## Lecture 6 - Topics

• The relativistic point particle: Action, reparametrizations, and equations of motion

Reading: Zwiebach, Chapter 5

Continued from last time.

$$\frac{\partial \mathcal{P}^t}{\partial t} + \frac{\partial \mathcal{P}^x}{\partial x} = 0$$

$$\mathcal{P}^t = \mu_0 \partial y / \partial t$$
$$\mathcal{P}^x = -T_0 \partial y / \partial x$$

Similar to  $\partial_{\mu}J^{\mu}=0,\ \partial\rho/\partial t+\nabla\cdot\vec{J}=0,\ Q=\int dx\rho$ 

Free BC (Neumann BC):

$$\begin{split} \mathcal{P}^x(t,x_*) &= 0 \\ P_y &= \int_0^a \mu_0 dx (\partial y/\partial t) = \int_{0^a} dx \mathcal{P}^t \\ \partial P_y/\partial t &= \int_0^a dx \partial \mathcal{P}^t/\partial t = -\int_0^a dx \partial \mathcal{P}^x/\partial x = -[\mathcal{P}^x(t,x=a) - \mathcal{P}^x(t,x=0)] \end{split}$$

Conservation of momentum?

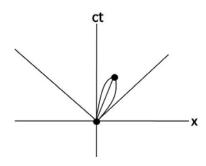
## Free Relativistic Particle

Non-relativistic Action:

$$S = \int dt \left(\frac{1}{2}mv^2\right)$$

Calculation: dv/dt = 0

Relativistic Particles:



Everyone should agree on action. It's a Lorentz invar.

$$-ds^2=-\eta_{\mu\nu}dx^\mu dx^\nu$$
 
$$ds=cdt\sqrt{1-v^2/c^2}=cd\tau$$
 
$$s=-mc^2\int_{\mathcal{P}}\frac{ds}{c}=-mc\int_{\mathcal{P}}ds$$
 So:  $s=-mc^2\int_{t_i}^{t_f}dt\sqrt{1-v^2/c^2}$ 

Check:

Lagrangian:

$$\begin{split} L &= -mc^2 \sqrt{1 - \frac{v^2}{c^2}} \\ &= -mc^2 (1 - \frac{1}{2} \frac{v^2}{c^2} - \ldots) & \text{Taylor Expansion} \\ &= \underbrace{-mc^2}_{\text{rest energy}} + \underbrace{\frac{1}{2} mv^2}_{\text{kinetic energy}} \end{split}$$

Momentum:

$$\begin{split} \vec{P} &= \frac{\partial L}{\partial \vec{v}} \\ &= -mc^2 \cdot \frac{\frac{1}{2} \frac{-2\vec{v}}{c^2}}{\sqrt{1 - \frac{v^2}{c^2}}} \\ &= \frac{mv^2}{\sqrt{1 - \frac{v^2}{c^2}}} \end{split}$$

Hamiltonian:

$$H = \vec{p} \cdot \vec{v} - L = \dots = \frac{mc^2}{\sqrt{1 - \frac{v^2}{c^2}}}$$

## Parameterization

Have parameterization  $x^{\mu}(\tau)$  (the  $x^{\mu}$ 's are functions of  $\tau$ )

$$ds^2 = -\eta_{\mu\nu} dx^\mu dx^\nu$$

$$ds = \sqrt{-\eta_{\mu\nu} \left(\frac{dx^{\mu}}{d\tau}\right) \left(\frac{dx^{\nu}}{d\tau}\right)} d\tau$$

$$s = -mc \int_{t_{i}}^{t_{f}} \sqrt{-\eta_{\mu\nu} \frac{dx^{\mu}}{d\tau} \frac{dx^{\nu}}{d\tau}} d\tau$$

$$\tau'(\tau):$$

$$\frac{dx^{\mu}}{d\tau} = \frac{dx^{\mu}}{d\tau'} \frac{d\tau'}{d\tau}$$

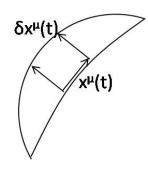
$$s = -mc \int_{t_{i}}^{t_{f}} \sqrt{-\eta_{\mu\nu} \left(\frac{dx^{\mu}}{d\tau'} \frac{dx^{\nu}}{d\tau'}\right)} \frac{d\tau'}{d\tau} d\tau$$

$$= -mc \int_{t_{i}}^{t_{f}} \sqrt{-\eta_{\mu\nu} \left(\frac{dx^{\mu}}{d\tau'} \frac{dx^{\nu}}{d\tau'}\right)} d\tau'$$

So using a different parameter,  $\tau'$  (instead of  $\tau$ ) gets same action s. s is reparameterization-invariant.

