

1. Estimate of Dirichlet Green Tensor

We need the following slight generalization of Van der Corput lemma for the oscillatory integral [3, P.152].

Lemma 1.1 *Let $-\infty < a < b < \infty$, and u is a C^k function u in (a, b) .*

1. *If $|u'(t)| \geq 1$ for $t \in (a, b)$ and u' is monotone in (a, b) , then for any $\phi(t)$ in (a, b) with integrable derivatives*

$$\left| \int_a^b e^{i\lambda u(t)} \phi(t) dt \right| \leq 3\lambda^{-1} \left[|\phi(b)| + \int_a^b |\phi'(t)| dt \right].$$

2. *For all $k \geq 2$, if $|u^{(k)}(t)| \geq 1$ for $t \in (a, b)$, then for any $\phi(t)$ in (a, b) with integrable derivatives*

$$\left| \int_a^b e^{i\lambda u(t)} \phi(t) dt \right| \leq 12k\lambda^{-1/k} \left[|\phi(b)| + \int_a^b |\phi'(t)| dt \right].$$

Proof. The assertion can be proved by extending the Van der Corput lemma in [3]. Here we omit the details. \square

We recall following lemma, see e.g. [6]:

Lemma 1.2 *Let $F(\rho, a) = \int_0^a t^{\alpha-1} f(t) e^{-i\rho t} dt$ where $0 < a \leq +\infty$, $0 < \alpha < 1$, $\rho > 0$ and $t^{\alpha-1} f \in L^1(0, a)$, then we have*

$$|F(\rho, a)| \leq C \left(\frac{1}{\rho^\alpha} f(0) + \frac{1}{\rho} (a^{\alpha-1} f(a) + |t^{\alpha-1} f|_{L^1(0, a)}) \right) \quad (1.1)$$

Proof. Put

$$g_0(t) = t^{\alpha-1} e^{-i\rho t} \quad (1.2)$$

and define

$$g_1(t) = - \int_t^{t-i\infty} x^{\alpha-1} e^{-i\rho x} dx \quad (1.3)$$

where the path of integration is the vertical line $x = t - iy, y \geq 0$. It is easy to show that $g_1(t)' = g_0(t)$. Substituting $x = t - iy$ into $g_1(t)$, we have

$$g_1(t) = i \int_0^\infty (t - iy)^{\alpha-1} e^{-i\rho t} e^{-\rho y} dy \quad (1.4)$$

Upon integration by parts, we have

$$\begin{aligned} F(\rho, a) &= \int_0^a f(t) dg_1(t) \\ &= e^{-i\frac{\alpha\pi}{2}} f(0) \Gamma(\alpha) \frac{1}{\rho^\alpha} + f(a) g_1(a) - \int_0^a f'(t) g_1(t) dt \\ &= e^{-i\frac{\alpha\pi}{2}} f(0) \Gamma(\alpha) \frac{1}{\rho^\alpha} - i \int_0^\infty e^{-\rho y} dy \int_0^a f'(t) (t - iy)^{\alpha-1} e^{-i\rho t} dt \end{aligned}$$

Let

$$h(y) = \int_0^a f'(t) (t - iy)^{\alpha-1} e^{-i\rho t} dt$$

and observe that

$$|h(y)| \leq \int_0^a |f'(t)|(t^2 + y^2)^{\frac{\alpha-1}{2}} dt$$

□

Lemma 1.3 *Let $F(\rho, a) = \int_0^a t^{-1/2} f(t) e^{-i\rho t} dt$ where $0 < a \leq +\infty$ and $\rho > 0$, then we have*

$$|F(\rho, a) - e^{-i\frac{\pi}{4}} f(0) \Gamma(1/2) \frac{1}{\rho^{1/2}}| \quad (1.5)$$

$$\leq C \left(\int_0^\infty e^{-\rho y} dy \int_0^a |f'(t)|(t^2 + y^2)^{-\frac{1}{4}} dt + \frac{1}{\rho} a^{-1/2} f(a) \right) \quad (1.6)$$

Proof. Put

$$g_0(t) = t^{-1/2} e^{-i\rho t} \quad (1.7)$$

and define

$$g_1(t) = - \int_t^{t-i\infty} x^{-1/2} e^{-i\rho x} dx \quad (1.8)$$

where the path of integration is the vertical line $x = t - iy, y \geq 0$. It is easy to show that $g_1(t) = g_0(t)$. Substituting $x = t - iy$ into $g_1(t)$, we have

$$g_1(t) = i \int_0^\infty (t - iy)^{-1/2} e^{-i\rho t} e^{-\rho y} dy \quad (1.9)$$

Upon integration by parts, we have

$$\begin{aligned} F(\rho, a) &= \int_0^a f(t) dg_1(t) \\ &= e^{-i\frac{\pi}{4}} f(0) \Gamma(1/2) \frac{1}{\rho^{1/2}} + f(a) g_1(a) - \int_0^a f'(t) g_1(t) dt \\ &= e^{-i\frac{\pi}{4}} f(0) \Gamma(1/2) \frac{1}{\rho^{1/2}} + i f(a) \int_0^\infty (a - iy)^{-1/2} e^{-i\rho t} e^{-\rho y} dy \\ &\quad - i \int_0^\infty e^{-\rho y} dy \int_0^a f'(t) (t - iy)^{-1/2} e^{-i\rho t} dt \end{aligned}$$

Let

$$h(y) = \int_0^a f'(t) (t - iy)^{-1/2} e^{-i\rho t} dt$$

and observe that

$$|h(y)| \leq \int_0^a |f'(t)|(t^2 + y^2)^{-\frac{1}{4}} dt$$

It is easy to see that

$$|g_1(a)| \leq a^{-1/2} \int_0^\infty e^{-\rho y} dy \leq C \frac{1}{\rho}$$

□

Lemma 1.4 Assume that $0 < \kappa := \sin \phi_\kappa < 1$, $0 < \phi_\kappa < \pi/2$, $0 \leq \phi \leq \pi/2$ and $-\pi/2 < t_1 < t_2 < \pi/2$ satisfy that $\kappa^2 = \sin^2(\phi + t_1) = \sin^2(\phi + t_2)$. Let $f(\theta)$:

$$f(t, \phi) := F(\sin(t + \phi), \cos(t + \phi), (\kappa^2 - \sin^2(t + \phi))^{1/2}) \quad (1.10)$$

be a function in $C^\infty(([-\pi/2, \pi/2] \setminus \{t_1, t_2\}) \times [0, \pi/2])$. Moreover, there exists $\epsilon > 0$ such that $f(\theta)$ can be represented as

$$f(t, \phi) = g_1(t, \phi) + g_2(t, \phi)(\kappa^2 - \sin^2(t + \phi))^{1/2})^{N/2} \quad (1.11)$$

where $g_1, g_2 \in C^\infty((\bigcup_{i=1,2} (t_i - \epsilon, t_i + \epsilon)) \times [0, \pi/2])$ and $N = \pm 1$. Then for any $\rho \geq 1$, we have

$$\begin{aligned} |I(\rho, \phi) &:= \int_{-\pi/2}^{\pi/2} f(\theta) e^{i\rho \cos \theta} d\theta - \frac{N+1}{2} \left(\frac{2\pi}{\rho}\right)^{1/2} f(0) e^{i\rho - i\pi/4}| \\ &\leq C \frac{1}{\rho^{(2+N)/4}} \end{aligned} \quad (1.12)$$

Proof. The proof will be split into two parts about whether ϕ equal to ϕ_κ .

If $\phi \neq \phi_\kappa$, there exists $0 < \delta < \pi/4$ such that

$$|(\kappa^2 - \sin^2(t + \phi))^{1/2}| > \frac{1}{2} |(\kappa^2 - \sin^2 \phi)^{1/2}| \quad (1.13)$$

for any $t \in (-\delta, \delta)$. Let $\chi_\delta \in C_0^\infty(-\pi/2, \pi/2)$ be the cut-off function with that $0 \leq \chi_\delta \leq 1$, $\chi_\delta = 1$ in $(-\delta/2, \delta/2)$ and $\chi_\delta = 0$ in $(-\pi/2, \pi/2) \setminus (-\delta, \delta)$. Then we can divide I into two parts such that

$$\begin{aligned} I &= \int_{-\delta}^{\delta} f(t) \chi_\delta(t) e^{i\rho \cos t} dt + \int_{-\pi/2}^{\pi/2} f(t) (1 - \chi_\delta(t)) e^{i\rho \cos t} dt \\ &=: I_1 + I_2 \end{aligned}$$

Substituting $t(s) = 2 \arcsin s/2$ for t in I_1 , we can obtain

$$I_1 = \int_{\mathbb{R}} f(t(s)) \chi_\delta(t(s)) \frac{1}{\sqrt{1 - s^2/4}} e^{i\rho} e^{-i\rho s^2/2} ds \quad (1.14)$$

$$= \int_{\mathbb{R}} h_\delta(s) e^{i\rho} e^{-i\rho s^2/2} ds \quad (1.15)$$

It is easy to see that $h_\delta(s) \in C_0^4(\mathbb{R})$. By the lemma of the stationary phase for quadratic term in [2], we have

$$I_1 = e^{i\rho} \int_{\mathbb{R}} h_\delta(s) e^{-i\frac{\rho}{2}s^2} ds = e^{i\rho} \int_{\mathbb{R}} \widehat{h}_\delta(y) \alpha(-y) dy \quad (1.16)$$

where

$$\alpha(y) = \left(\frac{1}{2\pi\rho}\right)^{1/2} e^{-i\pi/4} e^{\frac{i}{2\rho}y^2} \quad (1.17)$$

$$= \left(\frac{1}{2\pi\rho}\right)^{1/2} e^{-i\pi/4} (1 + O(\frac{y^2}{\rho})) \quad (1.18)$$

Consequently

$$I_1 = \left(\frac{1}{2\pi\rho}\right)^{1/2} e^{i\rho - i\pi/4} \int_{\mathbb{R}} \widehat{h}_\delta(y) (1 + \frac{1}{\rho} O(y^2)) dy \quad (1.19)$$

Moreover, $\int_{\mathbb{R}} \widehat{h}_\delta(y) dy = 2\pi h_\delta(0)$ and $|\int_{\mathbb{R}} \widehat{h}_\delta(y) y^2 dy| < C$ since $\widehat{h}_\delta(y) = O(1/y^4)$. Now, it turns to estimate I_2 .

When $N = 1$, using integration by parts, we have

$$|I_2| = \left| \int_{(-\frac{\pi}{2}, \frac{\pi}{2}) \setminus (-\frac{\delta}{2}, \frac{\delta}{2})} f(t)(1 - \chi_\delta(t)) / \sin t \, de^{i\rho \cos t} \right| \quad (1.20)$$

$$(1.21)$$

$$\leq C \frac{1}{\rho} + \left| \int_{(-\frac{\pi}{2}, \frac{\pi}{2}) \setminus (-\frac{\delta}{2}, \frac{\delta}{2})} (f(t)(1 - \chi_\delta(t)) / \sin t)' e^{i\rho \cos t} dt \right| \quad (1.22)$$

$$\leq C \frac{1}{\rho} \quad (1.23)$$

From above analysis, we obtain

$$\left| I(\rho, \phi) - \left(\frac{2\pi}{\rho} \right)^{1/2} f(0) e^{i\rho - i\pi/4} \right| \leq C(\phi) \frac{1}{\rho} \quad (1.24)$$

When $N = -1$, we can not use integration by parts again since $f'(\theta)$ is not integrable. However, for any $0 < \lambda_1 < 1$ and $1 < \lambda_2 < 1/\kappa$, there exists $0 < \sigma < \epsilon$, such that $\chi := ((t_1 - \sigma, t_1 + \sigma) \cup (t_2 - \sigma, t_2 + \sigma)) \cap (-\delta, \delta) = \emptyset$, dependent on λ_1, λ_2 and

$$\lambda_1 \kappa < |\sin(t + \phi)| < \lambda_2 \kappa. \quad (1.25)$$

for any $t \in \chi$.

We only analysis the integral on $\chi_1 = (t_1 - \sigma, t_1 + \sigma) \cap [-\pi/2, \pi/2]$ here, which denoted by I_{χ_1} , the proof of I_{χ_2} is similar. It is easy to see that $\sin(t + \phi)$ is monotonic in χ_1 . Without loss of generality, we assume that $\sin(t_1 - \sigma + \phi) < \kappa < \sin(t_1 + \sigma + \phi)$. Let $\sin(t + \phi) = \kappa \sin \theta$ and the implicit mapping from θ to t is denoted by $t(\theta)$ while the inverse mapping by $\theta(t)$, taking the interval χ_1 onto $L_\theta : \theta_1 \rightarrow \pi/2 \rightarrow \pi/2 - i\theta_2$ where $\sin(t_1 - \sigma + \phi) = \kappa \sin \theta_1, \sin(t_1 + \sigma + \phi) = \kappa \sin(\pi/2 - i\theta_2)$. By substituting $t(\theta)$ into I_{χ_1} , we have

$$I_{\chi_1} = \int_{t_1 - \sigma}^{t_1 + \sigma} \frac{f(t)(\kappa^2 - \sin^2(t + \phi))^{1/2}}{(\kappa^2 - \sin^2(t + \phi))^{1/2}} e^{i\rho \cos t} \quad (1.26)$$

$$= \int_{L_\theta} \frac{\kappa f(t(\theta)) \cos \theta}{(1 - \kappa^2 \sin^2 \theta)^{1/2}} e^{i\rho(\cos(t(\theta)))} d\theta \quad (1.27)$$

$$= \int_{L_\theta} \frac{\kappa g_1(t(\theta)) \cos \theta + g_2(t(\theta))}{(1 - \kappa^2 \sin^2 \theta)^{1/2}} e^{i\rho(\cos(t(\theta)))} d\theta \quad (1.28)$$

$$:= \int_{L_\theta} \frac{h(\theta)}{(1 - \kappa^2 \sin^2 \theta)^{1/2}} e^{i\rho(\cos(t(\theta)))} d\theta \quad (1.29)$$

Observe that $h(\theta)$ and $\partial h / \partial \theta$ are integrable on the path L_θ by (1.11). A simple computation show that

$$\frac{dt(\theta)}{d\theta} = \frac{\kappa \cos \theta}{\cos(t + \phi)} \quad \frac{d^2 t(\theta)}{d\theta^2} = \frac{\kappa^2 \cos^2 \theta \sin(t + \phi) - \kappa \sin \theta \cos^2(t + \phi)}{\cos^3(t + \phi)}$$

Then we can obtain

$$\frac{d \cos t}{d\theta} = \frac{-\kappa \sin t \cos \theta}{\cos(t + \phi)}$$

$$\begin{aligned}
\frac{d^2 \cos t}{d\theta^2} &= \frac{d^2 \cos t}{dt^2} \left(\frac{dt}{d\theta} \right)^2 + \frac{d \cos t}{dt} \frac{d^2 t}{d\theta^2} \\
&= \frac{-\kappa^2 \cos^2 \theta \cos t}{\cos^2(t + \phi)} + \frac{\kappa \sin \theta \cos^2(t + \phi) \sin t - \kappa^2 \cos^2 \theta \sin(t + \phi) \sin t}{\cos^3(t + \phi)} \\
&= \frac{-\kappa^2 \cos^2 \theta \cos \phi + \kappa \sin \theta \cos^2(t + \phi) \sin t}{\cos^3(t + \phi)} \\
&= \frac{(\sin^2(t + \phi) - \kappa^2) \cos \phi + \cos^2(t + \phi) \sin(t + \phi) \sin t}{\cos^3(t + \phi)}
\end{aligned}$$

Since $|\sin t| > |\sin \delta|$ and $1 - \lambda_2^2 \kappa^2 < \cos^2(t + \phi) < 1 - \lambda_1^2 \kappa^2$ for $t \in \chi_1$, it follows that $\theta = \pi/2$ is the only stationary point of $\cos(t(\theta))$ and

$$\left| \frac{d^2 \cos t}{d\theta^2}(\pi/2) \right| = \frac{(1 - \kappa^2)\kappa}{(1 - \kappa^2)^{3/2}} |\sin t| > \frac{(1 - \kappa^2)\kappa}{(1 - \kappa^2)^{3/2}} \sin \delta \quad (1.30)$$

Therefore, we can choose appropriate λ_1, λ_2 such that

$$\left| \frac{d^2 \cos t}{d\theta^2} \right| > \frac{(1 - \kappa^2)\kappa}{(1 - \kappa^2)^{3/2}} \sin \delta \quad (1.31)$$

for any $\theta \in \theta(\chi_1)$. According to lemma (6.1), we obtain $|I_{\chi_1}| \leq C \frac{1}{\rho^{1/2}}$, and also $|I_{\chi_2}| \leq C \frac{1}{\rho^{1/2}}$. Using integration by parts, we get

$$\left| \int_{[-\pi/2, \pi/2] \setminus ((-\delta, \delta) \cup \chi)} f(t)(1 - \chi_\delta(t)) e^{i\rho \cos t} dt \right| \leq C \frac{1}{\rho}$$

Consequently, for $N = -1$ and $\phi \neq \phi_\kappa$, we get $|I(\rho, \phi)| \leq \frac{1}{\rho^{1/2}}$.

