现代数学物理方法

第二章, 群论基础

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Importance of Young tableaux:

- Each different tableaux of n-boxes represents a different conjugacy class of S_n .
- **②** The Young tableaux are in *one-to-one* correspondence with the irreducible representations of S_n .

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The numbers of group elements in these conjugacy classes are:

$$\frac{3!}{3!} = 1, \qquad \frac{3!}{2} = 3, \qquad \frac{3!}{3} = 2.$$

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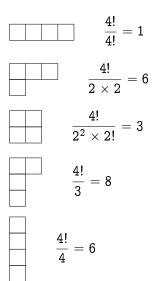
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Obviously, it is (167425)(3).

e.g.,

1 2

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$$\boxed{1 | 2} \iff [e + (12)] | 12 \rangle$$

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$$\frac{1 \mid 2}{3} \quad \longleftrightarrow \\
[e - (13)][e + (12)] \mid 123 \rangle$$

e.g.,

$$\begin{array}{c|c}
\hline 1 & 2 & \longrightarrow & [e + (12)] & |12\rangle & = & |12\rangle + & |21\rangle \\
\hline \hline 1 & 2 & \longrightarrow & \\
\hline 3 & & & \\
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\end{array}$$

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This state spans another 1-dimensional irreducible subspace which defines the so-called alternate representation D_1 of S_3 :

$$D_{1}[e] |\Psi_{1}\rangle = D_{1}[(123)] |\Psi_{1}\rangle = D_{1}[(132)] |\Psi_{1}\rangle = |\Psi_{1}\rangle$$

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Therefore, \Box is associated with a 2-d irreducible representation of S_3 .

Explanation:

The state related to the Young tableau

2	1
3	

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$$|\psi_{213}\rangle = [e - (23)][e + (12)]|213\rangle$$

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= $[e - (23) + (12) - (132)]|213\rangle$

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$$\begin{aligned} |\psi_{213}\rangle &= [e - (23)][e + (12)] |213\rangle \\ &= [e - (23) + (12) - (132)] |213\rangle \\ &= |213\rangle - |312\rangle + |123\rangle - |132\rangle \end{aligned}$$

The state related to the Young tableau

is determined as follows:

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angle &= [e-(23)][e+(12)]\,|213
angle \ &= [e-(23)+(12)-(132)]\,|213
angle \ &= |213
angle - |312
angle + |123
angle - |132
angle \end{array}$$

Recall that,

$$|\psi_{21}\rangle = |123\rangle + |213\rangle - |321\rangle - |231\rangle$$

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Hence,

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angle-|\psi_{22}
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Standard Young tableaux:

• In a standard Young tableau, the filled numbers increase within a row from left to right and within a column from top to down.

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$$\begin{array}{|c|c|c|}\hline 1 & 2 \\\hline 3 & & |\psi_{21}\rangle = |123\rangle + |213\rangle - |321\rangle - |231\rangle \\\hline \hline 1 & 3 \\\hline 2 & & |\psi_{22}\rangle = |132\rangle + |312\rangle - |231\rangle - |321\rangle \\\hline \end{array}$$

Standard Young tableaux:

- In a standard Young tableau, the filled numbers increase within a row from left to right and within a column from top to down.
- For a given Young tableau, the number of the standard Young tableaux is the same as the dimensions of the corresponding irreducible representation.

The standard Young tableaux of S_3 are as follows:

1 2 3

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				1	
1	2	3];	2	
				3	

The standard Young tableaux of S_3 are as follows:

	1	1 2	1 2
1 2 3;	2;	1 2,	1 3
	3	3	4

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• Go back to the construction of the 2-d irreducible representation of S_3 .

