

# Report for Electromagnetics Class Projects

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## 1 Deduction $Z_2$ from $\Gamma, T$

Firstly, in the coaxial line, there are the following expressions for electric and magnetic fields:

$$E = \frac{V \cdot \hat{\rho}}{\rho \cdot \ln(b/a)}$$

$$H = \frac{I \cdot \hat{\phi}}{2\pi\phi}$$

Because there is no surface current and surface charge distribution on the interfaces of  $z = z_0$  and  $z = z_0 + d$ , the electric field and magnetic field are continuous on the two interfaces. Therefore, the voltage and current are continuous on the two interfaces.

Contact the voltage and current expressions in the coaxial line, in medium 1 ( $z < z_0$ ) we have

$$\tilde{V}_1(z) = V_0^i(e^{-\gamma_1 z} + \Gamma e^{\gamma_1 z})$$

$$\tilde{I}_1(z) = \frac{V_0^i}{Z_1}(e^{-\gamma_1 z} - \Gamma e^{\gamma_1 z})$$

in medium 2 ( $z_0 < z < z_0 + d$ ), assume the magnitudes of voltage as  $B, C$ , it becomes to

$$\tilde{V}_2(z) = B e^{-\gamma_2 z} + C e^{\gamma_2 z}$$

$$\tilde{I}_2(z) = \frac{B}{Z_2} e^{-\gamma_2 z} - \frac{C}{Z_2} e^{\gamma_2 z}$$

in medium 3, it only has transmission wave, then it becomes to

$$\tilde{V}_3(z) = T e^{-\gamma_1 z}$$

$$\tilde{I}_3(z) = \frac{T}{Z_1} e^{-\gamma_1 z}$$

Then we have

$$\tilde{V}_1(z_0) = \tilde{V}_2(z_0)$$

$$\tilde{I}_1(z_0) = \tilde{I}_2(z_0)$$

$$\tilde{V}_2(z_0 + d) = \tilde{V}_3(z_0 + d)$$

$$\tilde{I}_2(z_0 + d) = \tilde{I}_3(z_0 + d)$$

Since only the ratio to the incident wave mode value is concerned, we set  $V_0^i = 1$ , then the four equations become to

$$\begin{aligned} e^{-\gamma_1 z_0} + \Gamma e^{\gamma_1 z_0} &= B e^{-\gamma_2 z_0} + C e^{\gamma_2 z_0} \\ \frac{e^{-\gamma_1 z_0} - \Gamma e^{\gamma_1 z_0}}{Z_1} &= \frac{B e^{-\gamma_2 z_0} - C e^{\gamma_2 z_0}}{Z_2} \\ B e^{-\gamma_2 (d+z_0)} + C e^{\gamma_2 (d+z_0)} &= T e^{-\gamma_1 (d+z_0)} \\ \frac{B e^{-\gamma_2 (d+z_0)} - C e^{\gamma_2 (d+z_0)}}{Z_2} &= \frac{T e^{-\gamma_1 (d+z_0)}}{Z_1} \end{aligned}$$

Following the tutorial, with the help of matlab, we can use  $\Gamma, T$  to express  $Z_2, \gamma_2$

$$\Gamma = \left\{ \begin{array}{l} -\frac{e^{-2\gamma_1 z_0} (\sigma_2 - \sigma_4 - \sigma_1 + \sigma_3)}{\sigma_2 - \sigma_4 + \sigma_1 - \sigma_3 + 2Z_1 Z_2 e^{\gamma_2 (d+z_0)} e^{-\gamma_2 z_0} + 2Z_1 Z_2 \sigma_5 e^{\gamma_2 z_0}} \\ \text{where} \\ \sigma_1 = Z_2^2 e^{\gamma_2 (d+z_0)} e^{-\gamma_2 z_0} \\ \sigma_2 = Z_1^2 e^{\gamma_2 (d+z_0)} e^{-\gamma_2 z_0} \\ \sigma_3 = Z_2^2 \sigma_5 e^{\gamma_2 z_0} \\ \sigma_4 = Z_1^2 \sigma_5 e^{\gamma_2 z_0} \\ \sigma_5 = e^{-\gamma_2 (d+z_0)} \end{array} \right.$$

$$T = \left\{ \begin{array}{l} \frac{4Z_1 Z_2 e^{\gamma_1 (d+z_0)} e^{-\gamma_1 z_0}}{Z_1^2 \sigma_2 e^{-\gamma_2 z_0} - Z_1^2 \sigma_1 e^{\gamma_2 z_0} + Z_2^2 \sigma_2 e^{-\gamma_2 z_0} - Z_2^2 \sigma_1 e^{\gamma_2 z_0} + 2Z_1 Z_2 \sigma_2 e^{-\gamma_2 z_0} + 2Z_1 Z_2 \sigma_1 e^{\gamma_2 z_0}} \\ \text{where} \\ \sigma_1 = e^{-\gamma_2 (d+z_0)} \\ \sigma_2 = e^{\gamma_2 (d+z_0)} \end{array} \right.$$