We now turn to the case of $\phi = \phi_\kappa$. By (1.11), we can define χ_ϵ similarly and also decompose I into I_1 and I_2 . Using the same argument above, we can easily carry out that: for $N = 1$, we have $|I_2| \leq C \frac{1}{\rho}$; for $N = -1$, we have $|I_2| \leq C \frac{1}{\rho^{1/2}}$. Finally, it remains to analysis I_1 . By (1.11), we have

$$\begin{aligned}
I_1 &= \int_{-\epsilon}^{\epsilon} g_1 \chi_\epsilon + g_2 \chi_\epsilon (\sin^2 \phi_\kappa - \sin^2(t + \phi_\kappa))^{N/2} e^{i\rho \cos t} dt \\
&= \int_{-\epsilon}^{\epsilon} g_1 \chi_\epsilon + g_2 \chi_\epsilon (-2(\sin \phi_\kappa + \sin(t + \phi_\kappa)) \cos \frac{2\phi_\kappa + t}{2} \sin t/2)^{N/2} e^{i\rho \cos t} dt \\
&= \int_{\mathbb{R}} g_1 \chi_\epsilon + g_2 \chi_\epsilon ((\sin \phi_\kappa + \sin(t + \phi_\kappa)) \cos \frac{2\phi_\kappa + t}{2})^{N/2} (-2 \sin t/2)^{N/2} e^{i\rho \cos t} dt
\end{aligned}$$

Also, substituting $t(s) = 2 \arcsin s/2$ for t in I_1 , it follows that

$$I_1 = \int_{\mathbb{R}} h_1(s) e^{-i\rho \frac{s^2}{2}} + h_2(s) (-s)^{N/2} e^{-i\rho \frac{s^2}{2}} \quad (1.32)$$

$$= I_{11} + I_{12} \quad (1.33)$$

where

$$\begin{aligned}
h_1(s) &= g_1(t(s)) \chi_\epsilon(t(s)) \sqrt{1 - s^2/4} e^{i\rho} \\
h_2(s) &= g_2 \chi_\epsilon((\sin \phi_\kappa + \sin(t + \phi_\kappa)) \cos \frac{2\phi_\kappa + t}{2})_{t=t(s)}^{N/2} \sqrt{1 - s^2/4} e^{i\rho}
\end{aligned}$$

and $h_1(s), h_2(s) \in C_c^\infty(\mathbb{R})$. Using stationary phase lemma similarly, if $N = 1$,

$$I_{11} = \left(\frac{2\pi}{\rho}\right)^{1/2} g_1(0) e^{i\rho - i\pi/4} + O\left(\frac{1}{\rho}\right) \quad (1.34)$$

$$= \left(\frac{2\pi}{\rho}\right)^{1/2} f(0) e^{i\rho - i\pi/4} + O\left(\frac{1}{\rho}\right) \quad (1.35)$$

if $N = -1$, we get $|I_{11}| \leq C \frac{1}{\rho^{1/2}}$. For I_{12} , we have

$$I_{12} = \int_0^\infty (ih_2(s) + h_2(-s)) s^{N/2} e^{-i\rho s^2/2} ds \quad (1.36)$$

$$= \frac{1}{2} \int_0^\infty (ih_2(\sqrt{s}) + h_2(-\sqrt{s})) s^{N/4-1/2} e^{-i\rho s/2} ds \quad (1.37)$$

By lemma (1.2), we get $|I_{12}| \leq C \frac{1}{\rho^{(N+2)/4}}$. \square

2. Some draft about Green Tensor Analysis

Let substitute $\xi = k \sin \theta$ into integral and shift the variable, we have

$$I(y) = \int_{\mathbb{R}} f(\xi) e^{i\xi y_1 + \mu(\xi) y_2} d\xi = \int_{\mathbb{R}} f(\xi) e^{i\xi(y_1 - z_1) + \mu(\xi)(y_2 - z_2)} e^{i\xi z_1 + \mu(\xi) z_2} d\xi \quad (2.1)$$

$$= k \int_L f(k \sin \theta) \cos \theta e^{ik|y-z| \cos(\theta-\eta)} e^{i|z| \cos(\theta-\phi)} d\theta \quad (2.2)$$

$$= k \int_{L_\phi} f(k \sin(\theta + \phi)) \cos(\theta + \phi) e^{ik|y-z| \cos(\theta+\phi-\eta)} e^{i|z| \cos \theta} d\theta \quad (2.3)$$

$$= k \int_L f(k \sin(\theta + \phi)) \cos(\theta + \phi) e^{ik|y-z| \cos(\theta+\phi-\eta)} e^{i|z| \cos \theta} d\theta \quad (2.4)$$

where $y_1, y_2 > 0$, $\sin \phi = \frac{z_1}{|z|}$, $\cos \phi = \frac{z_2}{|z|}$, $0 < \phi < \pi/2$ and $\sin \eta = \frac{y_1 - z_1}{|y - z|}$, $\cos \eta = \frac{y_2 - z_2}{|y - z|}$, $0 < \eta < \pi$. It is easy to see that $\phi + \eta < \pi$. Roughly, using stationary phase lemma, we obtain:

$$I(y) = f(k \sin \phi) k \cos \phi e^{ik|y-z| \cos(\phi-\eta)} \left(\frac{2\pi}{|z|}\right)^{1/2} e^{i|z| - i\frac{\pi}{4}} (1 + O(\frac{1}{|z|})) \quad (2.5)$$

$$\cos(a + ib) = \frac{e^b + e^{-b}}{2} \cos a + i \frac{e^{-b} - e^b}{2} \sin a \quad (2.6)$$

$$\sin(a + ib) = \frac{e^b + e^{-b}}{2} \sin a + i \frac{e^b - e^{-b}}{2} \cos a \quad (2.7)$$

When $\theta \in (-a - \pi/2, -a - \pi/2 + i\infty)$, let $\theta = -a - \pi/2 + it$, where $t > 0$, $0 \leq a \leq \phi$, then

$$\begin{aligned} & -\text{Im}(|z| \cos \theta + |y - z| \cos(\theta + \phi - \eta)) \\ &= |z| \sin(a + \pi/2) + |y - z| \sin(a + \pi/2 - \phi + \eta) \end{aligned} \quad (2.8)$$

$$= |z| \cos a + |y - z| \cos(a - \phi + \eta) \quad (2.9)$$

$$= |z| \cos a + \cos a |y - z| (\cos \phi \cos \eta + \sin \phi \sin \eta) \quad (2.10)$$

$$+ \sin a |y - z| (\sin \phi \cos \eta - \cos \phi \sin \eta) \quad (2.11)$$

$$= |z| \cos a + \cos a((y_2 - z_2) \cos \phi + (y_1 - z_1) \sin \phi) \quad (2.12)$$

$$+ \sin a((y_2 - z_2) \sin \phi - (y_1 - z_1) \cos \phi) \quad (2.13)$$

$$= y_1 \sin(\phi - a) + y_2 \cos(\phi - a) > 0 \quad (2.14)$$

Now, Using Cauchy Integral Theorem, we have

$$I(y) = k \int_L f(k \sin(\theta + \phi)) \cos(\theta + \phi) e^{\mathbf{i}k|y-z| \cos(\theta+\phi-\eta)} e^{\mathbf{i}|z| \cos \theta} d\theta \quad (2.15)$$

Let $L_1 = (-\pi/2, -\pi/2 + \mathbf{i}\infty)$ and $\theta = -\pi/2 + \mathbf{i}t, t > 0$, then

$$I_1(y) = k \int_{L_1} f(k \sin(\theta + \phi)) \cos(\theta + \phi) e^{\mathbf{i}k|y-z| \cos(\theta+\phi-\eta)} e^{\mathbf{i}|z| \cos \theta} d\theta \quad (2.16)$$

$$= \quad (2.17)$$

$$I(y) = f(k \sin \phi) k \cos \phi e^{\mathbf{i}k|y-z| \cos(\phi-\eta)} \left(\frac{2\pi}{|z|}\right)^{1/2} e^{\mathbf{i}|z| - \mathbf{i}\frac{\pi}{4}} \quad (2.18)$$

$$+ \frac{kz_2}{|z|} O\left(\left(\frac{1}{k|z|}\right)^{3/4} + \frac{1}{k|y|}\right) + \frac{kz_1}{|z|} O\left(\left(\frac{1}{k|z|}\right)^{5/4} + \left(\frac{1}{k|y|}\right)^2\right) \quad (2.19)$$

It is easy to see

$$\int_{-d}^d \frac{k}{(k|x-z|)^\alpha} \frac{1}{(k|x-y|)^\beta} dx_1 \leq C \left(\frac{1}{(kz_2)^{\alpha+\beta-1}} + \frac{1}{(ky_2)^{\alpha+\beta-1}} \right) \quad (2.20)$$

where $z, y \in \mathbb{R}_+^2$, $x \in \Gamma_0$ and $\alpha + \beta > 0$.

$$e^{\mathbf{i}\mu y_2 + \mathbf{i}\xi(x_1 - y_1)} = e^{\mathbf{i}\mu y_2 - \mathbf{i}y_2 / \tan \phi} = e^{\mathbf{i}y_2(\mu - \xi / \tan \phi)} \quad (2.21)$$

Another method

$$\int_{-\pi/2}^{\pi/2} f(k \sin(\theta + \psi)) k \cos(\theta + \psi) e^{\mathbf{i}k|x-y| \cos \theta} \quad (2.22)$$

$$= \int_{-\pi/2}^{\pi/2} f(k \sin(\theta + \psi)) k \cos(\theta + \psi) e^{\mathbf{i}k|x-y| \cos(\theta+\psi-\psi)} \quad (2.23)$$

$$= \int_{-\pi/2}^{\pi/2} f(k \sin(\theta + \psi)) k \cos(\theta + \psi) e^{\mathbf{i}ky_2 \cos(\theta+\psi) + \mathbf{i}k|x_1 - y_1| \sin(\theta+\psi)} \quad (2.24)$$

$$= \int_{-\pi/2}^{\pi/2} f(k \sin(\theta + \psi)) k \cos(\theta + \psi) \quad (2.25)$$

$$e^{\mathbf{i}k(y_2 - z_2) \cos(\theta+\psi) + \mathbf{i}k(|x_1 - y_1| - |x_1 - z_1|) \sin(\theta+\psi) + \mathbf{i}k|z| \cos(\theta+\psi-\phi)} \quad (2.26)$$

3. Finite Aperture Point Spread Function

If $x \in \Gamma_0$ and $z, y \in \mathbb{R}_+^2$, by lemma (??) we have

$$\begin{aligned} G(x, y) &= \frac{\mathbf{i}k_s}{\mu\sqrt{2\pi}} \frac{1}{\delta(\xi)} \begin{pmatrix} \mu_s \beta & \xi \beta \\ 2\xi \mu_s \mu_p & 2\xi^2 \mu_p \end{pmatrix}_{\xi=k_s \frac{x_1 - y_1}{|x - y|}} \frac{y_2}{|x - y|} \frac{1}{(k_s |x - y|)^{1/2}} e^{\mathbf{i}k_s |x - y| - \mathbf{i}\frac{\pi}{4}} \\ &+ \frac{\mathbf{i}k_p}{\mu\sqrt{2\pi}} \frac{1}{\delta(\xi)} \begin{pmatrix} 2\xi^2 \mu_s & -2\xi \mu_s \mu_p \\ -\xi \beta & \mu_p \beta \end{pmatrix}_{\xi=k_p \frac{x_1 - y_1}{|x - y|}} \frac{y_2}{|x - y|} \frac{1}{(k_p |x - y|)^{1/2}} e^{\mathbf{i}k_p |x - y| - \mathbf{i}\frac{\pi}{4}} \end{aligned} \quad (3.1)$$

$$\begin{aligned}
& +O\left(\frac{y_2}{|x-y|} \frac{1}{(k_s|x-y|)^{3/4}} + \frac{|x_1-y_1|}{|x-y|} \frac{1}{(k_s|x-y|)^{5/4}}\right) \\
& := \mathcal{G}_s(x, y) + \mathcal{G}_p(x, y) + O\left(\frac{y_2}{|x-y|} \frac{1}{(k_s|x-y|)^{3/4}} + \frac{|x_1-y_1|}{|x-y|} \frac{1}{(k_s|x-y|)^{5/4}}\right) \\
T_D(x, z) &= \frac{k_s}{\sqrt{2\pi}} \frac{1}{\gamma(\xi)} \begin{pmatrix} \mu_s \mu_p & \xi \mu_p \\ \xi \mu_s & \xi^2 \end{pmatrix}_{\xi=k_s \frac{x_1-z_1}{|x-z|}} \frac{z_2}{|x-z|} \frac{1}{(k_s|x-z|)^{1/2}} e^{i k_s |x-z| - i \frac{\pi}{4}} \\
& + \frac{k_p}{\sqrt{2\pi}} \frac{1}{\gamma(\xi)} \begin{pmatrix} \xi^2 & -\xi \mu_p \\ -\xi \mu_s & \mu_p \mu_s \end{pmatrix}_{\xi=k_p \frac{x_1-z_1}{|x-z|}} \frac{z_2}{|x-z|} \frac{1}{(k_p|x-z|)^{1/2}} e^{i k_p |x-z| - i \frac{\pi}{4}} \quad (3.2) \\
& + O\left(\frac{k_s z_2}{|x-z|} \frac{1}{(k_s|x-z|)^{3/4}} + \frac{k_s |x_1-z_1|}{|x-z|} \frac{1}{(k_s|x-z|)^{5/4}}\right) \\
& := \mathcal{T}_s(x, z) + \mathcal{T}_p(x, z) + O\left(\frac{k_s z_2}{|x-z|} \frac{1}{(k_s|x-z|)^{3/4}} + \frac{k_s |x_1-z_1|}{|x-z|} \frac{1}{(k_s|x-z|)^{5/4}}\right)
\end{aligned}$$

Now we consider the finite aperture point spread function $J_d(z, y)$:

$$\int_{-d}^d (T_D(x_1, 0; z_1, z_2))^T \overline{G(x_1, 0; y_1, y_2)} dx_1 \quad (3.3)$$

Recall following standard asymptotic expansion:

$$|x-y| = |x-z| + \widehat{x-z} \cdot (z-y) + O\left(\frac{|y-z|^2}{|x-z|}\right) \quad (3.4)$$

$$|y|^{-\alpha} = |z|^{-\alpha} \left(1 + \frac{|y|-|z|}{|z|}\right)^{-\alpha} = |z|^{-\alpha} \left(1 + O\left(\frac{|y-z|}{|z|}\right)\right) \quad (3.5)$$

$$e^{it} = 1 + O(t) \quad (3.6)$$

$$|a^{1/2} - b^{1/2}| \leq C \sqrt{|a-b|} \quad (3.7)$$

where $x, y, z \in \mathbb{R}^2$, $t, a, b \in \mathbb{R}$ and $\alpha > 0$.

Lemma 3.1 For any $z, y \in \mathbb{R}_+^2$, $J_d(z, y) = F(z, y) + O\left((1 + \frac{|y-z|}{z_2}) \left(\frac{1}{k_s z_2}\right)^{1/4} + \frac{(k_s |y-z|)^2}{k_s z_2} + \left(\frac{|y-z|}{z_2}\right)^{1/2}\right)$, where

$$F(z, y) = -\frac{\mathbf{i}}{2\pi\mu} \int_{\theta_1^d}^{\theta_2^d} f_s(\theta) \begin{pmatrix} \sin^2 \theta & \sin \theta \cos \theta \\ \sin \theta \cos \theta & \cos^2 \theta \end{pmatrix} e^{i k_s (z_1 - y_1) \cos \theta + i k_s (z_2 - y_2) \sin \theta} d\theta \quad (3.8)$$

$$-\frac{\mathbf{i}}{2\pi\mu} \int_{\theta_1^d}^{\theta_2^d} f_p(\theta) \begin{pmatrix} \cos^2 \theta & -\sin \theta \cos \theta \\ -\sin \theta \cos \theta & \sin^2 \theta \end{pmatrix} e^{i k_p (z_1 - y_1) \cos \theta + i k_p (z_2 - y_2) \sin \theta} d\theta \quad (3.9)$$

and

$$\begin{aligned}
f_s(\theta) &= \frac{\sin \theta ((\kappa^2 - \cos^2 \theta)^{1/2} (1 - 2 \cos^2 \theta) + 2(\kappa^2 - \cos^2 \theta)^{1/2} \cos^2 \theta)}{(\cos^2 \theta + \sin \theta (\kappa^2 - \cos^2 \theta)^{1/2}) ((1 - 2 \cos^2 \theta)^2 + 4 \cos^2 \theta \sin \theta (\kappa^2 - \cos^2 \theta)^{1/2})} \\
f_p(\theta) &= \frac{\sin \theta (1/\kappa^2 - \cos^2 \theta)^{1/2}}{(\cos^2 \theta + \sin \theta (1/\kappa^2 - \cos^2 \theta)^{1/2}) ((1/\kappa^2 - 2 \cos^2 \theta)^2 + 4 \cos^2 \theta \sin \theta (1/\kappa^2 - \cos^2 \theta)^{1/2})}
\end{aligned}$$

where $0 < \theta_1^d < \pi/2 < \theta_2^d < \pi$ and $z_2 = (d + z_1) \tan \theta_1^d = (z_1 - d) \tan \theta_2^d$.

Proof.

$$\begin{aligned}
& \frac{y_2}{|x-y|} \frac{1}{(k_s|x-y|)^{3/4}} + \frac{|x_1-y_1|}{|x-y|} \frac{1}{(k_s|x-y|)^{5/4}} \\
&= \left(\frac{z_2}{|x-z|} \frac{1}{(k_s|x-z|)^{3/4}} + \frac{|x_1-z_1|}{|x-z|} \frac{1}{(k_s|x-z|)^{5/4}} \right) (1 + O(\frac{|y-z|}{|x-z|})) \\
& \quad |\mu_i(k_j \frac{x_1-y_1}{|x-y|}) - \mu_i(k_j \frac{x_1-z_1}{|x-z|})| \\
& \leq Ck_j \sqrt{\left| \frac{x_1-y_1}{|x-y|} - \frac{x_1-z_1}{|x-z|} \right|} \leq Ck_j \left(\frac{|y-z|}{|x-z|} \right)^{1/2}
\end{aligned}$$

where $i, j = s, p$. By above, we can obtain

$$\mathcal{G}_s(x, y) = \mathcal{G}_s(x, z) e^{\widehat{\mathbf{i}k_s x - z} \cdot (z-y)} + O\left(\frac{(k_s|y-z|)^2}{(k_s|x-z|)^{3/2}}\right) + O\left(\frac{(k_s|y-z|)^{1/2}}{k_s|x-z|}\right) \quad (3.10)$$

$$\mathcal{G}_p(x, y) = \mathcal{G}_p(x, z) e^{\widehat{\mathbf{i}k_p x - z} \cdot (z-y)} + O\left(\frac{(k_p|y-z|)^2}{(k_p|x-z|)^{3/2}}\right) + O\left(\frac{(k_p|y-z|)^{1/2}}{k_p|x-z|}\right) \quad (3.11)$$

For $l > 1$, a simple computation show that

$$\int_{-d}^d \frac{k_s}{(k_s|x-z|)^l} dx_1 = \frac{1}{(k_s z_2)^{l-1}} \int_{\frac{-d-z_1}{z_2}}^{\frac{d-z_1}{z_2}} \frac{1}{(1+t^2)^{l/2}} dt \leq C \frac{1}{(k_s z_2)^{l-1}} \quad (3.12)$$

