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博士学位论文



基于 ATLAS 探测器上 ZZ 玻色子到全轻子
通道的衰变事例对标准模型 ZZ 过程和寻找
重共振态衰变到双 Z 玻色子过程的研究

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**Studies of Standard Model ZZ
Production and Search for heavy ZZ
resonances in Purely Leptonic
Decay with ATLAS Detector**

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摘要

论文介绍了本人在粒子物理领域基于大型强子对撞机（LHC）上 ATLAS 实验做的研究工作。大型强子对撞机是当今世界上最大的、能量最高的对撞机，是建立在理论和实验之间的重要桥梁。而 ATLAS 实验是 LHC 上的一个通用粒子探测器实验，同时也是体积最大的探测器。基于 ATLAS 实验在 LHC 上收集到的亮度为 139 fb^{-1} 能量为 13 TeV 的质子-质子对撞数据，本文重点介绍了两个 Z 玻色子衰变到四轻子末态过程的一系列研究。包括，标准模型（SM）下 ZZ 到四轻子过程截面的测量、矢量玻色子散射（VBS）过程在 ZZ 到四轻子末态的观测，和寻找重共振态衰变到 ZZ 到四轻子末态的过程。

ZZ 到四轻子过程截面的测量结果为 $\sigma_{ZZjj}^{tot} = 1.27 \pm 0.12(\text{stat}) \pm 0.02(\text{theo}) \pm 0.07(\text{exp}) \pm 0.01(\text{bkg}) \pm 0.03(\text{lumi})$ ，总体相对误差为 11%。在误差范围内，该结果和标准模型预言值 $1.14 \pm 0.04(\text{stat}) \pm 0.20(\text{theo})$ 相吻合。同时，在两个 Z 玻色子伴随着两个喷注（jets）末态的电弱相互作用过程的寻找中，我们观测到偏离本底假说超过 5 倍标准差 (5.5σ) 的明显偏差。在此基础上，本文也介绍了对于下一代高亮度大型强子对撞机（HL-LHC）在两个 Z 玻色子伴随着两个喷注（jets）末态的电弱相互作用过程的模拟预言。

另一方面，本文介绍了在一对 Z 玻色子衰变至四轻子末态过程中寻找重共振态的实验。根据不同的信号模型，寻找的粒子质量区间设置在 200 GeV 到 2000 GeV 之间。基于该测量结果，没有证据可以证明重共振态的存在。因此，研究给出了基于不同信号模型的截面上限，包括在不同衰变宽度假说下自旋为 0 的共振态，以及基于 Randall–Sundrum 模型的自旋为 2 的引力子（graviton）。在该分析中，我们认为，信号主要可通过 gluon-gluon Fusion (ggF) 和 Vector Boson Fusion (VBF) 过程产生。在自旋为 0 的窄衰变宽度模型下，我们对 ggF 和 VBF 两个过程都进行了研究。而对于大宽度模型，由于在质量很高的区间分辨率很差以及 VBF 过程的统计量太小等客观原因，只对 ggF 过程进行了研究。对于自旋为 2 的模型，实验给出了 Randall–Sundrum 模型的引力子的理论质量下限，为 1500 GeV 。

28

ABSTRACT

29 This dissertation presents my research in the field of Particle Physics with the ATLAS
30 experiment at the Large Hadron Collider (LHC). The LHC is the world's largest and most
31 powerful collider, and it was built as a bridge between the theories and the experiment.
32 The ATLAS experiment is a general-purpose particle detector experiment with the largest
33 volume at the LHC. This dissertation focus on the studies with two Z bosons production
34 decaying into $\ell\ell\ell'\ell'$ final state, where ℓ stands for electron or muon, using 139 fb^{-1} of
35 13 TeV proton-proton (pp) collision data collected by ATLAS experiment at the LHC.
36 The ZZ production in $\ell\ell\ell'\ell'$ channel provides a most clean and sensitive tool to test
37 the Standard Model (SM) at the energy frontier and to study the *Higgs* physics. Studies
38 including the measurement on SM $ZZjj$ production cross section, the observation of
39 Vector Boson Scattering (VBS) process as well as the searches of heavy resonances in
40 ZZ production decaying into $\ell\ell\ell'\ell'$ final state are reported in this dissertation.

41 The fiducial cross section for SM $ZZjj$ production is measured to be $\sigma_{ZZjj}^{tot} = 1.27 \pm$
42 $0.12(\text{stat}) \pm 0.02(\text{theo}) \pm 0.07(\text{exp}) \pm 0.01(\text{bkg}) \pm 0.03(\text{lumi})[\text{pb}]$ with a total relative
43 uncertainty of 11% for the $\ell\ell\ell'\ell'$ final state, and found to be compatible with the SM
44 prediction of $1.14 \pm 0.04(\text{stat}) \pm 0.20(\text{theo})[\text{pb}]$. The electroweak production of two jets
45 in association with a Z -boson pair (EW- $ZZjj$) is observed with a significant deviation
46 from the background-only hypothesis corresponding to a statistical significance of 5.5σ .
47 Following with the observation, the prospect study for the EW- $ZZjj$ production at the
48 High luminosity LHC (HL-LHC) using 3000 fb^{-1} simulated pp collision data at a centre-
49 of-mass energy of 14 TeV is presented, with a expected significance of around 7σ .

50 A search for heavy resonances decaying into a pair of Z bosons to $\ell\ell\ell'\ell'$ final state is
51 also conducted in this dissertation. Different mass ranges for the hypothetical resonances
52 are considered, depending on the signal models and spanning between 200 GeV and
53 2000 GeV. Data is found to agree with a background-only hypothesis, thus, the results
54 are interpreted as upper limits on production cross section for sevaral different models,
55 including heavy Higgs like (spin-0) narrow-width approximation (NWA) and large-width
56 approximation (LWA), as well as the Randall–Sundrum model with a graviton excitation
57 spin-2 resonance (RSG). The signal is assumed to generate dominatly via gluon-gluon
58 Fusion (ggF) production mode and Vector Boson Fusion (VBF) production mode. Both
59 ggF and VBF channels are studied in NWA, while for LWA, only ggF channel is studied
60 due to worse resolution in higher mass region and the lack of statistic for VBF process. In

Abstract

- 61 addition, mass of RS Graviton is constrained, $m(G_{KK}) < 1500$ GeV is excluded at 95%
- 62 CL by $ZZ \rightarrow \ell\ell\ell'\ell'$ analysis.

63

64

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590

Dedicated to my mother!

591

献给我的母亲！

592

593

594

Chapter 1 Introduction

595 The goal of particle physics is to understand how our universe works at its most fundamental level. It can be accomplished by pursuing the mysteries of the basic construction of
596 matter and energy, probing the interactions between elementary particles, and exploring
597 the basic nature of space and time itself.

599 **Elementary particles**

600 From around the 6th century BC, ancient Greek philosophers Leucippus, Democritus, and Epicurus brought up a philosophical idea that everything is composed of “uncuttable” elementary particles. In the 19th century, John Dalton, through his work
601 on stoichiometry, concluded that each element of nature was composed of a single, unique
602 type of particle. The particle was named as “atom” after the Greek word atomos, with
603 the meaning of “indivisible”. However this Dalton’s atom theory was strongly chal-
604 lenged later. Near the end of 19th century, physicists discovered that Dalton’s atoms
605 are not, in fact, the fundamental particles of nature, but conglomerates of even smaller
606 particles. Electron was discovered by J. J. Thomson in 1897, and then its charge was
607 carefully measured by Robert Andrews Millikan and Harvey Fletcher in their “oil drop
608 experiment” of 1909. In early 20th-century, Rutherford’s “gold foil experiment” showed
609 that the atom is mainly empty space, with almost all its mass concentrated in a tiny posi-
610 tively charge atomic nucleus. Then the discoveries of anti-particles (the positron in 1932)
611 and other particles (e.g. the muon in 1936) show that more discoveries could be expected
612 in future experiments.

615 Starting from 1950s, more accelerator facilities were put into service. Throughout
616 the 1950s and 1960s, a bewildering variety of particles were found in collisions of parti-
617 cles from increasingly high-energy beams. It was referred to informally as the “particle
618 zoo”. In 1964, the quark model was independently proposed by physicists Murray Gell-
619 Mann and George Zweig, and experimentally confirmed of their existence in mid-1970s.
620 In 1970s, the establishment of quantum chromodynamics (QCD) postulated the funda-
621 mental strong interaction, experienced by quarks and mediated by gluons.

622 The well-known Standard model (SM) was developed in stages throughout the latter
623 half of the 20th century. Since then, confirmation of the top quark (1995), the tau neu-
624 trino (2000), and the Higgs boson (2012) have added further credence to the Standard
625 Model. Now, the quarks, leptons and gauge bosons are the elementary constituents in a
626 framework of Standard Model of particle physics, which theoretically describes three of

627 the four known fundamental forces (the electromagnetic, weak, and strong interactions,
628 and not including the gravitational force) in the universe, as well as classifies all
629 known elementary particles.

630 **Higgs mechanics and electroweak symmetry breaking**

631 In 1961, Sheldon Glashow, Steven Weinberg and Abdus Salam together brought for-
632 ward a unified electroweak theory to combine the electromagnetic and weak interactions.
633 In the standard model, if under the condition that the energy is high enough but elec-
634 troweak symmetry is unbroken, all elementary particles are massless. But measurements
635 show the fact that the W and Z bosons actually have masses. Later on, the Higgs mechan-
636 ics resolves this conundrum. The simplest description of the mechanism adds a Higgs
637 field that permeates all space to the Standard Model. Below some extremely high energy,
638 the field causes spontaneous symmetry breaking during interactions. All massive parti-
639 cles in the Standard Model, including the W and Z bosons, interact with Higgs boson to
640 acquire their mass.

641 Over the past few decades, with the combination of electroweak theory, Higgs me-
642 chanics and strong interactions has been widely accepted. But the Higgs boson, which is
643 essential to explain the mechanics of the property “mass” for gauge bosons and fermions,
644 had been the final missing piece in the Standard Model of particle physics at that time. The
645 mass of Higgs boson was not specifically predicted, and it has been searched in several
646 large experiments (eg. LEP at CERN, Tevatron at Fermilab, and LHC at CERN) with
647 different energy. In 2012, the discovery of Higgs boson was finally announced by the
648 ATLAS and CMS collaborations at the Large Hadron Collider (LHC) with its mass round
649 125 GeV. Peter Higgs and Francois Englert were award the 2013’s Nobel Prize in Physics
650 for their theoretical discovery of a mechanism that contributes to our understanding of the
651 origin of mass of subatomic particles.

652 **Contents of this thesis**

653 This dissertation is organized as follows. Section 2 briefly introduces the Standard
654 Model of particle physics, the Higgs mechanism related to the dissertation and the LHC
655 phenomenology. Section 3 gives an overview of the LHC and the ATLAS detector. The
656 detector simulation and the reconstruction of physics objects are described in section 4.
657 And then section 5 focuses on the Standard model ZZ production cross section measure-
658 ment in $ZZ \rightarrow \ell\ell\ell'\ell'$ channel, where ℓ stands for electron or muon, and the observation
659 of its electroweak component as well as its further prospects in High luminosity LHC
660 (HL-LHC). Section 6 present the search of possible heavy resonances in $ZZ \rightarrow \ell\ell\ell'\ell'$
661 channel. In the end, section 7 gives the summary and outlook for future physics in LHC.

662

Chapter 2 Theory

663

2.1 The Standard Model of Particle Physics

664 The standard model (SM) reflects our current understanding of elementary particles
 665 and several basic interactions. It is a gauge quantum field theory containing the internal symmetries of the unitary product group $SU(3) \times SU(2) \times U(1)$, in which the color
 666 group $SU(3)$ presents the strong interaction, and $SU(2) \times U(1)$ describes the electroweak
 667 interactions. Over the past decades, the SM has been widely tested through various experiments with extremely high precision.
 668

670

2.1.1 Elementary particles in the Standard Model

671 The elementary particles in SM can be classified into 3 class: *fermions*, *gauge bosons* and the *Higgs boson* as shown in Figure 2.1.

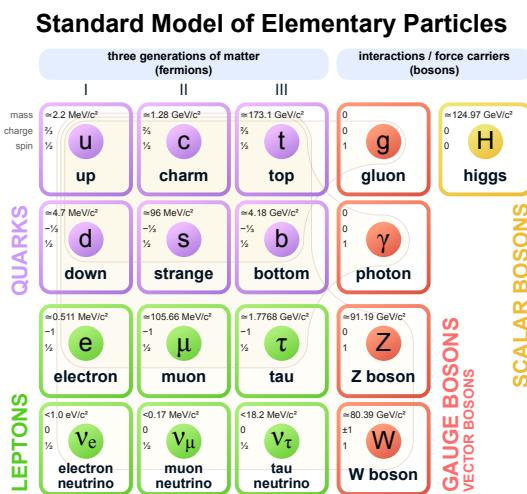


Fig. 2.1 The elementary particles of the Standard Model.

672

673 **Fermions** The Standard Model includes 12 elementary particles of spin- $\frac{1}{2}$ obeying the
 674 Fermi-Dirac statistics, known as fermions. They are classified into two types: *leptons* and
 675 *quarks* according to their interactions. The *leptons* include three generations: electron
 676 (e) and electron neutrino (ν_e); muon (μ) and muon neutrino (ν_μ); tau (τ) and tau neutrino
 677 (ν_τ). The e , μ and τ carry electric charge of -1 and three neutrinos are electrically neutral.
 678 All the leptons can participate in electroweak interactions. Also there are three generations
 679 of *quarks*: up (u) and down (d); charm (c) and strange (s); top (t) and bottom (b). The
 680 defining property of the quarks is that they carry color charge (while leptons don't), and
 681 hence interact via the strong interaction, letting them to be strongly bound from one to

another, forming color-neutral composite particles (known as hadrons) containing either a quark and an antiquark (mesons) or three quarks (baryons). In the meantime, u , c and t -quark carry electric charge of $2/3$, and d , s and b -quark carry electric charge of $-1/3$. Hence they interact via all three interactions described in SM. Each fermion also has a corresponding antiparticles.

Gauge bosons act as force carriers that mediate the strong, weak, and electromagnetic interactions in SM. They are spin-1 particles obeying the Bose-Einstein statistics. There are three types of gauge bosons:

- The eight massless *gluons* mediate the strong interactions between color charged particles (quarks).
- The massless *photons* mediate the electromagnetic force between electrically charged particles.
- The W^+ , W^- and Z bosons mediate the weak interactions between particles of different flavors (all quarks and leptons). All these three bosons are massive, the W^\pm carries an electric charge of $+1$ and -1 and couples to the electromagnetic interaction while Z boson is electrically neutral.

Figure 2.2 shows the Feynman diagrams of corresponding interactions in SM.

Higgs boson is a massive scalar elementary particle with spin-0. It plays a unique role in the SM by explaining the origin of masses of massive gauge bosons (W^\pm and Z) and fermions. And it is the last discovered particle in SM.

2.1.2 Electroweak theory

The electroweak interaction is the unified description of two of the four known fundamental interactions of nature: electromagnetism and the weak interaction. It is based on the gauge group of $SU(2)_L \times SU(1)_Y$, in which L is the left-handed fields and Y is the weak hypercharge^[1]. It follows the Lagrangian of

$$L_{EW} = L_{gauge} + L_{Higgs} + L_{fermion} + L_{Yukawa} \quad (2.1)$$

L_{gauge} is the **gauge term** part

$$L_{gauge} = -\frac{1}{4} W_{\mu\nu}^i W^{\mu\nu i} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} \quad (2.2)$$

where W_μ^i and B_μ present the $SU(2)_L$ and $SU(1)_Y$ gauge fields respectively, with the

Standard Model Interactions (Forces Mediated by Gauge Bosons)

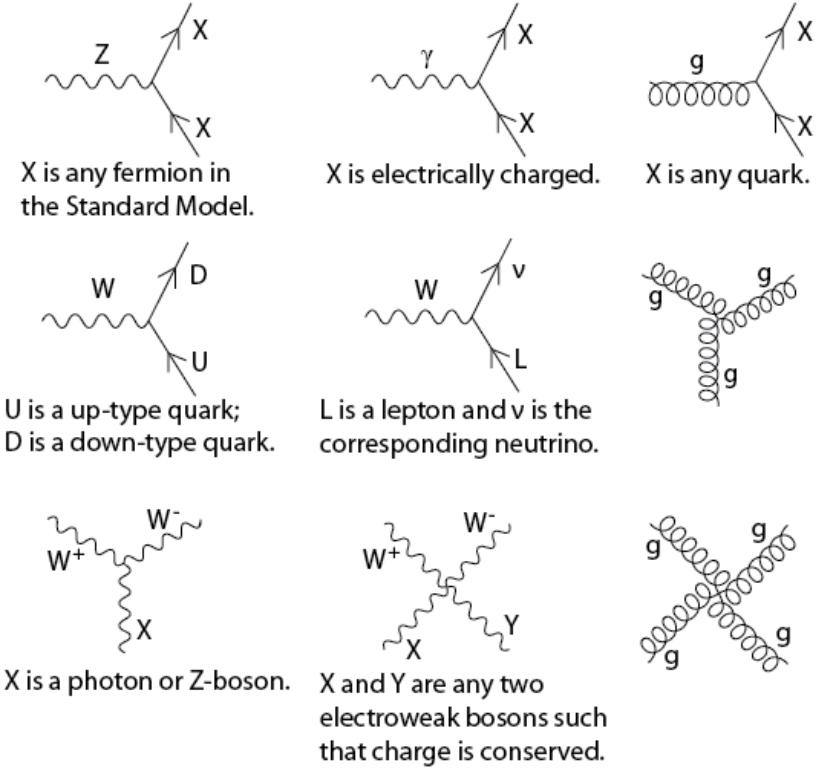


Fig. 2.2 The Feynman diagrams of interactions that form the basis of the standard model.

709 corresponding field strength tensors of

$$\begin{aligned} B_{\mu\nu} &= \partial_\mu B_\nu - \partial_\nu B_\mu \\ W_{\mu\nu}^i &= \partial_\mu W_\nu^i - \partial_\nu W_\mu^i - g \epsilon_{ijk} W_\mu^j W_\nu^k \end{aligned} \quad (2.3)$$

In the equations above, g is the $SU(2)_L$ gauge coupling and ϵ_{ijk} is the totally antisymmetric tensor. The gauge Lagrangian has three and four-point self interactions of W^i , which result in triple and quartic gauge boson couplings.

713 The second term of the Lagrangian is the **scalar part**:

$$L_{Higgs} = (D^\mu \phi)^\dagger D_\mu \phi - V(\phi) \quad (2.4)$$

714 where $\phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$ is a complex Higgs scalar, and $V(\phi)$ is the Higgs potential which is
 715 restricted into the form of

$$V(\phi) = +\mu^2 \phi^\dagger \phi + \lambda (\phi^\dagger \phi)^2 \quad (2.5)$$

⁷¹⁶ due to the combination of $SU(2)_L \times SU(1)_Y$ invariance and renormalizability. In Eq. 2.5,

717 μ is a mass-dependent parameter and λ is the quartic Higgs scalar coupling, which rep-
 718 resents a quartic self-interaction between the scalar fields. When $\mu^2 < 0$, there will be
 719 spontaneous symmetry breaking (more details in section 2.1.3). To maintain vacuum sta-
 720 bility, $\lambda > 0$ is required. And in Eq. 2.4, the gauge covariant derivative is defined as

$$D_\mu \phi = \left(\partial_\mu + ig \frac{\tau^i}{2} W_\mu^i + \frac{ig'}{2} B_\mu \right) \phi \quad (2.6)$$

721 in which τ^i represents the Pauli matrices, and g' is the $U(1)_Y$ gauge coupling. The square
 722 of the covariant derivative results in three and four -point interactions between the gauge
 723 and scalar fields.