Based on that, we can find out the expression of  $Z_2, \gamma_2$  using  $\Gamma, T$ , that is

$$Z_2 = \left\{ \begin{array}{l} \left( \begin{array}{l} \frac{Z_1 (\sigma_6 + T^2 e^{2\gamma_1 z_0} + \sigma_1 - \Gamma^2 \sigma_6 e^{4\gamma_1 z_0})}{\sigma_4} - \sigma_2 \\ \frac{Z_1 (\sigma_6 + T^2 e^{2\gamma_1 z_0} - \sigma_1 - \Gamma^2 \sigma_6 e^{4\gamma_1 z_0})}{\sigma_4} - \sigma_2 \end{array} \right) \\ \text{where} \\ \sigma_1 = \sqrt{-(T e^{\gamma_1 z_0} - \sigma_5 + \sigma_3)(\sigma_5 + T e^{\gamma_1 z_0} + \sigma_3)(\sigma_5 + T e^{\gamma_1 z_0} - \sigma_3)(\sigma_5 - T e^{\gamma_1 z_0} + \sigma_3)} \\ \sigma_2 = \frac{Z_1 (-\sigma_6 e^{4\gamma_1 z_0} \Gamma^2 + e^{2\gamma_1 z_0} T^2 + \sigma_6)}{\sigma_4} \\ \sigma_3 = \Gamma \sigma_5 e^{2\gamma_1 z_0} \\ \sigma_4 = \sigma_6 e^{4\gamma_1 z_0} \Gamma^2 - 2\sigma_6 e^{2\gamma_1 z_0} \Gamma - e^{2\gamma_1 z_0} T^2 + \sigma_6 \\ \sigma_5 = e^{\gamma_1 (d+z_0)} \\ \sigma_6 = e^{2\gamma_1 (d+z_0)} \end{array} \right.$$

$$\gamma_2 = \left\{ \begin{array}{l} \left( \frac{\log \left( \frac{e^{-\gamma_1 (d+z_0)} e^{-\gamma_1 z_0} (\sigma_5 + T^2 e^{2\gamma_1 z_0} + \sigma_1 - \sigma_2)}{2T} \right)}{\log \left( \frac{e^{-\gamma_1 (d+z_0)} e^{-\gamma_1 z_0} (\sigma_5 + T^2 e^{2\gamma_1 z_0} - \sigma_1 - \sigma_2)}{2T} \right)} \right) \\ \frac{d}{d} \\ \text{where} \\ \sigma_1 = \sqrt{-(\sigma_4 - \sigma_6 + \sigma_3)(\sigma_6 + \sigma_4 + \sigma_3)(\sigma_6 + \sigma_4 - \sigma_3)(\sigma_6 - \sigma_4 + \sigma_3)} \\ \sigma_2 = \Gamma^2 \sigma_5 e^{4\gamma_1 z_0} \\ \sigma_3 = \Gamma \sigma_6 e^{2\gamma_1 z_0} \\ \sigma_4 = T e^{\gamma_1 z_0} \\ \sigma_5 = e^{2\gamma_1 (d+z_0)} \\ \sigma_6 = e^{\gamma_1 (d+z_0)} \end{array} \right.$$

## 2 Deduction other parameters

From the lots, we have

$$\begin{aligned} a &= 0.002m \\ b &= 0.003m \\ d &= 0.55m \\ \epsilon_1 &= 1 \\ f &= 10^9 Hz \end{aligned}$$

Then we have

$$\begin{aligned} \gamma_1 &= j * \omega * \sqrt{\mu_0 \epsilon_1 \epsilon_0} \\ Z_1 &= \frac{60 \ln(b/a)}{\sqrt{\epsilon_1 \epsilon_0}} \end{aligned}$$

