theta|^{1/2}} x^{-1/2}(x^{2}+y^{2})^{-\frac{1}{4}} dx$$

$$\leq C \int_{0}^{\infty} e^{-\rho y} dy \int_{0}^{(2\sin\frac{\delta}{2})^{2}} \frac{1}{|\sqrt{x}-2|\sin\theta||^{1/2}} x^{-1/2}(x^{2}+y^{2})^{-\frac{1}{4}} dx$$

$$\leq C \int_{0}^{\infty} e^{-\rho y} dy \int_{0}^{2\sin\frac{\delta}{2}} \frac{1}{|x-2\sin|\theta||^{1/2}} (x^{4}+y^{2})^{-\frac{1}{4}} dx$$

$$\leq C \int_{0}^{\infty} e^{-\rho y} dy \int_{-2\sin|\theta|}^{2\sin\frac{\delta}{2}} \frac{1}{|x|^{1/2}} ((x+2\sin|\theta|)^{4}+y^{2})^{-\frac{1}{4}} dx$$

$$\leq C \int_{0}^{\infty} e^{-\rho y} dy \int_{-2\sin|\theta|}^{2\sin\frac{\delta}{2}} \frac{1}{|x|^{1/2}} ((x+2\sin|\theta|)^{4}+y^{2})^{-\frac{1}{4}} dx$$

$$\leq C \int_{0}^{\infty} e^{-\rho y} dy \left(\int_{0}^{\sqrt{2\sin\frac{\delta}{2}}} (x^{8}+y^{2})^{-\frac{1}{4}} dx + \int_{0}^{\sqrt{2\sin|\theta|}} ((x^{2}-2\sin|\theta|)^{4}+y^{2})^{-\frac{1}{4}} dx \right)$$

$$\leq C \int_{0}^{\infty} e^{-\rho y} dy \left(\int_{0}^{\sqrt{2\sin\frac{\delta}{2}}} (x^{8}+1)^{-\frac{1}{4}} dx + y^{-\frac{1}{4}} \int_{0}^{\sqrt{2\sin|\theta|}} (2\sin|\theta|-x^{2})^{-\frac{1}{2}} dx \right)$$

$$\leq C \int_{0}^{\infty} e^{-\rho y} dy \left(\int_{0}^{\infty} (x^{8}+1)^{-\frac{1}{4}} dx + y^{-\frac{1}{4}} \int_{0}^{\sqrt{2\sin|\theta|}} (2\sin|\theta|-x^{2})^{-\frac{1}{2}} dx \right)$$

$$\leq C \int_{0}^{\infty} y^{-\frac{1}{4}} e^{-\rho y} dy \left(\int_{0}^{\infty} (x^{8}+1)^{-\frac{1}{4}} dx + \int_{0}^{1} (1-x^{2})^{-\frac{1}{2}} dx \right) \leq C \int_{0}^{\frac{1}{2}} y^{-\frac{1}{4}} dx$$

Using the same agrument, we can also carry out

$$|I_{12} - \frac{1}{2}\sqrt{\frac{2\pi}{\rho}}f(0)e^{\mathbf{i}\rho - \mathbf{i}\frac{\pi}{4}}| \le C\frac{1}{\rho^{3/4}}$$
(5.9)

It remains to estimate I_2 . Note that there exits m > 0 such that $|\sin t| \ge m$ for any $t \in [-\pi/2, \pi/2] \setminus (-\delta/2, \delta/2)$. Upon integration by parts and representation (5.2) again, we have

$$|I_{12}| \le C\rho^{-1} (1 + \left| \int_{[-\pi/2, \pi/2] \setminus (-\delta/2, \delta/2)} \frac{\partial (f(t)(1 - \chi_{\delta}(t)))}{\partial t} \frac{1}{\sin t} dt \right|)$$

$$\le C\rho^{-1} (1 + \int_{-\pi/2}^{\pi/2} \left| \frac{\partial (f(t)(1 - \chi_{\delta}(t)))}{\partial t} \right| dt)$$

$$\le C\rho^{-1} (1 + \int_{-\pi/2}^{\pi/2} |(\kappa^2 - \sin^2(t + \phi))^{-1/2}| dt)$$

$$\le C\rho^{-1} (1 + \int_{-\pi/2}^{\pi/2} |(\kappa^2 - \sin^2 t)^{-1/2}| dt)$$

$$\le C\rho^{-1}$$

This completes the proof.

6. cross term of psf, 17.11.15

We need the following slight generalization of Van der Corput lemma for the oscillatory integral [2, P.152].