724 The third term of the Lagrangian is the **fermion part**

$$\begin{aligned} L_{fermion} = & \sum_{m=1}^F (\bar{q}_{mL}^0 \gamma_\mu D_\mu q_{mL}^0 + \bar{l}_{mL}^0 \gamma_\mu D_\mu l_{mL}^0 + \bar{u}_{mR}^0 \gamma_\mu D_\mu u_{mR}^0 \\ & + \bar{d}_{mR}^0 \gamma_\mu D_\mu d_{mR}^0 + \bar{e}_{mR}^0 \gamma_\mu D_\mu e_{mR}^0 + \bar{v}_{mR}^0 \gamma_\mu D_\mu v_{mR}^0) \end{aligned} \quad (2.7)$$

725 In Eq. 2.7, m is the family index of fermions, F is the number of families. The subscripts
 726 $L(R)$ stand for the left (right) chiral projection $\psi_{L(R)} \equiv (1 \mp \gamma_5) \psi / 2$.

$$q_{mL}^0 = \begin{pmatrix} u_m^0 \\ d_m^0 \end{pmatrix}_L \quad l_{mL}^0 = \begin{pmatrix} v_m^0 \\ e_m^{-0} \end{pmatrix}_L \quad (2.8)$$

727 are the $SU(2)$ doublets of left-hand quarks and leptons, while u_{mR}^0 , d_{mR}^0 , e_{mR}^{-0} and v_{mR}^0 are
 728 the right-hand singlets.

729 The last term in Eq. 2.1 is **Yukawa term**

$$\begin{aligned} L_{Yukawa} = & - \sum_{m,n=1}^F [\Gamma_{mn}^u \bar{q}_{mL}^0 \tilde{\phi} u_{nR}^0 + \Gamma_{mn}^d \bar{q}_{mL}^0 \phi d_{nR}^0 \\ & + \Gamma_{mn}^e \bar{l}_{mL}^0 \phi e_{nR}^0 + \Gamma_{mn}^v \bar{l}_{mL}^0 \tilde{\phi} v_{nR}^0] + h.c. \end{aligned} \quad (2.9)$$

730 the matrices Γ_{mn} refer to the Yukawa couplings between single Higgs doublet (ϕ) and the
 731 various flavors of quarks (m) and leptons (n).

732 2.1.3 Higgs mechanism and Electroweak symmetry breaking

733 As shown in previous subsection, the Lagrangian L_{gauge} does not involve any mass
 734 term due to the requirement of gauge invariance. So all the W and B bosons should be
 735 massless. But experimental observations show that the gauge bosons are massive. There-
 736 fore, the gauge invariance must be broken spontaneously. The Higgs field is introduced
 737 to break the $SU(2)_L \times U(1)_Y$ symmetry and gauge bosons and fermions can interact with

738 Higgs filed to acquire their masses. And this specific process is named *Higgs mechanism*
 739 in SM.

740 The Higgs field ϕ is a doublet and can be written in a Hermitian basis as

$$\phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi_1 - i\phi_2 \\ \phi_3 - i\phi_4 \end{pmatrix} \quad (2.10)$$

741 where $\phi_i = \phi_i^+$ stand for four Hermitian field. In this new basis, the Higgs potential in
 742 Eq. 2.5 can be expressed as:

$$V(\phi) = \frac{1}{2}\mu^2 \left(\sum_{i=1}^4 \phi_i^2 \right) + \frac{1}{4}\lambda \left(\sum_{i=1}^4 \phi_i^2 \right)^2 \quad (2.11)$$

743 To simplify the situation, the axis in this four-dimensional space can be chosen to satisfied
 744 $\langle 0 | \phi_i | 0 \rangle = 0$ for $i = 1, 2, 4$, and $\langle 0 | \phi_3 | 0 \rangle = v$. Thus,

$$V(\phi) \rightarrow V(v) = \frac{1}{2}\mu^2 v^2 + \frac{1}{4}\lambda v^4 \quad (2.12)$$

745 The minimization of this potential depends on the sign of μ^2 as shown in figure 2.3. When
 746 $\mu^2 > 0$ the minimum occurs at $v = 0$, namely the vacuum is empty space and $SU(2)_L \times$
 747 $U(1)_Y$ symmetry is unbroken. In the case of $\mu^2 < 0$, the $v = 0$ symmetric point is no
 748 longer stable and the minimum occurs at nonzero value of $v = (-\mu^2/\lambda)^{1/2}$ which breaks
 the $SU(2)_L \times U(1)_Y$ symmetry. Thus, the classical vacuum ϕ_0 of Higgs doublet can be

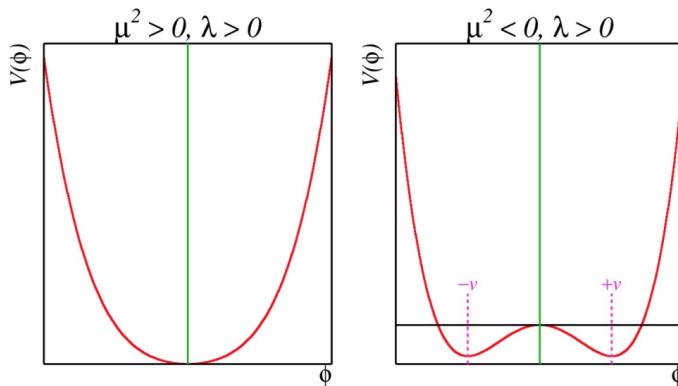


Fig. 2.3 The Higgs potential $V(\phi)$ with $\mu^2 > 0$ (left) and $\mu^2 < 0$ (right).

749

750 expressed by

$$\phi_0 = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix} \quad (2.13)$$

751 And to quantize around the classical vacuum in a general form:

$$\phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H \end{pmatrix} \quad (2.14)$$

752 Where H is a Hermitian field for physical Higgs scalar. In this gauge, the Lagrangian

753 L_{Higgs} in Eq. 2.4 takes a simple form

$$\begin{aligned} L_{Higgs} &= (D^\mu \phi)^\dagger D_\mu \phi - V(\phi) \\ &= M_W^2 W^{\mu+} W_\mu^- \left(1 + \frac{H}{v}\right)^2 + \frac{1}{2} M_Z^2 Z^\mu Z_\mu \left(1 + \frac{H}{v}\right)^2 \\ &\quad + \frac{1}{2} (\partial_\mu H)^2 - V(\phi) \end{aligned} \quad (2.15)$$

754 where the W and Z fields are

$$\begin{aligned} W^\pm &= \frac{1}{\sqrt{2}} (W^1 \mp iW^2) \\ Z &= -\sin\theta_W B + \cos\theta_W W^3 \end{aligned} \quad (2.16)$$

755 Therefore, in Eq. 2.15 spontaneous symmetry breaking brings out masses for the W and

756 Z gauge bosons

$$\begin{aligned} M_W &= \frac{gv}{2} \\ M_Z &= \sqrt{g^2 + g'^2} \frac{v}{2} = \frac{M_W}{\cos\theta_W} \end{aligned} \quad (2.17)$$

757 where θ_W is the weak angle defined as

$$\sin\theta_W = \frac{g'}{\sqrt{g^2 + g'^2}} \quad \cos\theta_W = \frac{g}{\sqrt{g^2 + g'^2}} \quad \tan\theta_W = \frac{g'}{g} \quad (2.18)$$

758 Then another gauge boson photon remains massless with the field of

$$A = \cos\theta_W B + \sin\theta_W W^3 \quad (2.19)$$

759 After the symmetry breaking, the Higgs potential in unitary gauge can be written into

$$V(\phi) = -\frac{\mu^4}{4\lambda} - \mu^4 H^2 + \lambda v H^3 + \frac{\lambda}{4} H^4 \quad (2.20)$$

760 The first term in V is a constant, while the second term denotes a (tree-level) mass of

761 Higgs boson

$$M_H = \sqrt{-2\mu^2} = \sqrt{2\lambda} v \quad (2.21)$$

762 Due to the unknown of quartic Higgs coupling λ , the Higgs mass is not predicted. The

763 third and fourth terms in the Higgs potential V denote the induced cubic and quartic in-
 764 teractions of the Higgs scalar.

765 Through the Higgs mechanism, fermions can also acquire their masses. In the unitary
 766 gauge, Yukawa Lagrangian (L_{Yukawa}) can be written as a simple form of^[2]

$$L_{Yukawa} = - \left(1 + \frac{H}{v} \right) (m_d \bar{d}d + m_u \bar{u}u + m_l \bar{l}l) \quad (2.22)$$

767 in which $m_f = \frac{y_f v}{\sqrt{2}}$ for $f = d, u, l$.

768 2.2 Phenomenology of Large Hadron Collider

769 The Large Hadron Collider (LHC) was built as a bridge between the theories and the
 770 experiment. Physicists hope that the LHC can help to answer some of the fundamental
 771 open questions in physics, concerning the basic laws of interactions and forces among the
 772 elementary particles, the deep structure of space and time, and in particular the interrela-
 773 tion between quantum mechanics and general relativity. This section will talk about firstly
 774 the general introduction of Physics inside hadronic collision, then followed by two impor-
 775 tant LHC phenomenologies of the Higgs physics and Diboson physics that are related
 776 closely to this dissertation.

777 2.2.1 Physics at hadronic collision

778 Protons are not the elementary particle, which actually are composed of quarks and
 779 gluons. So in proton-proton (pp) collision at the LHC, it is not protons themselves interact
 780 but quarks and gluons. Scattering processes can then be further classified into either *hard*
 781 or *soft* processes according to the momentum transfer during the interaction^[3]. QCD, as
 782 an underlying theory for both processes, its approach and level of understandings in two
 783 cases are quite different. For hard process, eg. Higgs, vector bosons and jets production,
 784 the rates and event properties can be precisely predicted based on perturbation theory.
 785 However, for soft processes like total cross-section, the underlying events, the rates and
 786 properties are dominated by non-perturbative QCD effects that are less understood. For
 787 many hard processes, the hard interactions are accompanied by soft ones. An example of
 788 the hadronic collision is illustrated in figure 2.4. and the typical features are summarized
 789 as below:

- 790 • **Parton Distribution Function (PDF):** $f_i(x, Q^2)$ gives the probability of finding a
 791 parton with flavor i (quark or gluon), carrying a momentum fraction of x and with
 792 the scale of momentum transfer Q in a proton. Parton distribution function cannot

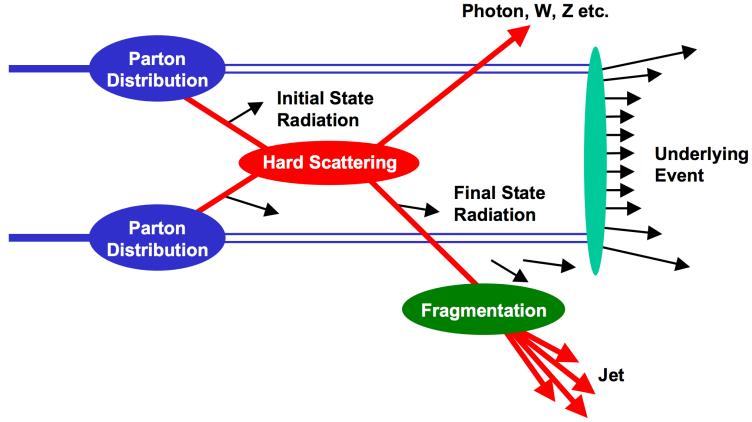


Fig. 2.4 Schematic view of a hadron-hadron collision^[4].

793 be fully calculated by perturbative QCD because of the inherent non-perturbative
 794 nature of partons. There are many different sets of PDFs that are determined by
 795 fits to data from experimental observables in various processes. As an example,
 796 figure 2.5 shows the *PDF4LHC15 NLO PDFs*, which are based on the combination
 797 of the *CT14*, *MMHT14* and *NNPDF3.1* NNLO PDF sets^[5].

- 798 • **Fragmentation and hadronization:** The processes to produce final state particles
 799 (or jets) from the partons produced in hard scattering.
- 800 • **Initial/Final state radiation:** The incoming and outgoing partons that carry color
 801 charge can emit QCD radiation, which gives rise to additional jets. Also the charged
 802 incoming and outgoing particles can emit Quantum Electrodynamics (QED) radia-
 803 tions with photons.
- 804 • **Underlying events:** Products from soft processes (not come from the primary hard
 805 scattering) as the remnants of scattering interactions.

806 Cross section of hard scattering

807 According to *QCD factorization theorems*^[6], the perturbative calculations can be ap-
 808 plied to many important hard processes involving hadrons. The basic problem addressed
 809 by factorization theorems is how to calculate high energy cross sections. Consider the
 810 process of scattering between two hardons A and B to produce a final state X, the cross
 811 section σ can be obtained by summing over all the subprocess cross section $\hat{\sigma}$ ^[7]

$$\sigma_{AB} = \int dx_a dx_b f_{q/A}(x_a) f_{q/B}(x_b) \hat{\sigma}_{ab \rightarrow X} \quad (2.23)$$

812 where $f_{q/A}(x_q)$ is the parton distribution functions of parton q . Taking into account the

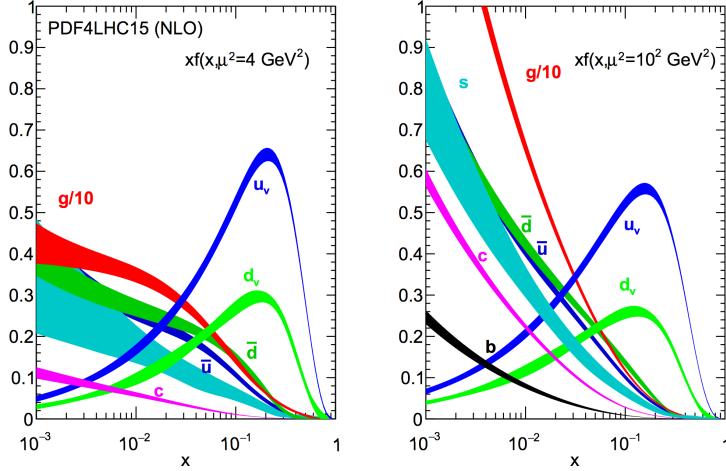


Fig. 2.5 The PDF4LHC15 NLO PDFs at a low scale $\mu^2 = Q^2 = 4 \text{ GeV}^2$ (left) and at $\mu^2 = Q^2 = 100 \text{ GeV}^2$ (right) as a function of x .

813 leading order correction:

$$\sigma_{AB} = \int dx_a dx_b f_{a/A}(x_a Q^2) f_{b/B}(x_b Q^2) \hat{\sigma}_{ab \rightarrow X} \quad (2.24)$$

814 where Q^2 represents large momentum scale that characterizes the hard scattering. Later
815 on, since the finite corrections were not universal and had to be calculated separately for
816 each process, the perturbative $O(\alpha_S^n)$ corrections to the leading logarithm cross section
817 in Eq. 2.24 need to be applied, one can get:

$$\sigma_{AB} = \int dx_a dx_b f_{a/A}(x_a \mu_F^2) f_{b/B}(x_b \mu_F^2) \hat{\sigma}_{ab \rightarrow X}(\alpha_S, \mu_R, \mu_F) \quad (2.25)$$

818 in which μ_F is *factorization scale* which can represent the scale that separates the long-
819 and short-distance physics, and μ_R is the *renormalization scale* for QCD running coupling.
820 $\hat{\sigma}_{ab \rightarrow X}$ is the parton-level hard scattering cross section that can be calculated perturbatively
821 in QCD with the form of

$$\hat{\sigma}_{ab \rightarrow X}(\alpha_S, \mu_R, \mu_F) = (\alpha_S)^n \left[\hat{\sigma}^{(0)} + (\alpha_S/2\pi) \hat{\sigma}^{(1)}(\mu_R, \mu_F) + (\alpha_S/2\pi)^2 \hat{\sigma}^{(2)}(\mu_R, \mu_F) + \dots \right] \quad (2.26)$$

822 where $\hat{\sigma}^{(0)}$ stands for the leading-order (LO) partonic cross section, while $\hat{\sigma}^{(1)}$ and $\hat{\sigma}^{(2)}$ are
823 the next-to-leading-order (NLO) and next-to-next-to-leading-order (NNLO) cross section.

824 μ_R and μ_F depend on the order of truncation in Eq. 2.26. In principle, if cross section is
825 calculated to all orders, it is invariant under changes in these parameters. The choices of μ_R
826 and μ_F are arbitrary. To avoid unnaturally large logarithms reappearing in the perturbation
827 series, it is sensible to choose μ_R and μ_F values of the order of the typical momentum
828 scales of the hard scattering process and $\mu_R = \mu_F$ is also often assumed. Take Drell–Yan

process as an example, the standard choice is $\mu_R = \mu_F = m_{ll}$, where m_{ll} is the invariant mass of di-lepton pair.

2.2.2 Higgs physics at the LHC

One important physics purpose of the LHC is searching for the Higgs boson, which was the last missing part in the SM. This section will discuss both the production and decay modes of the SM Higgs boson in proton-proton collision.

Higgs productions

The Higgs boson can be produced through several processes. There are 4 main production modes at the LHC: gluon-gluon Fusion (ggF), vector boson Fusion (VBF), associated production with vector-bosons (VH) (also called the Higgs Strahlung) and associated production with a pair of top/anti-top quarks (ttH)^[8]. Figure 2.6 shows the corresponding Feynman diagrams of each process (at LO). For pp collisions, the cross section of pro-

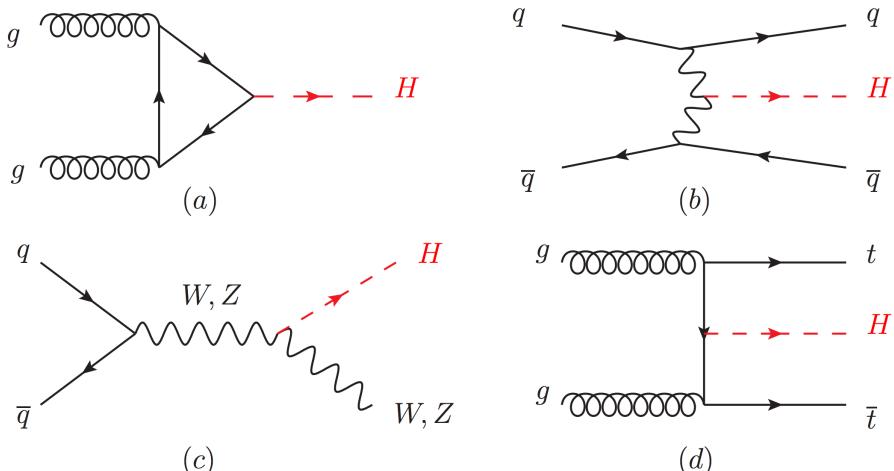


Fig. 2.6 Feynman diagrams of the Higgs production modes: (a) ggF ; (b) VBF ; (c) VH ; (d) ttH .

duction of Higgs boson is a function of centre-of-mass energy \sqrt{s} . Figure 2.7 depicts the cross section of SM Higgs, whose mass is 125 GeV, for several different production modes when centre-of-mass energy varying from 6 to 15 TeV. Figure 2.8 shows the prospect of production cross section as a function of Higgs mass from 10 to 2000 GeV for pp collision at the centre-of-mass energy of 13 TeV and 14 TeV^[9].

Higgs decays

The Higgs boson can interact with gauge bosons and fermions through gauge coupling and the Yukawa coupling as introduced in section 2.1.3. Figure 2.9 depicts the Feynman diagrams of various possible Higgs decay channels. The branching ratio of Higgs boson decaying into different final states as a function of Higgs mass is shown in figure 2.10.

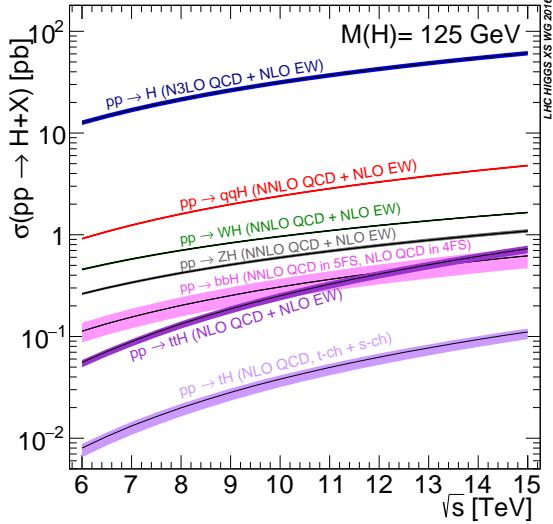


Fig. 2.7 The SM Higgs boson production cross sections for various production modes as a function of the centre-of-mass energy for pp collision.