Due to the property of coaxial lines, we have

$$Z = \frac{60 \ln(b/a)}{\sqrt{\epsilon}}$$

where  $\epsilon$  can be the complex permittivity. Therefore we have

$$\epsilon_2 = \frac{\mu_0}{Z_2^2 (2\pi / \ln(b/a))^2}$$

## 3 Calculation and Simulation 1

With the first group of coefficients, we have

$$\begin{aligned} \Gamma_1 &= -0.18367707842865 - 0.336458601842074i \\ T_1 &= -0.73358668120551 - 0.561167203468699i \end{aligned}$$

$$Z_2 = 8.103846546579724 - 0.000090357978102i \, \Omega$$

$$\gamma_2 = 0.000049836674385 - 51.364551044460249i \, m^{-1}$$

Actually, there are two solutions for  $Z_2$  and  $\gamma_2$ , however, the real part of the other solution for  $Z_2$  is negative, which is impossible in real world, so we discard that solution.

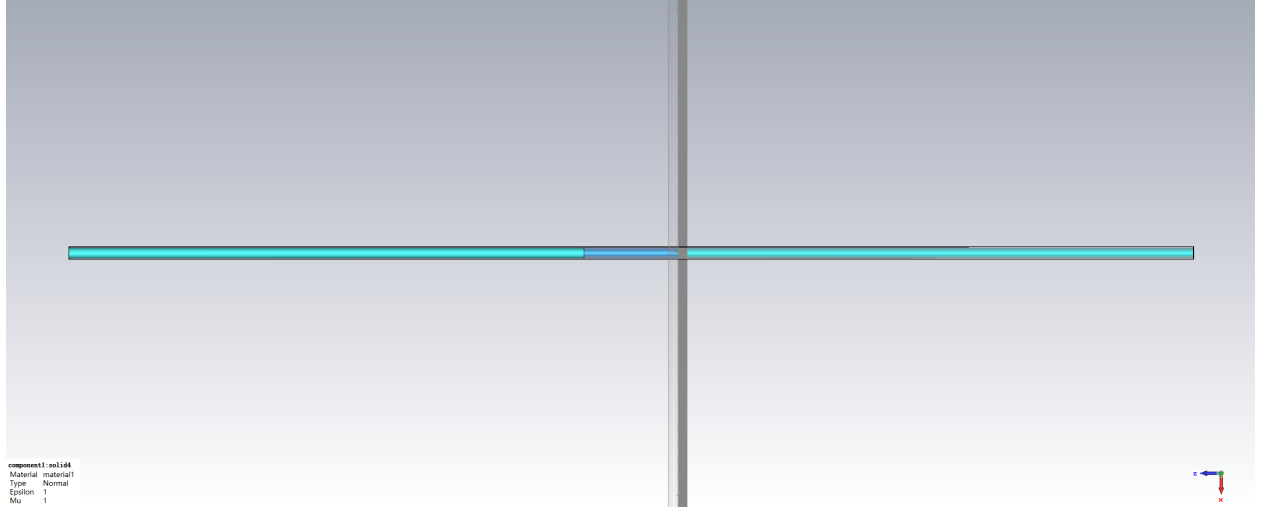
Then by further calculating, we have

$$\epsilon_2 = 7.968465274556406 * 10^{-11} + 1.776969508829155 * 10^{-15}i$$

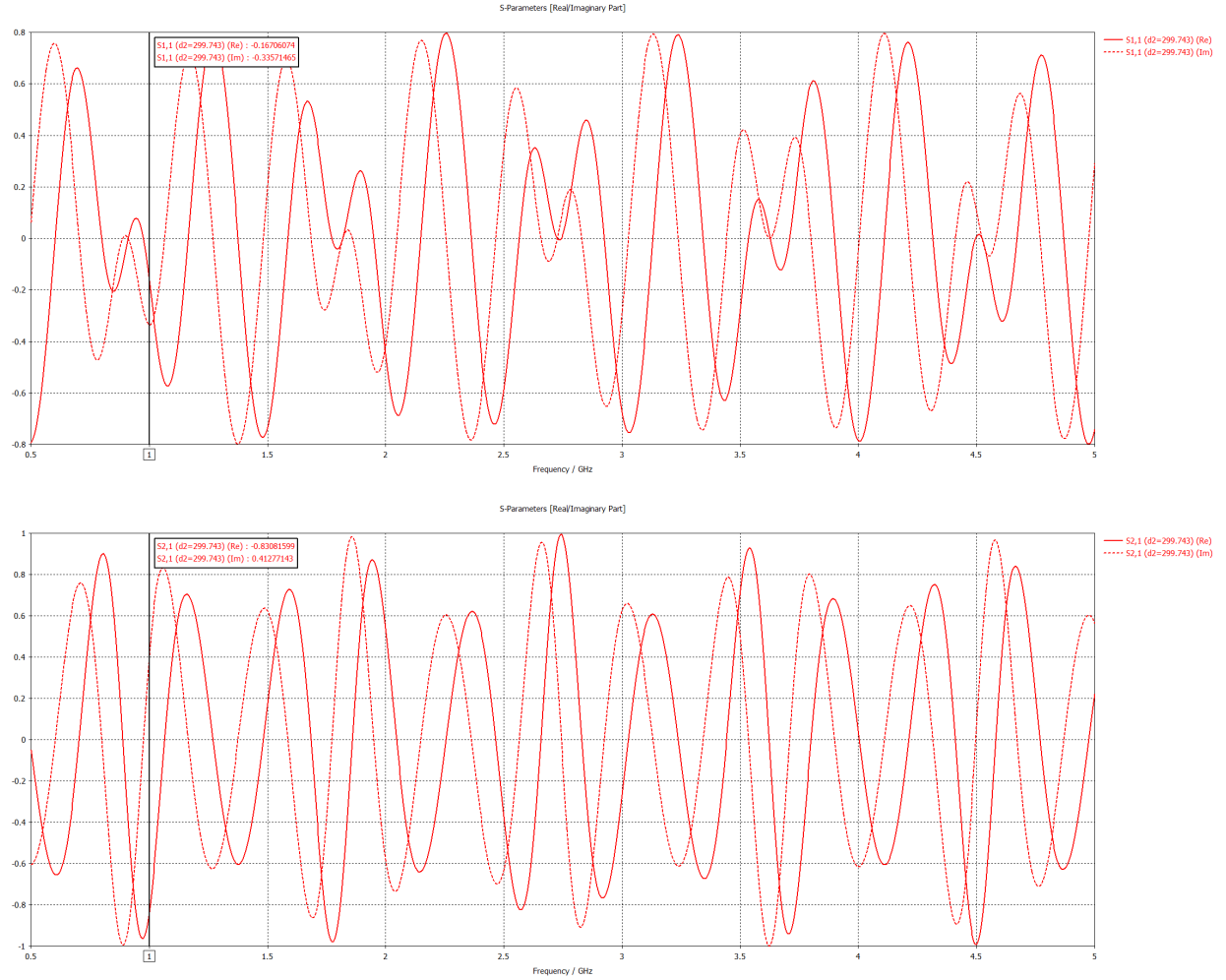
$$\epsilon_{r2} = \frac{real(\epsilon_2)}{\epsilon_0} \approx 8.999656929862034$$

$$\sigma = -imag(\epsilon_2) * \omega \approx -1.116502870918147 * 10^{-5} \approx 0S/m$$