Lemma 6.1 Let $-\infty < a < b < \infty$, and u is a C^k function u in (a,b).

1. If $|u'(t)| \ge 1$ for $t \in (a,b)$ and u' is monotone in (a,b), then for any $\phi(t)$ in (a,b) with integrable derivatives

$$\left| \int_{a}^{b} e^{i\lambda u(t)} \phi(t) dt \right| \leq 3\lambda^{-1} \left[|\phi(b)| + \int_{a}^{b} |\phi'(t)| dt \right].$$

2. For all $k \geq 2$, if $|u^{(k)}(t)| \geq 1$ for $t \in (a,b)$, then for any $\phi(t)$ in (a,b) with integrable derivatives

$$\left| \int_a^b e^{\mathbf{i}\lambda u(t)} \phi(t) dt \right| \le 12k\lambda^{-1/k} \left[|\phi(b)| + \int_a^b |\phi'(t)| dt \right].$$

Proof. The assertion can be proved by extending the Van der Corptut lemma in [2]. Here we omit the details.

Lemma 6.2 For $0 < \kappa < 1$, let $F(\lambda) = \int_0^{\kappa} f(t)e^{i\lambda(\sqrt{1-t^2}-\tau\sqrt{\kappa^2-t^2}+\alpha t)}dt$, where $\tau \ge c_0 > 0$ and $\alpha \in \mathbb{R}$, then we have

$$|F(\lambda)| \le C(\kappa)\lambda^{-\frac{1}{2N_*}} \left[|f(\kappa)| + \int_0^{\kappa} |f'(t)| dt \right]$$

where $N_* = \min\{N | \kappa^{2N-1} < c_0, N \in \mathbb{Z}_+ \}.$

Proof. Put $\phi(t) = -\sqrt{1-t^2}$ and $\psi(t,\tau) = \tau \phi(t/\kappa) - \phi(t) + \alpha t$. For easy of notations, we denote the *n*-th partial derivative of g(t) with respect to t by $g^{(n)}(t)$. Then, it is to see that, for n > 1

$$\psi^{(n)}(t,\tau) = \frac{\tau}{\kappa^{n-1}} \phi^{(n)}(\frac{t}{\kappa}) - \phi^{(n)}(t)$$

A standard computation show that

$$\phi^{(1)}(t) = \frac{t}{\sqrt{1 - t^2}}$$

$$\phi^{(2)}(t) = \frac{1}{(1 - t^2)^{3/2}}$$

$$\phi^{(3)}(t) = \frac{3t}{(1 - t^2)^{5/2}}$$

Moreover, for $n \geq 3$, we have

$$\phi^{(n)}(t) = \frac{p_n(t)}{(1-t^2)^{n-1/2}} \tag{6.1}$$

where $p_n = \sum_{k=0}^{n-2} a_k^n t^k$ is a (n-2)-th polynomial such that its coefficients satisfy the following recursion formula:

$$\begin{aligned} a_{n-1}^{n+1} &= (n+1)a_{n-2}^n, \quad a_{n-2}^{n+1} &= (n+2)a_{n-3}^n \\ a_k^{n+1} &= (k+1)a_{k+1}^n + (2n-k)a_{k-1}^n \quad \text{for } 1 \leq k \leq n-3 \\ a_0^{n+1} &= a_1^n \end{aligned}$$

Since the polynmial coefficients are all positive, it is obvious that for $n \ge 1$, $\phi^{(n)}(t)$ is a monotone increasing positive function. Using the recursion formula, it follows that

$$\phi^{(n)}(0) = \begin{cases} 0 & \text{n is odd,} \\ (n-1)!!(n-3)!! & \text{n is even.} \end{cases}$$
 (6.2)

where (2k-1)!! is double factorial and n>3. We are now in the position to proof the inequality. Since $0<\kappa<1$, obersev that

$$\psi^{(2N_*+1)}(t,\tau) \ge \frac{\tau}{\kappa^{2N_*}} \phi^{(2N_*+1)}(t) - \phi^{(2N_*+1)}(t) > 0$$

Therefore, $\psi^{(2N_*)}(t,\tau)$ is monotone increasing in $[0,\kappa)$. By (6.2), we get

$$\psi^{(2N_*)}(t,\tau) \ge \psi^{(2N_*)}(0,\tau) \ge \psi^{(2N_*)}(0,c_0) = C(2N_*)(\frac{c_0}{\kappa^{2N_*-1}} - 1) > 0$$
(6.3)

The lemma is now a direct consequence of lemma (6.1).