Quick calculation to find equation of motion from  $s=-mc\int\sqrt{1-\frac{v^2}{c^2}}dt$ . Should get derivative of rel. momentum with respect to time = 0.

$$S = -mc \int dS$$
$$\delta S = -mc \int \delta(dS)$$



$$dS^{2} = -\eta_{\mu\nu} dx^{\mu} dx^{\nu}$$
$$(dS)^{2} = -\eta_{\mu\nu} \frac{dx^{\mu}}{d\tau} \frac{dx^{\nu}}{d\tau} (d\tau)^{2}$$
$$2(dS) \cdot \delta(dS) = -2\eta_{\mu\nu} \delta\left(\frac{dx^{\mu}}{d\tau}\right) \frac{dx^{\nu}}{d\tau} (d\tau)^{2}$$

$$\delta(dS) = -\eta_{\mu\nu} \frac{d}{d\tau} (\delta x^{\mu}) \frac{dx^{\nu}}{ds} d\tau$$

Must vary with  $dx^{\mu}/d\tau$  and  $dx^{\nu}/d\tau$ , but since  $\eta_{\mu\nu}$  is symmetric sufficient to vary just  $dx^{\mu}/d\tau$  and multiply by 2.

$$\delta(dS) = -\frac{d}{d\tau} (\delta x^{\mu}) \frac{dx_{\mu}}{dS} d\tau$$

$$\begin{split} \delta S &= \int_{\tau_i}^{\tau_f} \frac{d(\delta x^\mu)}{d\tau} \bigg( mc \frac{dx^\mu}{ds} \bigg) d\tau \\ &= \int_{\tau_i}^{\tau_f} \bigg[ \frac{d}{d\tau} (\delta x^\mu P_\mu) - \delta x^\mu \frac{dP_\mu}{d\tau} \bigg] d\tau \end{split}$$

$$\delta x^{\mu}(\tau_i) = \delta x^{\mu}(\tau_f) = 0$$

$$dS = -\int_{\tau_i}^{\tau_f} \left( \delta x^{\mu}(\tau) \frac{dP_{\mu}}{d\tau} \right) d\tau$$

Equation of Motion:

$$\frac{dP_{\mu}}{d\tau} = 0$$

This means that  $P_{\mu}$  constant on world-line. Constant as a function of any parameter!

$$\underbrace{\frac{dP_{\mu}}{dt}}_{0} = \underbrace{\frac{dP_{\mu}}{d\tau}}_{0} \cdot \underbrace{\frac{d\tau}{dt}}_{\neq 0}$$

Therefore:  $\frac{d}{d\tau} \left( \frac{dx^{\mu}}{ds} \right) = 0$ ,  $\frac{d^2}{ds^2} (dx^{\mu}) = 0$  (if  $\tau = s$ . Okay because  $\tau$  is arbitrary.)

But can't assign  $s=\tau\colon d^2x^\mu/d\tau^2 \neq =0.$ 

$$\frac{d}{ds} \left( \frac{dx^{\mu}}{d\tau} \right) \neq 0$$

## Coupling to Electromagnetism

Lorentz Force Equation:

$$\begin{split} \frac{dP_{\mu}}{dS} &= \frac{q}{c} \cdot F_{\mu\nu} \frac{dx^{\nu}}{ds} \\ \frac{dP_{\mu}}{d\tau} &= \frac{q}{c} \cdot F_{\mu\nu} \frac{dx^{\nu}}{d\tau} \\ S &= -mc \int_{P} dS + \frac{q}{c} \int_{\mathcal{P}} A_{\mu}(x(\tau)) \frac{dx^{\mu}}{d\tau} d\tau \end{split}$$

A: Nevitz-Schwartz Tensor