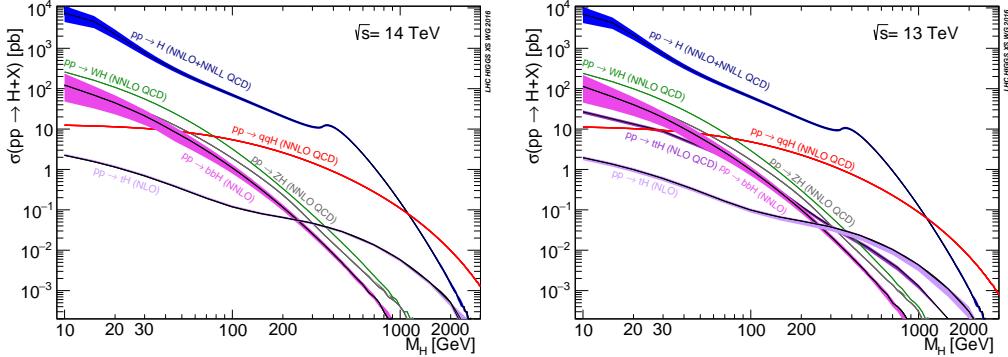


Fig. 2.8 Higgs boson production cross section for various production modes as a function of the Higgs mass for $\sqrt{s} = 13$ TeV (left) and 14 TeV (right) for pp collision.

2.2.3 Diboson physics

The study of diboson physics is another important test for SM of particle physics in electroweak sector, while the Vector Boson Scattering (VBS) is a key process for probing the mechanism of the electroweak symmetry breaking (EWSB). In the meantime, the non-resonant diboson productions are crucial backgrounds for Higgs studies at the LHC, which make the precise measurement of their cross section becomes very important.

Diboson productions

About 90% of diboson productions at hadron collider is from quark-antiquark annihilation, while others are from gluon initiated process. Figure 2.11 shows the tree-level Feynman diagrams of diboson production. Then figure 2.12 illuminates the total production cross-section presented by ATLAS as a function of centre-of-mass energy \sqrt{s} from 7 to 13 TeV for several diboson processes compared to some other major processes in

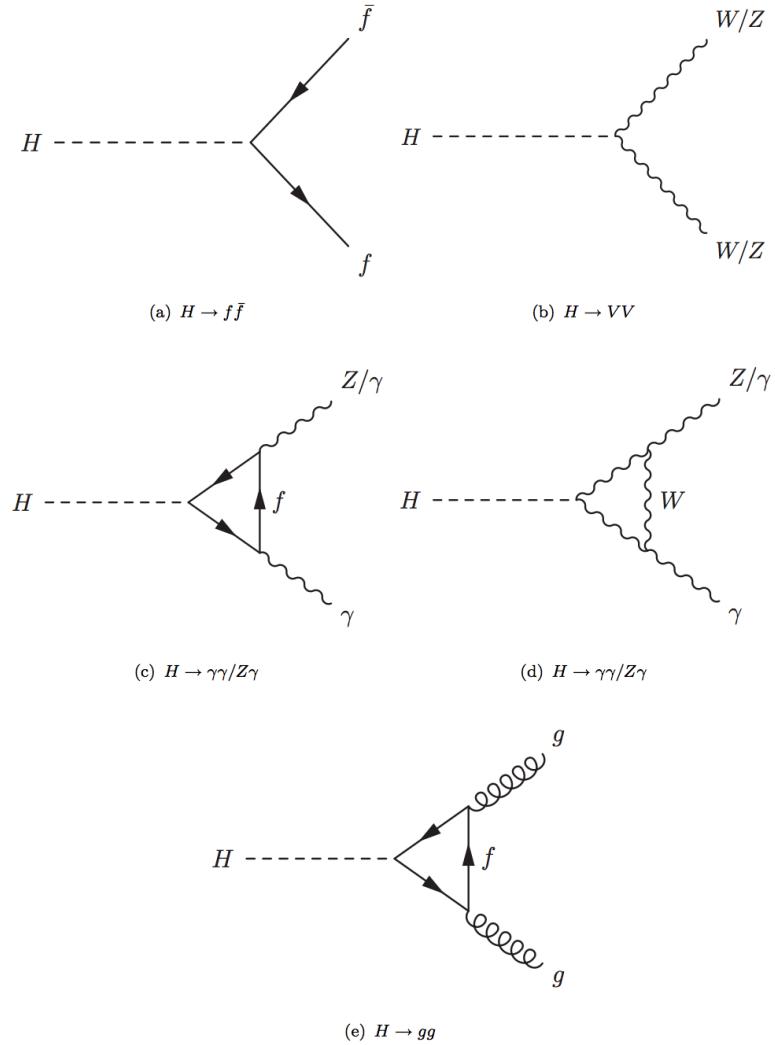
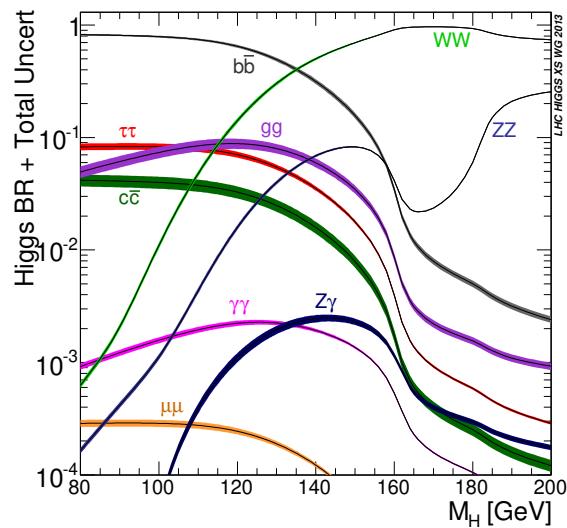


Fig. 2.9 SM Higgs decay channels.


 Fig. 2.10 Branching ratio of Higgs decays in various channels as a function of Higgs mass^[10].

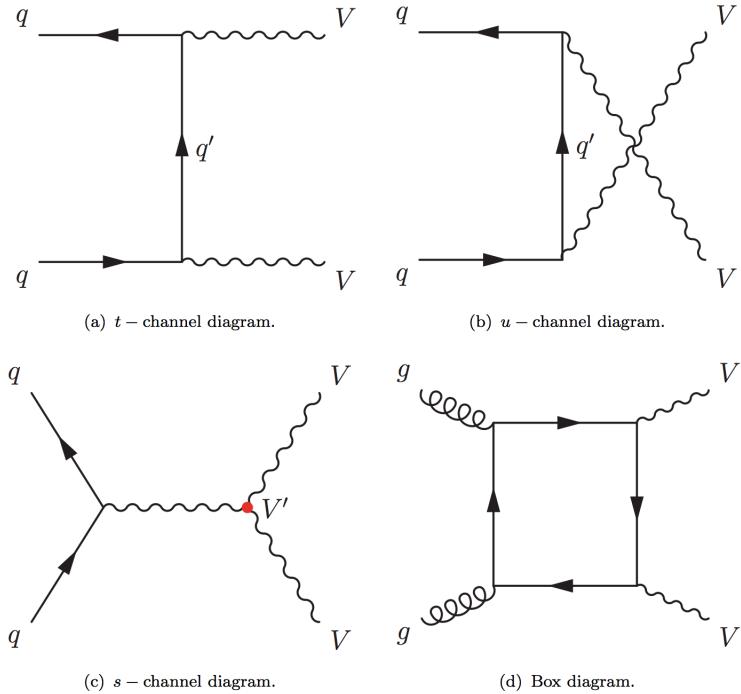


Fig. 2.11 The tree-level Feynman diagrams of diboson production at the LHC.

hadron collision. The cross section for diboson processes are calculated at next-to-next leading order (NNLO).

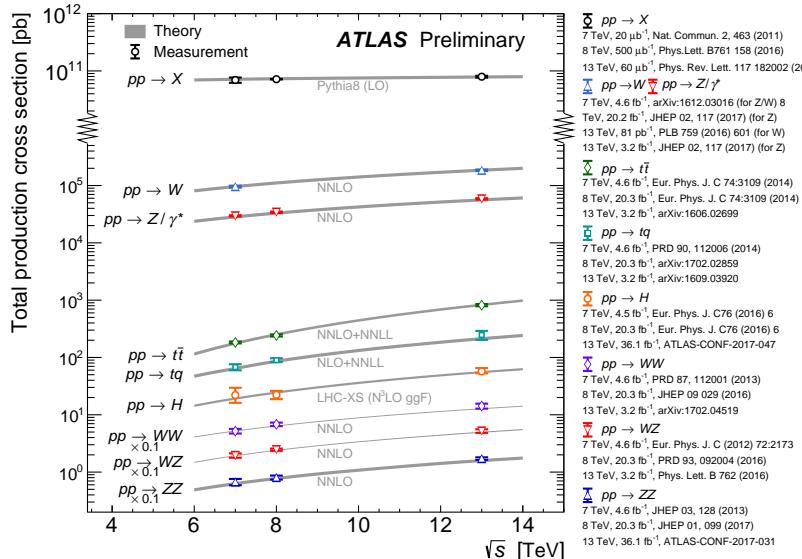


Fig. 2.12 Total production cross-section presented by ATLAS as a function of centre-of-mass energy \sqrt{s} from 7 to 13 TeV for some selected processes, the diboson measurements are scaled by a factor 0.1 to allow a presentation without overlaps.

864

Vector boson scattering

865 The $SU(2)_L \times U(1)_Y$ structure in SM predicts self-interactions between electroweak
866 gauge bosons. Those self-couplings can involve either three or four gauge bosons at a sin-
867 gle vertex, known as triple gauge coupling (TGC) or quartic gauge couplings (QGC), re-

869 spectsively. Vector boson scattering (VBS) is carried out by four electroweak vector bosons,
 870 namely Z , W^\pm and photon (γ) as the Feynman diagrams shown in figure 2.13. And the
 871 vertexes include either those self-interactions or the interactions with the Higgs boson are
 described in figure 2.14.

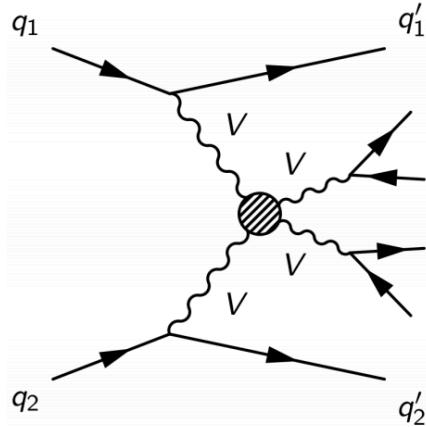


Fig. 2.13 Feynman diagrams of the vector boson scattering.

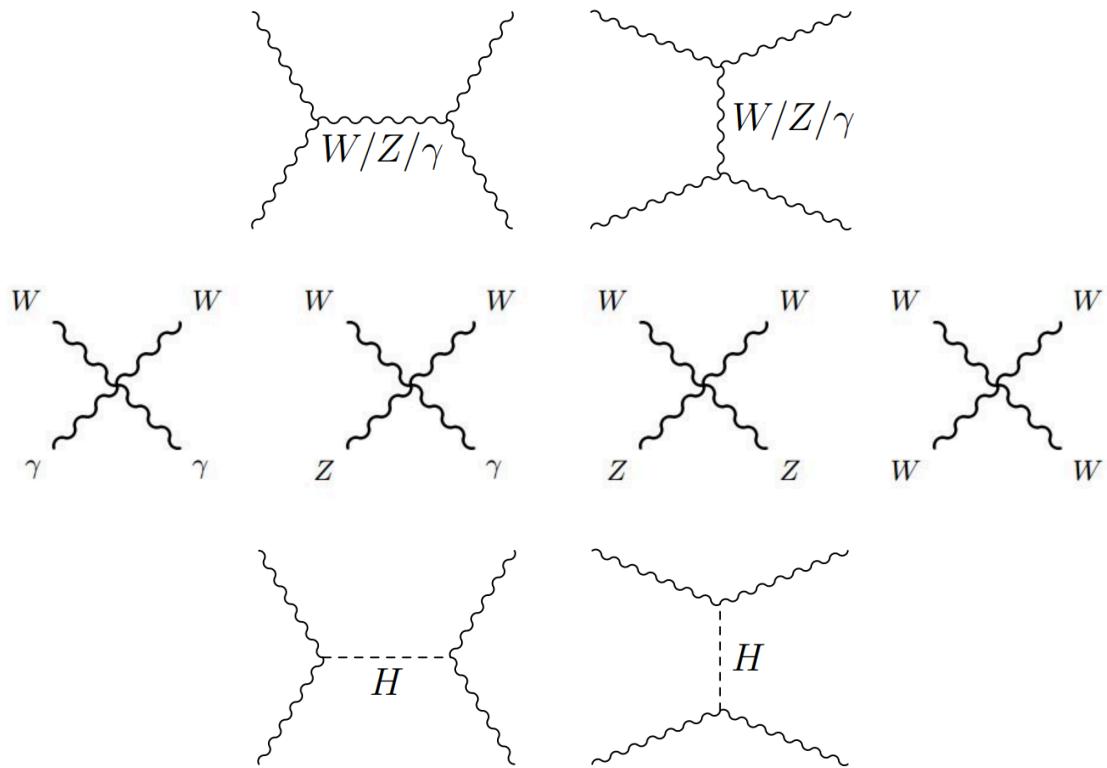


Fig. 2.14 Feynman diagrams of vertexes involving QGC, TGC and Higgs.

873 The amplitudes of leading-order (LO) VBS can be expressed as^[11]:

$$\begin{aligned} iM_{TGC}^{s-channel} &= -i \frac{g_1^2}{4m_W^4} [s(t-u) - 3m_W^2(t-u)] \\ iM_{TGC}^{t-channel} &= -i \frac{g_1^2}{4m_W^4} \left[(s-u)t - 3m_W^2(s-u) + \frac{8m_W^2}{s} u^2 \right] \end{aligned} \quad (2.27)$$

874

$$iM_{QGC} = i \frac{g_1^2}{4m_W^4} \left[s^2 + 4st + t^2 - 4m_W^2(s+t) - \frac{8m_W^2}{s} ut \right] \quad (2.28)$$

875

$$\begin{aligned} iM_{Higgs} &= -i \frac{C_v^2 g_1^2}{4m_W^2} \left[\frac{(s-2m_W^2)^2}{s-m_H^2} + \frac{(t-2m_W^2)^2}{t-m_H^2} \right] \\ &\simeq -i \frac{C_v^2 g_1^2}{4m_W^2} (s+t) \end{aligned} \quad (2.29)$$

876 Combining s- and t-channel of TGC in Eq. 2.27 and the QGC term in Eq. 2.28:

$$iM_{TGC} + iM_{QGC} = i \frac{g_1^2}{4m_W^2} (s+t) + O((s/m_W^2)^0) \quad (2.30)$$

877 In Eq. 2.30, the amplitude grows as a function of centre-of-mass energy (\sqrt{s}), which
 878 violates the unitarity in the TeV region. Considering the Higgs term in Eq. 2.29 can
 879 perfectly cancel out this growing, and the remaining term $O((s/m_W^2)^0)$ only depends on
 880 the total amplitude in SM.

881 In conclusion, the Higgs boson acts as "moderator" to unitarize high-energy longitudi-
 882 nal vector boson scattering as introducing the Higgs restores the unitarity of total am-
 883 plitude in high energy region.

884 **Chapter 3 The Large Hadron Collider and the ATLAS 885 Detector**

886 3.1 The Large Hadron Collider

887 Located near the French-Swiss border at the European Organization for Nuclear Re-
888 search (CERN), the Large Hadron Collider (LHC) is the world's largest and most pow-
889 erful particle collider. It's the proton-proton collider with the centre-of-mass energy up
890 to 14 TeV. The beams inside the LHC are made to collide at four locations around its
891 27-kilometer accelerator ring, corresponding to the positions of four particle detectors -
892 ATLAS, CMS, ALICE and LHCb. With its unprecedented energy, the LHC is designed
893 to observe physics that involve highly massive particles, which have never been observed
894 in previous accelerators with lower energies.

895 3.1.1 Operation history and machine layout

896 **Operation history**

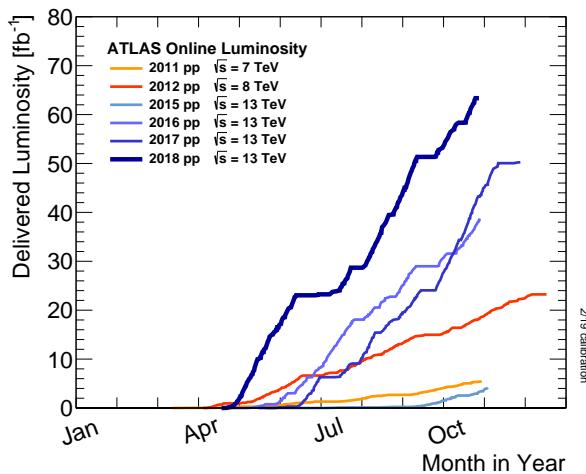
897 The LHC^[12-15] is a two-ring-superconducting-hadron accelerator and collider lies in
898 a tunnel 27 kilometres in circumference and as deep as 175 metres underground. It's
899 designed to provide proton-proton (pp) collisions at the centre-of-mass energy (\sqrt{s}) up
900 to 14 TeV with a unprecedented luminosity of $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$. In the meantime, it can also
901 collide heavy (Pb) ions with an energy of 2.8 TeV per nucleon and a peak luminosity of
902 $10^{27} \text{ cm}^{-2} \text{ s}^{-1}$. Table 3.1 shows the main design parameters of the LHC for proton-proton
903 collisions.

904 The LHC was built from 1998 to 2008. It started its first beam in September 2008, but
905 then was interrupted by a quench incident only after a few days running. Then it resumed
906 the operation in November 2009 with a low energy beams. From March 2010, physics runs
907 took place at the centre-of-mass energy of 7 TeV. Later on, this energy was increased in
908 2012 to $\sqrt{s} = 8 \text{ TeV}$, with an integrated luminosity of 20.3 fb^{-1} , and this period is called
909 "run-1". After run-1, the LHC was shut down for two years for hardware maintenance and
910 upgrade, starting from February 2013.

911 The second operation period with higher centre-of-mass energy at 13 TeV started
912 from 2015 called "run-2". And it continued to the end of 2018 with total integrated lumi-
913 nosity reaching about 147 fb^{-1} for ATLAS experiment. Figure 3.1 shows the cumulative
914 luminosity as a function of time in month delivered to ATLAS experiment during stable
915 beams in years from 2011 to 2018.

Table 3.1 Summary of design parameters of the LHC for pp collisions.

Circumference	26.7 km
Beam energy at collision	7 TeV
Beam energy at injection	0.45 TeV
Dipole field at 7 TeV	8.33 T
Luminosity	$10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
Beam current	0.56 A
Protons per bunch	1.1×10^{11}
Number of bunches	2808
Nominal bunch spacing	24.95 ns
Normalized emittance	$3.75 \mu\text{m}$
Total crossing angle	$300 \mu\text{rad}$
Energy loss per turn	6.7 keV
Critical synchrotron energy	44.1 eV
Radiated power per beam	3.8 kW
Stored energy per beam	350 MJ
Stored energy in magnets	11 GJ
Operating temperature	1.9 K

**Fig. 3.1 Cumulative luminosity as a function of time in years from 2011 to 2018 for ATLAS detector.**

916

Machine layout

917

The layout of CERN accelerator complex is shown in figure 3.2. The protons are accelerated by a series of machines before being injected into the main ring. At beginning, the 50 MeV protons are produced in the linear particle accelerator LINAC2, and further accelerated to 1.4 GeV in Proton Synchrotron Booster (PSB). The protons are then injected

921 into the Proton Synchrotron (PS) to gain the energy of 26 GeV and further accelerated to
 922 450 GeV in Super Proton Synchrotron (SPS). At the end, they are injected into the main
 923 ring, and can reach a maximum energy of 7 TeV.