Let

$$\mathcal{G}_\alpha(x, y) = \frac{\mathbf{i}}{\sqrt{2\pi\mu}} g_\alpha\left(\frac{x_1-y_1}{|x-y|}, \kappa\right) \frac{1}{(k_\alpha|x-y|)^{1/2}} e^{\mathbf{i}k_\alpha|x-y| - \mathbf{i}\frac{\pi}{4}} \quad (3.13)$$

$$\mathcal{T}_\alpha(x, y) = \frac{k_\alpha}{\sqrt{2\pi}} t_\alpha\left(\frac{x_1-z_1}{|x-z|}, \kappa\right) \frac{1}{(k_s|x-z|)^{1/2}} e^{\mathbf{i}k_\alpha|x-z| - \mathbf{i}\frac{\pi}{4}} \quad (3.14)$$

where $\alpha = s, p$. Now, by substituting (3.10-3.11) into $J_d(z, y)$ and using inequality (3.12), we have

$$\begin{aligned}
J_d(z, y) &= \frac{-\mathbf{i}}{2\pi\mu} \int_{-d}^d t_s\left(\frac{x_1-z_1}{|x-z|}, \kappa\right) \overline{g_s\left(\frac{x_1-z_1}{|x-z|}, \kappa\right)} \frac{e^{\widehat{\mathbf{i}k_s x - z} \cdot (y-z)}}{|x-z|} \\
& \quad + t_p\left(\frac{x_1-z_1}{|x-z|}, \kappa\right) \overline{g_p\left(\frac{x_1-z_1}{|x-z|}, \kappa\right)} \frac{e^{\widehat{\mathbf{i}k_p x - z} \cdot (y-z)}}{|x-z|} dx_1 \quad (3.15)
\end{aligned}$$

$$- \frac{\mathbf{i}}{2\pi\mu} \int_{-d}^d t_p\left(\frac{x_1-z_1}{|x-z|}, \kappa\right) \overline{g_s\left(\frac{x_1-z_1}{|x-z|}, \kappa\right)} \frac{e^{\widehat{\mathbf{i}k_s x - z} \cdot (y-z)}}{|x-z|} \quad (3.16)$$

$$+ t_s\left(\frac{x_1-z_1}{|x-z|}, \kappa\right) \overline{g_p\left(\frac{x_1-z_1}{|x-z|}, \kappa\right)} \frac{e^{\widehat{\mathbf{i}k_p x - z} \cdot (y-z)}}{|x-z|} dx_1 \quad (3.17)$$

$$+ O\left((1 + \frac{|y-z|}{z_2}) \left(\frac{1}{k_s z_2}\right)^{1/4} + \frac{(k_s|y-z|)^2}{k_s z_2} + \left(\frac{|y-z|}{z_2}\right)^{1/2}\right) \quad (3.18)$$

$$:= F(z, y) + R(z, y) \quad (3.19)$$

$$+ O\left((1 + \frac{|y-z|}{z_2}) \left(\frac{1}{k_s z_2}\right)^{1/4} + \frac{(k_s|y-z|)^2}{k_s z_2} + \left(\frac{|y-z|}{z_2}\right)^{1/2}\right) \quad (3.20)$$

We denote $\widehat{x-z} = x-z/|x-z| = (\cos(\phi+\pi), \sin(\phi+\pi))$, then taking the substitution $x_1 = z_1 - z_2 \cot \phi$, we obtain

$$F(z, y) = \frac{-\mathbf{i}}{2\pi\mu} \int_{\theta_1^d}^{\theta_2^d} A_s(\phi, \kappa) e^{\mathbf{i}k_s(z_1-y_1) \cos \phi + \mathbf{i}k_s(z_2-y_2) \sin \phi} \quad (3.21)$$

$$+ \frac{-\mathbf{i}}{2\pi\mu} \int_{\theta_1^d}^{\theta_2^d} A_p(\phi, \kappa) e^{\mathbf{i}k_p(z_1-y_1) \cos \phi + \mathbf{i}k_p(z_2-y_2) \sin \phi} \quad (3.22)$$

$$R(z, y) = \frac{-\mathbf{i}}{2\pi\mu} \int_{\theta_1^d}^{\theta_2^d} B_s(\phi, \kappa) e^{\mathbf{i}k_s(z_1-y_1) \cos \phi + \mathbf{i}k_s(z_2-y_2) \sin \phi + (k_p-k_s)|x-z|} \quad (3.23)$$

$$+ \frac{-\mathbf{i}}{2\pi\mu} \int_{\theta_1^d}^{\theta_2^d} B_p(\phi, \kappa) e^{\mathbf{i}k_p(z_1-y_1) \cos \phi + \mathbf{i}k_p(z_2-y_2) \sin \phi + (k_s-k_p)|x-z|} \quad (3.24)$$

It is easy to see that $|R(z, y)| \leq C \frac{|z-y|}{z_2}$. \square

Let

$$g(x_1) = \frac{1}{((x_1 - z_1)^2 + z_2^2)^{3/4} ((x_1 - y_1)^2 + y_2^2)^{1/4}}$$

$$\phi(x_1) = ((x_1 - z_1)^2 + z_2^2)^{1/2} - ((x_1 - y_1)^2 + y_2^2)^{1/2}$$

Then, we have

$$g'(x_1) = -g(x_1) \left[\frac{3(x_1 - z_1)}{2((x_1 - z_1)^2 + z_2^2)} + \frac{(x_1 - y_1)}{2((x_1 - y_1)^2 + y_2^2)} \right]$$

$$\phi'(x_1) = \frac{x_1 - z_1}{((x_1 - z_1)^2 + z_2^2)^{1/2}} - \frac{x_1 - y_1}{((x_1 - y_1)^2 + y_2^2)^{1/2}}$$

$$= \frac{\frac{(x_1 - z_1)^2}{(x_1 - z_1)^2 + z_2^2} - \frac{(x_1 - y_1)^2}{(x_1 - y_1)^2 + y_2^2}}{\frac{x_1 - z_1}{((x_1 - z_1)^2 + z_2^2)^{1/2}} + \frac{x_1 - y_1}{((x_1 - y_1)^2 + y_2^2)^{1/2}}}$$

$$= \frac{(x_1 - z_1)^2 y_2^2 - (x_1 - y_1)^2 z_2^2}{\left(\frac{x_1 - z_1}{((x_1 - z_1)^2 + z_2^2)^{1/2}} + \frac{x_1 - y_1}{((x_1 - y_1)^2 + y_2^2)^{1/2}} \right) ((x_1 - z_1)^2 + z_2^2) ((x_1 - y_1)^2 + y_2^2)}$$

$$\phi''(x_1) = \frac{z_2^2}{((x_1 - z_1)^2 + z_2^2)^{3/2}} - \frac{y_2^2}{((x_1 - y_1)^2 + y_2^2)^{3/2}}$$

Using integration by parts, we can obtain

$$\int_{-d}^d g(x_1) e^{\mathbf{i}\phi(x_1)} dx_1$$

$$= \left(\frac{g(d)}{\phi'(d)} e^{\mathbf{i}\phi(d)} - \frac{g(-d)}{\phi'(-d)} e^{\mathbf{i}\phi(-d)} \right) - \int_{-d}^d \frac{g'(x_1)}{\phi'(x_1)} - \frac{g(x_1)\phi''(x_1)}{(\phi'(x_1))^2} dx_1$$

Assume that

$$|y_1| \leq c_0 d \quad |z_1| \leq c_0 d \quad h \leq y_2, z_2 \leq c_1 h \quad d \leq c_2 h$$

where $0 < c_0 < 1$. Let define $0 < \theta_y, \theta_z < \pi$ such that

$$\cos \theta_y = \frac{x_1 - y_1}{((x_1 - y_1)^2 + y_2^2)^{1/2}}$$

$$\cos \theta_z = \frac{x_1 - z_1}{((x_1 - z_1)^2 + z_2^2)^{1/2}}$$

By mean value theorem and the law of sines, we get

$$\begin{aligned}
|\phi'(x_1)| &= |\cos \theta_z - \cos \theta_y| = |\sin \theta'| |\theta_z - \theta_y| \\
&\geq \frac{h}{(1+c_0)d} |\sin(\theta_z - \theta_y)| \\
&= \frac{h}{(1+c_0)d} \frac{|z-y|}{|x-y|} \sin \theta_{|x-y|} \\
&= \frac{h}{(1+c_0)d} \frac{|z-y|}{|x-z|} \sin \theta_{|x-z|} \\
&\geq \frac{h^2}{(1+c_0)^2 d^2} \frac{|z-y|}{|x-y|} \\
\text{or} \quad &\geq \frac{h^2}{(1+c_0)^2 d^2} \frac{|z-y|}{|x-z|}
\end{aligned}$$

Then we have

$$\begin{aligned}
\left| \frac{g(x_1)}{\phi'(x_1)} \right| &\leq \frac{(1+c_0)^2 d^2}{h^2} \frac{1}{|z-y||x-y|^{1/2}|x-z|^{1/2}} \\
&\leq C \frac{d^2}{h^3} \frac{1}{|z-y|}
\end{aligned}$$

Moreover, by mean value theorem again, we have

$$\begin{aligned}
|\phi''(x_1)| &= \left| \frac{\sin^2 \theta_z}{|x-z|} - \frac{\sin^2 \theta_y}{|x-y|} \right| \\
&= \left| \frac{2 \sin \theta' \cos \theta'}{|x-y'|} (\theta_z - \theta_y) - \frac{\sin^2 \theta'}{|x-y'|^2} (|x-z| - |x-y|) \right| \\
&\leq \pi \frac{|\sin(\theta_z - \theta_y)|}{h} + \frac{|z-y|}{h^2} \\
&\leq \pi \frac{|\sin \theta_{|x-z|}| |z-y|}{h|x-z|} + \frac{|z-y|}{h^2} \\
&\leq C \frac{|z-y|}{h^2}
\end{aligned}$$

Now, it is easy to see that

$$\begin{aligned}
&\left| \int_{-d}^d \frac{g'(x_1)}{\phi'(x_1)} - \frac{g(x_1)\phi''(x_1)}{(\phi'(x_1))^2} dx_1 \right| \\
&\leq C \frac{d^3}{h^4} \frac{1}{|z-y|} + C \frac{d^3}{h^3} \frac{1}{|z-y|} \frac{d^2}{h^3}
\end{aligned}$$

Based on the above analysis, we can obtain

$$\left| \int_{-d}^d z_2 g(x_1) e^{i\phi(x_1)} \right| \leq C \left(\left(\frac{d}{h} \right)^2 + \left(\frac{d}{h} \right)^3 + \left(\frac{d}{h} \right)^5 \right) \frac{1}{|z-y|}$$

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$$\begin{aligned}
\sin \phi_\kappa - \sin(t + \phi) &= -2 \cos\left(\frac{\phi_\kappa + \phi + t}{2}\right) \sin\left(\frac{t + \phi - \phi_\kappa}{2}\right) \\
\sin\left(\frac{t + \phi - \phi_\kappa}{2}\right) &= \sin \frac{t}{2} \cos\left(\frac{\phi - \phi_\kappa}{2}\right) + \cos \frac{t}{2} \sin\left(\frac{\phi - \phi_\kappa}{2}\right)
\end{aligned}$$

Some think, substituting $t = 2 \arcsin s/2$ into following integral

$$\begin{aligned} & \int_0^\infty \chi(t)(\sin \phi_\kappa - \sin(t + \phi))^{1/2} e^{-i\rho \cos t} \\ &= \int_0^\infty \chi(t(s))(-s \cos(\frac{\phi - \phi_\kappa}{2}) - \sqrt{4 - s^2} \sin(\frac{\phi - \phi_\kappa}{2}))^{1/2} e^{-i\rho s^2/2} \\ &= \int_0^\infty \chi(t)(-\sqrt{t} \cos(\frac{\phi - \phi_\kappa}{2}) - \sqrt{4 - t} \sin(\frac{\phi - \phi_\kappa}{2}))^{1/2} t^{-1/2} e^{-i\rho t/2} \end{aligned}$$

Let

$$\begin{aligned} f(t) &= t^{-1/2} e^{-i\rho t/2} \\ g(t) &= - \int_t^{t-i\infty} x^{-1/2} e^{-i\rho x/2} dx \\ &= i \int_0^\infty (t - ix)^{-1/2} e^{-i\rho t - \rho x} dx \end{aligned}$$

It is easy to see that $g'(t) = f(t)$. Then we have

$$\begin{aligned} &= \int_0^\infty \chi(t)(-\sqrt{t} \cos(\frac{\phi - \phi_\kappa}{2}) - \sqrt{4 - t} \sin(\frac{\phi - \phi_\kappa}{2}))^{1/2} t^{-1/2} e^{-i\rho t/2} \\ &= \chi(0)(-2 \sin(\frac{\phi - \phi_\kappa}{2}))^{1/2} g(0) \\ &\quad - \int_0^\infty (\chi(t)(-\sqrt{t} \cos(\frac{\phi - \phi_\kappa}{2}) - \sqrt{4 - t} \sin(\frac{\phi - \phi_\kappa}{2}))^{1/2})' g(t) dt \end{aligned}$$

We get

$$\begin{aligned} g(x) &= \int_0^\infty \chi(t)(-\sqrt{t} \cos(\frac{\phi - \phi_\kappa}{2}) - \sqrt{4 - t} \sin(\frac{\phi - \phi_\kappa}{2}))^{1/2} t^{-1/2} (t - ix)^{-1/2} e^{-i\rho t} dt \\ R(\rho) &= \int_0^\infty g(x) e^{-\rho x} dx \end{aligned}$$

Because $\chi(t)$ has compact support $(-\delta, \delta)$, we obtain

$$gg(x) = \int_0^\delta (\sqrt{t} \cos(\theta) - \sqrt{4 - t} \sin \theta)^{-1/2} t^{-1/2} (t^2 + x^2)^{-1/4} dt$$

where $\theta = \frac{\phi - \phi_\kappa}{2}$. For $x > 0$, Put $L(x)$:

$$\begin{aligned} & \int_0^a \frac{1}{t^{3/4}} \frac{1}{(t^2 + x^2)^{1/4}} dt \\ &= 4 \int_0^a \frac{1}{(t^2 + x^2)^{1/4}} dt^{1/4} \\ &= 4 \int_0^{a^{1/4}} \frac{1}{(t^8 + x^2)^{1/4}} dt \\ &= 4x^{-1/4} \int_0^{(\frac{a}{x})^{1/4}} \frac{1}{(t^8 + 1)^{1/4}} dt \\ &= 4x^{-1/4} \int_0^{(\frac{a}{x})^{1/4}} \frac{1}{(t^8 + 1)^{1/4}} dt \\ &\leq 4x^{-1/4} \int_0^\infty \frac{1}{(t^8 + 1)^{1/4}} dt \end{aligned}$$

Back to analysis $gg(x)$, we have

$$\begin{aligned}
gg(x) &\leq \int_0^\delta \left| \frac{\sqrt{t} + 2|\sin \theta|}{t - 4\sin^2 \theta} \right|^{1/2} t^{-1/2} (t^2 + x^2)^{-1/4} dt \\
&= \int_0^\delta \left| \frac{1}{\sqrt{t} - 2|\sin \theta|} \right|^{1/2} t^{-1/2} (t^2 + x^2)^{-1/4} dt \\
&= 2 \int_0^{\sqrt{\delta}} \left| \frac{1}{t - 2|\sin \theta|} \right|^{1/2} (t^4 + x^2)^{-1/4} dt \\
&= 2 \int_{-2|\sin \theta|}^{\sqrt{\delta} - 2|\sin \theta|} |t|^{-1/2} ((t + 2|\sin \theta|)^4 + x^2)^{-1/4} dt \\
&\leq 4 \int_0^{\delta^{1/4}} (t^8 + x^2)^{-1/4} dt + 4 \int_0^{\sqrt{2|\sin \theta|}} ((t^2 - 2|\sin \theta|)^4 + x^2)^{-1/4} dt \\
&\leq Cx^{-1/4} (1 + \int_0^{\sqrt{2|\sin \theta|}} ((t^2 - 2|\sin \theta|)^4/x + x)^{-1/4} dt) \\
&\leq Cx^{-1/4} (1 + \int_0^{\sqrt{2|\sin \theta|}} (t^2 - 2|\sin \theta|)^{-1/2} dt) \\
&= Cx^{-1/4} (1 + \int_0^1 (1 - t^2)^{-1/2} dt) \leq Cx^{-1/4}
\end{aligned}$$

Immediately, we can obtain

$$|g(x)| \leq Cx^{-1/4}$$

It follows that

$$R(\rho) \leq \int_0^\infty x^{-1/4} e^{-\rho x} \leq C\rho^{-3/4}$$

5. stationary of phase lemma

Lemma 5.1 Assume that $0 < \kappa := \sin \phi_\kappa < 1$, $0 < \phi_\kappa < \pi/2$, $0 \leq \phi \leq \pi/2$. Let

$$f(t, \phi) := F(\sin(t + \phi), \cos(t + \phi), (\kappa^2 - \sin^2(t + \phi))^{1/2}) \quad (5.1)$$

be a complexed function in $C([-\pi/2, \pi/2] \times [0, \pi/2])$. Moreover, its partial derivative with respect to t can be represented as

$$\frac{\partial f(t, \phi)}{\partial t} = g(t, \phi) (\kappa^2 - \sin^2(t + \phi))^{-1/2} \quad (5.2)$$

where $g(t, \phi)$ is uniformly bounded. Then for any $\rho \geq 1$, we have

$$\begin{aligned}
&\left| I(\rho, \phi) := \int_{-\pi/2}^{\pi/2} f(t) e^{i\rho \cos t} dt - \left(\frac{2\pi}{\rho} \right)^{1/2} f(0) e^{i\rho - i\pi/4} \right| \\
&\leq C \frac{1}{\rho^{3/4}}
\end{aligned} \quad (5.3)$$

Proof. Solving the following equation:

$$\kappa^2 - \sin^2(t + \phi) = 0$$

we have, if $0 < \phi < \pi/2 - \phi_\kappa$,

$$t_1(\phi) = \phi_\kappa - \phi \quad t_2(\phi) = -\phi_\kappa - \phi$$

and if $\pi/2 - \phi_\kappa \leq \phi < \pi/2$,

$$t_1(\phi) = \phi_\kappa - \phi \quad t_2(\phi) = \pi - \phi_\kappa - \phi$$

Since $|t_2(\phi)| < \phi_\kappa$ or $|t_2(\phi)| < \pi/2 - \phi_\kappa$, we now define $\delta := \min(\frac{\phi_\kappa}{2}, \frac{\pi/2 - \phi_\kappa}{2})$ and it is easy to see that

$$\kappa + \sin(t + \phi) \neq 0 \quad (5.4)$$

$$\cos\left(\frac{t + \phi + \phi_\kappa}{2}\right) \neq 0 \quad (5.5)$$

for any $(t, \phi) \in [-\delta, \delta] \times [0, \pi/2]$. Let $\chi_\delta \in C_0^\infty(-\pi/2, \pi/2)$ be the cut-off function with that $0 \leq \chi_\delta \leq 1$, $\chi_\delta = 1$ in $(-\delta/2, \delta/2)$ and $\chi_\delta = 0$ in $(-\pi/2, \pi/2) \setminus (-\delta, \delta)$. Then we can divide I into two parts such that

$$\begin{aligned} I &= \int_{-\delta}^{\delta} f(t) \chi_\delta(t) e^{i\rho \cos t} dt + \int_{-\pi/2}^{\pi/2} f(t) (1 - \chi_\delta(t)) e^{i\rho \cos t} dt \\ &=: I_1 + I_2 \end{aligned}$$

