

广相复习

1.在假定空间平坦及均匀各向同性的情况下,时空度规可写为

$$ds^2 = -dt^2 + a^2(t)(dx^2 + dy^2 + dz^2)$$

$a(t)$ 为宇宙尺度因子,是时间的函数。在只存在暗能量的情况下,引力场方程为 $R_{uv} = -\Lambda g_{uv}$,其中, $\Lambda > 0$ 为常数。给定条件 $a(T) = a_0$,求解 $a(t)$,即尺度因子随时间的演化。

解:涉及到只求里奇张量时,可以不用求黎曼张量,具体解法如下:

$$R_{\mu\lambda\nu}^{\lambda} = -\Gamma_{\mu\lambda,\nu}^{\lambda} + \Gamma_{\mu\nu,\lambda}^{\lambda} - \Gamma_{\mu\lambda}^{\sigma}\Gamma_{\sigma\nu}^{\lambda} + \Gamma_{\mu\nu}^{\sigma}\Gamma_{\sigma\lambda}^{\lambda}$$

$$g_{tt} = -1, g_{ii} = a^2$$

$$\Gamma_{ti}^i = \frac{\dot{a}}{a}, \Gamma_{ii}^t = a\dot{a}$$

其中不为0的分量有: R_{tt}, R_{ii}

$$R_{tt} = R_{txt}^x + R_{tyt}^y + R_{tzt}^z = \frac{3\ddot{a}}{a}$$

$$R_{ii} = -(\ddot{a}a + 2\dot{a}^2)$$

$$\text{对应分量相同: } \frac{3\ddot{a}}{a} = \Lambda, -(\ddot{a}a + 2\dot{a}^2) = -\Lambda a^2, a = a_0 e^{\sqrt{\Lambda}(t-T)/3}$$

2.对于三位度规球: $ds^2 = d\theta^2 + \sin^2\theta d\varphi^2 + \cos^2\theta d\psi^2$,求联络,曲率张量,里奇张量,标量曲率。

解:首先求得度规为: $g_{\theta\theta} = 1, g_{\varphi\varphi} = \sin^2\theta, g_{\psi\psi} = \cos^2\theta$,然后计算出不为0的联络为 $\Gamma_{\theta\varphi}^{\varphi} = \frac{\cos\theta}{\sin\theta}, \Gamma_{\theta\psi}^{\psi} = -\frac{\sin\theta}{\cos\theta}, \Gamma_{\varphi\varphi}^{\theta} = -\sin\theta \cos\theta, \Gamma_{\psi\psi}^{\theta} = \sin\theta \cos\theta$

$$R_{\lambda\mu\nu}^{\rho} = -\Gamma_{\lambda\mu,\nu}^{\rho} + \Gamma_{\lambda\nu,\mu}^{\rho} - \Gamma_{\lambda\mu}^{\sigma}\Gamma_{\sigma\nu}^{\rho} + \Gamma_{\lambda\nu}^{\sigma}\Gamma_{\sigma\mu}^{\rho}$$

根据对称性和 $\mu \neq \nu$ 不为0的黎曼张量分量为: $R_{\varphi\theta\varphi\theta}, R_{\psi\theta\psi\theta}, R_{\varphi\psi\varphi\psi}, R_{\varphi\psi\varphi\theta}, R_{\varphi\psi\psi\theta}, R_{\varphi\theta\psi\theta}$

这里的上下指标变换需要乘以度规!!!

$$R_{\theta\psi\theta}^{\psi} = 1, R_{\theta\varphi\theta}^{\varphi} = 1, R_{\psi\varphi\psi}^{\varphi} = \cos^2\theta$$

$$R_{\psi\varphi\theta}^{\varphi} = 0, R_{\theta\psi\theta}^{\varphi} = 0, R_{\psi\psi\theta}^{\varphi} = 0$$

将所有的指标写成下指标的形式:

$$R_{\varphi\theta\varphi\theta} = \sin^2\theta, R_{\psi\theta\psi\theta} = \cos^2\theta, R_{\varphi\psi\varphi\psi} = \sin^2\theta \cos^2\theta$$

$$R_{\varphi\psi\varphi\theta} = R_{\varphi\psi\psi\theta} = R_{\varphi\theta\psi\theta} = 0$$

根据里奇张量的公式：

$$\begin{aligned} R_{\theta\theta} &= g^{\varphi\varphi} R_{\varphi\theta\theta\varphi} + g^{\psi\psi} R_{\psi\theta\theta\psi} = -2 \\ R_{\psi\psi} &= g^{\theta\theta} R_{\theta\psi\psi\theta} + g^{\varphi\varphi} R_{\varphi\psi\psi\varphi} = -2 \cos^2 \theta \\ R_{\varphi\varphi} &= g^{\theta\theta} R_{\theta\varphi\varphi\theta} + g^{\psi\psi} R_{\psi\varphi\varphi\psi} = -2 \sin^2 \theta \end{aligned}$$

因此标量曲率为：注意这里还是要乘以度规!!!