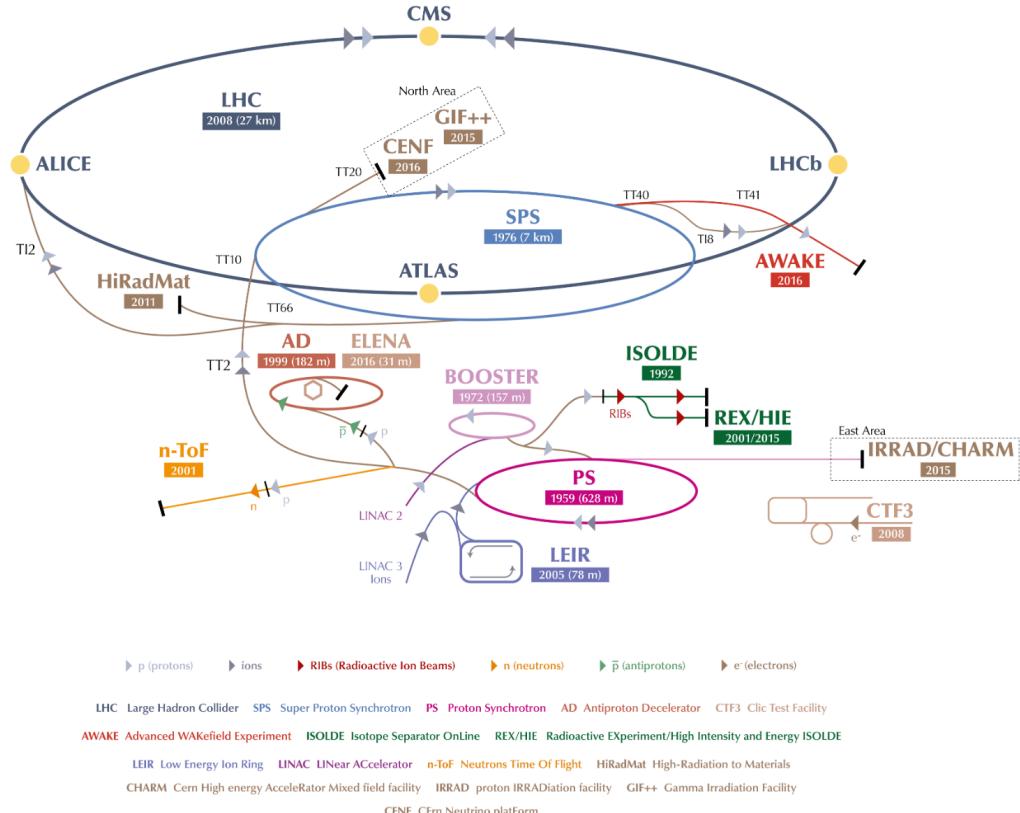


Fig. 3.2 CERN accelerator complex^[16].

924 The collisions can occur in 4 points, with corresponding 4 major detector experiments
 925 that are briefly described as follows:

- **ATLAS:** A Toroidal LHC ApparatuS, one of the two general-purpose particle detector experiments and detector with largest volume at the LHC. It is designed to search for the Higgs boson, test the stardand model of particle physics and search for possible beyond SM physics.
- **CMS:** Compact Muon Solenoid, another large general-purpose particle physics detector, with the same physics goal (also cross check) as ATLAS.
- **ALICE:** A Large Ion Collider Experiment, it is optimized to study heavy-ion (Pb-Pb nuclei) collisions at a centre-of-mass energy of 2.76 TeV per nucleon pair.
- **LHCb:** Large Hadron Collider beauty, it is a specialized b-physics experiment, designed primarily to measure the parameters of CP violation in the interactions of b-hadrons.

3.1.2 Luminosity and pile-up

Luminosity

In beam-beam collisions, the event rate for a process is given by^[15]:

$$N = \mathcal{L}\sigma \quad (3.1)$$

where σ is the cross section of the process, and \mathcal{L} is the luminosity. For the studies of rare events, \mathcal{L} must be as high as possible. The luminosity only depends on the beam parameters, and can be written as:

$$\mathcal{L} = \frac{N_b^2 n f_r \gamma}{4\pi \epsilon_n \beta^*} \quad (3.2)$$

where N_b denotes the number of particles per bunch, n is the number of bunches per beam, f_r is the revolution frequency, γ represents relativistic γ factor, ϵ_n is the normalized transverse emittance and β^* denotes the β function at the collision point. To reduce the beam-beam interaction effects, the bunches must have a crossing angle, which produces a geometrical luminosity reduction factor F :

$$F = 1/\sqrt{1 + \left(\frac{\theta_c \sigma_Z}{2\sigma^*}\right)} \quad (3.3)$$

where θ_c denotes the crossing angle at the interaction point, σ_Z is the root mean square (RMS) bunch length and σ^* is the transverse RMS beam size at crossing point.

The luminosity expressed in Eq. 3.2 is normally the instantaneous luminosity. In fact the running conditions usually vary with time, so the luminosity can change as well. To take into account the time dependence, integrated luminosity is invited, by integraling the instantaneous luminosity over time:

$$L = \int \mathcal{L}(t) dt \quad (3.4)$$

The unit of integrated luminosity we commonly use is b^{-1} that satisfying $1b^{-1} = 10^{24} cm^{-2}$. Figure 3.3 shows integrated luminosity as a function of time delivered to ATLAS (green), recorded by ATLAS (yellow), and certified to be good quality data (blue) during run-2 pp collisions. For most physics analysis, the data with good quality (require to satisfy *Good Run List*) is used.

Pile-up

In collisions, multiple interactions can happen in one single bunch crossing, which is called “pile-up”. The variable $\langle \mu \rangle$, representing the average number of interactions per

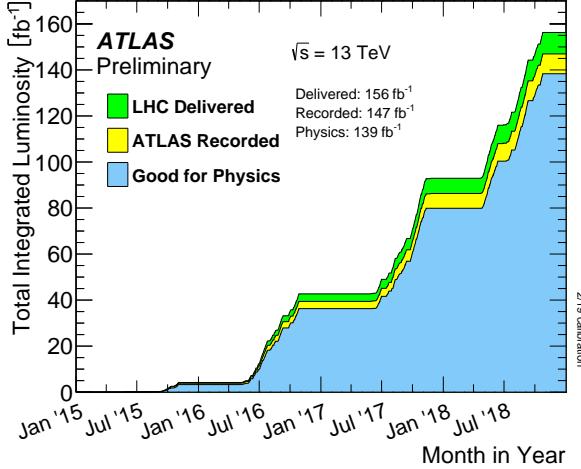


Fig. 3.3 Integrated luminosity in ATLAS.

bunch crossing, is defined to describe pile-up effect:

$$\langle \mu \rangle = \frac{L_{tot}\sigma}{f_r n_{bunch}} \quad (3.5)$$

where L_{tot} is the instantaneous luminosity, σ is the inelastic cross section, f_r is the LHC revolution frequency and n_{bunch} is the number of colliding bunches. Normally, with increasing luminosity, the pile-up becomes more significant. Figure 3.4 shows the luminosity-weighted distribution of the mean number of interactions per crossing for pp collision data from 2015 to 2018 (full run-2), the challenge of pile-up increased in each year.

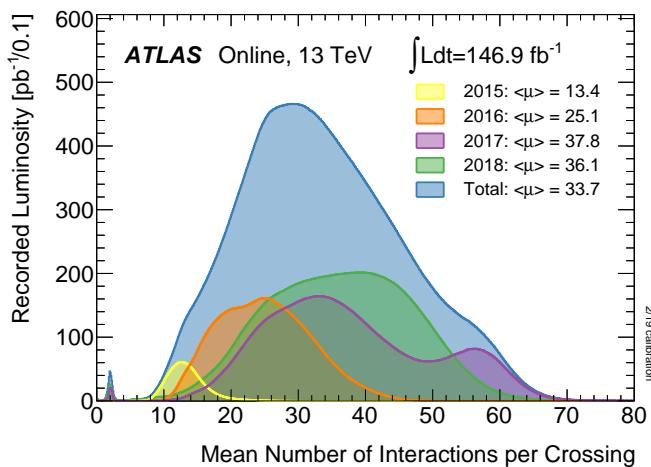


Fig. 3.4 Number of Interactions per Crossing from 2015-2018 in ATLAS.

969 3.2 ATLAS detector

970 3.2.1 Detector overview

971 ATLAS (A Toroidal LHC ApparatuS) is the largest volume detector ever constructed
 972 for a particle collider. It is a cylinder with 46 meters long, 25 meters in diameter, and sits
 973 in a cavern 100 meters below ground. The detector contains about 3000 km of cables and
 974 it weights 7000 tonnes.

975 This paragraph briefly summarizes the coordinate system and nomenclature used to
 976 describe the ATLAS detector^[17]. As depicted in figure 3.5, we define the nominal inter-
 977 action point as the origin of the coordinate system, the beam direction as the z -axis and
 978 the x - y plane is transverse to the beam direction. The positive x -axis is defined to be the
 979 direction pointing to the center of the LHC ring, while the positive y -axis is pointing up-
 980 wards. There are two sides of detector A and C, in which A (C) -side is in the positive

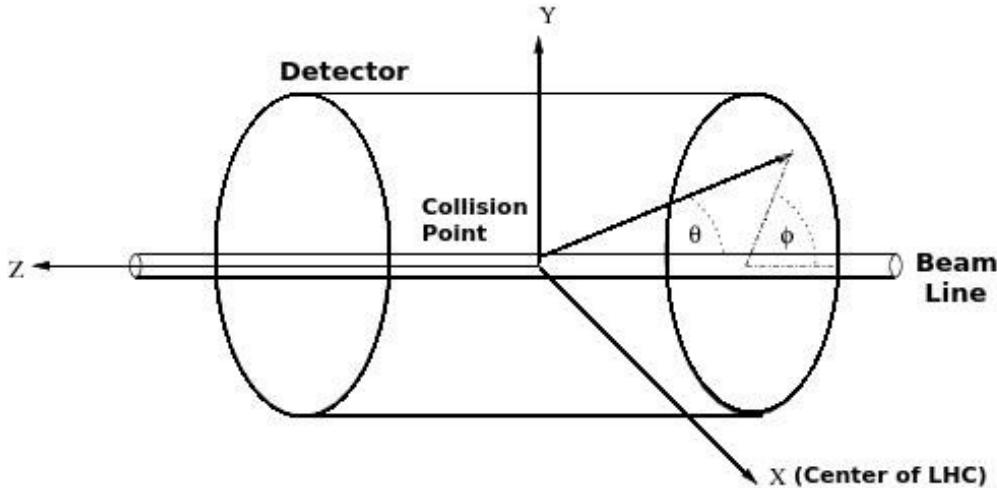


Fig. 3.5 Coordinate system used by the ATLAS experiment at the LHC^[18].

980
 981 (negative) z direction. The azimuthal angle ϕ is measured as usual around the beam axis,
 982 while the polar angle θ is the angle from the beam axis. In physics analysis, we usually
 983 use the pseudorapidity instead of θ angle, which is designed as $\eta = -\ln \left[\tan \left(\frac{\theta}{2} \right) \right]$.

984 For massive objects (eg. jets), the rapidity $y = \frac{1}{2} \ln \left[\frac{E+p_z}{E-p_z} \right]$ is used. In addition, the
 985 transverse momentum p_T , transverse energy E_T and the missing transverse energy E_T^{miss}
 986 are defined in x - y plane. The ΔR , a commonly used distance measurement, is defined in
 987 the pseudorapidity-azimuthal angle space as $\Delta R = \sqrt{\Delta\eta^2 + \Delta\phi^2}$.

988 The overall ATLAS layout is shown in figure 3.6, which is forward-backward sym-
 989 metric with respect to the interaction point. The magnet configuration comprises a thin
 990 superconducting solenoid surrounding the inner-detector cavity, and three large super-
 991 conducting toroids (one barrel and two end-caps) arranged with an eight-fold azimuthal

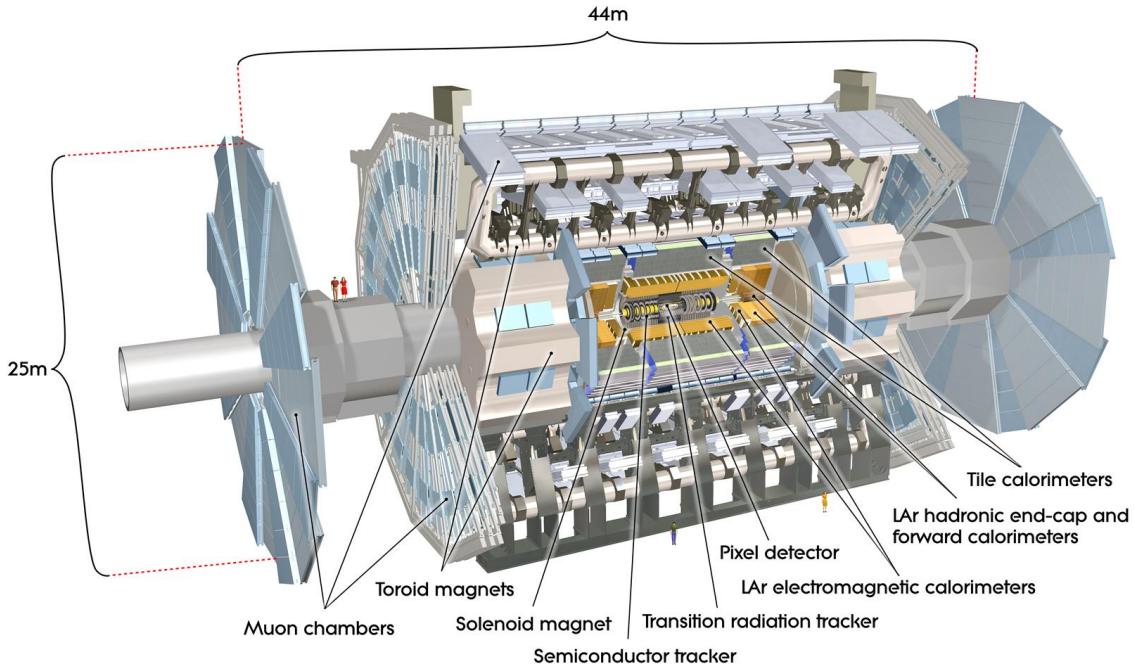


Fig. 3.6 Cut-away view of the ATLAS detector^[19].

992 symmetry around the calorimeters.

993 **The inner detector**, which is the innermost part of ATLAS, is immersed in a 2 T
 994 solenoidal magnetic field. It's used for pattern recognition, momentum and vertex mea-
 995 surements and electron identification, with the combination of tracking system.

996 **The calorimeter** is outside the solenoid, for electromagnetic and hadronic en-
 997 ergy measurements. The high granularity liquid-argon (LAr) electromagnetic sampling
 998 calorimeters is used to measure energy and position with range up to $|\eta| < 3.2$ for electrons
 999 and photons. For hadron, a scintillator-tile calorimeter is used in the range of $|\eta| < 1.7$,
 1000 and the liquid-argon hadronic endcap calorimeters (HEC) is used in end-cap region. And
 1001 then the LAr forward calorimeters provide both electromagnetic and hadronic energy mea-
 1002 surements with the coverage in forward region up to $|\eta| = 4.9$.

1003 **The muon spectrometer** is the outermost layer. It's a air-core toroid system, with
 1004 a long barrel and two inserted end-cap magnets that provides strong bending power in a
 1005 large volume within a light and open structure. A set of chambers measuring the tracks
 1006 of muons with high spatial precision and accurate time-resolution are used. Multiple-
 1007 scattering effects are minor, and excellent muon momentum resolution can be achieved.

1008 3.2.2 Physics requirement

1009 As mentioned previously, ATLAS is one of two general-purpose particle detector ex-
1010 periment at the LHC. It's designed to take advantage of the unprecedeted energy at the
1011 LHC, as the discovery of Higgs boson is one of its benchmark. Lots of precise tests
1012 and measurements of SM physics are ongoing with ATLAS experiment. while, in the
1013 meantime, ATLAS is also designed to observe the phenomena that involve highly massive
1014 particles, which can also explore the possibility of extra dimensions proposed by several
1015 models in TeV region. To fulfil many diverse physics goals, a set of general requirements
1016 are needed:

- 1017 • The high-speed and radiation-hard electronics are required due to the experimental
1018 conditions at the LHC.
- 1019 • High detector granularity is needed to reduce the overlapping events and handle the
1020 particle fluxes.
- 1021 • Large acceptance in pseudorapidity and azimuthal angle coverage is needed.
- 1022 • For inner detector, good charged-particle momentum resolution and reconstruction
1023 efficiency are crucial. And the vertex detectors close to the interaction region are
1024 required to be able to observe secondary vertices for offline tagging of τ -lepton and
1025 b -jets.
- 1026 • Good electromagnetic (EM) calorimetry for electron and photon, as well as full-
1027 coverage hadronic calorimetry for accurate jet and missing transverse energy mea-
1028 surements, are essentially required, since these measurements form the basis of
1029 many studies.
- 1030 • Good muon spectrometer is also required for muon identification and momentum
1031 measurement over a wide range of momenta.
- 1032 • Highly efficient but with sufficient background rejection triggers are also needed
1033 and extremely important for objects with low transverse-momentum.

1034 More detailed descriptions of each sub-system will be given in the following subsec-
1035 tions.

1036 3.2.3 Magnet system

1037 A strong magnetic field is required for precise measurement of charged particle mo-
1038 mента. The ATLAS detector uses two large superconducting magnet systems, a hybrid
1039 system of a central superconducting solenoid and three outer superconducting toroids, to
1040 bend charged particles^[20]. The total magnet system is 22 m in diameter and 26 m in length
1041 as shown in figure 3.7.

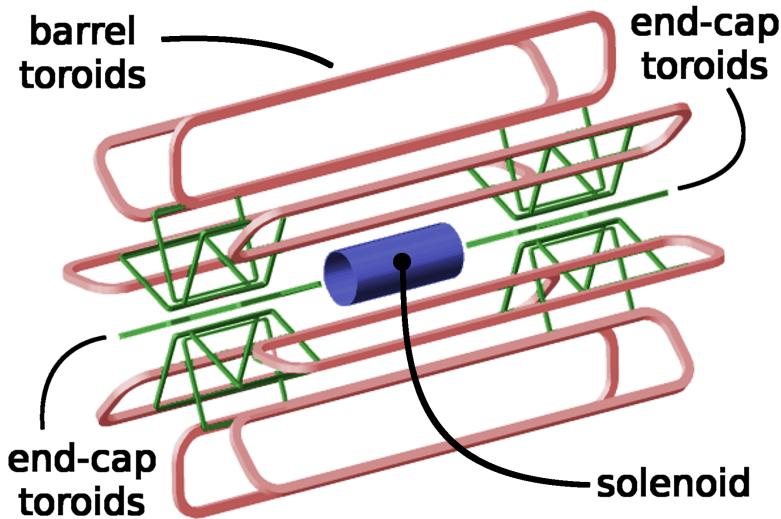


Fig. 3.7 Schematic diagram of the ATLAS magnet system.

1042 The central solenoid produces two Tesla (T) magnetic field surrounding the inner De-
 1043 tector. When obtaining such high field strength, at the same time, the solenoid needs to
 1044 be thin in order to reduce the material in front of the calorimeter.

1045 The outer toroid system comprises one barrel superconducting toroid and two end-
 1046 caps. The barrel one is composed of eight coils encased in individual racetrack-shaped,
 1047 stainless-steel vacuum vessels and produces the magnetic field in the cylindrical volume
 1048 surrounding the calorimeters. Each end-cap toroid consists of a single cold mass built up
 1049 from eight flat, square coil units and eight keystone wedges and provides a magnetic field
 1050 of approximately 1 T for the muon detectors in the end-cap regions.