Substituting $t(s) = 2 \arcsin s/2$ for t in I_1 , we can obtain

$$I_1 = \int_{-2 \sin \frac{\delta}{2}}^{2 \sin \frac{\delta}{2}} f(t(s)) \chi_\delta(t(s)) \frac{1}{\sqrt{1 - s^2/4}} e^{i\rho} e^{-i\rho s^2/2} ds \quad (5.6)$$

$$= \int_0^{2 \sin \frac{\delta}{2}} (f(t(s)) \chi_\delta(t(s)) + f(-t(s)) \chi_\delta(-t(s))) \frac{1}{\sqrt{1 - s^2/4}} e^{i\rho} e^{-i\rho s^2/2} ds \quad (5.7)$$

$$:= I_{11} + I_{12} \quad (5.8)$$

Taking substitution $s = \sqrt{x}$, we get

$$I_{11} = \frac{1}{2} \int_0^{(2 \sin \frac{\delta}{2})^2} f(t(\sqrt{x})) \chi_\delta(t(\sqrt{x})) \frac{1}{\sqrt{1 - x/4}} x^{-1/2} e^{i\rho} e^{-i\rho x/2} dx$$

Observe that

$$\begin{aligned} \sin \phi_\kappa - \sin(t + \phi) &= -2 \cos\left(\frac{\phi_\kappa + \phi + t}{2}\right) \sin\left(\frac{t + \phi - \phi_\kappa}{2}\right) \\ \sin\left(\frac{t + \phi - \phi_\kappa}{2}\right) &= \sin \frac{t}{2} \cos\left(\frac{\phi - \phi_\kappa}{2}\right) + \cos \frac{t}{2} \sin\left(\frac{\phi - \phi_\kappa}{2}\right) \\ &:= \sin \frac{t}{2} \cos \theta + \cos \frac{t}{2} \sin \theta \end{aligned}$$

where $\theta = \frac{\phi - \phi_\kappa}{2}$. By lemma (1.3) and using representation (5.2), inequality (5.4-5.5), it follows that

$$\begin{aligned} &|I_{11} - \frac{1}{2} \sqrt{\frac{2\pi}{\rho}} f(0) e^{i\rho - i\frac{\pi}{4}}| \\ &\leq \int_0^\infty e^{-\rho y} dy \int_0^{(2 \sin \frac{\delta}{2})^2} \left| \frac{\partial(f(t(\sqrt{x})) \chi_\delta(t(\sqrt{x})) \frac{1}{\sqrt{1 - x/4}})}{\partial x} \right| (x^2 + y^2)^{-\frac{1}{4}} dx \end{aligned}$$

$$\begin{aligned}
&\leq C \int_0^\infty e^{-\rho y} dy \int_0^{(2 \sin \frac{\delta}{2})^2} |\sqrt{x} \cos \theta + \sqrt{4-x} \sin \theta|^{-1/2} x^{-1/2} (x^2 + y^2)^{-\frac{1}{4}} dx \\
&\leq C \int_0^\infty e^{-\rho y} dy \int_0^{(2 \sin \frac{\delta}{2})^2} \frac{(\sqrt{x} |\cos \theta| + \sqrt{4-x} |\sin \theta|)^{1/2}}{|x - 4 \sin^2 \theta|^{1/2}} x^{-1/2} (x^2 + y^2)^{-\frac{1}{4}} dx \\
&\leq C \int_0^\infty e^{-\rho y} dy \int_0^{(2 \sin \frac{\delta}{2})^2} \frac{1}{|\sqrt{x} - 2 \sin \theta|^{1/2}} x^{-1/2} (x^2 + y^2)^{-\frac{1}{4}} dx \\
&\leq C \int_0^\infty e^{-\rho y} dy \int_0^{2 \sin \frac{\delta}{2}} \frac{1}{|x - 2 \sin |\theta||^{1/2}} (x^4 + y^2)^{-\frac{1}{4}} dx \\
&\leq C \int_0^\infty e^{-\rho y} dy \int_{-2 \sin |\theta|}^{2 \sin \frac{\delta}{2} - 2 \sin |\theta|} \frac{1}{|x|^{1/2}} ((x + 2 \sin |\theta|)^4 + y^2)^{-\frac{1}{4}} dx \\
&\leq C \int_0^\infty e^{-\rho y} dy \int_{-2 \sin |\theta|}^{2 \sin \frac{\delta}{2}} \frac{1}{|x|^{1/2}} ((x + 2 \sin |\theta|)^4 + y^2)^{-\frac{1}{4}} dx \\
&\leq C \int_0^\infty e^{-\rho y} dy \left(\int_0^{\sqrt{2 \sin \frac{\delta}{2}}} (x^8 + y^2)^{-\frac{1}{4}} dx + \int_0^{\sqrt{2 \sin |\theta|}} ((x^2 - 2 \sin |\theta|)^4 + y^2)^{-\frac{1}{4}} dx \right) \\
&\leq C \int_0^\infty e^{-\rho y} dy \left(y^{-\frac{1}{4}} \int_0^\infty (x^8 + 1)^{-\frac{1}{4}} dx + y^{-\frac{1}{4}} \int_0^{\sqrt{2 \sin |\theta|}} (2 \sin |\theta| - x^2)^{-\frac{1}{2}} dx \right) \\
&\leq C \int_0^\infty y^{-\frac{1}{4}} e^{-\rho y} dy \left(\int_0^\infty (x^8 + 1)^{-\frac{1}{4}} dx + \int_0^1 (1 - x^2)^{-\frac{1}{2}} dx \right) \leq C \frac{1}{\rho^{3/4}}
\end{aligned}$$

Using the same argument, we can also carry out

$$|I_{12} - \frac{1}{2} \sqrt{\frac{2\pi}{\rho}} f(0) e^{i\rho - i\frac{\pi}{4}}| \leq C \frac{1}{\rho^{3/4}} \quad (5.9)$$

It remains to estimate I_2 . Note that there exists $m > 0$ such that $|\sin t| \geq m$ for any $t \in [-\pi/2, \pi/2] \setminus (-\delta/2, \delta/2)$. Upon integration by parts and representation (5.2) again, we have

$$\begin{aligned}
|I_{12}| &\leq C \rho^{-1} \left(1 + \left| \int_{[-\pi/2, \pi/2] \setminus (-\delta/2, \delta/2)} \frac{\partial(f(t)(1 - \chi_\delta(t)))}{\partial t} \frac{1}{\sin t} dt \right| \right) \\
&\leq C \rho^{-1} \left(1 + \int_{-\pi/2}^{\pi/2} \left| \frac{\partial(f(t)(1 - \chi_\delta(t)))}{\partial t} \right| dt \right) \\
&\leq C \rho^{-1} \left(1 + \int_{-\pi/2}^{\pi/2} |(\kappa^2 - \sin^2(t + \phi))^{-1/2}| dt \right) \\
&\leq C \rho^{-1} \left(1 + \int_{-\pi/2}^{\pi/2} |(\kappa^2 - \sin^2 t)^{-1/2}| dt \right) \\
&\leq C \rho^{-1}
\end{aligned}$$

This completes the proof. \square

6. cross term of psf, 17.11.15

We need the following slight generalization of Van der Corput lemma for the oscillatory integral [3, P.152].

Lemma 6.1 *Let $-\infty < a < b < \infty$, and u is a C^k function u in (a, b) .*

1. *If $|u'(t)| \geq 1$ for $t \in (a, b)$ and u' is monotone in (a, b) , then for any $\phi(t)$ in (a, b) with integrable derivatives*

$$\left| \int_a^b e^{i\lambda u(t)} \phi(t) dt \right| \leq 3\lambda^{-1} \left[|\phi(b)| + \int_a^b |\phi'(t)| dt \right].$$

2. *For all $k \geq 2$, if $|u^{(k)}(t)| \geq 1$ for $t \in (a, b)$, then for any $\phi(t)$ in (a, b) with integrable derivatives*

$$\left| \int_a^b e^{i\lambda u(t)} \phi(t) dt \right| \leq 12k\lambda^{-1/k} \left[|\phi(b)| + \int_a^b |\phi'(t)| dt \right].$$

Proof. The assertion can be proved by extending the Van der Corput lemma in [3]. Here we omit the details. \square

Lemma 6.2 *For $0 < \kappa < 1$, let $F(\lambda) = \int_0^\kappa f(t) e^{i\lambda(\sqrt{1-t^2} - \tau\sqrt{\kappa^2-t^2} + \alpha t)} dt$, where $\tau \geq c_0 > 0$ and $\alpha \in \mathbb{R}$, then we have*

$$|F(\lambda)| \leq C(\kappa) \lambda^{-\frac{1}{2N_*}} \left[|f(\kappa)| + \int_0^\kappa |f'(t)| dt \right]$$

where $N_* = \min\{N | \kappa^{2N-1} < c_0, N \in \mathbb{Z}_+\}$.

Proof. Put $\phi(t) = -\sqrt{1-t^2}$ and $\psi(t, \tau) = \tau\phi(t/\kappa) - \phi(t) + \alpha t$. For easy of notations, we denote the n -th partial derivative of $g(t)$ with respect to t by $g^{(n)}(t)$. Then, it is to see that, for $n > 1$

$$\psi^{(n)}(t, \tau) = \frac{\tau}{\kappa^{n-1}} \phi^{(n)}\left(\frac{t}{\kappa}\right) - \phi^{(n)}(t)$$

A standard computation show that

$$\begin{aligned} \phi^{(1)}(t) &= \frac{t}{\sqrt{1-t^2}} \\ \phi^{(2)}(t) &= \frac{1}{(1-t^2)^{3/2}} \\ \phi^{(3)}(t) &= \frac{3t}{(1-t^2)^{5/2}} \end{aligned}$$

Moreover, for $n \geq 3$, we have

$$\phi^{(n)}(t) = \frac{p_n(t)}{(1-t^2)^{n-1/2}} \tag{6.1}$$

where $p_n = \sum_0^{n-2} a_k^n t^k$ is a $(n-2)$ -th polynomial such that its coefficients satisfy the following recursion formula:

$$\begin{aligned} a_{n-1}^{n+1} &= (n+1)a_{n-2}^n, \quad a_{n-2}^{n+1} = (n+2)a_{n-3}^n \\ a_k^{n+1} &= (k+1)a_{k+1}^n + (2n-k)a_{k-1}^n \quad \text{for } 1 \leq k \leq n-3 \\ a_0^{n+1} &= a_1^n \end{aligned}$$

Since the polynomial coefficients are all positive, it is obvious that for $n \geq 1$, $\phi^{(n)}(t)$ is a monotone increasing positive function. Using the recursion formula, it follows that

$$\phi^{(n)}(0) = \begin{cases} 0 & n \text{ is odd,} \\ (n-1)!!(n-3)!! & n \text{ is even.} \end{cases} \quad (6.2)$$

where $(2k-1)!!$ is double factorial and $n > 3$. We are now in the position to proof the inequality. Since $0 < \kappa < 1$, obersev that

$$\psi^{(2N_*+1)}(t, \tau) \geq \frac{\tau}{\kappa^{2N_*}} \phi^{(2N_*+1)}(t) - \phi^{(2N_*+1)}(t) > 0$$

Therefore, $\psi^{(2N_*)}(t, \tau)$ is monotone increasing in $[0, \kappa)$. By (6.2), we get

$$\psi^{(2N_*)}(t, \tau) \geq \psi^{(2N_*)}(0, \tau) \geq \psi^{(2N_*)}(0, c_0) = C(2N_*) \left(\frac{c_0}{\kappa^{2N_*-1}} - 1 \right) > 0 \quad (6.3)$$

The lemma is now a direct consequence of lemma (6.1). \square

7. Other exponential decay term: 17.11.16 on G1

The parameterization of hyperbolic curve passing $(\pm 1, 0)$ is:

$$\xi_1 = \pm \sqrt{t^2 + 1} \quad \xi_2 = t$$

where $t \in \mathbb{R}$. Substituting $\xi = \xi_1 + \mathbf{i}\xi_2$ into $\mu(\xi) := (1 - \xi^2)^{1/2}$ and $\mu_\kappa(\xi) := (\kappa^2 - \xi^2)^{1/2}$, we get

$$\begin{aligned} \text{Im } \mu(\xi) &= \text{Im} (1 - (\xi_1^2 - \xi_2^2 + \mathbf{i}2\xi_1\xi_2))^{1/2} \\ &= \text{Im} (-2t\sqrt{t^2 + 1}\mathbf{i})^{1/2} = t^{1/2}(t^2 + 1)^{1/4} \end{aligned}$$

$$\begin{aligned} \text{Im } \mu_\kappa(\xi) &= \text{Im} (\kappa^2 - (\xi_1^2 - \xi_2^2 + \mathbf{i}2\xi_1\xi_2))^{1/2} \\ &= \text{Im} (\kappa^2 - 1 - 2t\sqrt{t^2 + 1}\mathbf{i})^{1/2} \\ &= \sqrt{\frac{\sqrt{(1 - \kappa^2)^2 + 4t(t^2 + 1)} + 1 - \kappa^2}{2}} \\ &\geq t^{1/2}(t^2 + 1)^{1/4} \end{aligned}$$

where we only consider the branch, denoted by Γ^+ , in the first quadrant here. For $a > 0, b > 0$, we have

$$|e^{\mathbf{i}\xi a + \mathbf{i}\mu(\xi)b + \mathbf{i}\mu_\kappa(\xi)c}| \leq e^{-ta - t^{1/2}(t^2 + 1)^{1/4}b - t^{1/2}(t^2 + 1)^{1/4}c} \leq e^{-t(b+c)}$$

Lemma 7.1 For $\xi \in \Gamma_0$, let $f(\xi)$ is a complex valued function in $L^1(\Gamma^+)$ such that $|f(\xi)| \leq C(1 + \xi^k)$, $k \in \mathbb{Z}_+$. Then we have

$$\begin{aligned} |I(a, b, c) &:= \int_{\Gamma^+} f(\xi) e^{\mathbf{i}\xi a + \mathbf{i}\mu(\xi)b + \mathbf{i}\mu_\kappa(\xi)c} d\xi| \\ &\leq C \left(\frac{1}{b+c} + \frac{1}{(b+c)^k} \right) \end{aligned}$$

Proof.

$$\frac{d\xi(t)}{dt} = \frac{t}{\sqrt{t^2 + 1}} + \mathbf{i}$$

Substituting $\xi(t)$ into $I(a, b, c)$, we have

$$\begin{aligned} |I(a, b, c)| &= \left| \int_0^\infty |f(\xi(t)) \frac{d\xi(t)}{dt} e^{\mathbf{i}\xi(t)a + \mathbf{i}\mu(\xi(t))b + \mathbf{i}\mu_\kappa(\xi(t))c} dt \right| \\ &\leq C \int_0^\infty (1 + t^k) e^{-t(b+c)} dt \\ &\leq C \left(\frac{1}{b+c} + \frac{1}{(b+c)^k} \right) \end{aligned}$$

□

Lemma 7.2 *Let $f(\xi)$ is a bounded complex valued function in $L^1((\kappa, 1))$. Then we have*

$$\begin{aligned} |I(a, b)| &:= \int_\kappa^1 |f(\xi) e^{\mathbf{i}\xi a + \mathbf{i}\mu_\kappa(\xi)b} d\xi| \\ &\leq C \frac{1}{b} \end{aligned}$$

Proof. It is simple to see that

$$\begin{aligned} |I(a, b)| &\leq C \int_\kappa^1 e^{-b\sqrt{\xi^2 - \kappa^2}} d\xi \\ &\leq C \int_0^{\sqrt{1-\kappa^2}} \frac{t}{\sqrt{t^2 + \kappa^2}} e^{-bt} dt \\ &\leq C \frac{1}{b} \end{aligned}$$

□

8. about principle of arguement

Put

$$\begin{aligned} \delta_\pm(t) &= (\kappa - 2t^2)^2 \mp \mathbf{i}4t^2\sqrt{1-t^2}\sqrt{t^2-\kappa} \\ &:= f_1(t) \mp \mathbf{i}f_2(t) \end{aligned}$$

where $0 < \kappa < 1$ and we have

$$\delta'_\pm(t) = f'_1(t) \mp \mathbf{i}f'_2(t)$$

It is easy to see $f_2(1) = f_2(\kappa) = 0$ and $f_1(t) > 0$ for any $\kappa \leq t \leq 1$. Then

$$\begin{aligned} &\int_\kappa^1 \frac{\delta'_+(t)}{\delta_+(t)} - \frac{\delta'_-(t)}{\delta_-(t)} dt \\ &= 2\mathbf{i} \int_\kappa^1 \operatorname{Im} \left(\frac{\delta'_+(t)}{\delta_+(t)} \right) dt \end{aligned}$$

$$\begin{aligned}
&= 2\mathbf{i} \int_{\kappa}^1 \operatorname{Im} \frac{(f_1'(t) - \mathbf{i}f_2'(t))f_1(t) + \mathbf{i}f_2(t)}{(f_1(t) - \mathbf{i}f_2(t))(f_1(t) + \mathbf{i}f_2(t))} dt \\
&= 2\mathbf{i} \int_{\kappa}^1 \frac{f_1'(t)f_2(t) - f_1(t)f_2'(t)}{f_1^2(t) + f_2^2(t)} dt \\
&= 2\mathbf{i} \int_{\kappa}^1 \frac{f_1^2(t)}{f_1^2(t) + f_2^2(t)} \frac{f_1'(t)f_2(t) - f_1(t)f_2'(t)}{f_1^2(t)} dt \\
&= -2\mathbf{i} \int_{\kappa}^1 \frac{f_1^2(t)}{f_1^2(t) + f_2^2(t)} d \frac{f_2(t)}{f_1(t)} \\
&= -2\mathbf{i} \arctan \frac{f_2(t)}{f_1(t)} \Big|_{\kappa}^1 = 0
\end{aligned}$$

Notic that, the condition only used above are $f_2(1) = f_2(\kappa) = 0$ and $f_1(t) > 0$.

