# Conformal Field Theory

# 1. Conformal Algebra

The generators of infinitesimal conformal transformations on functions of  $x^{\mu}$  are :

$$P_{\mu} = -i\partial_{\mu}; \quad L_{\mu\nu} = i(x_{\mu}\partial_{\nu} - x_{\nu}\partial_{\mu});$$

$$D = -ix^{\mu}\partial_{\mu}; \quad K_{\mu} = -i(2x_{\mu}x^{\nu}\partial_{\nu} - x^{\nu}x_{\nu}\partial_{\mu})$$

Verify the commutation relations:

$$[D, P_{\mu}] = iP_{\mu}, \qquad [D, K_{\mu}] = -iK_{\mu}.$$

### 2. Special Conformal Transformations

We want to construct finite special conformal transformation, starting from the infinite-simal ones :

$$x^{\mu'} = 2(b^{\nu}x_{\nu})x^{\mu} - b^{\mu}x^2$$
  $(x^2 \equiv x^{\mu}x_{\mu}).$ 

1. Show that the *inversion* 

$$I: x^{\mu} \to x^{\mu'} = \frac{x^{\mu}}{x^2}$$

is a conformal transformation. Can we obtain I from an infinitesimal transformation?

2. Show that the infinitesimal special conformal generators K are related to the translation generators P by

$$K_{\mu} = -IP_{\mu}I$$

[use the action of the transformations on the coordinates]

3. From the previous result and the fact that  $I^2 = 1$ , show that

$$e^{ib^{\mu}K_{\mu}} = Ie^{-ib^{\mu}P_{\mu}}I.$$

Use this identity to find the form of the finite special conformal transformation with parameter  $b^{\mu}$ .

#### 3. Noether currents

Consider a classical field theory for a scalar field  $\phi(x)$  and an infinitesimal transformation, with parameter  $\epsilon$ , which we assume linear in  $\epsilon$  (but not necessarily in  $\phi$ ):

$$\delta \phi \equiv \phi'(x) - \phi(x) = X(\phi)\epsilon$$

If such a transformation leaves the action invariant, or equivalently if the functional variation of the Lagrangian is a total divergence,

$$\delta \mathcal{L} = \epsilon(\partial^{\mu} F_{\mu}),$$

then Noether's theorem ensures the existence of a current  $J^{\mu}$ , conserved on the equations of motion, and given by

$$J^{\mu} = \frac{\partial \mathcal{L}}{\partial (\partial_{\mu} \phi)} X(\phi) - F^{\mu}.$$

Below we consider the two kinds of transformations corresponding to translation and dilatation of the coordinates, and the corresponding action on a scalar field of conformal weight  $\Delta$ ,

$$x'_{\mu} = x_{\mu} - \epsilon_{\mu}, \quad \phi'(x') = \phi(x)$$

and

$$x'_{\mu} = x_{\mu} - \epsilon x_{\mu}, \quad \phi'(x') = (1 + \Delta \epsilon)\phi(x)$$

- 1. Write the functional variation  $\delta\phi(x)$  of the field under the two transformations
- 2. Write functional variation of the Lagrangian, and verify that it is a total divergence in both cases (assuming the classical theory is scale invariant). This will lead you to identify the quantity  $F^{\mu}$ . [Recall that for a CFT, the Lagrangian is a scalar field of weight d]
- 3. Construct the Noether currents for translation and dilatation using the general formula, and check they are given by the following expressions:

$$T_{\mu\nu} = \frac{\partial \mathcal{L}}{\partial(\partial^{\mu}\phi)} \partial_{\nu}\phi - \eta_{\mu\nu}\mathcal{L}, \qquad J_{\mu}^{D} = x^{\nu}T_{\mu\nu} + \Delta\phi \frac{\partial \mathcal{L}}{\partial(\partial^{\mu}\phi)}.$$

## 4. Improved stress tensor and conformal coupling

Consider the theory of a real scalar field in d dimensions with action

$$S = -\int d^d x \left[ \frac{1}{2} \partial^\mu \phi \partial_\mu \phi + \frac{\lambda}{n} \phi^n \right]$$

#### 3.1. Canonical stress tensor

- i. What must be the conformal weight  $\Delta$  of the field  $\phi$  such that the kinetic term has the appropriate scaling dimension?
- ii. Write down the explicit form of the canonical (i.e. Noether) stress tensor  $T_{\mu\nu}$ , of its trace, and of the dilatation Noether current  $J_D^{\mu}$ . For which value(s) of n is the dilatation current conserved (on a solution of the equations of motion)? Is the stress tensor traceless in this case?