1051 3.2.4 Inner detector

1052 The inner detector, as shown in figure 3.8, is the detector closest to beam pipe. It's
 1053 used to measure the position of charged particle tracks in high precision together with
 1054 good momentum resolution, among which the measurement of primary and secondary
 1055 vertices and electron identification are especially important. Due to the extremely high
 1056 luminosity produced by the LHC, the precise measurements of vertex and momentum
 1057 becomes tough and fine-granularity detectors are crucial. The inner detector consists of
 1058 three subdetectors described as below:

1059 **Pixel detector**

1060 The pixel detector^[22] is the innermost part of ATLAS tracking system. With finest
 1061 granularity of materials, it has the best spatial resolution and 3-dimensional space-point
 1062 measurement in inner detector. ATLAS Pixel Detector for the LHC run-2 is composed of
 1063 4 layers of barrel pixel detector and two end-caps with three pixel disks each, as shown in

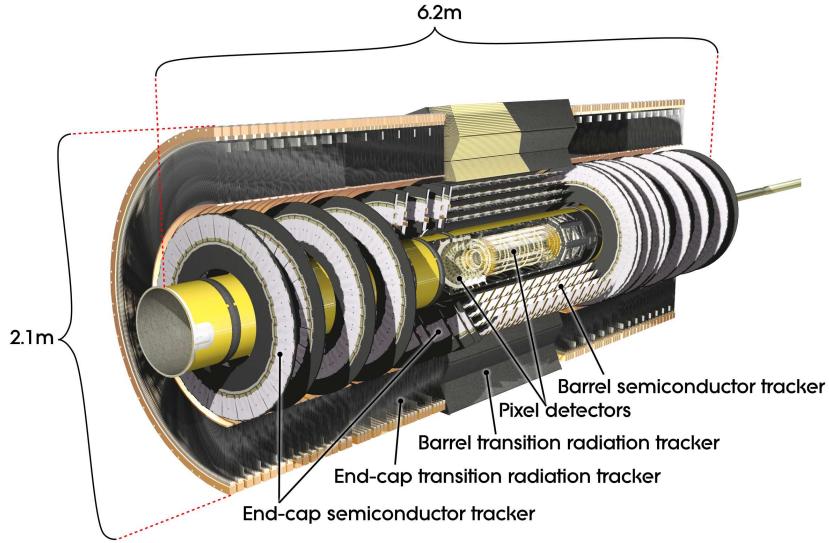


Fig. 3.8 Schematic diagram of the ATLAS inner detector^[21].

figure 3.9. There are three outer layers that originally installed for run-1 and one additional layer called Insertable B-Layer (IBL) that newly constructed in run-2^[23]. Now the 4-layer pixel detector has very good reconstruction of primary and secondary vertices, which is even crucial for long-lived particles like τ -lepton and b-quark.

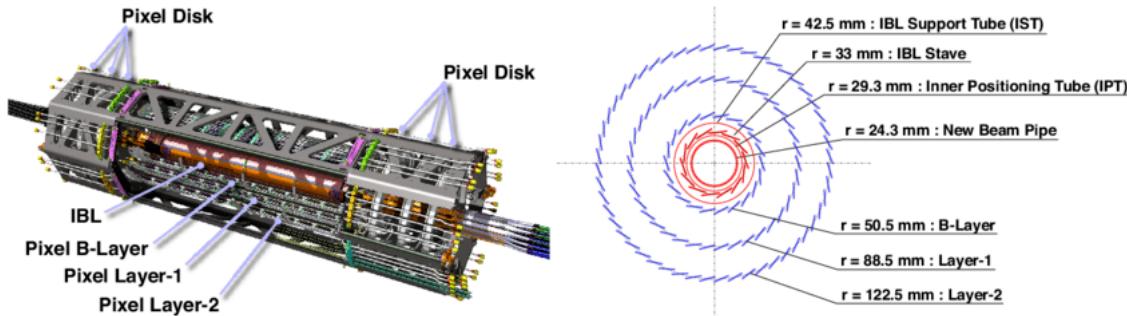


Fig. 3.9 Schematic diagram of the ATLAS 4-Layer Pixel Detector.

1067

1068 Semiconductor Tracker

1069 The Semiconductor Tracker (SCT)^[24] installed outside the pixel detector is the middle
1070 component of the inner detector. It has similar function as pixel detector but with long
1071 and narrow strips instead of small pixels, which makes a much larger coverage than pixel
1072 detector. The SCT consists of 4088 modules, and contains four concentric layers in barrel
1073 (2112 modules) and nine disks in each of the two end-caps (1976 modules) as shown in
1074 figure 3.10. And it measures particles over a large area with 6.3 million readout channels
1075 and a total area of 61 square meters. The SCT is the most critical part of the inner detector
1076 for 2D track hit reconstruction. In barrel, the hit precision is 17 μm in the $r\text{-}\phi$ coordinate
1077 and 580 μm in z coordinate. In end-caps, the precision is 17 μm in the $z\text{-}\phi$ coordinate and
1078 580 μm in r coordinate.

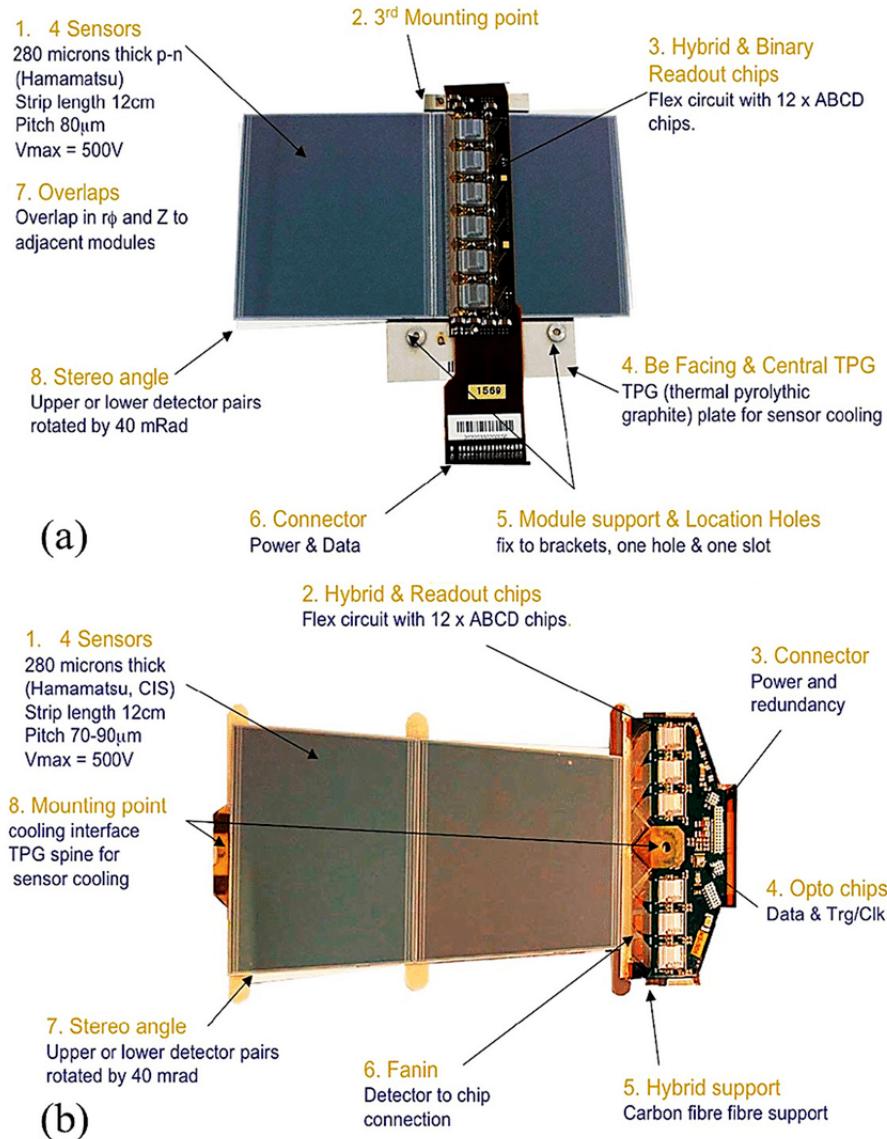


Fig. 3.10 SCT (a) barrel module and (b) end-cap^[25].

Transition radiation tracker

The transition radiation tracker (TRT)^[26] is the outermost part of inner detector, which has a very different design comparing to the two previously described sub-detectors. It can be separated into three parts: one barrel and two end-cap regions with the $|\eta|$ coverage up to 2.0. There are 73 barrel layers and 224 end-cap layers (112 in each) with 372000 straws in total, and about 351000 readout channels for TRT. The TRT provides better z resolution but much worse $r\phi$ resolution (about 130 μ m) comparing to the pixel detector and SCT per straw. But the straw hits still make significant contributions to momentum measurement, since its lower precision per point (compared to silicon) can be compensated by the large number of measurements and long track length.

3.2.5 Calorimeters

The calorimeters are designed to measure the energy from particles by absorbing them. They are located outside the solenoidal magnet that surrounds the inner detector. The ATLAS calorimeters are comprised of a number of sampling calorimeters with full ϕ -symmetry and the pseudorapidity range of $|\eta| < 4.9$. Figure 3.11 shows the layout of the ATLAS calorimeter system. There are two basic calorimeter systems: an inner electromagnetic (EM) calorimeter and an outer hadronic calorimeter. The EM calorimeter is designed for precise measurements of electrons and photons with fine granularity; while the hadronic one has relative coarser granularity but satisfies the physics requirements for jets reconstructions and E_T^{miss} measurements. Two different sampling techniques are used, the EM calorimeter is purely based on liquid-argon (LAr) technology, while the hadronic one use both LAr and scintillating tiles calorimeters. More details are described as below:

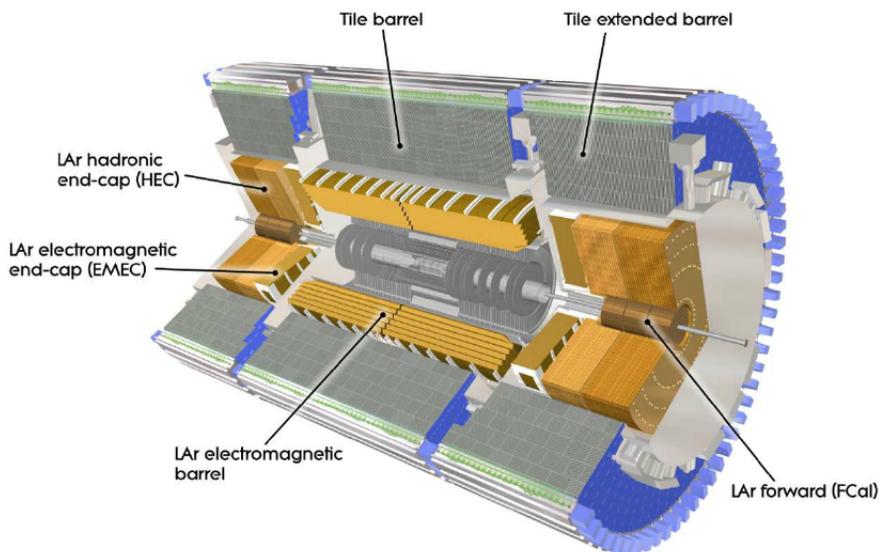


Fig. 3.11 Cut-away view of the ATLAS calorimeters. The LAr calorimeters are seen inside the scintillator-based tile hadronic calorimeters^[27].

Liquid Argon calorimeter

The LAr calorimeter uses liquid-argon as active medium. The LAr sampling calorimeter technique with “accordion-shaped” electrodes is used for all electromagnetic calorimetry covering the pseudorapidity range of $|\eta| < 3.2$; and for hadronic calorimetry with range from $|\eta| = 1.4$ to the acceptance limit $|\eta| = 4.9$ ^[28]. Figure 3.12 depicts a segment of the barrel calorimeter, which had “accordion-shaped” electrodes and absorber. For barrel EM calorimeter, the absorbing material is lead-liquid argon, while the hadronic end-cap calorimeter uses copper plates. In addition, the forward calorimeter is split into

1110 three parts, an EM sector in which copper is used as absorbing material and two hadronic sectors using tungsten outside the EM sector.

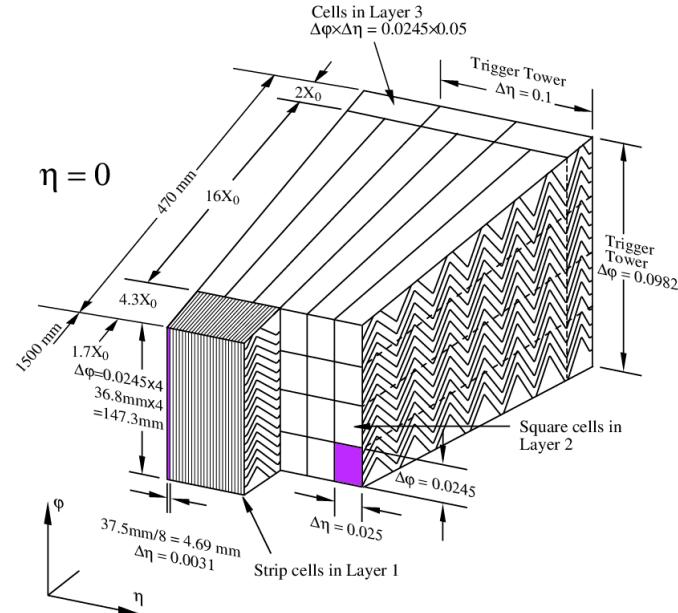


Fig. 3.12 Schematic diagram of a LAr EM calorimeter barrel module.

1111

1112 Tile calorimeter

1113 Tile calorimeter is a sampling calorimeter using scintillating plates as active medium
 1114 and steel as absorber. It consists of three sections: the central barrel with the pseudorapidity range of $|\eta| < 1.0$ and two extended barrels with $0.8 < |\eta| < 1.7$. Figure 3.13 shows
 1115 the design of one tile calorimeter module. It's used for energy reconstruction of jets and
 1116 E_T^{miss} measurement by combining the measurements with the end-cap and forward LAr
 1117 hadronic calorimeter.
 1118

1119

3.2.6 Muon spectrometer

1120 Muon spectrometer^[30] is the outermost part of the ATLAS detector with an extremely
 1121 large tracking system. It measures a large range of muon momentum, and the accuracy is
 1122 about 3% at 100 GeV and 10% at 1 TeV. The muon spectrometer comprises three main
 1123 parts: a magnetic field produced by three toroidal magnets; a set of chambers measur-
 1124 ing the tracks of muons with high spatial precision; and triggering chambers with accu-
 1125 rate time-resolution. Figure 3.14 shows the schematic of ATLAS muon spectrometer that
 1126 consists of four types of muon chambers (*MDT*, *CSC*, *RPC*, *TGC*) as well as the magnet
 1127 systems (barrel and end-cap toroid).

1128

More details of four chambers are given as below:

1129

- **Monitored Drift Tubes (MDT).** MDTs provide the precise momentum measure-

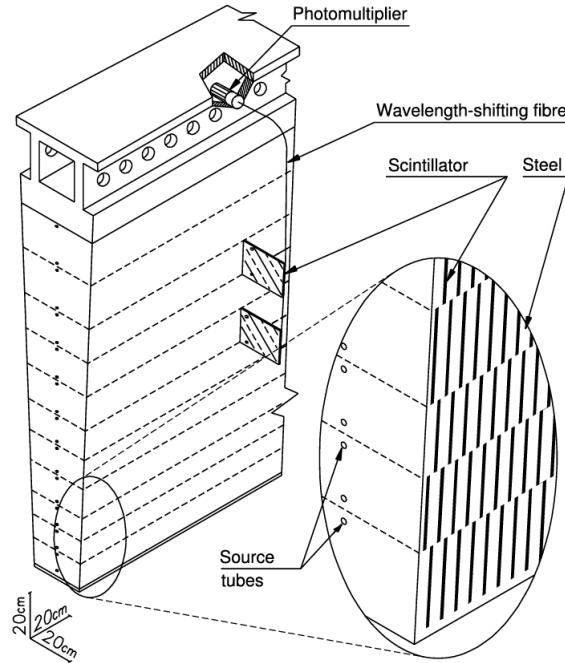


Fig. 3.13 Schematic diagram of tile calorimeter module^[29].

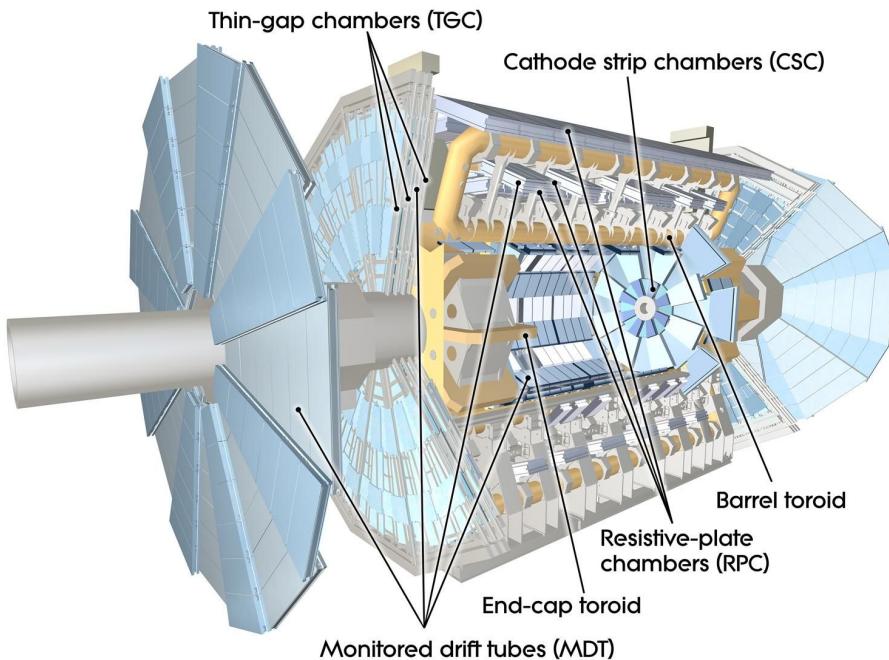


Fig. 3.14 Cut-away view of the ATLAS muon spectrometer^[31].

ment with the $|\eta|$ range up to 2.7, except in the innermost end-cap layer where the coverage is limited to $|\eta| < 2.0$. The chambers comprises three or four layers of drift tubes, with a diameter of 29.970 mm, operated with Ar/CO₂ gas (93/7) at 3 bar. The average resolution can reach 80 μm per tube and 30 μm per chamber.

- **Cathode strip chambers (CSC).** CSCs are used in the forward region of $2 < |\eta| < 2.7$ in the innermost tracking layers, due to their good time resolution and high

1136 rate capability. The CSCs are multi-wire proportional chambers (MWPC) with the
1137 cathode planes segmented into strips in orthogonal directions, which allows both
1138 coordinates to be measured from the induced-charge distribution. The resolution of
1139 a chamber is about $40 \mu\text{m}$ for bending plane and 5 mm for the transverse plane.

- 1140 • **Resistive plate chambers (RPC).** The RPCs serve as fast triggers in the barrel
1141 region of $|\eta| < 1.05$ due to the high rate capability and good spatial and time res-
1142 olution. It is a gaseous parallel electrode-plate detector without any wires. There
1143 are three concentric cylindrical layers around the beam axis, as three trigger sta-
1144 tions. Each station consists of two independent layers to measure the transverse
1145 coordinates of η and ϕ .
- 1146 • **Thin gap chambers (TGC).** TGCs are used as trigger system for the end-cap region
1147 of $1.5 < |\eta| < 2.4$, and works based on the same principle as multi-wire propor-
1148 tional chambers. In addition, they can also provide the second azimuthal coordinate
1149 to complement the measurement of MDT in bending direction.

1150

3.2.7 Trigger system

1151 Trigger system in ATLAS is a very essential component, which is responsible for
1152 deciding whether to keep a given collision event for later study or not. In the LHC run-
1153 2, higher energy, luminosity and pile-up lead to a large increase of event rate by up to
1154 a factor of five, which causes to a even larger challenge and more strict requirement of
1155 trigger system.