9. Fundamental solution of Elastic wave

$$G(x; y) = \frac{1}{\omega^2} (\nabla \times \nabla \cdot (g_s(x; y)\mathbb{I}) - \nabla \nabla g_p(x; y)) \quad (9.1)$$

$$= \frac{1}{\omega^2} (k_s^2 g_s(x, y) + \nabla \nabla (g_s(x; y) - g_p(x; y))) \quad (9.2)$$

where y is the Dirac source, $g_p(x; y)$ or $g_s(x; y)$ is the fundamental solution of the scalar Helmholtz equation with wavenumbers $k_p = \omega/c_p$ or $k_s = \omega/c_s$.

$$g_{\alpha} = \frac{\mathbf{i}}{4} H_0^{(1)}(k_{\alpha}|x - y|) \quad (9.3)$$

where $H_0^{(1)}(t)$ is the Hankel function of the first type and order zero. By straight calculation using $H_1^{(1)}(t) = -dH_0^{(1)}(t)/dt$ and $dH_1^{(1)}(t)/dt = H_0^{(1)}(t) - H_1^{(1)}(t)/t$, we have

$$\begin{aligned}
G_{ij}(x; y) &= \frac{\mathbf{i}}{4} \left\{ \left(\frac{k_s^2}{\omega^2} H_0^{(1)}(k_s|x - y|) - \frac{1}{\omega^2} \frac{k_s H_1^{(1)}(k_s|x - y| - k_p H_1^{(1)}(k_p|x - y|)}{|x - y|} \right) \delta_{ij} \right. \\
&\quad \left. + \frac{1}{\omega^2} \left[\left(\frac{2k_s H_1^{(1)}(k_s|x - y| - 2k_p H_1^{(1)}(k_p|x - y|)}{|x - y|} - (k_s^2 H_0^{(1)}(k_s|x - y|) - k_p^2 H_0^{(1)}(k_p|x - y|)) \right) \frac{(x_i - y_i)(x_j - y_j)}{|x - y|^2} \right] \right\}
\end{aligned}$$

The definition of hankal function is $H_k^{(1)}(t) = J_k(t) + \mathbf{i}Y_k(t)$ where

$$J_k(t) = \sum_{p=0}^{\infty} \frac{(-1)^p}{p!(k+p)!} (t/2)^{k+2p}$$

Specially

$$\begin{aligned}
J_0(t) &= \sum_{p=0}^{\infty} \frac{(-1)^p}{p!p!} (t/2)^{2p} = 1 + \dots \\
J_1(t) &= \sum_{p=0}^{\infty} \frac{(-1)^p}{p!(1+p)!} (t/2)^{1+2p} = \frac{t}{2} + \dots
\end{aligned}$$

and

$$Y_k(t) = \frac{1}{\pi} \{\ln t^2 - 2 \ln 2 + 2C_{euler}\} J_k(t) - \frac{1}{\pi} \sum_{p=0}^{k-1} \frac{(k-1-p)!}{p!} (2/t)^{k-2p} \\ - \frac{1}{\pi} \sum_{p=0}^{\infty} \frac{(-1)^p}{p!(k+p)!} (t/2)^{k+2p} \{\psi(p+k) + \psi(p)\}$$

Speceilly

$$Y_0(t) = \frac{1}{\pi} \{\ln t^2 - 2 \ln 2 + 2C_{euler}\} J_0(t) - \frac{1}{\pi} \sum_{p=1}^{\infty} \frac{(-1)^p}{p!p!} (t/2)^{2p} \{2\psi(p)\} \\ Y_1(t) = \frac{1}{\pi} \{\ln t^2 - 2 \ln 2 + 2C_{euler}\} J_1(t) - \frac{1}{\pi} \frac{2}{t} - \frac{t}{2\pi} \\ - \frac{1}{\pi} \sum_{p=1}^{\infty} \frac{(-1)^p}{p!(1+p)!} (t/2)^{1+2p} \{\psi(p+1) + \psi(p)\}$$

Thus, we have

$$H_0^{(1)}(kr) = 1 + \mathbf{i} \frac{2}{\pi} (C_{euler} + \ln k - \ln 2) + \mathbf{i} \frac{1}{\pi} \ln r^2 + o(kr) \\ H_1^{(1)}(kr) = \frac{kr}{2} + \mathbf{i} \frac{1}{\pi} (C_{euler} + \ln k - \ln 2 - \frac{1}{2}) kr - \mathbf{i} \frac{1}{\pi} \frac{2}{kr} + \mathbf{i} \frac{1}{\pi} \ln r^2 \frac{kr}{2} + o(k^2 r^2) \\ \frac{\mathbf{i}}{4} H_0^{(1)}(kr) = \frac{\mathbf{i}}{4} - \frac{1}{2\pi} (C_{euler} + \ln k - \ln 2) - \frac{1}{4\pi} \ln r^2 + o(kr) \\ \frac{\mathbf{i}}{4} H_1^{(1)}(kr) = \mathbf{i} \frac{kr}{8} - \frac{1}{4\pi} (C_{euler} + \ln k - \ln 2 - \frac{1}{2}) kr + \frac{1}{4\pi} \frac{2}{kr} - \frac{1}{4\pi} \ln r^2 \frac{kr}{2} + o(k^2 r^2)$$

We also need to define the surface traction $T_x^n(\cdot)$ on the normal direction \mathbf{n} ,

$$T_x^n u(x) := \sigma \cdot \mathbf{n} = 2\mu \frac{\partial u}{\partial n} + \lambda n \operatorname{div} u + \mu \mathbf{n} \times \operatorname{curl} u$$

where

$$\sigma(u) = \begin{pmatrix} (\lambda + 2\mu) \partial u_1 / \partial x_1 + \lambda \partial u_2 / \partial x_2 & \mu \partial u_1 / \partial x_2 + \mu \partial u_2 / \partial x_1 \\ \mu \partial u_1 / \partial x_2 + \mu \partial u_2 / \partial x_1 & (\lambda + 2\mu) \partial u_2 / \partial x_2 + \lambda \partial u_1 / \partial x_1 \end{pmatrix}$$

A simple computation show that

$$\frac{\partial^3 H_0^{(1)}(k|x-y|)}{\partial x_i^2 \partial x_j} = (1 + 2\delta_{ij})(-k^2 H_0^{(1)}(kr) \frac{r_j}{r^2} + 2k H_1^{(1)}(kr) \frac{r_j}{r^3}) \\ + k^3 H_1^{(1)}(kr) \frac{r_i^2 r_j}{r^3} + 4k^2 H_0^{(1)}(kr) \frac{r_i^2 r_j}{r^4} - 8k H_1^{(1)}(kr) \frac{r_i^2 r_j}{r^5}$$

where $r = |x - y|$ and $r_i = x_i - y_i$.

$$\frac{\mathbf{i}}{4} H_0^{(1)}(kr) = -\frac{1}{2\pi} (\ln \frac{kr}{2} + C_{euler}) (1 - (\frac{kr}{2})^2 + \dots) + \frac{1}{4\pi} (2(\frac{kr}{2})^2 + \dots) + \frac{\mathbf{i}}{4} (1 - (\frac{kr}{2})^2 + \dots) \\ = -\frac{1}{2\pi} (\ln \frac{kr}{2}) (1 + O(r^2)) - \frac{1}{2\pi} C_{euler} + \frac{\mathbf{i}}{4} + O(r^2) \\ \frac{\mathbf{i}}{4} H_1^{(1)}(kr) = -\frac{1}{2\pi} (\ln \frac{kr}{2} + C_{euler}) (\frac{kr}{2} - \frac{1}{2} (\frac{kr}{2})^3 + \dots) + \frac{1}{4\pi} (\frac{kr}{2} + O(r^3)) + \frac{\mathbf{i}}{4} (\frac{kr}{2} - \frac{1}{2} (\frac{kr}{2})^3 + \dots) + \frac{1}{2\pi} \frac{1}{kr} \\ = -\frac{1}{4\pi} (\ln \frac{kr}{2}) (kr + O(r^3)) - \frac{kr}{4\pi} C_{euler} + \frac{kr}{8\pi} + \frac{\mathbf{i}kr}{8} + \frac{1}{2\pi} \frac{1}{kr} + O(r^3)$$

$$A(kr) := \frac{\mathbf{i}}{4}(k^2 H_0^{(1)}(kr) - 2k H_1^{(1)}(kr)/r) = \frac{k^2}{4\pi}(\ln \frac{kr}{2})(\frac{kr}{2})^2 - \frac{1}{\pi r^2} - \frac{k^2}{4\pi} + O(r^2)$$

and

$$A_{sp}(r) = A(k_s r) - A(k_p r) = \frac{k_s^2}{4\pi}(\ln \frac{kr}{2})(\frac{kr}{2})^2 - \frac{k_p^2}{4\pi}(\ln \frac{kr}{2})(\frac{kr}{2})^2 - (\frac{k_s^2}{4\pi} - \frac{k_p^2}{4\pi}) + O(r^2)$$

Let $g^{jkk} = \frac{\partial^3 g}{\partial x_j \partial x_k^2}$, $d = g_s - g_p$, thus

$$\begin{aligned} g^{ij} &= (1 + 2\delta_{ij})(-A(kr)\frac{r_j}{r^2}) + \frac{\mathbf{i}k^3}{4}H_1^{(1)}(kr)\frac{r_i^2 r_j}{r^3} + 4A(kr)\frac{r_i^2 r_j}{r^4} \\ &= (1 + 2\delta_{ij})(\frac{1}{\pi r^3} + \frac{k^2}{4\pi r})\frac{r_j}{r} - (\frac{4}{\pi r^3} + \frac{k^2}{2\pi r})\frac{r_i^2 r_j}{r^3} + O(r \ln r) \\ d^{ij} &= (1 + 2\delta_{ij})(-A_{sp}\frac{r_j}{r^2}) + (\frac{\mathbf{i}k_s^3}{4}H_1^{(1)}(k_s r) - \frac{\mathbf{i}k_p^3}{4}H_1^{(1)}(k_p r))\frac{r_i^2 r_j}{r^3} + 4A_{sp}\frac{r_i^2 r_j}{r^4} \\ &= (1 + 2\delta_{ij})(\frac{k_s^2}{4\pi r} - \frac{k_p^2}{4\pi r})\frac{r_j}{r} - (\frac{k_s^2}{2\pi r} - \frac{k_p^2}{2\pi r})\frac{r_i^2 r_j}{r^3} + O(r \ln r) \\ &= (1 + 2\delta_{ij})\frac{(\lambda + \mu)\omega^2}{\mu(\lambda + 2\mu)}\frac{1}{4\pi r}\frac{r_j}{r} - \frac{(\lambda + \mu)\omega^2}{\mu(\lambda + 2\mu)}\frac{1}{2\pi r}\frac{r_i^2 r_j}{r^3} + O(r \ln r) \\ k^2 g^i &= -\frac{\mathbf{i}}{4}H_1^{(1)}(kr)\frac{kr_i}{r} = -\frac{k^2}{2\pi r}\frac{r_i}{r} + O(r \ln r) \\ d^{ij} + d^{jjj} &= 2\frac{(\lambda + \mu)\omega^2}{\mu(\lambda + 2\mu)}\frac{1}{4\pi r}\frac{r_j}{r} + O(r \ln r) \end{aligned}$$

Then we have [5, p43]

$$\begin{aligned} \sigma(Ge_1)n &= \frac{1}{\omega^2} \begin{pmatrix} (\lambda + 2\mu)(k_s^2 g_s^1 + d^{111}) + \lambda d^{122} & \mu(k_s^2 g_s^2 + d^{112}) + \mu d^{112} \\ \mu(k_s^2 g_s^2 + d^{112}) + \mu d^{112} & (\lambda + 2\mu)d^{122} + \lambda(k_s^2 g_s^1 + d^{111}) \end{pmatrix} n \\ &= \frac{\mu}{2\pi(\lambda + 2\mu)} \left(\left(\frac{2(\lambda + \mu)r_1^2}{\mu r^2} + 1 \right) \left(-\frac{r_1 n_1}{r^2} - \frac{r_2 n_2}{r^2} \right) - \begin{pmatrix} 0 \\ \frac{r_2 n_1 - r_1 n_2}{r^2} \end{pmatrix} \right) + O(r \ln r) \end{aligned}$$

and

$$\begin{aligned} \sigma(Ge_2)n &= \frac{1}{\omega^2} \begin{pmatrix} (\lambda + 2\mu)d^{112} + \lambda(k_s^2 g_s^2 + d^{222}) & \mu(k_s^2 g_s^1 + d^{122}) + \mu d^{122} \\ \mu(k_s^2 g_s^1 + d^{122}) + \mu d^{122} & (\lambda + 2\mu)(k_s^2 g_s^2 + d^{222}) + \lambda d^{112} \end{pmatrix} n \\ &= \frac{\mu}{2\pi(\lambda + 2\mu)} \left(\left(\frac{2(\lambda + \mu)r_1 r_2}{\mu r^2} \right) \left(-\frac{r_1 n_1}{r^2} - \frac{r_2 n_2}{r^2} \right) - \begin{pmatrix} \frac{r_1 n_2 - r_2 n_1}{r^2} \\ 0 \end{pmatrix} \right) + O(r \ln r) \end{aligned}$$

Now Let u be represented as single potential:

$$u = \int_{\partial D} G(x, y) \phi(y) ds(y) \quad (9.4)$$

with Neumann boundary condition

$$T_x u(x) = f(x) \quad \text{on } \partial D \quad (9.5)$$

Then we obtain corresponding integral equation

$$\mathbf{P.V.} \int_{\partial D} T_x G(x, y) \phi(y) ds(y) - \frac{1}{2} \phi(x) = f(x) \quad (9.6)$$

where $x \in \partial D$. We describe the necessary parametrization of the integral equation in the two-dimensional case. We assume that the boundary curve ∂D possesses a regular analytic and 2π -periodic parametric representation of the form

$$x(t) = (x_1(t), x_2(t))$$

in counterclockwise orientation satisfying $|x'(t)| > 0$ for all t . Let [4]

$$\begin{aligned} T(x, y) &= (T_x(N_1(x, y))n, T_x(N_1(x, y))n) \\ T_0(x, y) &= -\frac{\mu}{2\pi(\lambda + 2\mu)} \begin{pmatrix} 0 & \frac{r_1 n_2 - r_2 n_1}{r^2} \\ \frac{r_2 n_1 - r_1 n_2}{r^2} & 0 \end{pmatrix} \\ T_1(x, y) &= T(x, y) - T_0(x, y) \end{aligned}$$

Then by above analysis we have

$$\int_0^{2\pi} T_1(x(t), x(\tau))\phi(x(\tau))|x'(\tau)|dt + \mathbf{P.V.} \int_0^{2\pi} T_0(x(t), x(\tau))\phi(x(\tau))|x'(\tau)|dt - \frac{1}{2}\phi(x(t)) = f(x(t))$$

In particular, using expansion above, we can deduce the diagonal terms:

$$\lim_{\tau \rightarrow t} \frac{-r_1 n_1 - r_2 n_2}{r^2} = \lim_{\tau \rightarrow t} \frac{(x_1(\tau) - x_1(t))x_2'(t) - (x_2(\tau) - x_2(t))x_1'(t)}{|x(t) - x(\tau)|^2|x'(t)|} = \frac{x_1''(t)x_2'(t) - x_2''(t)x_1'(t)}{2|x'(t)|^3}$$

and

$$\begin{aligned} &\lim_{\tau \rightarrow t} \frac{(r_1 n_2 - r_2 n_1)|x'(\tau)|(\tau - t)}{r^2} \\ &= \lim_{\tau \rightarrow t} \frac{(x_1(\tau) - x_1(t))x_1'(t) + (x_2(\tau) - x_2(t))x_2'(t))|x'(\tau)|(\tau - t)}{|x(t) - x(\tau)|^2|x'(t)|} = 1 \end{aligned}$$

According to above analysis, it is easy to see that T_1 has no singularity. Therefore, the numerical formulation for T_1 requires only straightforward application of simple quadrature formula. we choose an equidistant set of knots $t_j := \pi j/n, j = 0, \dots, 2n-1$ and divide $[0, 2\pi)$ into n equivalent interval $I_i = [2j\pi/n, 2(j+1)\pi/n]$ where $[0, 2\pi) = \bigcup_{j=1}^{j=n-1} I_i$. Each I_i is described by 3 nodes of the intrinsic variable $\xi (-1 \leq \xi \leq 1)$ and the

quadratic shape functions are:

$$\begin{aligned} A_{-1}(\xi) &= \frac{\xi(\xi - 1)}{2} \\ A_0(\xi) &= 1 - \xi^2 \\ A_1(\xi) &= \frac{\xi(\xi + 1)}{2} \end{aligned}$$

let

$$f(x, y) = \frac{(x_1 - y_1)n_2^x - (x_2 - y_2)n_1^x}{|x - y|^2}$$

and it is easy to see that $(n_x^1, n_x^2) = (x_2'(t), -x_1'(t))/|x'(t)|$. Now, we can represent the integral in variable $\xi \in [-1, 1]$. For $x \in \partial D$ we have:

$$\begin{aligned} &\mathbf{P.V.} \int_{\partial D} f(x, y)\phi(y)ds(y) \\ &= \sum_{x \notin I_i} \int_{I_i} f(x, x(t))\phi(x(t))|x'(t)|dt + \mathbf{P.V.} \sum_{x \in I_i} \int_{I_i} f(x, x(t))\phi(x(t))|x'(t)|dt \end{aligned}$$

if $x = x(2j\pi/n)$, above integral becomes to

$$\begin{aligned}
& \sum_{i \neq j, j-1} \frac{\pi}{n} \int_{-1}^1 \frac{\sum_{k=1,2} (x_k^i(\xi) - x_k \frac{2j\pi}{n})) x'_k(\frac{2j\pi}{n})}{|x(\frac{2j\pi}{n}) - x^i(\xi)|^2} \phi(x^i(\xi)) \frac{|x'(\frac{(2i+1+\xi)\pi}{n})|}{|x'(\frac{2j\pi}{n})|} d\xi \\
& + \text{P.V.} \sum_{i=j, j-1} \frac{\pi}{n} \int_{-1}^1 \frac{\sum_{k=1,2} (x_k^i(\xi) - x_k \frac{2j\pi}{n})) x'_k(\frac{2j\pi}{n})}{|x(\frac{2j\pi}{n}) - x^i(\xi)|^2} \phi(x^i(\xi)) \frac{|x'(\frac{(2i+1+\xi)\pi}{n})|}{|x'(\frac{2j\pi}{n})|} d\xi \\
& \approx \sum_{i \neq j, j-1} \sum_{l=-1}^1 \phi(x^i(l)) \frac{\pi}{n} \int_{-1}^1 \frac{\sum_{k=1,2} (x_k^i(\xi) - x_k \frac{2j\pi}{n})) x'_k(\frac{2j\pi}{n})}{|x(\frac{2j\pi}{n}) - x^i(\xi)|^2} A_l(\xi) \frac{|x'(\frac{(2i+1+\xi)\pi}{n})|}{|x'(\frac{2j\pi}{n})|} d\xi \\
& + \phi(x(\frac{2j\pi}{n})) \text{P.V.} \frac{\pi}{n} \left(\int_{-1}^1 \frac{\sum_{k=1,2} (x_k^j(\xi) - x_k^j(-1)) x'_k(\frac{2j\pi}{n})}{|x^j(-1) - x^j(\xi)|^2} A_{-1}(\xi) \frac{|x'(\frac{(2j+1+\xi)\pi}{n})|}{|x'(\frac{2j\pi}{n})|} d\xi \right. \\
& \left. + \int_{-1}^1 \frac{\sum_{k=1,2} (x_k^{j-1}(\xi) - x_k^{j-1}(1)) x'_k(\frac{2j\pi}{n})}{|x^{j-1}(1) - x^{j-1}(\xi)|^2} A_1(\xi) \frac{|x'(\frac{(2j-1+\xi)\pi}{n})|}{|x'(\frac{2j\pi}{n})|} d\xi \right)
\end{aligned}$$

