ICFP M2 – SOFT MATTER PHYSICS Tutorial 2. The n = 0 theorem

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One of the important achievements of de Gennes was to draw a connexion between critical phenomena and polymer statistics. Here, we would like to understand the essence of this link – without questioning the mathematical legitimity of the invoked limit. We consider the following magnetic system: a d-dimension periodic lattice at temperature T, at each site i of which is located one n-component classical spin \mathbf{S}_i . Two spins \mathbf{S}_i and \mathbf{S}_j interact through the exchange energy $-J_{i,j} \mathbf{S}_i \cdot \mathbf{S}_j$, with $J_{i,j} > 0$ to promote alignment. We assume $J_{i,j} = k_{\mathrm{B}}T_J$ (where k_{B} is the Boltzmann constant) to be a finite constant when i and j correspond to neighbouring sites, and we impose $J_{i,j}$ to be zero otherwise. Finally, for practical purposes, we choose the following normalization: $\mathbf{S}_i^2 = n$.

I Preliminaries

- 1 Express the Hamiltonian H of the problem.
- **2** Write the associated canonical partition function Z as an integral over the solid angles Ω_i of all the spins \mathbf{S}_i .
- 3 By introducing $\Omega = \prod_i \Omega_i$, rewrite Z as an average (noted $\langle \cdot \rangle_0$) over all possible orientations of the spins.
- 4 Recall the definition of the average $\langle \cdot \rangle$ in the canonical ensemble, and relate it to $\langle \cdot \rangle_0$.

II The n=0 characteristic function

We consider a given spin $\mathbf{S}_i \equiv \mathbf{S}$ as a uniformly distributed random vector.

- 1 Write the characteristic function $f(\mathbf{k})$, of wave vector \mathbf{k} , associated to \mathbf{S} .
- **2** Explain how f allows to generate the successive moments of the distribution of **S**. Relate in particular $\langle S_{\alpha}S_{\beta}\rangle_0$ to f, where the α and β indexes indicate components of **S**.
- **3** By invoking the normalization of **S**, demonstrate that f satisfies the following partial differential equation:

$$\Delta f + nf = 0 , \qquad (1)$$

where Δ is the Laplacian operator in Fourier space.

ICFP M2 – Soft Matter Physics

4 By symmetry of the problem, how does the orientation of \mathbf{k} matter? Simplify Eq. (2) into the following ordinary differential equation:

$$f'' + \frac{n-1}{k}f' + nf = 0 , (2)$$

where the prime denotes the derivative with respect to the norm $k = ||\mathbf{k}||$ of \mathbf{k} .

- **5** What are the boundary conditions at k = 0?
- **6** When n = 0, solve for f and conclude about the moments of the distribution of **S**.

III Loop expansion of the n = 0 partition function

1 Show that the partition function has the exact expansion at n = 0:

$$\frac{Z}{\Omega} = \lim_{n \to 0} \left\langle \prod_{(i < j)} \left[1 + \frac{T_J}{T} \sum_{\alpha} S_{i\alpha} S_{j\alpha} + \frac{1}{2} \left(\frac{T_J}{T} \right)^2 \sum_{\alpha, \beta} S_{i\alpha} S_{i\beta} S_{j\alpha} S_{j\beta} \right] \right\rangle_0$$
(3)

where the notation (i < j) indicates that i < j together with sites i and j being nearest neighbours.

- **2** Propose a diagrammatic representation, involving bonds connecting sites on a lattice as well as "charges", of each term in the bracket $[\cdot]$.
- 3 Show that only closed loops made of the previous elementary diagrams contribute to Z. What about loops with self-intersections?
- 4 Calculate Z using those considerations, and comment the obtained result.

IV The n=0 theorem

- 1 Express the correlator $\langle S_{p\delta}S_{q\delta}\rangle$, for two given sites p and q and a component index δ , as an average $\langle \cdot \rangle_0$ over all possible orientations of all the spins.
- **2** As in Eq. (3), provide an exact expansion of $\langle S_{p\delta}S_{q\delta}\rangle$ in the n=0 case.
- **3** What are the diagrams contributing to $\langle S_{p\delta} S_{q\delta} \rangle$?
- 4 Let $W_N(p,q)$ be the number of self-avoiding random walks linking sites p and q in exactly N steps. Show that for n=0:

$$\langle S_{p\delta} S_{q\delta} \rangle = \sum_{N} W_N(p,q) \left(\frac{T_J}{T}\right)^N .$$
 (4)

Comment this central result and link it with polymer physics.