1156 The trigger system in run-2 consists of a hardware-based first level trigger (Level-
1157 1) and a software-based high level trigger (HLT)^[32]. As depicted in figure 3.15, in
1158 Level-1, the inputs from coarse granularity calorimeter and muon detector information
1159 together with some other subsystems are sent to the Central Trigger Processor to deter-
1160 mine Regions-of-Interest (RoIs) in the detector. The event rate can be reduced by Level-1
1161 triggers from 30 MHz to 100 kHz. After that, the RoI information from Level-1 is sent to
1162 HLT, in which more sophisticated selection algorithms are run for regional reconstruction.
1163 The HLT reduces the rate from Level-1 from 100 kHz to about 1 kHz on average. At the
1164 end, the events that accepted by HLT are transferred to local storage at experimental site
1165 for offline reconstruction. Details about Level-1 and HLT trigger systems are described
1166 as below:

1167

Level-1 trigger

1168 Substantial upgrades have been delivered in ATLAS Level-1 trigger system for run-2
1169 data taking. The upgrades took place in both hardware and detector readout, allow the

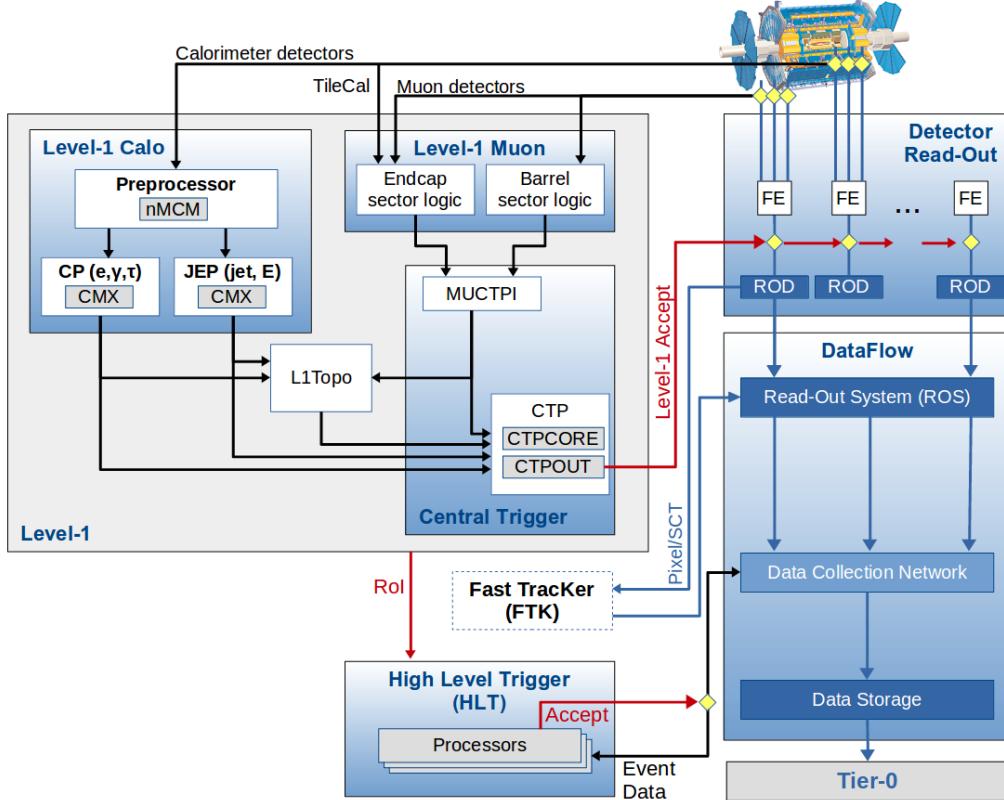


Fig. 3.15 Schematic diagram of the ATLAS trigger and data acquisition system in run-2.

trigger rate increasing from 70 kHz (in run-1) to 100 kHz (in run-2). There are two major parts of Level-1 triggers, including Level-1 calorimeter (L1calo) trigger and Level-1 muon (L1mu) trigger.

Level-1 Calorimeter trigger uses the reduced granularity information from the electromagnetic and hadronic calorimeters to search for electrons, photons, taus and jets and missing transverse energy (E_T^{miss}). It can identify an Region-of-Interest (RoI) as a 2×2 trigger tower cluster in the EM calorimeter as shown in figure 3.16, and 4×4 or 8×8 trigger tower for Jet RoIs. One important upgrade was that, the new FPGA-based (field-programmable gate array) Multi-Chip Modules are used to replace the ASICs (application-specific integrated circuits) included in the modules used in run-1, which allows the use of auto-correlation filters to suppress pile-up.

The Level-1 Muon trigger system includes one barrel section (RPC) and two end-cap section (TGC), which provides fast trigger signals from the muon detectors for the Level-1 trigger decision. By requiring a coincidence with hits from the innermost muon chambers, it can reduce the L1_MU15 rate by about 50% in the region of $1.3 < |\eta| < 1.9$ with only a loss of around 2% signal efficiency. In addition, the coverage was extended by around 4% due to installing new chambers in the feet region of the muon detector.

1187 High Level Trigger

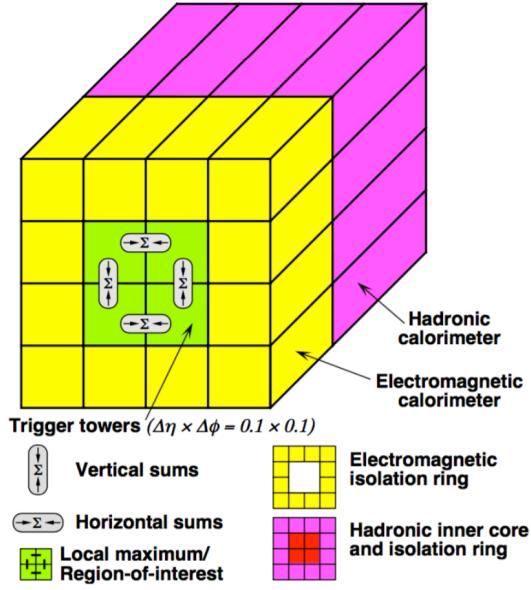


Fig. 3.16 An examples of L1 calorimeter trigger tower for electron and photon triggers^[33].

1188 In run-1, the Event Filter computer clusters and Level-2 trigger system were separated,
 1189 while now in run-2, they have been merged into a single HLT event processing. The new
 1190 arrangement helps to reduce the complexity and duplication of algorithm, which leads to
 1191 a more flexible high level trigger system. During the long-shutdown between the LHC
 1192 run-1 and run-2, lots of reoptimizations have been done for trigger reconstruction algo-
 1193 rithms as well as the offline analysis selections, which can improve the efficiency by more
 1194 than a factor of two in some cases like hadronic tau triggers. For some triggers, the HLT
 1195 processing performed within RoIs also allows to aggregate from RoIs to single objects.
 1196 This improvement reduces the CPU processing for events with overlapping RoIs, and the
 1197 average output rate has been increased from 400 Hz to 1 kHz.

1198 The HLT reconstruction algorithm can be divided into fast and precision online recon-
 1199 struction steps. As illuminated by figure 3.17, the initial fast reconstruction helps to reduce
 1200 the event rate, and to seed into precision reconstruction. Then the final online precision
 1201 reconstruction is improved and uses offline-like algorithms as much as possible. In partic-
 1202 ular, multivariate analysis techniques (based on machine learning) have been introduced
 1203 online in many aspects.

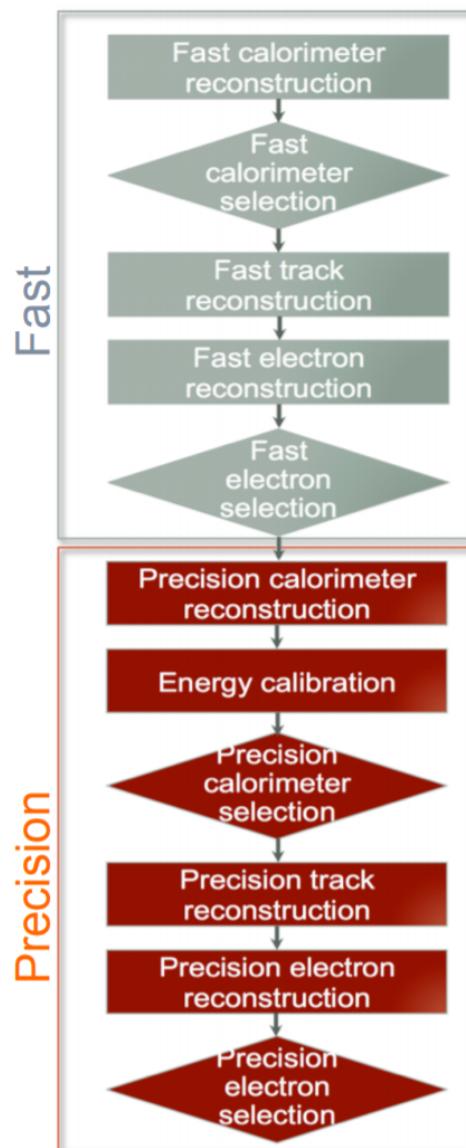


Fig. 3.17 The HLT trigger algorithm sequence^[33].

1204 Chapter 4 Simulation and Event Reconstruction for 1205 the ATLAS Experiment

1206 In current LHC pp collision, bunches of protons collide every 25 nanoseconds (ns),
1207 which gives a large challenge to event reconstruction and selections. To predict and model
1208 each process, Monte Carlo simulations of physics events are essential for high-energy
1209 physics experiments. This section will briefly discuss the event simulation and recon-
1210 struction programs based on the ATLAS software framework.

1211 4.1 Event simulation

1212 The ATLAS simulation program is integrated into the ATLAS software framework
1213 called *Athena*^[34], which uses Python as an object-oriented scripting and interpreter lan-
1214 guage to configure and load C++ algorithms and objects. Figure 4.1 shows the overview
1215 of ATLAS simulation data flow^[35]. In the diagrams, the square-cornered boxes represents
algorithms and applications to be run and round-cornered boxes denote data objects.

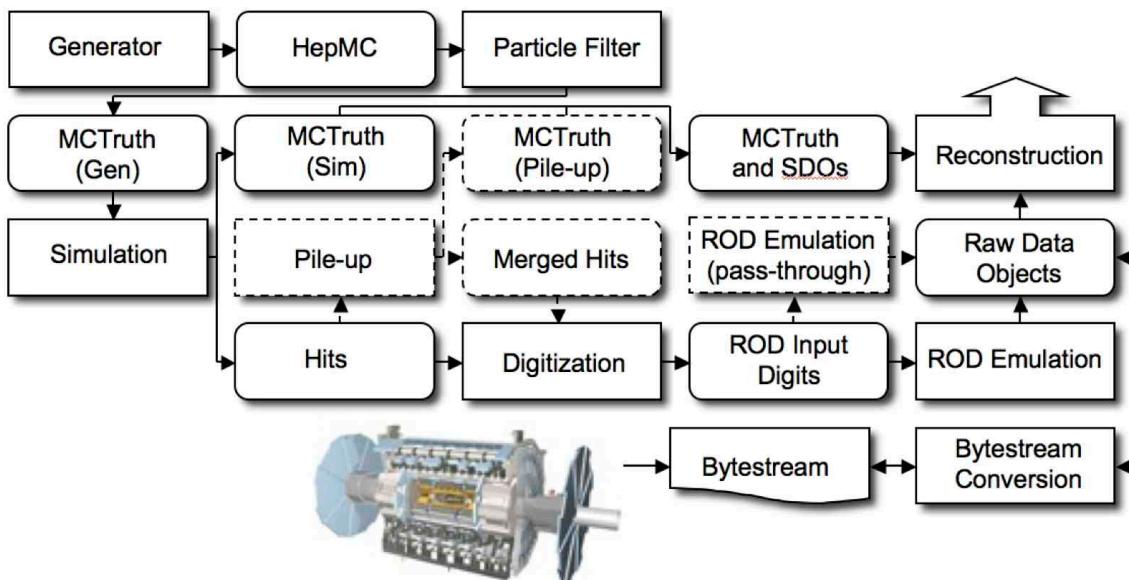


Fig. 4.1 The flow of the ATLAS simulation software.

1216
1217 First of all, events are produced by MC generators in standard HepMC format and
1218 then read into the simulation. During the simulation, particles are propagated through the
1219 full ATLAS detector whose configurations can be set by users via GEANT4 toolkit. The
1220 energies deposited in the sensitive regions of the detector are recorded as *hits* that contains
1221 the total energy deposition, position and time, and are written to a simulation hit file. In

the meantime, the events in “truth” format are also recorded to contain the history of the interactions from the generator, including incoming and outgoing particles. Simulated Data Objects (SDOs) are created from truth, which are maps between hits in sensitive portions of the detector and truth information of particles in simulation. The files are then sent to digitization, with constructs “digits” inputs and be written into Raw Data Object (RDO) file used for reconstruction.

In conclusion, there are three main parts of framework: *Generation*, *Simulation* and *Digitization*. More details are given as below:

Event generation

As shown in figure 4.2^[36], at hardon colliders, multiple scattering and rescattering effects arise, which needs to be simulated by Monte Carlo (MC) event generators to reflect the full complexity of those event structures. Several MC event generators can be used to generate events in HepMC format. The events can be filtered at generation time with some certain requirements (eg. decay channel or missing energy above a certain threshold). The generator is responsible for any prompt decays (e.g. W or Z bosons) and stores any “stable” particle expected to propagate through a part of the detector. During the generation steps, any interactions with detector are ignored and only immediate decays are considered.

There are several MC generators that have been widely used with general purpose, including SHERPA^[37], HERWIG++^[38], POWHEGBox^[39], MC@NLO^[40] and PYTHIA8^[41].

Simulation

GEANT4 is used as standard simulation toolkit for the ATLAS experiment, which transports physics particles through the detector’s geometry. During the generation level, the entire connected chain of the HepMC event is stored as the Monte Carlo truth. Only the stable particles are read into GEANT4 for further simulation and selection, and transformations can be applied to these events to select certain processes. During the simulation, many secondary tracks can be produced, therefore only information from the interactions of interest are stored, including the incoming particles, step sequence, vertex as well as outgoing particles. The output of GEANT4 is called *hit file*, which contains metadata describing the configuration of the simulation during the run, all requested truth information and a collection of hits for each subdetector.

Since the standard ATLAS detector simulation cost very large computing resources to accurately model the complex detector geometry and physics descriptions, some fast simulation programss are developed according to different user purpose. Some popular fast-sim toolkits include *Fast G4 Simulation*^[42], *ATLFAST-I*^[43] and *ATLFAST-II*^[44].

Digitization

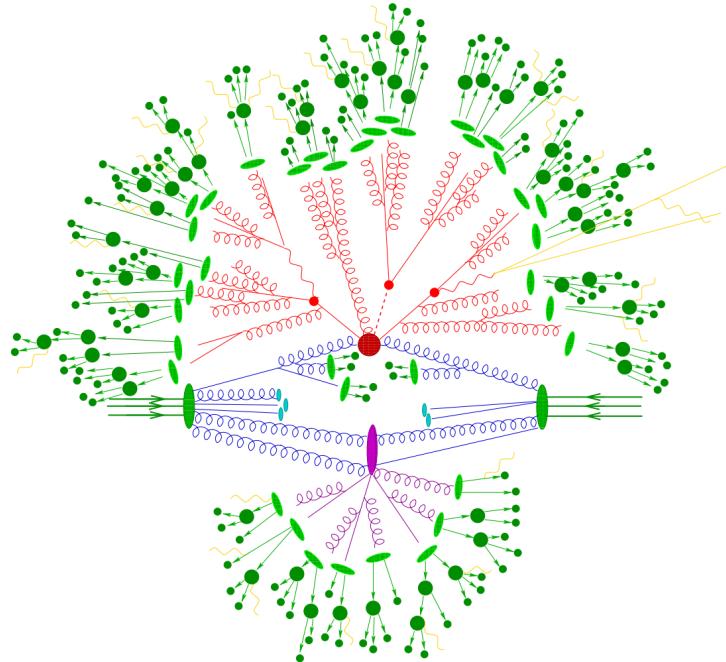


Fig. 4.2 Sketch of a hardon-hardon collision simulated by MC event generator. The red blob in center denotes the hard collision, surrounded by tree-like structures representing Bremsstrahlung which is simulated by Parton Showers. The purple blob stands for a secondary hard scattering event. The light green blobs indecate the parton-to-hardon transitions and the dark green blobs represents hadron decays. The yellow lines are soft photon radiations.

1257 The hit outputs from simulated events, including hard scattering signal, minimum
 1258 bias, beam halo, beam gas and cavern background events, are then sent into digitization
 1259 procedure, converted into detector response called “digits”. Before converted into detector
 1260 signal as “digits” formart, each type of event can be overlaid at a user-specified rate. Those
 1261 overlay, called “pile-up”, can be done during digitization to save the CPU time. At this
 1262 stage, the detector noise and the first level trigger that implemented with hardware on
 1263 the real detector are added into events. The digitization firstly constructs “digits” inputs
 1264 to the readout drivers (RODs) in the detector electronics. Then the ROD functionality
 1265 is emulated, and the output digits are written out as Raw Data Object (RDO) file. In
 1266 addition, the digitization algorithms can also produce Simulated Data Objects (SDOs),
 1267 which contain information about all the particles, noise and the amount of energy that
 1268 contributed to the signal. Then all information are sent into reconstruction level described
 1269 in next subsection.

4.2 Event reconstruction

The data flow of ATLAS data processing is sketched in figure 4.3^[45]. Data from detector is firstly filtered by online trigger system before sending to the *Tier-0* (*T0*) for initial processing by offline reconstruction software based on Athena. A small amount of data named “express stream” is processed in almost real time in *T0* for online data quality monitoring. In addition, some other dedicated data streams are sent out at trigger level for detector alignment and calibration. These calibration and alignment information are then used for bulk reconstruction in *T0*. At the end of the reconstruction chain, the data are delivered into *Tier-1* (*T1*) and *Tier-2* (*T2*) centers around the world for further analysis and production of simulated data. *T1* centers are also responsible for data reprocessing by re-running data reconstruction with improved calibration and alignment constants and with improved reconstruction algorithms.

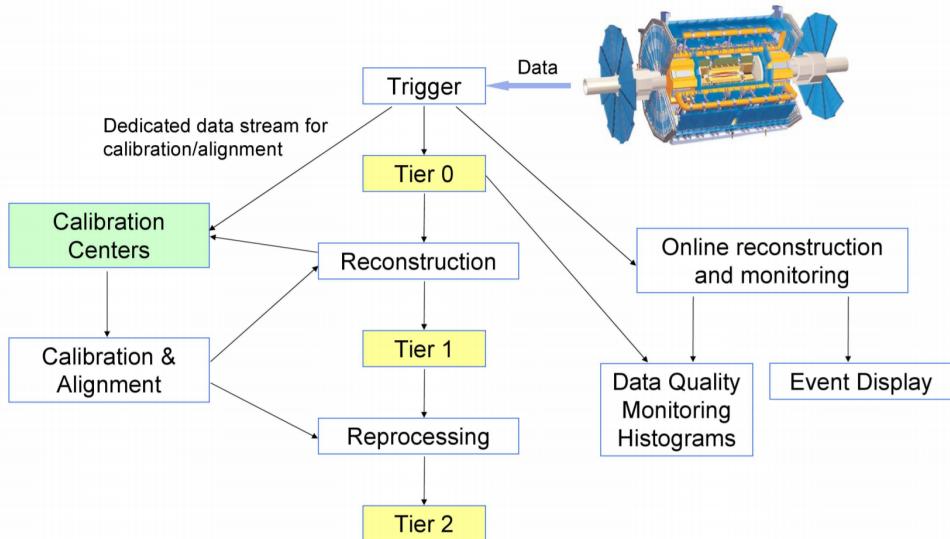


Fig. 4.3 The flowchart of the ATLAS data processing.

This section describes the reconstruction of some important physics objects in ATLAS experiment: tracks, vertices, electrons, muons, jets, and missing energies.

4.2.1 Track

The ATLAS detector is composed of two independent tracking systems: the Inner Detector (ID) close to the interaction point, and the Muon Spectrometer (MS) located in the outermost region. The reconstructed charged-particle trajectories in the ID and MS are referred to as ID tracks and MS tracks respectively. The challenge of ID reconstruction is that it needs to handle high track density that imposes a large number of combinatorial track candidates, while the MS reconstruction is however largely limited by the huge

amount of inert material, the large background and the highly inhomogeneous magnetic field^[46]. More details of these two types of track are given as below:

Inner detector track

Figure 4.4 sketches the ID system used for detecting charge-particle tracks. The ID track reconstructions contains two sequences: *inside-out* track reconstruction and *outside-in* one.