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$$\begin{array}{c|ccccc}
\hline 1 & 2 & 3 \\
\hline 3 & 5 & \hline 2 \\
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\end{array}$$

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\end{array}$$

$$D_2\lceil(12)
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 $D_2[(12)]\ket{\Psi_{22}}$

$$D_2[(12)]\ket{\Psi_{22}} = D_2[(12)]igg\{\ket{132}+\ket{312}-\ket{231}-\ket{321}igg\}$$

$$egin{aligned} D_2 ig[(12) ig] ig| \Psi_{22}
angle &= D_2 ig[(12) ig] igg\{ ig| 132
angle + ig| 312
angle - ig| 231
angle - ig| 321
angle igg\} \ &= igg\{ ig| 231
angle + ig| 321
angle - ig| 132
angle - ig| 312
angle igg\} \end{aligned}$$

$$egin{aligned} D_2[(12)] \ket{\Psi_{22}} &= D_2[(12)] igg\{ \ket{132} + \ket{312} - \ket{231} - \ket{321} igg\} \ &= igg\{ \ket{231} + \ket{321} - \ket{132} - \ket{312} igg\} \ &= -\ket{\Psi_{22}} \end{aligned}$$

$$\begin{split} D_{2}[(12)] \ket{\Psi_{22}} &= D_{2}[(12)] \bigg\{ \ket{132} + \ket{312} - \ket{231} - \ket{321} \bigg\} \\ &= \bigg\{ \ket{231} + \ket{321} - \ket{132} - \ket{312} \bigg\} \\ &= -\ket{\Psi_{22}} \end{split}$$

By setting
$$|\psi_{21}\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$
 and $|\psi_{22}\rangle = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$,

$$\begin{split} D_{2}[(12)] \ket{\Psi_{22}} &= D_{2}[(12)] \bigg\{ \ket{132} + \ket{312} - \ket{231} - \ket{321} \bigg\} \\ &= \bigg\{ \ket{231} + \ket{321} - \ket{132} - \ket{312} \bigg\} \\ &= -\ket{\Psi_{22}} \end{split}$$

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$$\begin{split} D_2 \big[(12) \big] \, \big| \Psi_{22} \big\rangle \, &= D_2 \big[(12) \big] \bigg\{ \, \big| 132 \big\rangle + \big| 312 \big\rangle - \big| 231 \big\rangle - \big| 321 \big\rangle \, \bigg\} \\ &= \bigg\{ \, \big| 231 \big\rangle + \big| 321 \big\rangle - \big| 132 \big\rangle - \big| 312 \big\rangle \, \bigg\} \\ &= - \big| \Psi_{22} \big\rangle \end{split}$$

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$$D_2[(12)] = \left[egin{array}{cc} 1 & -1 \ 0 & -1 \end{array}
ight]$$

$$D_2[(23)]\ket{\psi_{21}}$$

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$$egin{aligned} D_2 ig[(23) ig] ig| \psi_{21} ig> &= D_2 ig[(23) ig] igg\{ ig| 123 ig> + ig| 213 ig> - ig| 321 ig> - ig| 231 iga> igg\} \ &= igg\{ ig| 132 ig> + ig| 312 ig> - ig| 231 ig> - ig| 321 ig> igg\} \end{aligned}$$

$$D_{2}[(23)] |\psi_{21}\rangle = D_{2}[(23)] \Big\{ |123\rangle + |213\rangle - |321\rangle - |231\rangle \Big\}$$

$$= \Big\{ |132\rangle + |312\rangle - |231\rangle - |321\rangle \Big\}$$

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$$\begin{split} D_{2}[(23)] \ket{\psi_{21}} &= D_{2}[(23)] \bigg\{ \ket{123} + \ket{213} - \ket{321} - \ket{231} \bigg\} \\ &= \bigg\{ \ket{132} + \ket{312} - \ket{231} - \ket{321} \bigg\} \\ &= \ket{\Psi_{22}} \\ D_{2}[(23)] \ket{\psi_{22}} &= D_{2}[(23)] \bigg\{ \ket{132} + \ket{312} - \ket{231} - \ket{321} \bigg\} \\ &= \bigg\{ \ket{123} + \ket{213} - \ket{321} - \ket{231} \bigg\} \\ &= \ket{\psi_{21}} \end{split}$$

Hence,

$$D_2[(23)] = \left[\begin{array}{cc} 0 & 1 \\ 1 & 0 \end{array} \right]$$

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$$= \begin{bmatrix} 1 & -1 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