### 3.2. Improved stress tensor

We want to construct an *improved* stress tensor  $\Theta_{\mu\nu}$  and an improved dilatation current  $\tilde{J}_D^{\mu}$  such that 1)  $\Theta_{\mu\nu}$  it is conserved 2) it is traceless iff the theory is conformal.

i. Define:

$$\Theta_{\mu\nu} = T_{\mu\nu} + \xi \left( \partial_{\mu} \partial_{\nu} - \eta_{\mu\nu} \partial^{\rho} \partial_{\rho} \right) \phi^{2}$$

Show that it is conserved and that the associated Noether charges  $P_{\mu}$  coincide with those associated to  $T_{\mu\nu}$ ,

$$P_{\mu} \equiv \int d^3x \, T_{0\mu} = \int d^3x \, \Theta_{0\mu}$$

- ii. Show, for an appropriate value of  $\xi$  (which depends on d),  $\Theta_{\mu\nu}$  is traceless (on solutions of the EOM) for the appropriate value of n (which you have found in part 1) for which the theory is conformally invariant.
- iii. Define the improved dilatation current  $\tilde{J}_{\mu}^{D}$  as

$$\tilde{J}^D_{\mu} = x^{\nu} \Theta_{\mu\nu}.$$

Show that tracelessness of  $\Theta_{\mu\nu}$  implies conservation of  $\tilde{J}_{\mu}^{D}$ , and that the associated dilatation charge D is the same as the one obtained from  $J_{\mu}^{D}$ , i.e.

$$D \equiv \int d^3x J_0^D = \int d^3x \, \tilde{J}_0^D.$$

The conserved charges constructed from  $\Theta_{\mu\nu}$  and  $\tilde{J}_{\mu}^{D}$  are the same as those obtained from  $T_{\mu\nu}$  and  $J_{\mu}^{D}$ , respectively. Therefore the improved conserved currents are associated to the same symmetries as the old (Noether) currents.

#### 3.3. Covariant version

Now consider the same scalar field theory minimally coupled to a background metric  $g_{\mu\nu}$ ,

$$S = -\int d^d x \sqrt{-g} \left[ \frac{1}{2} g^{\mu\nu} \partial_{\mu} \phi \partial_{\nu} \phi + \frac{\lambda}{n} \phi^n \right]$$

i. Show that the definition

$$T_{\mu\nu} = -\frac{2}{\sqrt{-g}} \frac{\delta S}{\delta g^{\mu\nu}}$$

gives back the canonical stress tensor (after setting  $g_{\mu\nu} = \eta_{\mu\nu}$ )

ii. Can you find a (covariant) term such that, when added to the action, gives the improved stress tensor of part 2 as

$$\Theta_{\mu\nu} = -\frac{2}{\sqrt{-g}} \frac{\delta S}{\delta g^{\mu\nu}} ?$$

[Hint: recall how the Ricci scalar changes under a variation of the metric. Remember to set  $g_{\mu\nu} = \eta_{\mu\nu}$  at the end of the calculation, since we want to reproduce the flat space result of part 2.].

The resulting action is the action of a *conformally coupled* scalar field. It is conformal in any background, i.e. it is invariant under Weyl transformations (which are *not* coordinate transformations)

$$g_{\mu\nu} \to e^{2\omega(x)} g_{\mu\nu}, \quad \phi \to e^{-\Delta\omega(x)} \phi$$

for an arbitrary function  $\omega(x)$ .