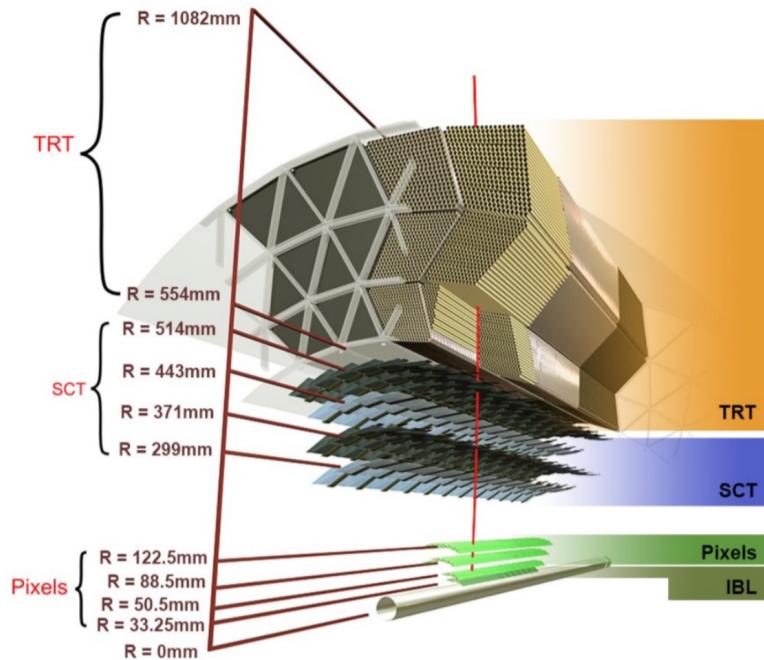


Fig. 4.4 Schematic view of the ATLAS inner detector showing all the corresponding components.

For inside-out tracking, it exploits the high granularity of the pixel and SCT detectors to discover prompt tracks originating from the interaction point. In first step, the track seeds are formed by combining the information of space-points in the three pixel layers and the first SCT layer. Then, these seeds are extended throughout the SCT to build track candidates. After that, these candidates are fitted with some quality cuts applied to remove the outlier clusters, reject the fake tracks and resolve ambiguities in the cluster-to-track association. The selected tracks are then further extended to TRT, and refitted with the full information from pixel, SCT and TRT detectors.

Another complementary approach, outside-in, searches for unused track segments start from TRT instead. These segments are then extended into the SCT and pixel detectors to improve the tracking efficiency for secondary tracks from conversions or decays of long-lived particles.

Muon spectrometer track

1310 The MS track reconstruction^[47] starts from searching hit patterns inside each muon
 1311 chamber to form segments. In each MDT chamber and nearby trigger chamber, a Hough
 1312 transform^[48] is used to search the hits lie on a certain trajectory in the bending plane of the
 1313 detector. The MDT segments are reconstructed by performing a linear fit to the hits found
 1314 in each layer. The RPC or TGC hits can be built by measuring the coordinate orthogonal to
 1315 the bending plane. And the segments of CSC can be built using a separate combinatorial
 1316 search in the η and ϕ detector planes.

1317 Then muon track candidates are built by fitting hits from segments in different layers
 1318 together. This task makes use of the algorithm by performing a segment-seeded combi-
 1319 natorial search, which starts by using the segments generated in the middle layers of the
 1320 detector where more trigger hits are available as seeds. The search is then extended to use
 1321 the segments as seeds from the inner and outer layers. The segments are selected based
 1322 on criteria of hit multiplicity and fit quality, and are matched using their relative positions
 1323 and angles. To build a track, at least two matching segments are required, except in the
 1324 barrel-endcap transition region where a single high-quality segment with η and ϕ infor-
 1325 mation can be used to build a track. At beginning, the same segment can be used to build
 1326 more than one track candidates. Later on, an overlap removal algorithm is performed to
 1327 select the best assignment to a single track, or decide whether allows the certain segment
 1328 to be shared between two tracks.

1329 The hits associated with each track candidate are then fitted using a global χ^2 fit. The
 1330 algorithm accepts the track candidate if its fitting χ^2 passes the selection criteria. Hits
 1331 with large contribution to χ^2 are removed and the track fit is repeated. In addition, the
 1332 algorithm performs a hit recovery procedure looking for additional hits consistent with
 1333 the candidate trajectory, and the track candidate is refit if additional hits are found.

1334 4.2.2 Primary vertex

1335 The reconstruction of primary vertex (PV) uses the reconstructed tracks introduced in
 1336 previous section as inputs. The tracks must satisfy the following criteria^[49]:

- 1337 • $p_T > 400$ MeV
- 1338 • $|\eta| < 2.5$
- 1339 • Number of silicon hits $\geq \begin{cases} 9 & \text{if } |\eta| \leq 1.65 \\ 11 & \text{if } |\eta| > 1.65 \end{cases}$
- 1340 • IBL hits + B-layer hits ≥ 1
- 1341 • A maximum of 1 shared module (1 shared pixel hit or 2 shared SCT hits)
- 1342 • Pixel holes = 0

1343 • SCT holes ≤ 1

1344 A candidate vertex is formed by requiring two tracks passing these selection criteria.

1345 The reconstruction of PV can be divided into two steps^[50]: vertex finding and vertex
1346 fitting. The first step is the pattern recognition process, namely the association of recon-
1347 structed tracks to vertex candidates. The latter one works on the reconstruction of the
1348 actual vertex position and its covariance matrix. More details are described as below:

1349 First of all, a set of tracks passing the selection criteria mentioned above is selected.
1350 Then a seed position for the first vertex is chosen. This seed position is determined by beam
1351 spot in the transverse plane. The starting point for x- and y- coordinates are directly from
1352 the centre of the beam spot, while the one for z-coordinate is calculated as the mode of
1353 z-coordinates of tracks at their respective points with closest approach to the reconstructed
1354 centre of the beam spot.

1355 After determining the seed position, the iterative primary vertex finding procedure
1356 starts. An vertex fitting algorithm is adopted to find the optimal vertex position by per-
1357 forming an iterative χ^2 minimization, in which the seed position is used as the start point
1358 and the reconstructed tracks are used as input measurements. For this fitting procedure, the
1359 input tracks are assigned weights to reflect their compatibility with the vertex estimation,
1360 and the vertex position is re-calculated based on these weighted tracks. Then the iterative
1361 procedure is repeated by re-calculating the track weight according to the new vertex posi-
1362 tion. After the last iteration, the final weight of each track used in vertex fit is estimated.
1363 And those incompatible tracks ($> 7 \sigma$) are then rejected from this vertex candidate and
1364 moved back to the unused pool for next vertex finding. Then iteration procedure describes
1365 above are repeated again by using the remaining tracks, until no un-associated tracks are
1366 left or no additional vertex can be found in remaining tracks.

1367 At the end, the vertices with at least two associated tracks passing through are treated
1368 as possible PV candidates. And the output of this vertex reconstruction algorithm is the in-
1369 formation of three dimensional vertex positions and their covariance matrices. In physics
1370 analysis, it's most often to choose the one with highest sum of transverse momentum
1371 ($\sum p_T^2$) as PV.

1372 4.2.3 Electron

1373 Many interesting physical processes are with the involvement of one or more elec-
1374 trons (or positrons) at the LHC. But these electrons can be subjected to large amount of
1375 backgrounds such as hadrons, non-prompt electrons from photon conversions and non-
1376 isolated electrons from heavy flavor hadon decays. It is therefore essential to efficiently

1377 reconstruct and identify electrons as well as, in the meantime, to keep high background
1378 rejection.

1379 In ATLAS, in central region, the electrons leave tracks in inner detector (ID) and
1380 deposit the energies in the electromagnetic (EM) calorimeter. Firstly the signals from
1381 calorimeter are used for L1 trigger system, and then combined with the information from
1382 ID tracks to reconstruct electron candidates that will be used for the high level trigger
1383 (HLT) decision algorithms^[51]. The backgrounds mentioned above can then be further
1384 suppressed by using several identification criteria. In addition, electrons are required to
1385 be isolated from other activities to be further distinguished from background.

1386 More details of electron *reconstruction*, *identification* and *isolation* are described as
1387 below.

1388 **Electron reconstruction**

1389 Several steps are proceeded for electron reconstruction in the central region of ATLAS
1390 detector ($|\eta| < 2.47$):

1391 1. **Seed-cluster reconstruction:** A sliding window with size of 3×5 in unit of
1392 $\Delta\eta^{tower} \times \Delta\phi^{tower} = 0.025 \times 0.025$ in $\eta \times \phi$ space is utilized to search for elec-
1393 tron cluster seeds with total cluster transverse energy greater than 2.5 GeV. Then
1394 a clustering algorithm^[52] is applied to form the clusters around the seeds, which
1395 can take advantage of removing the duplications. The kinematics of clusters are
1396 then reconstructed by using an extended window depending on the cluster posi-
1397 tion. The efficiency of cluster search is from about 95% at $E_T = 7GeV$ to 99% for
1398 $E_T \geq 15GeV$.

1399 2. **Track reconstruction:** The track reconstruction can be divided into two steps: pat-
1400 tern recognition and track fit. The standard pattern recognition in ATLAS uses pion
1401 hypothesis for energy loss caused by interactions with detector material. If a track
1402 seed with $p_T > 1$ GeV cannot be successfully extended to a full track required at
1403 least seven hits using this pion hypothesis, but still falls inside one of the EM clus-
1404 ter region of interest, as a second attempt, the pattern recognition using electron
1405 hypothesis is then used to allow larger energy loss. Depending on the pattern used
1406 in previous stage, the track candidates are then fitted with either the pion hypothesis
1407 or the electron hypothesis by using ATLAS Global χ^2 Track Fitter^[53]. If a track
1408 candidate fails the fit by using pion hypothesis, it can be refit with the electron hy-
1409 pothesis again. In this method, a specific electron-oriented algorithm is integrated
1410 into the ATLAS standard track reconstruction, which improves the performance for
1411 electron and as well as maintain minimal interference with the main track recon-

1412 struction.

1413 **3. Electron specific track fit:** Once the tracks are obtained, they are loosely matched
1414 to EM cluster using the distance in η and ϕ between the position of track (after ex-
1415 trapolation) in calorimeter's middle layer and the cluster barycentre. The matching
1416 conditions take into account the energy loss of bremsstrahlung and the number of
1417 precise hits in silicon detector.

1418 **4. Electron candidate reconstruction:** The electron candidate is reconstructed by
1419 matching the track candidate to EM cluster seed to eventually completes the electron
1420 reconstruction procedure. If more than one track satisfy the matching condition,
1421 one track is chosen as primary track based on the information of the cluster-track
1422 distance R, the number of pixel hits and the presence of a hit in the first silicon
1423 layer^[54]. In addition, the electron candidates are removed from electron pool if it's
1424 without any associated precise hit tracks, and moved into photon candidates pool.
1425 Then we reformed the electron clusters by using 3×7 (5×5) longitudinal towers
1426 of cells in barrel (end-caps) in EM calorimeter. The measured energy is calibrated
1427 to original electron energy based on MC simulated samples by using multivariate
1428 techniques (MVA).

1429 In addition, in physics analysis, to reduce the background from photon conversions
1430 and secondary particles, the track associated with electron is required to be compatible
1431 with the primary vertex of the hard collision. Practically, the impact parameters cuts such
1432 as $d_0/\sigma_{d_0} < 5$ and $z_0 \sin\theta < 0.5$ mm are usually applied, where d_0 is the closest distance
1433 of the track to the measured beam-line, z_0 is the distance along the beam-line between the
1434 point where d_0 is measured and the beam-spot position, and θ is polar angle of the track,
1435 σ_{d_0} denotes the estimated uncertainty of d_0 parameter. Figure 4.5 depicts the definition of
1436 each track impact parameter.

1437 **Electron identification**

1438 The electron identifications are applied to determine whether the reconstructed elec-
1439 tron candidate is more signal-like or background-like object. The identification algorithms
1440 make use of quantities of related variables from electron cluster and track measurements
1441 including calorimeter shower shapes, track properties, as well as variables measuring
1442 bremsstrahlung effects for distinguishing signal from background. Taking the advantage
1443 of new IBL in run-2, the number of hits in this innermost pixel layer is utilized for discrim-
1444 inating between electrons and converted photons. In addition, a likelihood method based
1445 on the TRT high-threshold hits is adopted to compensate the lower transition radiation
1446 absorption probability of the argon.

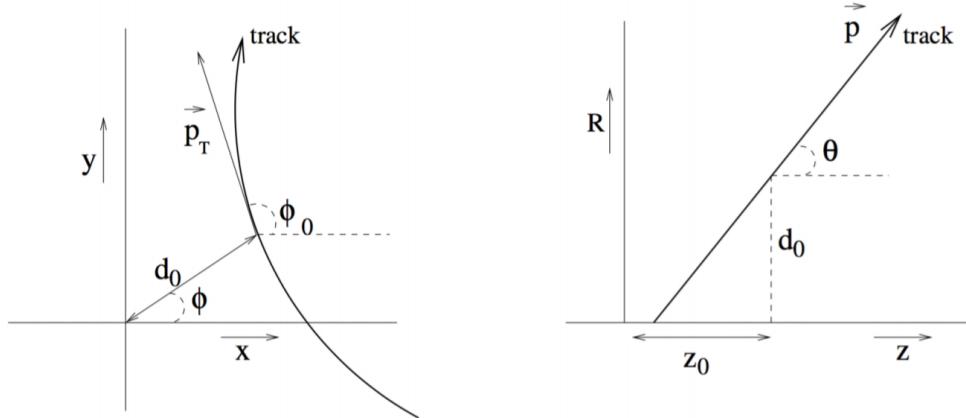


Fig. 4.5 Schematic of the impact parameters of a track in the transverse plane (left) and RZ-plane (right), as defined in the global ATLAS tracking frame^[55].

1447 The baseline identification algorithm introduced for ATLAS run-2 data analysis is the
 1448 likelihood-based (LH) method, which uses a MVA technique to simultaneously evaluate
 1449 several properties of electron candidates when making a decision. The LH method utilizes
 1450 the probability density functions (PDFs) of signal and background as the input discrimi-
 1451 nating variables. Based on these PDFs, it can calculate overall probabilities of the object
 1452 to be signal or background. Then the probabilities of signal and background are combined
 1453 together into a discriminant $d_{\mathcal{L}}$:

$$d_{\mathcal{L}} = \frac{\mathcal{L}_S}{\mathcal{L}_S + \mathcal{L}_B}, \quad \mathcal{L}_{S(B)}(\mathbf{x}) = \prod_{i=1}^n P_{s(b),i}(x_i) \quad (4.1)$$

1454 where \mathbf{x} denotes the vector of discriminating variables and $P_{s(b),i}(x_i)$ represents the value
 1455 of signal (background) PDF of the i^{th} variable as x_i .

1456 Three levels of working points (WPs) for electron identification are provided: *Loose*,
 1457 *Medium* and *Tight*, in order of increasing background rejection. Samples selected by a
 1458 looser WP are subsets of a tighter one, for example, the electrons passing Medium can
 1459 all be selected by Loose. The identification efficiency varies as function of transverse
 1460 energy (E_T) as shown in figure 4.6. For evaluations, the electron candidates from MC
 1461 simulation of $Z \rightarrow ee$ decays (di-jet) are used as signal (background). Depending on the
 1462 working point, the signal (background) efficiencies for reconstructed electron candidates
 1463 at $E_T = 25 GeV$ are in the range of 78 to 90% (0.3 to 0.8%), and increase (decrease) with
 1464 E_T .

1465 **Electron isolation**

1466 In addition to the identification criteria, most analyses have electron isolation require-
 1467 ment to further distinguish signal from background. To quantify the energy of particles

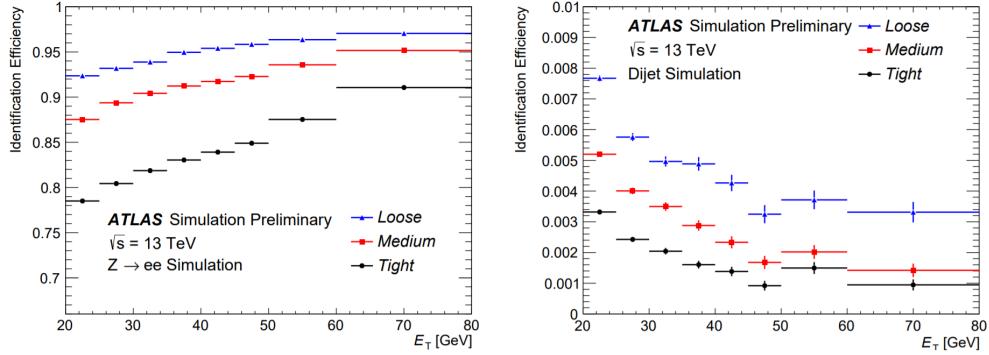


Fig. 4.6 The efficiencies of three electron identification WPs from $Z \rightarrow ee$ (left) events and hadrons misidentified as electrons estimated using di-jet MC samples (right).

around the electron candidate, the isolation variables can help to separate the prompt electron from other non-isolated electrons, like the electrons from converted photons or from heavy flavour hadron decays. There are two kinds of discriminating variables that have been designed:

- **Calorimeter-based variable:** $E_T^{topocone20}$. It's defined as the sum of transverse energies of topological clusters^[56], calibrated at EM scale within a cone of $\Delta R = 0.2$ around the candidate electron cluster. It only considers the clusters with positive reconstructed energy. In addition, a correction as a function of (E_T, η) values is applied to account for the electron energy leakage outside the cluster.

- **Track-based variable:** $p_T^{varcone20}$. It's calculated as the sum of transverse momentum of all satisfied tracks within a cone of $\Delta R = \min(0.2, 10\text{GeV}/E_T)$ around the candidate electron track. To calculate the sum, it requires the tracks are originating from the reconstruction PV of hard collision, and exclude the associated tracks of electron itself.

Based on the values of $E_T^{topocone20}/p_T$ and $p_T^{varcone20}/p_T$, a series of working points with different selection requirements are defined. The resulting WPs are divided into two kinds:

- Efficiency targeted working points: varying requirements to obtain a certain isolation efficiency, which can either be a constant or as a function of E_T .
- Fixed requirement working points: set the constant upper thresholds on isolation variables.

The distribution of two discriminating variables are shown in figure 4.7 for $ZZ \rightarrow ee$ events with $E_T > 27\text{GeV}$ and satisfying *Tight* requirement.

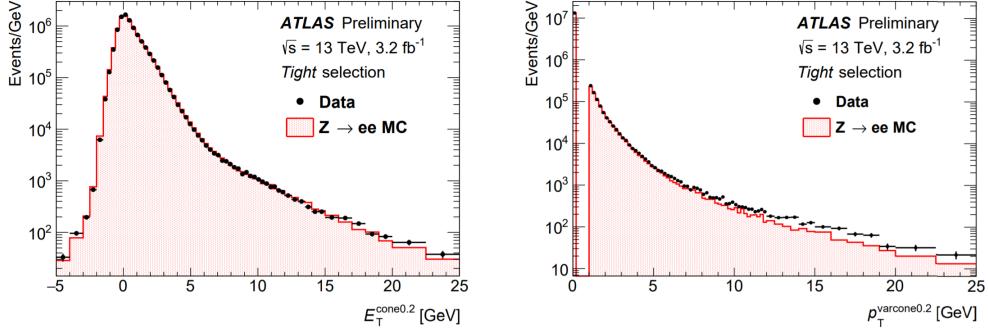


Fig. 4.7 Distributions of $E_T^{cone0.2}$ (left) and $p_T^{varcone0.2}$ (right) for electrons from $ZZ \rightarrow ee$ events in data and MC simulation. The simulated events (full histograms) are normalized to data.

4.2.4 Muon

Muons are distinctive signatures in final states of many physics analyses at the LHC including the Higgs analyses, SM measurements, BSM searches and so on. High performance of muon reconstruction and identifications are crucial. This section briefly describes some more details of the *reconstruction, identification and isolation* of muon.