Set the length of medium 1 of both two sides as  $d_2 = \lambda_1$ , which allows us ignore the phase difference caused by medium 1, then we have the following model



By simulation, we have thus we obtain



$$S_{11} = -0.16706074 - 0.33571465i \approx \Gamma_1$$

$$S_{21} = T_1 = -0.83081599 + 0.41277143i$$

$$\approx T_1 * e^{-\gamma_1 d} = -0.810675666421054 + 0.442564134294426i$$

which is within the allowable error range.

## 4 Calculation and Simulation 2

With the second group of coefficients, we have

$$\Gamma_2 = -0.195677924334674 + 0.116950135912389i$$

$$T_2 = -0.134265494843727 - 0.7238751863504i$$

However, the  $T$  we assumed in the equation is  $\frac{\tilde{V}^r(z_0 + d)}{\tilde{V}^i(z_0)}$ , therefore, we should use  $T_1 \cdot e^{-\gamma_1 d}$  instead of  $T_1$  to solve  $\epsilon_2$ .

$$Z_2 = 9.384395457555984 + 0.700022010466726i \, \Omega$$

$$\gamma_2 = -4.026968258633184 - 53.993480480585660i \, m^{-1}$$

Actually, there are two solutions for  $Z_2$  and  $\gamma_2$ , however, the real part of the other solution for  $Z_2$  is negative, which is impossible in real world, so we discard that solution.