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• In conclusion,

$$D_{2}[(123)] = D_{2}[(12)(23)] = D_{2}[(12)]D_{2}[(23)]$$

$$= \begin{bmatrix} 1 & -1 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

$$= \begin{bmatrix} -1 & 1 \\ -1 & 0 \end{bmatrix}$$

• In conclusion, the 2-d irreducible Rep. $D_2(S_3)$ is realized by,

$$D_2[e] = egin{bmatrix} 1 & 0 \ 0 & 1 \end{bmatrix}$$
 $D_2[(12)] = egin{bmatrix} 1 & -1 \ 0 & -1 \end{bmatrix}$ $D_2[(13)] = egin{bmatrix} 0 & 1 \ -1 & 1 \end{bmatrix}$ $D_2[(132)] = egin{bmatrix} 0 & 1 \ 1 & 0 \end{bmatrix}$ $D_2[(132)] = egin{bmatrix} 0 & -1 \ 1 & -1 \end{bmatrix}$

$$\chi_2[e]=2$$

$$\chi_2[e] = 2$$
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$$H = \sum_{g \in S_3} igl[D_2(g) igr]^\dagger D_2(g) = igl[egin{array}{cc} 8 & -4 \ -4 & 8 \ \end{array} igr]$$

$$\left[egin{array}{cc} 8 & -4 \ -4 & 8 \end{array}
ight] \left[egin{array}{cc} a \ b \end{array}
ight] = \lambda \left[egin{array}{cc} a \ b \end{array}
ight]$$

$$\left[\begin{array}{cc} 8 & -4 \\ -4 & 8 \end{array}\right] \left[\begin{array}{c} a \\ b \end{array}\right] = \lambda \left[\begin{array}{c} a \\ b \end{array}\right]$$

$$0 = \begin{vmatrix} 8 - \lambda & -4 \\ -4 & 8 - \lambda \end{vmatrix}$$

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$$|\lambda=4
angle=rac{1}{\sqrt{2}}e^{i\phi_1}\left[egin{array}{c}1\1\end{array}
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where ϕ_1 and ϕ_2 are two arbitrary real parameters (phases).

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where ϕ_1 and ϕ_2 are two arbitrary real parameters (phases). These two eigenvectors can be used to define a unitary matrix

$$u=\left[egin{array}{cc} rac{e^{i\phi_1}}{\sqrt{2}} & rac{e^{i\phi_2}}{\sqrt{2}} \ rac{e^{i\phi_1}}{\sqrt{2}} & -rac{e^{i\phi_2}}{\sqrt{2}} \end{array}
ight]$$

$$H = \left[\begin{array}{cc} 8 & -4 \\ -4 & 8 \end{array} \right]$$

$$H \hspace{.1in} = \left[egin{array}{ccc} 8 & -4 \ -4 & 8 \end{array}
ight] = u \left[egin{array}{ccc} 4 & 0 \ 0 & 12 \end{array}
ight] u^{\dagger}$$

$$egin{array}{ll} H & = \left[egin{array}{ccc} 8 & -4 \ -4 & 8 \end{array}
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ight] u^{\dagger} \ & = \left[egin{array}{ccc} rac{e^{i\phi_1}}{\sqrt{2}} & rac{e^{i\phi_2}}{\sqrt{2}} \ rac{e^{i\phi_2}}{\sqrt{2}} & -rac{e^{i\phi_2}}{\sqrt{2}} \end{array}
ight] \left[egin{array}{ccc} 4 & 0 \ 0 & 12 \end{array}
ight] \left[egin{array}{ccc} rac{e^{-i\phi_1}}{\sqrt{2}} & rac{e^{-i\phi_1}}{\sqrt{2}} \ rac{e^{-i\phi_2}}{\sqrt{2}} \end{array}
ight] \end{array}$$

$$H = \begin{bmatrix} 8 & -4 \\ -4 & 8 \end{bmatrix} = u \begin{bmatrix} 4 & 0 \\ 0 & 12 \end{bmatrix} u^{\dagger}$$

$$= \begin{bmatrix} \frac{e^{i\phi_1}}{\sqrt{2}} & \frac{e^{i\phi_2}}{\sqrt{2}} \\ \frac{e^{i\phi_1}}{\sqrt{2}} & -\frac{e^{i\phi_2}}{\sqrt{2}} \end{bmatrix} \begin{bmatrix} 4 & 0 \\ 0 & 12 \end{bmatrix} \begin{bmatrix} \frac{e^{-i\phi_1}}{\sqrt{2}} & \frac{e^{-i\phi_1}}{\sqrt{2}} \\ \frac{e^{-i\phi_2}}{\sqrt{2}} & -\frac{e^{-i\phi_2}}{\sqrt{2}} \end{bmatrix}$$