$$R = g^{\theta\theta} R_{\theta\theta} + g^{\psi\psi} R_{\psi\psi} + g^{\varphi\varphi} R_{\varphi\varphi} = -6$$

3.引入场强张量 $F^{\mu\nu}, F^{\mu\nu} = -F^{\nu\mu}, \mu, \nu = 0, 1, 2, 3$, 其中 $F^{0i} = E^i, F^{ij} = \varepsilon^{ijk} B_k, i, j, k = 1, 2, 3$, 证明麦克斯韦方程组

$$\begin{aligned} \varepsilon^{ijk} \partial_j B_k - \partial_0 E^i &= 4\pi J^i & \partial_i E^i &= 4\pi J^0 \\ \varepsilon^{ijk} \partial_j E_k + \partial_0 B^i &= 0 & \partial_i B^i &= 0 \end{aligned}$$

可写为

$$\partial_\mu F^{\mu\nu} = -4\pi J^\nu \quad \partial_\mu F_{\nu\lambda} + \partial_\nu F_{\lambda\mu} + \partial_\lambda F_{\mu\nu} = 0$$

并得出电荷守恒方程 $\partial_\mu J^\mu = 0$. (注意: $F_{\mu\nu} = \eta_{\mu\nu} \eta_{\rho\sigma} F^{\lambda\sigma}$, 所以 $F_{0i} = \eta_{00} F^{0i} = -F^{0i}, F_{ij} = F^{ij}$)

解: ijk都是表示的分量, $i, j, k = 1, 2, 3$

$$\begin{aligned} \text{由前两个式子可以得出: } \partial_j F^{ij} - \partial_0 F^{0i} &= 4\pi J_i \\ \partial_0 F^{i0} = 0, \partial_0 F^{00} &= 0 \\ \partial_i F^{0i} &= 4\pi J^0 \end{aligned}$$

由上述式子可以推出要证明的第一个式子

$$\text{乘以度规进行上下指标的置换: } -F_{0i} = E_i, B^i = \frac{1}{2} \varepsilon^{ijk} F_{jk}$$

$$\begin{aligned} \text{代入第三个式子 } \frac{1}{2} \varepsilon^{ijk} (2\partial_j F_{k0} + \partial_0 F_{jk}) &= 0 \\ \frac{1}{2} \varepsilon^{ijk} (\partial_k F_{0j} + \partial_j F_{k0} + \partial_0 F_{jk}) &= 0 \end{aligned}$$

$$\text{上式对jk求和可以得出: } \partial_k F_{0j} + \partial_j F_{k0} + \partial_0 F_{jk} = 0$$

$$\text{对第四个式子: } \partial_i F_{jk} + \partial_k F_{ij} + \partial_j F_{ki} = 0 \quad (1)$$

上面已经得出来对于一个指标一个0和没有0的情况, 由于三个指标都需要变成 $\mu\nu\lambda$, 所以还有两个0和三个0的情况, 上面是式子是对i 指标的情况, 其他的指标和i指标是类似的, 因此只需要对一个指标求出来了, 其他指标也就求出来了。

对最后一个式子的证明,利用对称性:

$$\begin{aligned} \partial_\nu F^{\nu\mu} &= -4\pi J^\mu \\ \partial_\mu \partial_\nu F^{\nu\mu} &= -4\pi \partial_\mu J^\mu \\ &= -\partial_\mu \partial_\nu F^{\mu\nu} \\ &= -\partial_\nu \partial_\mu F^{\nu\mu} \end{aligned}$$

4.电磁场和电流耦合的作用量为

$$S[A_\mu] = \int d^4x \left(-\frac{1}{4} F^{\mu\nu} F_{\mu\nu} + A_\mu J^\mu \right),$$

其中 $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$, $J^\mu = (\rho, \vec{J})$ 为给定的四维电流密度, 对 A_μ 作变分, 得出电磁场运动方程 $\partial_\mu F^{\mu\nu} = -J^\nu$

解: 注意这里的电磁场张量是反对称的, 并且这里有 $\nu\mu$ 代表求和

$$\delta(F^{\nu\mu} F_{\nu\mu}) = 2F^{\mu\nu} \delta F_{\nu\mu}$$

$$F^{\mu\nu} \delta(\partial_\mu A_\nu - \partial_\nu A_\mu) = F^{\mu\nu} \partial_\mu \delta A_\nu + F^{\nu\mu} \partial_\nu \delta A_\mu = 2F^{\mu\nu} \partial_\mu \delta A_\nu$$

$$\begin{aligned} \delta S[A_\nu] &= \int d^4x (-F^{\mu\nu} \partial_\mu \delta A_\nu + J^\nu \delta A_\nu) \\ &= \int d^4x (-\partial_\mu F^{\mu\nu} \delta A_\nu + \delta A_\nu \partial_\mu F^{\mu\nu} + J^\nu \delta A_\nu) \\ &= \int d^4x (\partial_\mu F^{\mu\nu} + J^\nu) \delta A_\nu = 0 \end{aligned}$$