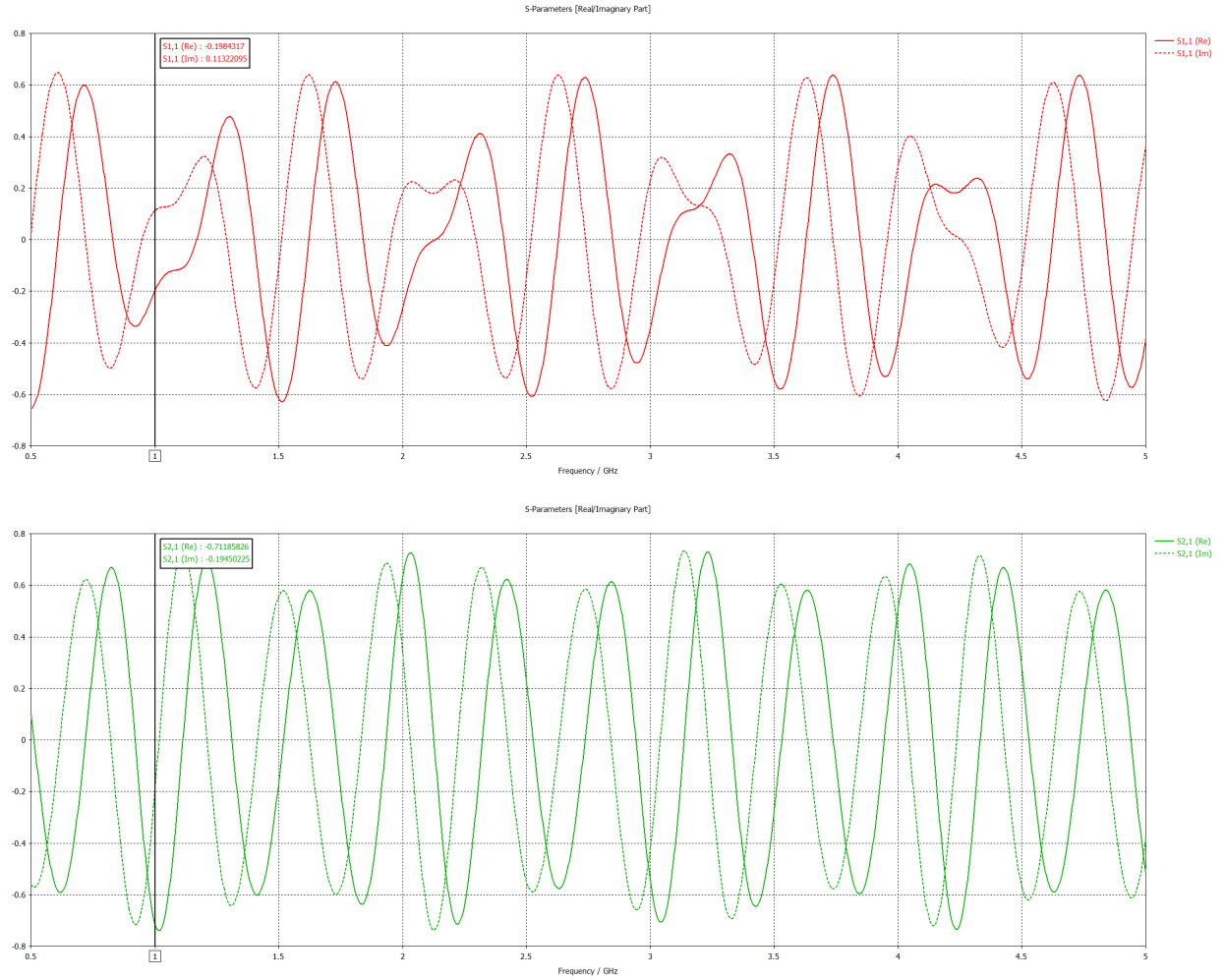
Then by further calculating, we have

$$\epsilon_2 = 5.843882975868956 * 10^{-11} - 8.767185499046495 * 10^{-12}i$$

$$\epsilon_{r2} = \frac{real(\epsilon_2)}{\epsilon_0} \approx 6.600134418482435$$

$$\sigma = -imag(\epsilon_2) * \omega \approx 0.055085851112927$$

By simulation, we have thus we obtain



$$S_{11} = -0.1984317 + 0.11322095i \approx \Gamma_2$$

$$S_{21} = -0.71185826 - 0.19450225i$$

$$\approx T_2 * e^{-\gamma_1 d} = -0.716040335476707 - 0.171197974549531i$$

which is within the allowable error range.