and if $x = x((2j+1)\pi/n)$, above integral becomes to

$$\begin{aligned}
& \sum_{i \neq j} \frac{\pi}{n} \int_{-1}^1 \frac{\sum_{k=1,2} (x_k^i(\xi) - x_k \frac{2j+1\pi}{n})) x'_k(\frac{2j+1\pi}{n})}{|x(\frac{2j+1\pi}{n}) - x^i(\xi)|^2} \phi(x^i(\xi)) \frac{|x'(\frac{(2i+1+\xi)\pi}{n})|}{|x'(\frac{2j+1\pi}{n})|} d\xi \\
& + \text{P.V.} \sum_{i=j, j-1} \frac{\pi}{n} \int_{-1}^1 \frac{\sum_{k=1,2} (x_k^i(\xi) - x_k \frac{2j+1\pi}{n})) x'_k(\frac{2j+1\pi}{n})}{|x(\frac{2j+1\pi}{n}) - x^i(\xi)|^2} \phi(x^i(\xi)) \frac{|x'(\frac{(2i+1+\xi)\pi}{n})|}{|x'(\frac{2j+1\pi}{n})|} d\xi \\
& \approx \sum_{i \neq j} \sum_{l=-1}^1 \phi(x^i(l)) \frac{\pi}{n} \int_{-1}^1 \frac{\sum_{k=1,2} (x_k^i(\xi) - x_k \frac{(2j+1)\pi}{n})) x'_k(\frac{(2j+1)\pi}{n})}{|x(\frac{(2j+1)\pi}{n}) - x^i(\xi)|^2} A_l(\xi) \frac{|x'(\frac{(2i+1+\xi)\pi}{n})|}{|x'(\frac{(2j+1)\pi}{n})|} d\xi \\
& + \phi(x(\frac{(2j+1)\pi}{n})) \text{P.V.} \frac{\pi}{n} \left(\int_{-1}^1 \frac{\sum_{k=1,2} (x_k^j(\xi) - x_k^j(0)) x'_k(\frac{(2j+1)\pi}{n})}{|x^j(0) - x^j(\xi)|^2} A_0(\xi) \frac{|x'(\frac{((2j+1)+\xi)\pi}{n})|}{|x'(\frac{(2j+1)\pi}{n})|} d\xi \right)
\end{aligned}$$

where $x^i(\xi) = x((2i+1+\xi)\pi/n)$. For simplicity, we denote a 3-D tensor M_{jil} :

$$M_{jil} = \frac{\pi}{n} \int_{-1}^1 \frac{\sum_{k=1,2} (x_k^i(\xi) - x_k \frac{j\pi}{n})) x'_k(\frac{j\pi}{n})}{|x(\frac{j\pi}{n}) - x^i(\xi)|^2} A_l(\xi) \frac{|x'(\frac{(2i+1+\xi)\pi}{n})|}{|x'(\frac{j\pi}{n})|} d\xi$$

Then, the general Matrix is

$$G_{ji} = \begin{cases} M_{j,i,-1} + M_{j,i-1,1} & j \text{ is even} \\ M_{j,i,0} & j \text{ is odd} \end{cases}$$

10. Traction Tenson of Neumann Green Function

$$\hat{\Phi}(\xi, x_2; y_2) = \frac{\mathbf{i}}{2\omega^2} \left[\begin{pmatrix} \mu_s & -\xi \frac{x_2-y_2}{|x_2-y_2|} \\ -\xi \frac{x_2-y_2}{|x_2-y_2|} & \frac{\xi^2}{\mu_s} \end{pmatrix} e^{\mathbf{i}\mu_s|x_2-y_2|} + \begin{pmatrix} \frac{\xi^2}{\mu_p} & \xi \frac{x_2-y_2}{|x_2-y_2|} \\ \xi \frac{x_2-y_2}{|x_2-y_2|} & \mu_p \end{pmatrix} e^{\mathbf{i}\mu_p|x_2-y_2|} \right]$$

$$\sigma(\Phi)(\xi, x_2; y_2) e_1 = \frac{\mu}{2\omega^2} \left[\begin{pmatrix} -2\xi\mu_s & \frac{x_2-y_2}{|x_2-y_2|} 2\xi^2 \\ -\frac{x_2-y_2}{|x_2-y_2|} \beta & \frac{\xi\beta}{\mu_s} \end{pmatrix} e^{\mathbf{i}\mu_s|x_2-y_2|} + \begin{pmatrix} -\frac{\alpha\xi}{\mu_p} & -\frac{x_2-y_2}{|x_2-y_2|} \alpha \\ -\frac{x_2-y_2}{|x_2-y_2|} 2\xi^2 & -2\xi\mu_p \end{pmatrix} e^{\mathbf{i}\mu_p|x_2-y_2|} \right]$$

$$\sigma(\Phi)(\xi, x_2; y_2)e_2 = \frac{\mu}{2\omega^2} \left[\begin{pmatrix} -\frac{x_2-y_2}{|x_2-y_2|}\beta & \frac{\xi\beta}{2\xi\mu_s} \\ 2\xi\mu_s & -\frac{x_2-y_2}{|x_2-y_2|}2\xi^2 \end{pmatrix} e^{\mathbf{i}\mu_s|x_2-y_2|} + \begin{pmatrix} -\frac{x_2-y_2}{|x_2-y_2|}2\xi^2 & -2\xi\mu_p \\ -\frac{\beta\xi}{\mu_p} & -\frac{x_2-y_2}{|x_2-y_2|}\beta \end{pmatrix} e^{\mathbf{i}\mu_p|x_2-y_2|} \right]$$

where $\alpha(\xi) = k_s^2 - 2\mu_p^2$, $\beta(\xi) = k_s^2 - 2\xi^2$.

$$\hat{N}(\xi, x_2; y_2) = \hat{\Phi}(\xi, x_2; y_2) - \hat{\Phi}(\xi, x_2; -y_2) + \hat{N}_c(\xi, x_2; y_2) \quad (10.1)$$

$$\begin{aligned} \hat{N}_c(\xi, x_2; y_2) = \frac{\mathbf{i}}{\omega^2\delta(\xi)} & \left\{ A(\xi)e^{\mathbf{i}\mu_s(x_2+y_2)} + B(\xi)e^{\mathbf{i}\mu_p(x_2+y_2)} \right. \\ & \left. + C(\xi)e^{\mathbf{i}\mu_s x_2 + \mu_p y_2} + D(\xi)e^{\mathbf{i}\mu_p x_2 + \mu_s y_2} \right\} \end{aligned} \quad (10.2)$$

where

$$\begin{aligned} A(\xi) &= \begin{pmatrix} \mu_s\beta^2 & -4\xi^3\mu_s\mu_p \\ -\xi\beta^2 & 4\xi_4\mu_p \end{pmatrix} & B(\xi) &= \begin{pmatrix} 4\xi^4\mu_s & \xi\beta^2 \\ 4\xi^3\mu_s\mu_p & \mu_p\beta^2 \end{pmatrix} \\ C(\xi) &= \begin{pmatrix} 2\xi^2\mu_s\beta & -2\xi\mu_s\mu_p\beta \\ -2\xi^3\beta & 2\xi^2\mu_p\beta \end{pmatrix} & D(\xi) &= \begin{pmatrix} 2\xi^2\mu_s\beta & 2\xi^3\beta \\ 2\xi\mu_s\mu_p\beta & 2\xi^2\mu_p\beta \end{pmatrix} \end{aligned}$$

The traction of integral part N_c :

$$\begin{aligned} \sigma(N_c(\xi, x_2; y_2))e_1 &= \frac{-\mu}{\omega^2\delta(\xi)} \left\{ A(\xi)e^{\mathbf{i}\mu_s(x_2+y_2)} + B(\xi)e^{\mathbf{i}\mu_p(x_2+y_2)} \right. \\ & \left. + C(\xi)e^{\mathbf{i}\mu_s x_2 + \mu_p y_2} + D(\xi)e^{\mathbf{i}\mu_p x_2 + \mu_s y_2} \right\} \end{aligned}$$

where

$$\begin{aligned} A(\xi) &= \begin{pmatrix} 2\xi\mu_s\beta^2 & -8\mu_s\mu_p\xi^4 \\ \beta^3 & -4\xi^3\mu_p\beta \end{pmatrix} & B(\xi) &= \begin{pmatrix} 4\xi^3\mu_s\alpha & \alpha\beta^2 \\ 8\xi^4\mu_s\mu_p & 2\xi\mu_p\beta^2 \end{pmatrix} \\ C(\xi) &= \begin{pmatrix} 4\xi^3\mu_s\beta & -4\xi^2\mu_s\mu_p\beta \\ 2\xi^2\beta^2 & -2\xi\mu_p\beta^2 \end{pmatrix} & D(\xi) &= \begin{pmatrix} 2\xi\mu_s\alpha\beta & 2\xi^2\alpha\beta \\ 4\xi^2\mu_s\mu_p\beta & 4\xi^3\mu_p\beta \end{pmatrix} \end{aligned}$$

In particular, for $y_2 = 0$, a more simpler form are deduced:

$$\sigma(N(\xi, x_2; y_2))e_1 = \frac{-1}{\delta(\xi)} \left\{ \begin{pmatrix} 2\xi\mu_s\beta & -4\mu_s\mu_p\xi^2 \\ \beta^2 & -2\xi\mu_p\beta \end{pmatrix} e^{\mathbf{i}\mu_s x_2} + \begin{pmatrix} 2\xi\mu_s\alpha & \alpha\beta \\ 4\xi^2\mu_s\mu_p & 2\xi\mu_p\beta \end{pmatrix} e^{\mathbf{i}\mu_p x_2} \right\}$$

and

$$\begin{aligned} \sigma(N_c(\xi, x_2; y_2))e_2 &= \frac{-\mu}{\omega^2\delta(\xi)} \left\{ A(\xi)e^{\mathbf{i}\mu_s(x_2+y_2)} + B(\xi)e^{\mathbf{i}\mu_p(x_2+y_2)} \right. \\ & \left. + C(\xi)e^{\mathbf{i}\mu_s x_2 + \mu_p y_2} + D(\xi)e^{\mathbf{i}\mu_p x_2 + \mu_s y_2} \right\} \end{aligned}$$

where

$$\begin{aligned} A(\xi) &= \begin{pmatrix} \beta^3 & -4\xi^3\mu_p\beta \\ -2\xi\mu_s\beta^2 & 8\mu_s\mu_p\xi^4 \end{pmatrix} & B(\xi) &= \begin{pmatrix} 8\xi^4\mu_s\mu_p & 2\xi\mu_p\beta^2 \\ 4\xi^3\mu_s\beta & \beta^3 \end{pmatrix} \\ C(\xi) &= \begin{pmatrix} 2\xi^2\beta^2 & -2\xi\mu_p\beta^2 \\ -4\xi^3\mu_s\beta & 4\xi^2\mu_s\mu_p\beta \end{pmatrix} & D(\xi) &= \begin{pmatrix} 4\xi^2\mu_s\mu_p\beta & 4\xi^3\mu_p\beta \\ 2\xi\mu_s\beta^2 & 2\xi^2\beta^2 \end{pmatrix} \end{aligned}$$

Similarly, for $y_2 = 0$, a more simpler form are deduced:

$$\sigma(N(\xi, x_2; y_2))e_2 = \frac{-1}{\delta(\xi)} \left\{ \begin{pmatrix} \beta^2 & -2\xi\mu_p\beta \\ -2\xi\mu_s\beta & 4\mu_s\mu_p\xi^2 \end{pmatrix} e^{\mathbf{i}\mu_s x_2} + \begin{pmatrix} 4\xi^2\mu_s\mu_p & 2\xi\mu_p\beta \\ 2\xi\mu_s\beta & \beta^2 \end{pmatrix} e^{\mathbf{i}\mu_p x_2} \right\}$$

where $\delta(\xi) = \beta^2 + 4\xi^2\mu_s\mu_p$

11. reflection of Plane wave

11.1. P-wave

We denote incident P-wave [1, p172] as

$$u^0 = A_0(\sin t_0, \cos t_0)^T e^{\mathbf{i}k_p(x_1 \sin t_0 + x_2 \cos t_0)} \quad (11.1)$$

and its stress as

$$\sigma(u^0) = \mathbf{i}k_p A_0(2\mu \sin t_0 \cos t_0, \lambda + 2\mu \cos^2 t_0)^T e^{\mathbf{i}k_p(x_1 \sin t_0 + x_2 \cos t_0)}$$

The reflected P-wave is represented as

$$\begin{aligned} u^1 &= A_1(\sin t_1, -\cos t_1)^T e^{\mathbf{i}k_p(x_1 \sin t_1 - x_2 \cos t_1)} \\ \sigma(u^1) &= \mathbf{i}k_p A_1(-2\mu \sin t_1 \cos t_1, \lambda + 2\mu \cos^2 t_1)^T e^{\mathbf{i}k_p(x_1 \sin t_1 + x_2 \cos t_1)} \end{aligned}$$

and reflected S-wave as

$$\begin{aligned} u^2 &= A_2(\cos t_2, \sin t_2)^T e^{\mathbf{i}k_s(x_1 \sin t_2 - x_2 \cos t_2)} \\ \sigma(u^2) &= \mathbf{i}k_s A_2(\mu(\sin^2 t_2 - \cos^2 t_2), -2\mu \sin t_2 \cos t_2)^T e^{\mathbf{i}k_s(x_1 \sin t_2 - x_2 \cos t_2)} \end{aligned}$$

We consider the clamped condition, then the total field on the $x_2 = 0$ vanish:

$$u^0(x_1, 0) + u^1(x_1, 0) + u^2(x_1, 0) = 0$$

for any $x_1 \in \mathbb{R}$. A simple computation show that

$$\begin{aligned} t_1 = t_0 \quad \text{and} \quad \frac{\sin t_2}{\sin t_0} &= \frac{k_p}{k_s} := \kappa \\ A_0 = \cos(t_0 - t_2) \quad A_1 &= \cos(t_0 + t_2) \quad A_2 = -\sin 2t_0 \end{aligned}$$

11.2. S-wave

Similarly, we denote incident S-wave as

$$u^0 = A_0(-\cos t_0, \sin t_0)^T e^{\mathbf{i}k_p(x_1 \sin t_0 + x_2 \cos t_0)} \quad (11.2)$$

$$\sigma(u^0) = \mathbf{i}k_s(\mu(\sin^2 t_0 - \cos^2 t_0), 2\mu \sin t_0 \cos t_0)^T e^{\mathbf{i}k_p(x_1 \sin t_0 + x_2 \cos t_0)} \quad (11.3)$$

The reflected P-wave is represented as

$$\begin{aligned} u^1 &= A_1(\sin t_1, -\cos t_1)^T e^{\mathbf{i}k_p(x_1 \sin t_1 - x_2 \cos t_1)} \\ \sigma(u^1) &= \mathbf{i}k_p A_1(-2\mu \sin t_1 \cos t_1, \lambda + 2\mu \cos^2 t_1)^T e^{\mathbf{i}k_p(x_1 \sin t_1 + x_2 \cos t_1)} \end{aligned}$$

and reflected S-wave as

$$\begin{aligned} u^2 &= A_2(\cos t_2, \sin t_2)^T e^{\mathbf{i}k_s(x_1 \sin t_2 - x_2 \cos t_2)} \\ \sigma(u^2) &= \mathbf{i}k_s A_2(\mu(\sin^2 t_2 - \cos^2 t_2), -2\mu \sin t_2 \cos t_2)^T e^{\mathbf{i}k_s(x_1 \sin t_2 - x_2 \cos t_2)} \end{aligned}$$

The result is

$$\begin{aligned} t_2 = t_0 \quad \text{and} \quad \frac{\sin t_1}{\sin t_0} &= \frac{k_s}{k_p} = \frac{1}{\kappa} \\ A_0 &= \cos(t_0 - t_1) \quad A_1 = \sin 2t_0 \quad A_2 = \cos(t_0 + t_1) \end{aligned}$$

12. scattering relation of elastic wave

The solution for the scattering of a plane P-wave u_p (or S-wave u_s) with incident direction d_0 at a plane $\Gamma := x \in \mathbb{R}^2 : x \cdot \nu = 0$ through the origin with normal vector ν is described by

$$u = u_p + u_{p,p} + u_{p,s} = A_0 d_0 e^{ik_{px} \cdot d} + A_1 d_1 e^{ik_{px} \cdot d_1} + A_2 d_2^\perp d_0^{iksx \cdot d_2} \quad (12.1)$$

$$u = u_s + u_{s,p} + u_{s,s} = A_0 d_0^\perp e^{iksx \cdot d} + A_1 d_1 e^{ik_{px} \cdot d_1} + A_2 d_2^\perp d^{iksx \cdot d_2} \quad (12.2)$$

where $d_i = (d_i^1, d_i^2)^T$ are unit vectors, $d_i^\perp = (d_i^2, -d_i^1)^T$ and A_i are corresponding amplitude. For fixed boundary, we have $u = 0$ for $x \in \Gamma$. After a standard computation, we get for P-wave:

$$d_1 = d_0 - 2\alpha\nu \quad (12.3)$$

$$d_2 = \kappa d_0 - \beta\nu \quad (12.4)$$

$$A_0 = \kappa(d, \nu)^2 - \kappa(d, \nu^\perp)^2 - \beta(d, \nu) \quad (12.5)$$

$$A_1 = \kappa - \beta(d, \nu) \quad (12.6)$$

$$A_2 = -2(d, \nu)(d, \nu^\perp) \quad (12.7)$$

where $\alpha = (d, \nu)$, $\beta = \kappa\alpha - \sqrt{\kappa^2\alpha^2 - \kappa^2 + 1}$ and $\kappa = k_p/k_s$. For S-wave:

$$d_1 = \kappa_1 d_0 - \gamma\nu \quad (12.8)$$

$$d_2 = d_0 - 2\alpha\nu \quad (12.9)$$

$$A_0 = \kappa_1(d, \nu)^2 - \kappa_1(d, \nu^\perp)^2 - \gamma(d, \nu) \quad (12.10)$$

$$A_1 = 2(d, \nu)(d, \nu^\perp) \quad (12.11)$$

$$A_2 = \kappa_1 - \gamma(d, \nu) \quad (12.12)$$

where $\gamma = \kappa_1\alpha - \sqrt{\kappa_1^2\alpha^2 - \kappa_1^2 + 1}$ and $\kappa_1 = 1/\kappa$. Thus the traction of $u(x)$ on the plane Γ can be obtained. For P-wave

$$\begin{aligned} \sigma(u) \cdot \nu &= [\mathbf{i}k_p A_0(\lambda\nu + 2\mu(d_0, \nu)d_0) + \mathbf{i}k_p A_1(\lambda\nu + 2\mu(d_1, \nu)d_1) \\ &+ \mathbf{i}k_s A_2\mu((d_2, \nu)d_2^\perp + (d_2^\perp, \nu)d_2)]e^{ik_{px} \cdot d} := \mathbf{i}k_p Rf_p(x, d, \nu)e^{ik_{px} \cdot d} \end{aligned}$$

For S-wave

$$\begin{aligned} \sigma(u) \cdot \nu &= [\mathbf{i}k_s A_0\mu((d_0, \nu)d_0^\perp + (d_0^\perp, \nu)d_0) + \mathbf{i}k_p A_1(\lambda\nu + 2\mu(d_1, \nu)d_1) \\ &+ \mathbf{i}k_s A_2\mu((d_2, \nu)d_2^\perp + (d_2^\perp, \nu)d_2)]e^{iksx \cdot d} := \mathbf{i}k_s Rf_s(x, d, \nu)e^{iksx \cdot d} \end{aligned}$$

Definition 12.1 For any unit vector $d \in \mathbb{R}^2$, let $u_p^i = d e^{ik_{px} \cdot d}$ or $u_s^i = d^\perp e^{iksx \cdot d}$ be the incident wave and $u_\alpha^s = u_\alpha^s(x; d)$ be the radiation solution of the Navier equation:

$$u_\alpha^s + \omega^2 u_\alpha^s = 0 \quad \text{in } \mathbb{R}^2 \setminus \bar{D} \quad (12.13)$$

$$u_\alpha^s = -u_\alpha^i \quad \text{on } \partial D \quad (12.14)$$

The scattering coecient $R(x;d)$ for $x \in \partial D$ is defined by the relation

$$\sigma(u_\alpha^s + u_\alpha^i) \cdot \nu = \mathbf{i}k_\alpha R_\alpha(x;d)e^{\mathbf{i}k_\alpha x \cdot d} \quad \text{on } \partial D$$

where $\alpha = p, s$.