7. Other exponential decay term: 17.11.16 on G1

The parameterization of hyperbolic curve passing $(\pm 1, 0)$ is:

$$\xi_1 = \pm \sqrt{t^2 + 1}$$
 $\xi_2 = t$

where $t \in \mathbb{R}$. Substituting $\xi = \xi_1 + \mathbf{i}\xi_2$ into $\mu(\xi) := (1 - \xi^2)^{1/2}$ and $\mu_{\kappa}(\xi) := (\kappa^2 - \xi^2)^{1/2}$, we get

$$\operatorname{Im} \mu(\xi) = \operatorname{Im} \left(1 - (\xi_1^2 - \xi_2^2 + \mathbf{i}2\xi_1\xi_2)\right)^{1/2}$$
$$= \operatorname{Im} \left(-2t\sqrt{t^2 + 1}\mathbf{i}\right)^{1/2} = t^{1/2}(t^2 + 1)^{1/4}$$

$$\operatorname{Im} \mu_{\kappa}(\xi) = \operatorname{Im} (\kappa^{2} - (\xi_{1}^{2} - \xi_{2}^{2} + \mathbf{i}2\xi_{1}\xi_{2}))^{1/2}$$

$$= \operatorname{Im} (\kappa^{2} - 1 - 2t\sqrt{t^{2} + 1}\mathbf{i})^{1/2}$$

$$= \sqrt{\frac{\sqrt{(1 - \kappa^{2})^{2} + 4t(t^{2} + 1)} + 1 - \kappa^{2}}{2}}$$

$$> t^{1/2}(t^{2} + 1)^{1/4}$$

where we only consider the branch, denoted by Γ^+ , in the first quadrant here. For a > 0, b > 0, we have

$$|e^{\mathbf{i}\xi a + \mathbf{i}\mu(\xi)b + \mathbf{i}\mu_{\kappa}(\xi)c}| \le e^{-ta - t^{1/2}(t^2 + 1)^{1/4}b - t^{1/2}(t^2 + 1)^{1/4}c} \le e^{-t(b + c)}$$

Lemma 7.1 For $\xi \in \Gamma_0$, let $f(\xi)$ is a complex valued function in $L^1(\Gamma^+)$ such that $|f(\xi)| \leq C(1+\xi^k)$, $k \in \mathbb{Z}_+$. Then we have

$$|I(a,b,c) := \int_{\Gamma^+} f(\xi) e^{\mathbf{i}\xi a + \mathbf{i}\mu(\xi)b + \mathbf{i}\mu_{\kappa}(\xi)c} d\xi|$$

$$\leq C\left(\frac{1}{b+c} + \frac{1}{(b+c)^k}\right)$$

Proof.

$$\frac{d\xi(t)}{dt} = \frac{t}{\sqrt{t^2 + 1}} + \mathbf{i}$$

Substituting $\xi(t)$ into I(a,b,c), we hvae

$$\begin{split} |I(a,b,c)| &= \Big| \int_0^\infty |f(\xi(t)) \frac{d\xi(t)}{dt} e^{\mathbf{i}\xi(t)a + \mathbf{i}\mu(\xi(t))b + \mathbf{i}\mu_\kappa(\xi(t))c} dt \Big| \\ &\leq C \int_0^\infty (1+t^k) e^{-t(b+c)} dt \\ &\leq C (\frac{1}{b+c} + \frac{1}{(b+c)^k}) \end{split}$$

Lemma 7.2 Let $f(\xi)$ is a bounded complex valued function in $L^1((\kappa, 1))$. Then we have

$$|I(a,b) := \int_{\kappa}^{1} |f(\xi)e^{\mathbf{i}\xi a + \mathbf{i}\mu_{\kappa}(\xi)b}d\xi|$$

$$\leq C\frac{1}{h}$$

Proof. It is simple to see that

$$\begin{split} |I(a,b)| &\leq C \int_{\kappa}^{1} e^{-b\sqrt{\xi^{2}-\kappa^{2}}} d\xi \\ &\leq C \int_{0}^{\sqrt{1-\kappa^{2}}} \frac{t}{\sqrt{t^{2}+\kappa^{2}}} e^{-bt} dt \\ &\leq C \frac{1}{b} \end{split}$$

References

- [1] Lawrence C Evans. Partial differential equations. 2nd ed. Marcel Dekker,, 2010.
- [2] Loukas Grafakos. Classical and modern Fourier analysis. Prentice Hall, 2004.
- [3] R. Wong, Werner Rheinboldt, and Daniel Siewiorek. Asymptotic Approximations of Integrals. 1989.