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$$\Omega_- = u \left[egin{array}{cc} \sqrt{4} & 0 \ 0 & \sqrt{12} \end{array}
ight] u^\dagger$$

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$$egin{array}{ll} \Omega &= u \left[egin{array}{ccc} \sqrt{4} & 0 \ 0 & \sqrt{12} \end{array}
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ight] \left[egin{array}{ccc} 2 & 0 \ 0 & 2\sqrt{3} \end{array}
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$$\begin{split} \Omega &= u \begin{bmatrix} \sqrt{4} & 0 \\ 0 & \sqrt{12} \end{bmatrix} u^\dagger \\ &= \begin{bmatrix} \frac{e^{i\phi_1}}{\sqrt{2}} & \frac{e^{i\phi_2}}{\sqrt{2}} \\ \frac{e^{i\phi_1}}{\sqrt{2}} & -\frac{e^{i\phi_2}}{\sqrt{2}} \end{bmatrix} \begin{bmatrix} 2 & 0 \\ 0 & 2\sqrt{3} \end{bmatrix} \begin{bmatrix} \frac{e^{-i\phi_1}}{\sqrt{2}} & \frac{e^{-i\phi_1}}{\sqrt{2}} \\ \frac{e^{-i\phi_2}}{\sqrt{2}} & -\frac{e^{-i\phi_2}}{\sqrt{2}} \end{bmatrix} \\ &= \begin{bmatrix} 1 + \sqrt{3} & 1 - \sqrt{3} \\ 1 - \sqrt{3} & 1 + \sqrt{3} \end{bmatrix} \end{split}$$

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The 2-dimensional unitary irreducible representation of S_3 is then constructed as,

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Explicitly,

$$D_2^{ ext{unitary}}ig[eig] = \left[egin{array}{cc} 1 & 0 \ 0 & 1 \end{array}
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$$D_2^{ ext{unitary}}ig[(12)ig] = \left[egin{array}{cc} \sqrt{3}/2 & -1/2 \ -1/2 & -\sqrt{3}/2 \end{array}
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$$\begin{split} &D_2^{\text{unitary}}[(13)] = \left[\begin{array}{cc} -\sqrt{3}/2 & -1/2 \\ -1/2 & \sqrt{3}/2 \end{array} \right] \\ &D_2^{\text{unitary}}[(23)] = \left[\begin{array}{cc} 0 & 1 \\ 1 & 0 \end{array} \right] \\ &D_2^{\text{unitary}}[(123)] = \left[\begin{array}{cc} -1/2 & \sqrt{3}/2 \\ -\sqrt{3}/2 & -1/2 \end{array} \right] \\ &D_2^{\text{unitary}}[(132)] = \left[\begin{array}{cc} -1/2 & -\sqrt{3}/2 \\ \sqrt{3}/2 & -1/2 \end{array} \right] \end{split}$$

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Warning:

The matrix forms of the 2-dimensional unitary irreducible representation of S_3 are still not unique, although they are equivalent to each other.

An alternative realization of this 2-d irreducible unitary representation for S_3 is,

$$D_2(e) = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$
 $D_2[(132)] = \begin{bmatrix} -1/2 & \sqrt{3}/2 \\ -\sqrt{3}/2 & -1/2 \end{bmatrix}$
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Problem:

• Find the group of all the discrete rotations that leave a regular tetrahedron invariant by labeling the four vertices and considering the rotations as permutations on the four vertices. This defines a four dimensional representation of a group. Find the conjugacy classes and the characters of the irreducible representations of this group.