In the case of high frequency approximation, the scattering coecient can be approximated by

$$R_\alpha(x;d) = \begin{cases} Rf_\alpha(x;d,\nu) & \text{if } x \in \partial D_d^- = \{x \in \partial D, \nu(x) \cdot d < 0\}, \\ 0 & \text{if } x \in \partial D_d^+ = \{x \in \partial D, \nu(x) \cdot d \geq 0\}. \end{cases}$$

13. Difference of solution of navier equation in full-space and half-space,0105

For any $0 < \varepsilon < 1$, we consider the problem

$$\Delta_\varepsilon u_1^\varepsilon + (1 + \mathbf{i}\varepsilon)\omega^2 u_1^\varepsilon = 0 \quad \text{in } \mathbb{R}_+^2 \setminus \bar{D} \quad (13.1)$$

$$u_1^\varepsilon = g \quad \text{on } \Gamma_D \quad (13.2)$$

$$\sigma(u_1^\varepsilon)e_2 = 0 \quad \text{on } \Gamma_0 \quad (13.3)$$

and

$$\Delta_\varepsilon u_2^\varepsilon + (1 + \mathbf{i}\varepsilon)\omega^2 u_2^\varepsilon = 0 \quad \text{in } \mathbb{R}^2 \setminus \bar{D} \quad (13.4)$$

$$u_2^\varepsilon = g \quad \text{on } \Gamma_D \quad (13.5)$$

Let $w^\varepsilon(x)$ be the solution of the problem:

$$\Delta_\varepsilon w^\varepsilon + (1 + \mathbf{i}\varepsilon)\omega^2 w^\varepsilon = 0 \quad \text{in } \mathbb{R}_+^2 \quad (13.6)$$

$$\sigma(w^\varepsilon)e_2 = -\sigma(u_2^\varepsilon)e_2 \quad \text{on } \Gamma_0 \quad (13.7)$$

Then $u_1^\varepsilon - u_2^\varepsilon - w^\varepsilon$ satisfies (13.1),(13.3) with the boundary condition $u_1^\varepsilon - u_2^\varepsilon - w^\varepsilon = -w^\varepsilon$ on Γ_D . Thus by the limiting absorption principle, we have

$$\|T_x^\nu(u_1^\varepsilon - u_2^\varepsilon)\|_{H^{-1/2}(\Gamma_D)} \leq C(\|w^\varepsilon\|_{H^{1/2}(\Gamma_D)} + |T_x^\nu(w^\varepsilon)|_{H^{-1/2}(\Gamma_D)}) \quad (13.8)$$

$$\leq C \max_{x \in D} (|w^\varepsilon(x)| + d_D |\nabla w^\varepsilon(x)|) \quad (13.9)$$

where C is independent of ε, ω . By the integral representation formula we have for any $z \in \Gamma_0$

$$u_2^\varepsilon(z) = \int_{\Gamma_D} (T_y^\nu \Phi^\varepsilon(y, z))^T u_2^\varepsilon(y) - \Phi^\varepsilon(z, y) (T_y^\nu u_2^\varepsilon(y)) ds(y) \quad (13.10)$$

which yields by using the integral representation again that for $x \in D$

$$w^\varepsilon(x) = \int_{\Gamma_0} N^\varepsilon(x, z) (T_z^{e_2} u_2^\varepsilon(z)) ds(z) \quad (13.11)$$

$$= \int_{\Gamma_D} ds(y) \int_{\Gamma_0} N^\varepsilon(x, z) (T_z^{e_2} ((T_y^\nu \Phi^\varepsilon(y, z))^T)) ds(z) \quad (13.12)$$

$$- \int_{\Gamma_D} v^\varepsilon(x, y) (T_y^\nu u_2^\varepsilon(y)) ds(y) \quad (13.13)$$

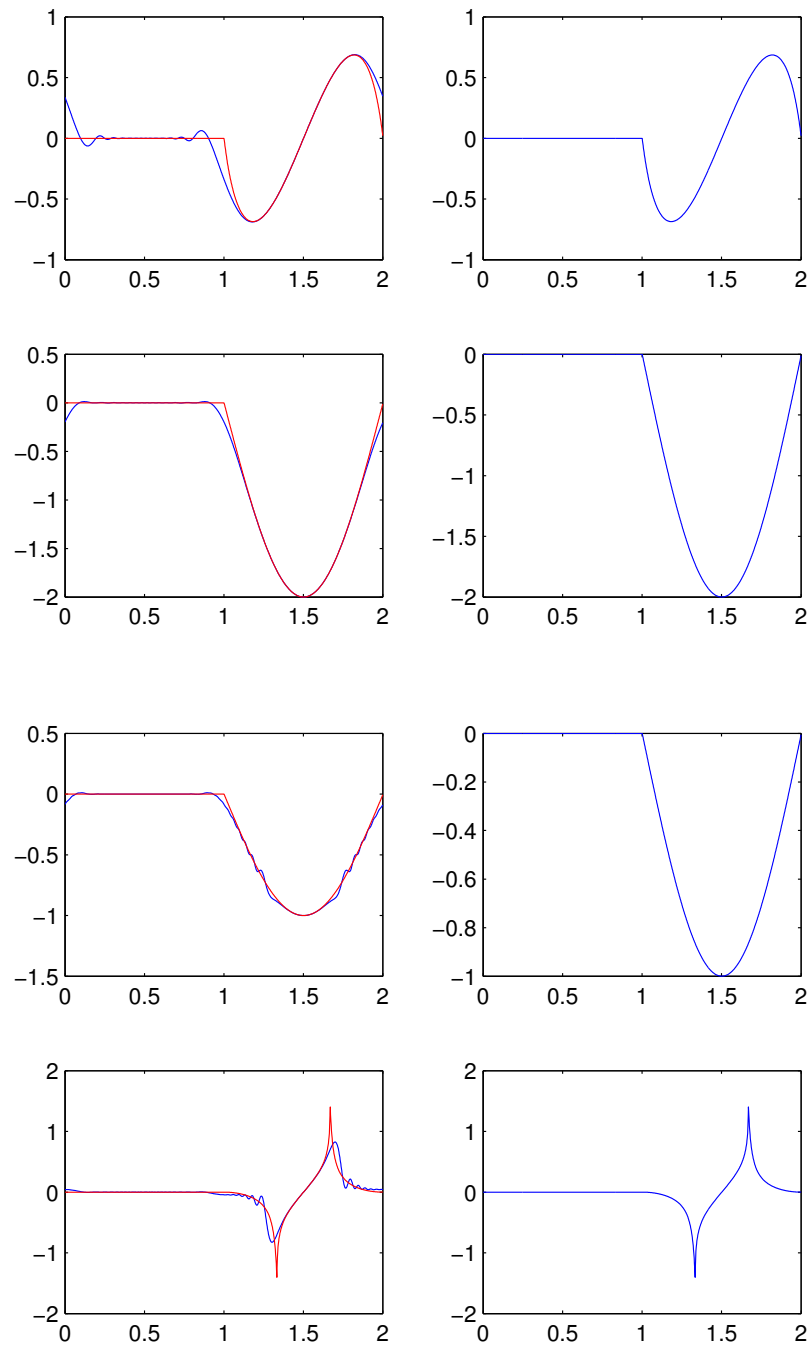


Figure 1. $\theta = 0\pi$

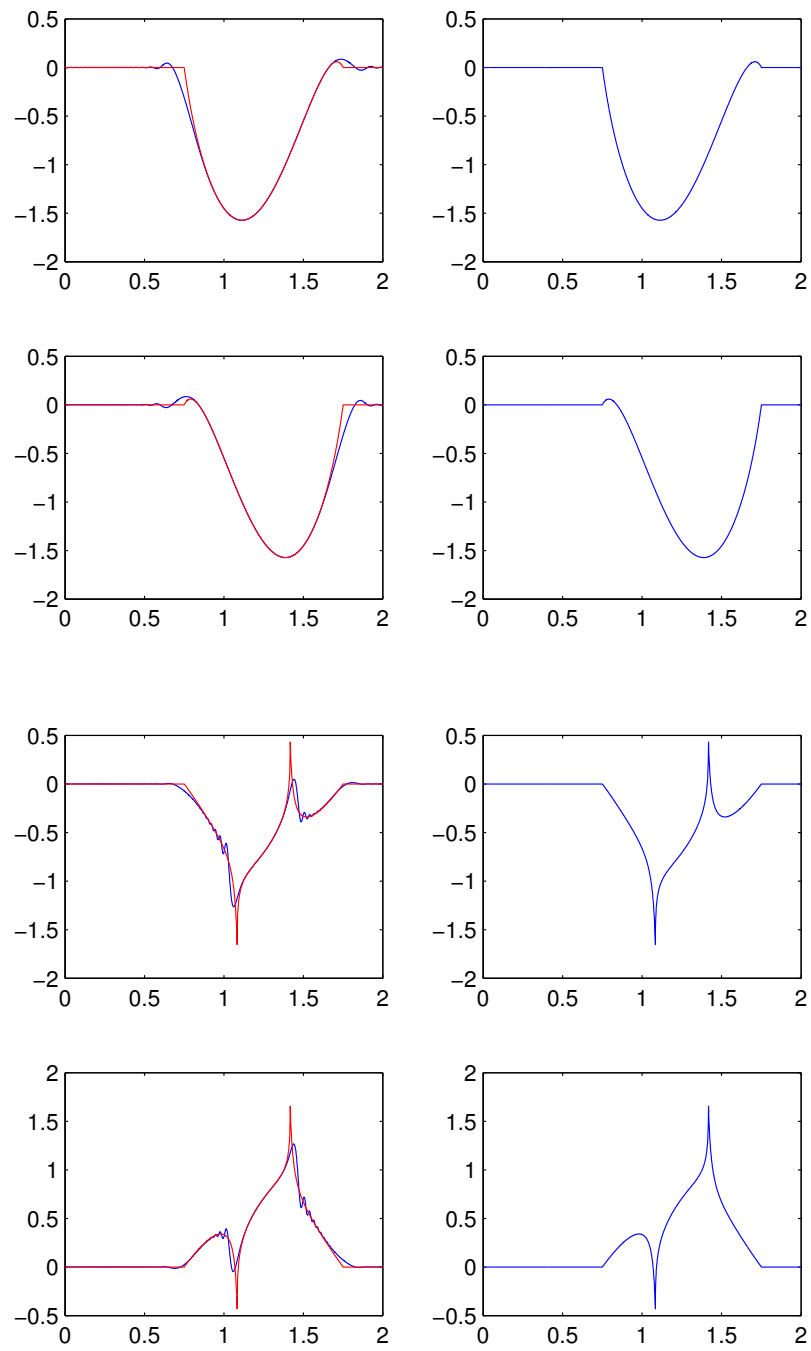


Figure 2. $\theta = \pi/4$

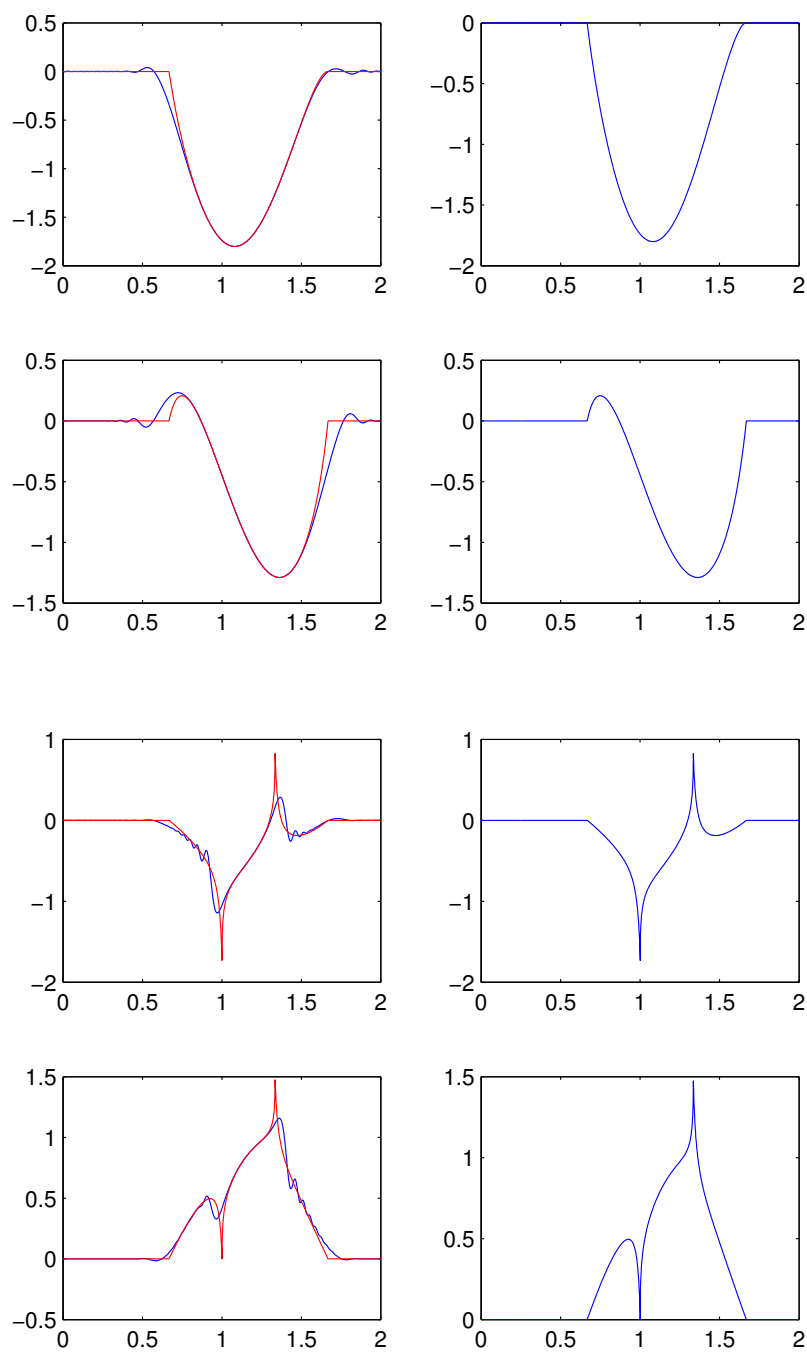


Figure 3. $\theta = \pi/3$

$$= \int_{\Gamma_D} ds(y) \int_{\Gamma_0} N^\varepsilon(x, z) (T_z^{e2}(\Phi^\varepsilon(y, z))^T (T_y^\nu)^T) ds(z) \quad (13.14)$$

$$- \int_{\Gamma_D} v^\varepsilon(x, y) (T_y^\nu u_2^\varepsilon(y)) ds(y) \quad (13.15)$$

$$= \int_{\Gamma_D} ds(y) \int_{\Gamma_0} N^\varepsilon(x, z) (T_y^\nu (T_z^{e2} \Phi^\varepsilon(z, y))^T)^T ds(z) \quad (13.16)$$

$$- \int_{\Gamma_D} v^\varepsilon(x, y) (T_y^\nu u_2^\varepsilon(y)) ds(y) \quad (13.17)$$

$$= \int_{\Gamma_D} (T_y^\nu (v^\varepsilon(x, y))^T)^T u_2^\varepsilon(y) - v^\varepsilon(x, y) (T_y^\nu u_2^\varepsilon(y)) ds(y) \quad (13.18)$$

where

$$v^\varepsilon(x, y) = \int_{\Gamma_0} N^\varepsilon(x, z) (T_z^{e2} \Phi^\varepsilon(z, y)) ds(z) \quad (13.19)$$

Since $\|T_x^\nu(u_2^\varepsilon)\|_{H^{-1/2}(\Gamma_D)} \leq C\|g\|_{H^{1/2}(\Gamma_D)}$, we obtain

$$|w^\varepsilon(x)| \leq C\|g\|_{H^{1/2}(\Gamma_D)} \max_{x \in D} (|v^\varepsilon(x, y)| + d_D |\nabla_y v^\varepsilon(x, y)|) \quad (13.20)$$

and

$$|\nabla w^\varepsilon(x)| \leq C\|g\|_{H^{1/2}(\Gamma_D)} \max_{x \in D} (|\nabla_x v^\varepsilon(x, y)| + d_D |\nabla_x \nabla_y v^\varepsilon(x, y)|) \quad (13.21)$$

By (13.8) and letting $\varepsilon \rightarrow 0^+$, we have

$$\|T_x^\nu(u_1 - u_2)\|_{H^{-1/2}(\Gamma_D)} \leq C\|g\|_{H^{1/2}(\Gamma_D)} \max_{x \in D} \lim_{\varepsilon \rightarrow 0^+} (|v^\varepsilon(x, y)| \quad (13.22)$$

$$+ d_D |\nabla_y v^\varepsilon(x, y)| + d_D |\nabla_x v^\varepsilon(x, y)| + d_D^2 |\nabla_x \nabla_y v^\varepsilon(x, y)|) \quad (13.23)$$

where u_1 is the scattering solution in the half-space and u_2 in the full-space. Now, it turns to estimate $v^\varepsilon(x, y)$. Applying the Fourier transformation to the first horizontal variable of $N^\varepsilon(z, x)$ and $T_z^{e2} \Phi^\varepsilon(z, y)$, we have

$$\mathcal{F}[N^\varepsilon](\xi, 0; x) = \frac{\mathbf{i}}{\mu \delta(\xi)} \left[\begin{pmatrix} 2\xi^2 \mu_s & -2\xi \mu_s \mu_p \\ -\xi \beta & \mu_p \beta \end{pmatrix} e^{\mathbf{i} \mu_p x_2} + \begin{pmatrix} \mu_s \beta & \xi \beta \\ 2\xi \mu_s \mu_p & 2\xi^2 \mu_p \end{pmatrix} e^{\mathbf{i} \mu_s x_2} \right] e^{-\mathbf{i} \xi x_1}$$

$$\mathcal{F}[T_z^{e2} \Phi^\varepsilon](\xi, 0; y) = \frac{\mu}{2\omega^2} \left[\begin{pmatrix} 2\xi^2 & -2\xi \mu_p \\ -\frac{\beta \xi}{\mu_p} & \beta \end{pmatrix} e^{\mathbf{i} \mu_p y_2} + \begin{pmatrix} \beta & \frac{\xi \beta}{\mu_s} \\ 2\xi \mu_s & 2\xi^2 \end{pmatrix} e^{\mathbf{i} \mu_s y_2} \right] e^{-\mathbf{i} \xi y_1}$$

Using Parseval identity combined with above two formula, we have

$$\lim_{\varepsilon \rightarrow 0^+} v^\varepsilon(x, y) = \lim_{\varepsilon \rightarrow 0^+} \int_{\mathbb{R}} \mathcal{F}[N^\varepsilon](\xi, 0; x)^T \mathcal{F}[T_z^{e2} \Phi^\varepsilon](-\xi, 0; y) d\xi$$

Lemma 13.1 *For any $x, y \in D$, let*

$$p(x, y) = \lim_{\varepsilon \rightarrow 0^+} p^\varepsilon(x, y) := \lim_{\varepsilon \rightarrow 0^+} \int_{\mathbb{R}} \frac{f(\mu_p^\varepsilon, \mu_s^\varepsilon, \xi)}{\delta^\varepsilon(\xi)} e^{\mathbf{i} \mu_\alpha^\varepsilon x_2 + \mathbf{i} \mu_\beta^\varepsilon y_2 + \mathbf{i} \xi (y_1 - x_1)} d\xi$$

where $f(a, b, c)$ is a homogeneous fifth order polynomial with respect to a, b, c and $\alpha = s, p$, $\beta = s, p$. Then there exists a constant $C > 0$ only dependent on κ such that

$$|p(x, y)| + k_s^{-1} |\nabla_x p(x, y)| + k_s^{-1} |\nabla_y p(x, y)| + k_s^{-2} |\nabla_x \nabla_y p(x, y)| \leq C((k_s h)^{-1/2} + e^{-\sqrt{k_R^2 - k_s^2} h})$$

uniformly for $x, y \in D$.

Proof. Without loss of generality, we assume $k_\alpha \leq k_\beta$. Then we can divide $p(x, y)$ into two parts:

$$\begin{aligned} p(x, y) &= \lim_{\varepsilon \rightarrow 0^+} \int_{I_1} + \int_{I_2} \frac{f(\mu_p^\varepsilon, \mu_s^\varepsilon, \xi)}{(k_\alpha^\varepsilon)^2 \delta^\varepsilon(\xi)} e^{i\mu_\alpha^\varepsilon x_2 + i\mu_\beta^\varepsilon y_2 + i\xi(y_1 - x_1)} d\xi \\ &= \int_{I_1} \frac{f(\mu_p, \mu_s, \xi)}{k_\alpha^2 \delta(\xi)} e^{i\mu_\alpha x_2 + i\mu_\beta y_2 + i\xi(y_1 - x_1)} d\xi \\ &\quad + \lim_{\varepsilon \rightarrow 0^+} \int_{I_2} \frac{f(\mu_p^\varepsilon, \mu_s^\varepsilon, \xi)}{(k_\alpha^\varepsilon)^2 \delta^\varepsilon(\xi)} e^{i\mu_\alpha^\varepsilon x_2 + i\mu_\beta^\varepsilon y_2 + i\xi(y_1 - x_1)} d\xi \\ &= p_1(x, y) + p_2(x, y) \end{aligned}$$

where $I_1 = (-k_\alpha, k_\alpha)$, $I_2 = (-2k_R + k_\alpha, k_\alpha) \cup (k_\alpha, 2k_R - k_\alpha)$ and $I_2 = R \setminus [-k_\alpha, k_\alpha]$. Substituting $\xi = k_\alpha t$ into $p_1(x, y)$, we get

$$p_1(x, y) = \int_{-1}^1 \frac{f(\mu_p(k_\alpha t), \mu_s(k_\alpha t), k_\alpha t)}{k_\alpha \delta(k_\alpha t)} e^{ik_\alpha x_2(\sqrt{1-t^2} + \tau\sqrt{\varsigma^2-t^2} + \gamma t)} dt$$

where $\tau = y_2/x_2$, $\varsigma = k_\beta/k_\alpha$ and $\gamma = (y_1 - x_1)/x_2$. It is easy to see that the phase function $\phi(t) = \sqrt{1-t^2} + \tau\sqrt{\varsigma^2-t^2} + \gamma t$ satisfies $|\phi''(t)| \geq 1/(1-t^2)^{3/2} \geq 1$ for $t \in (-1, 1)$. Then we can obtain $|p_1(x, y)| \leq C1/(k_s h)^{1/2}$ by lemma 6.1.

For $p_2(x, y)$, by changing the integration path and using same argument as in the proof of estimate for psf , we can easily obtain:

$$|p_2(x, y)| \leq C\left(\frac{1}{k_s h} + e^{-\sqrt{k_R^2 - k_s^2} h}\right)$$

This completes the proof of the estimate for $|p(x, y)|$. The other estimates can be proved by a similar argument. We omit the details \square

Lemma 13.2 For any $x, y \in D$, let

$$p(x, y) = \lim_{\varepsilon \rightarrow 0^+} p^\varepsilon(x, y) := \lim_{\varepsilon \rightarrow 0^+} \int_{\mathbb{R}} \frac{f(\mu_p^\varepsilon, \mu_s^\varepsilon, \xi)}{\delta^\varepsilon(\xi)} e^{i\mu_\alpha^\varepsilon x_2 + i\mu_\beta^\varepsilon y_2 + i\xi(y_1 - x_1)} d\xi$$

where $f(a, b, c)$ is a homogeneous fifth order polynomial with respect to a, b, c and $\alpha = s, p$, $\beta = s, p$. Then there exists a constant $C > 0$ only dependent on κ such that

$$|p(x, y)| + k_s^{-1} |\nabla_x p(x, y)| + k_s^{-1} |\nabla_y p(x, y)| + k_s^{-2} |\nabla_x \nabla_y p(x, y)| \leq C(1 + k_s d_D)((k_s h)^{-1/2} + e^{-\sqrt{k_R^2 - k_s^2} h})$$

uniformly for $x, y \in D$.

Proof. Without loss of generality, we assume $k_\alpha \leq k_\beta$. Then we can divide $p(x, y)$ into four parts:

$$\begin{aligned} p(x, y) &= \lim_{\varepsilon \rightarrow 0^+} \int_{I_1} + \int_{I_2} + \int_{I_3} \frac{f(\mu_p^\varepsilon, \mu_s^\varepsilon, \xi)}{(k_\alpha^\varepsilon)^2 \delta^\varepsilon(\xi)} e^{i\mu_\alpha^\varepsilon x_2 + i\mu_\beta^\varepsilon y_2 + i\xi(y_1 - x_1)} d\xi \\ &= \int_{I_1} + \text{PV} \int_{I_2} + \int_{I_3} \frac{f(\mu_p, \mu_s, \xi)}{k_\alpha^2 \delta(\xi)} e^{i\mu_\alpha x_2 + i\mu_\beta y_2 + i\xi(y_1 - x_1)} d\xi \\ &\quad + i\pi \left(\frac{f(\mu_p(k_R), \mu_s(k_R), k_R)}{k_\alpha^2 \delta'(k_R)} e^{i\mu_\alpha(k_R)x_2 + i\mu_\beta(k_R)y_2 + ik_R(y_1 - x_1)} \right) \end{aligned}$$

$$\begin{aligned}
& - \frac{f(\mu_p(k_R), \mu_s(k_R), -k_R)}{k_\alpha^2 \delta'(-k_R)} e^{i\mu_\alpha(k_R)x_2 + i\mu_\beta(k_R)y_2 - i k_R(y_1 - x_1)} \\
& : \quad = p_1(x, y) + p_2(x, y) + p_3(x, y) + p_4(x, y)
\end{aligned}$$

where $I_1 = (-k_\alpha, k_\alpha)$, $I_2 = (-2k_R + k_\alpha, k_\alpha) \cup (k_\alpha, 2k_R - k_\alpha)$ and $I_2 = R \setminus [-2k_R + k_\alpha, 2k_R - k_\alpha]$. Substituting $\xi = k_\alpha t$ into $p_1(x, y)$, we get

$$p_1(x, y) = \int_{-1}^1 \frac{f(\mu_p(k_\alpha t), \mu_s(k_\alpha t), k_\alpha t)}{k_\alpha \delta(k_\alpha t)} e^{i k_\alpha x_2 (\sqrt{1-t^2} + \tau \sqrt{\varsigma^2 - t^2} + \gamma t)} dt$$

where $\tau = y_2/x_2$, $\varsigma = k_\beta/k_\alpha$ and $\gamma = (y_1 - x_1)/x_2$. It is easy to see that the phase function $\phi(t) = \sqrt{1-t^2} + \tau \sqrt{\varsigma^2 - t^2} + \gamma t$ satisfies $|\phi''(t)| \geq 1/(1-t^2)^{3/2} \geq 1$ for $t \in (-1, 1)$. Then we can obtain $|p_1(x, y)| \leq C1/(k_s h)^{1/2}$ by lemma 6.1.

Let

$$g_\pm(\xi) = \frac{f(\mu_p, \mu_s, \xi)(\xi \pm k_R)}{\delta(\xi)} e^{i\mu_\alpha x_2 + i\mu_\beta y_2 + i\xi(y_1 - x_1)} d\xi$$

Then by the definition of cauchy principle value, we have

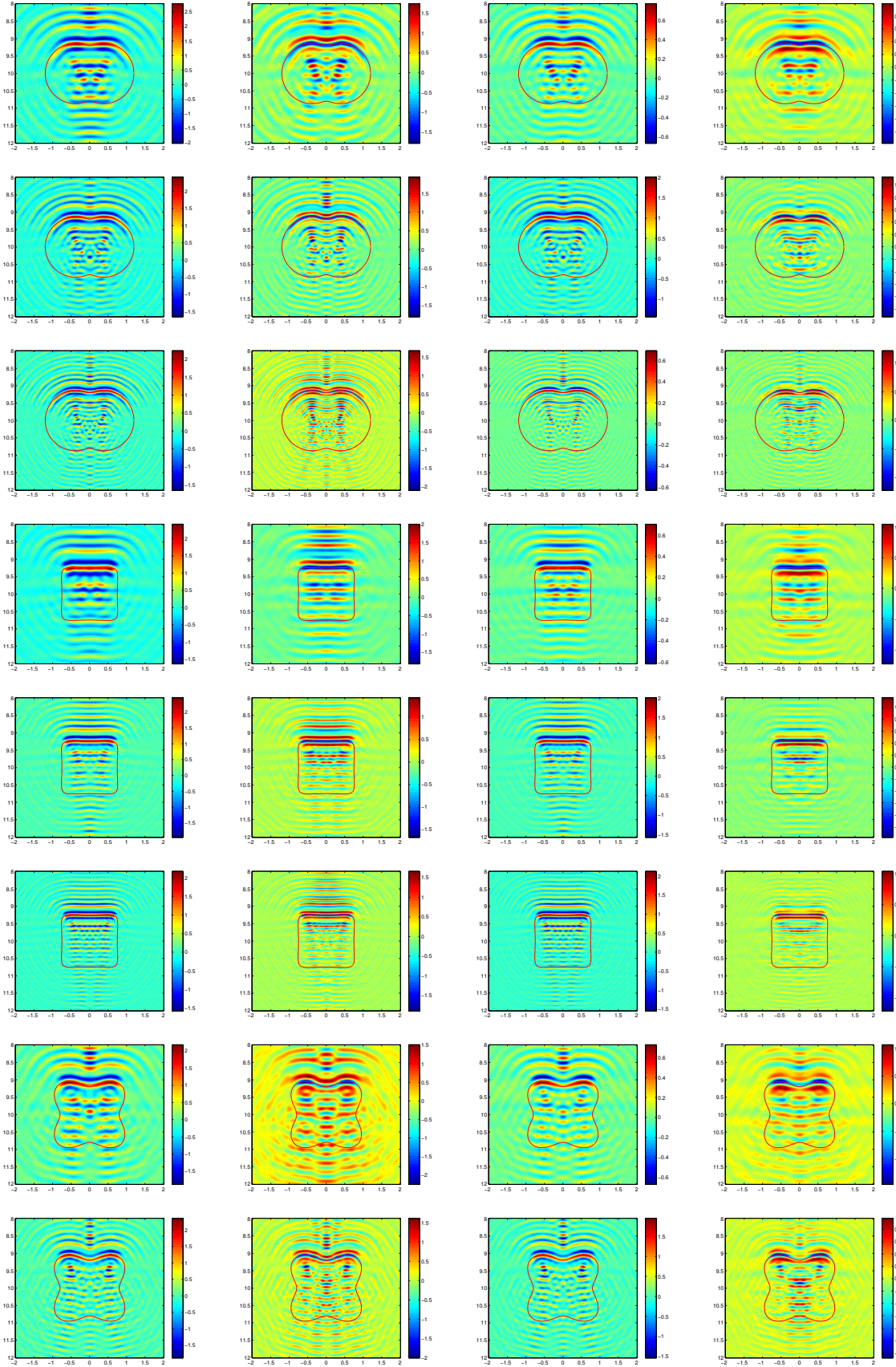
$$p_2(x, y) = \int_{-2k_R+k_\alpha}^{-k_\alpha} \frac{g_-(\xi) - g_-(-k_R)}{x + k_R} d\xi + \int_{k_\alpha}^{2k_R-k_\alpha} \frac{g_+(\xi) - g_+(k_R)}{x - k_R} d\xi$$

□

14. Some comment about phaseless imaging, 2018.01.24

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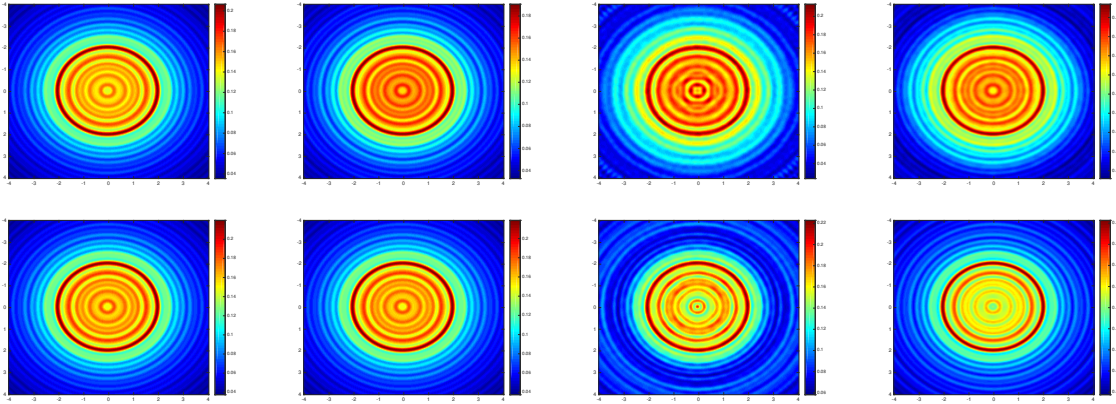


Figure 5. Circle; From left to right: vector imaging, scalar imaging, phaseless imaging128, phaseless imaging512; From up to down: $R=10$, $R=100$

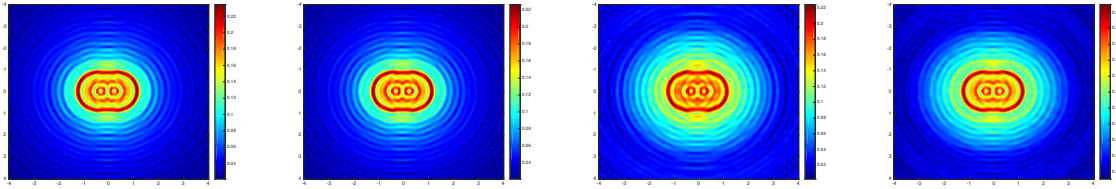


Figure 6. Peanut; From left to right: vector imaging, scalar imaging, phaseless imaging128, phaseless imaging512;

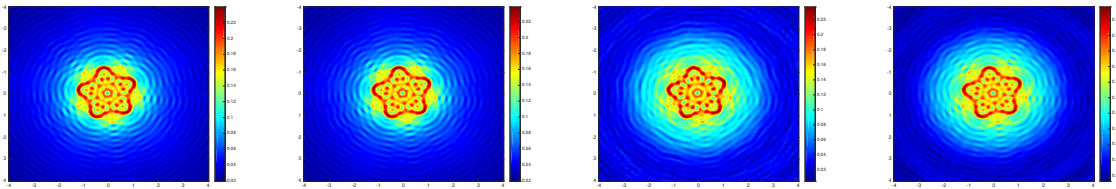


Figure 7. Peanut; From left to right: vector imaging, scalar imaging, phaseless imaging128, phaseless imaging512;

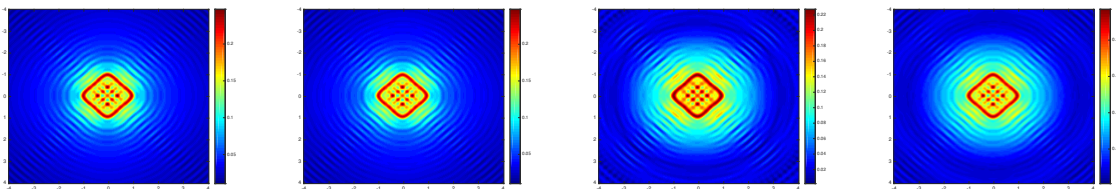


Figure 8. Peanut; From left to right: vector imaging, scalar imaging, phaseless imaging128, phaseless imaging512;