Muon reconstruction

Muon reconstruction is firstly performed in inner detector (ID) and muon spectrometer (MS) independently as given in section 4.2.1. The information from each individual detector is then combined together to form the muon tracks for physics analyses. The combined ID-MS reconstruction is developed according to several algorithm based on the information from ID, MS and calorimeters. Four different muon types are defined^[47]:

- **Combined (CB) muons:** a combined track is formed by using the reconstructed tracks performed independently in ID and MS with a global refit. To improve the fit quality, the hits from MS may be added to or removed from the track. The outside-in pattern recognition is utilized for the reconstruction of most muons, in which the muons are first reconstructed in MS and then extrapolated inward to match the ID track. In the meantime, the inside-out pattern is also used as a complementary method.
- **Segment-tagged (ST) muons:** a reconstructed track in ID is defined as muon, if it can be associated with at least one track segment in MDT or CSC chambers. These ST muons are used when they can only pass across one layer of MS chambers due to their low p_T or falling into regions with less MS acceptance.
- **Calorimeter-tagged (CT) muons:** a reconstructed track in ID is categorized as muon if it's matched to the energy deposit in calorimeter which is recognized with

1514 a minimum-ionizing particle. This CT muons have lowest purity amount all types
1515 of muons, but it covers the region where ATLAS muon spectrometer is only par-
1516 tially constructed. For the region of $|\eta| < 0.1$ and $15GeV < p_T < 100GeV$, the
1517 identification of CT muons are optimal.

- 1518 • **Extrapolated (ME) muons:** the muon is reconstructed based only on the MS track
1519 and a loose requirement of originating from the interaction point. In general, this
1520 type of muon needs to pass at least two (three) layers of MS chambers to provide
1521 a track measurement in barrel (forward) region. ME muons are designed to extend
1522 the acceptance for muon reconstruction into the region $2.5 < |\eta| < 2.7$ where ID
1523 doesn't cover.

1524 Before collecting those muons for physics analyses, overlap removals are performed
1525 between different muon types with the priority of CB > ST > CT, when two types of muons
1526 share the same ID track. Besides, the overlaps with ME muons are resolved by analyzing
1527 the track hit content, and selecting the track with better fit quality and larger number of
1528 hits.

1529 **Muon identification**

1530 After reconstruction, the muon identification is then performed to further discrimi-
1531 nate between signal and background, especially to suppress backgrounds from pion and
1532 kaon decays by requiring prompt muons with high efficiency and guaranteeing a robust
1533 momentum measurement. The muon identification is defined by using the fit quality of
1534 combined track. The variables utilized in judgement for CB tracks include:

- 1535 • *q/p significance*, the absolute difference between q/g (charge over momentum) of
1536 muons measured in ID and MS divided quadratic sum of their corresponding un-
1537 certainties;
- 1538 • ρ' , the absolute value of difference between the p_T (transverse momentum) mea-
1539 sured in ID and MS, divided by the p_T of combined track;
- 1540 • *Nomalized χ^2* of the combined track fit;
- 1541 • *Number of hits in ID and MS*

1542 In addition, some new variables used for *LowPt* muon working point what will be de-
1543 scribed later^[57]:

- 1544 • *Momentum balance significance (MBS)* is computed as momentum difference be-
1545 tween the ID and MS standalone measurements with respect to the uncertainty σ on
1546 energy lost in the calorimeter system.
- 1547 • *Scattering neighbor significance (SNS)* is defined to estimated the significance of a
1548 change in trajectory along the track, expected in the presence of a hadron decaying

1549 to a muon.

- 1550 • *Scattering curvature significance (SCS)* is defined as the normalized integral of the
1551 scattering angle significances, corrected for large kinks along the trajectory.

1552 Five selection working points are developed to satisfy the different needs for different
1553 physics goals: *LowPt*, *Loose*, *Medium*, *Tight* and *HighPt*. The *Tight*, *Medium*, *Loose* are
1554 subsets from the tighter one to looser one. More detailed definition of each working point
1555 is given as follow:

1556 • *Loose*: this working point is designed to maximize the reconstruction efficiency
1557 while keeping good-quality of muon tracks. And they are specifically developed
1558 for reconstructing the Higgs boson candidates from four-lepton final states. All
1559 four muon types are used for this selection level. The CB and ME muons passing
1560 Medium WP that will mentioned below are all included into Loose category. In
1561 addition, the CT and ST muons are restricted to $|\eta| < 0.1$ region. In the range of
1562 $|\eta| < 2.5$, around 97.5% Loose muons are CB muons, and about 1.5% are CT while
1563 remaining 1% are ST muons.

1564 • *Medium*: this working point is the default criteria of muon identification in ATLAS.
1565 This selection minimizes the systematic uncertainties of muon reconstruction and
1566 calibration. In this category, we only use CB and ME muons. For CB muons, at least
1567 3 hits in at least two layers of MDT are required, except $|\eta| < 0.1$ region, in which
1568 tracks with ≥ 1 MDT layer but ≤ 1 MDT hole layer are allowed. For ME muons, at
1569 least 3 MDT/CSC layers are required. Furthermore, a loose cut on the compatibility
1570 between measured momentum in ID and MS is applied to reduce the fake muons
1571 from hadrons misidentification. Besides, the q/p-significance is required to be less
1572 than 7.

1573 • *Tight*: this working point is used to maximize the purity of muons but with sacrifice
1574 of some selection efficiency. Only CB muons with hits in ≥ 2 stations of MS and
1575 passing Medium criteria are selected. In addition, the normalized χ^2 of combined
1576 track fit should be smaller than 8. Then, a two-dimensional cut of q/p-significance
1577 and ρ' is adopted as a function of muon p_T to ensure tighter background rejection
1578 for momentum below 20 GeV, in which the fake rate is usually higher.

1579 • *High- p_T* : this set of selections aims to maximize the momentum resolution for
1580 tracks with $p_T > 100\text{GeV}$ region. The selection is especially optimized for search-
1581 ing high-mass Z' and W' resonances. CB muons satisfying Medium selection and
1582 with ≥ 3 hits in 3 MS stations are chosen. The specific region in MS where align-
1583 ment is suboptimal are removed as a precaution.

- 1584 • *Low- p_T* : this type of muon is newly designed for physics analyses with ATLAS
 1585 software release version 21. It's designed to obtain a optimal muon identification
 1586 with very low transverse momentum of $3GeV < p_T < 5GeV$, which is crucial
 1587 for B-physics measurement in ATLAS. In this muon requirement, only CB muons
 1588 are used. In the range of $|\eta| < 1.3$, it requests muons hit at least one MS station; in
 1589 $1.3 < |\eta| < 1.55$, a least two MS stations are required; while in region of $|\eta| > 1.55$,
 1590 *Medium WP* is required. In addition, cuts are applied to suppress fakes as: $|MBS| <$
 1591 3.0 , $|SNS| < 3.0$ and $|SCS| < 3.0$.

1592 Figure 4.8 and 4.9 show the selection efficiency of different muon identification work-
 1593 ing points. For *Medium (Loose)*, *Tight* and *High- p_T* : $Z \rightarrow \mu\mu$ events with $p_T > 10GeV$
 1594 are used for measurement. In the top figure of figure 4.8, the efficiency of the Loose se-
 1595 lection (squares) is shown comparing to Medium one, where significant difference can be
 1596 observed in region of $|\eta| < 0.1$. For *LowPt*, $J/\Psi \rightarrow \mu\mu$ events with $3GeV < p_T < 10GeV$
 are used for measurement.

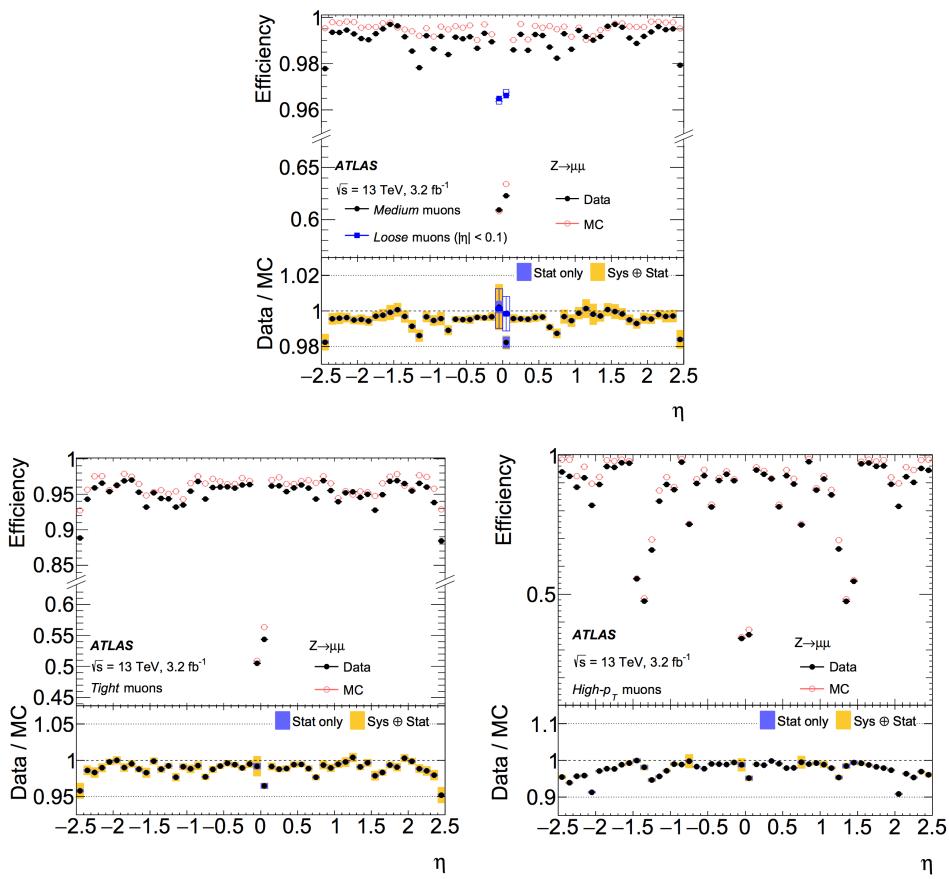


Fig. 4.8 Muon reconstruction efficiency as a function of η for: Medium (and Loose), Tight and High- p_T working points.

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1598

Muon isolation

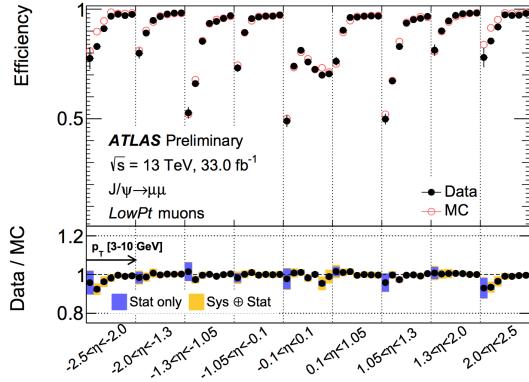


Fig. 4.9 Muon reconstruction efficiency for Low- p_T working point as a function of η .

Similar as electron, the muon isolation is used to further distinguish the prompt muon from non-prompt backgrounds. There are also two types of isolation variables for muon:

- **Calorimeter-based variable:** $E_T^{topocone20}$. It's defined as the sum of the transverse energy of topological clusters within a cone of size $\Delta R = 0.2$ around the candidate muon, after subtracting the contribution from the energy deposit of the muon itself and correcting for pile-up effects. The contributions from pile-up and underlying events are computed using the ambient energy-density technique^[58] and are corrected on an event-by-event basis.
- **Tracked-based variable:** $p_T^{varcone30}$. It's computed as the scalar sum of the transverse momenta of the tracks with $p_T > 1\text{GeV}$ in a cone size of $\Delta R = \min(10\text{GeV}/p_T^\mu, 0.3)$ around the candidate muon whose transverse momenta is p_T^μ after excluding the muon track itself. This p_T -dependent cone size can help to improve the performance for muons produced in the decay of particles with a large transverse momentum.

Then the isolation selections are applied based on *relative isolation variables*, which are computed as the ratio of the track- or calorimeter-based isolation variables to the transverse momentum of the muon. Figure 4.10 shows the distribution of those relative isolation variables by using $Z \rightarrow \mu\mu$ events and passing *Medium* identification criteria.

4.2.5 Jets

Jets are another important features for many physics analyses at the LHC, and especially the key signatures for vector boson fusion/scattering (VBF/VBS) processes. In ATLAS detector, jets are reconstructed as groups of topologically associated energy deposits in the calorimeters, tracks associated with charged particles measured in the inner tracking detector, or simulated particles. This section introduces the jet reconstruction, jet

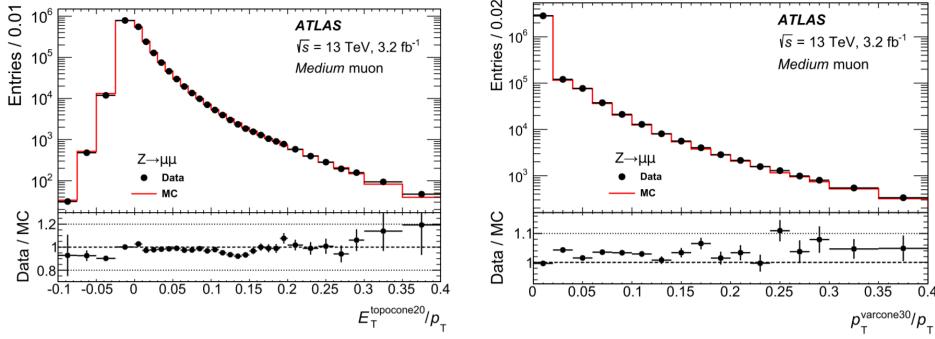


Fig. 4.10 Distributions of the calorimeter-based (right) and the track-based (left) relative isolation variables measured in $Z \rightarrow \mu\mu$ events.

1623 energy scale (JES) calibration and the b-jet tagging technical.

1624 Jet reconstruction

1625 Jets are reconstructed using anti- k_t algorithm^[59] and with radius parameter of $R = 0.4$
 1626 in most cases. The FASTJET software package^[60] is utilized for jet finding and reconstruc-
 1627 tion. A collection of four-vectors are used as inputs at each combination step in jet clus-
 1628 tering, the total four-momentum is therefore computed as the sum of four-vector of all its
 1629 constituents. There are three types of jets in ATLAS:

- 1630 • *Truth jets*: the inputs to jet algorithm are simulated particles.
- 1631 • *Track jets*: the inputs are charged tracks measured from inner detector.
- 1632 • *Calorimeter jets*: the inputs are energy deposits in calorimeters.

Figure 4.11 shows the schematic of ATLAS jet reconstruction.

Jet reconstruction

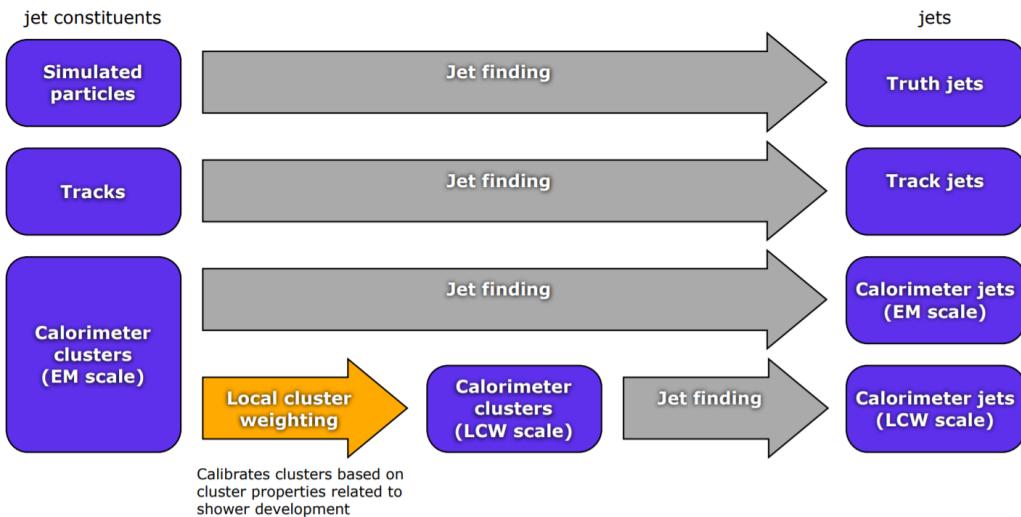


Fig. 4.11 A overview schematic of ATLAS jet reconstruction^[61].

1633

1634 The *truth jets* are reconstructed using anti- k_t algorithm with $R = 0.4$ by using final-
 1635 state, stable particles from MC simulation as inputs. It requires the candidate particles with

1636 lifetime $c_\tau > 10$ mm and excludes the particles from pile-up. Truth jets with $p_T > 7 GeV$
 1637 and $|\eta| < 4.5$ are then used for jets calibration described later.

1638 The *track jets* are reconstructed from charged particles within the full acceptance of
 1639 inner detector ($|\eta| < 2.5$). The track reconstruction has been introduced in section 4.2.1.
 1640 Reconstructed jets with $p_T > 500 MeV$ and associated with primary vertex are then se-
 1641 lected. Tracks are assigned to jets using ghost association^[58], a procedure that treats se-
 1642 lected tracks as four-vectors of infinitesimal magnitude during the jet reconstruction and
 1643 assigns them to the jet which they are clustered with. In addition, muon track segments
 1644 are used as a compensation for those uncaptured jet energy carried by energetic particles
 1645 passing through the calorimeters without being completely absorbed. Similar to the ID
 1646 track, muon segments are assigned to jets using the method of ghost association mentioned
 1647 above as well.

1648 The *calorimeter jets* are reconstructed using a set of three-dimensional, positive-
 1649 energy topological clusters (topo-clusters) made of calorimeter cell energies as input to
 1650 the anti- k_t algorithm^[62]. Topo-clusters are built from near-by calorimeter cells that con-
 1651 tains a significant energy above a noise threshold, which is estimated from measurements
 1652 of calorimeter electronic noise and simulated pile-up noise. Those calorimeter cell ener-
 1653 gies are measured at electromagnetic energy scale (EM scale) corresponding to the energy
 1654 deposited by electromagnetically interacting particles. And jets passing a p_T threshold of
 1655 7 GeV are reconstructed with the anti- k_t algorithm.

1656 Jet energy scale calibration

1657 Figure 4.12 depicts an overview of ATLAS jet calibration scheme for EM-scale
 1658 calorimeter jets. In this procedure, the jet energies are scaled to truth jets, which is recon-
 1659 structed at the particle-level. Each step of the calibration corrects the full four-momentum
 unless otherwise stated, scaling the jet p_T , energy, and mass.

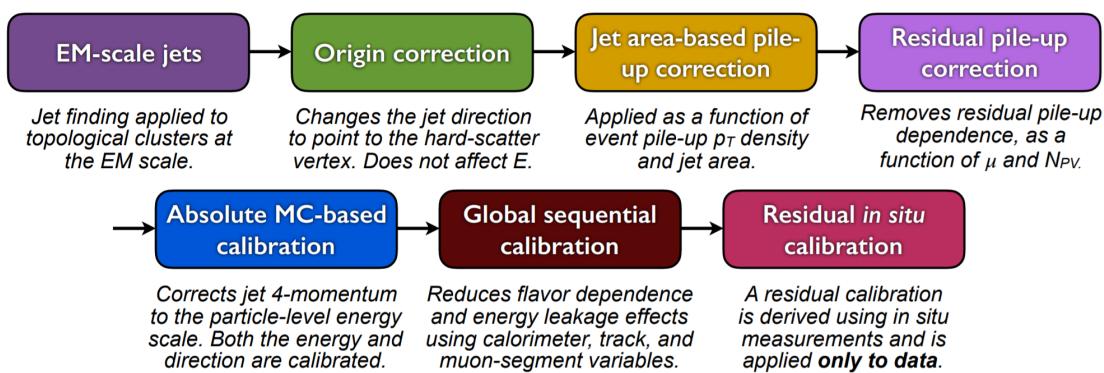


Fig. 4.12 A overview schematic of ATLAS jet calibration^[62].

1660

1661 First of all, the origin correction recompute the four-momentum of jets to point them to

the hard-scatter primary vertex instead of the centre of detector, and in the meantime keep the jet energy unchanged. This correction improves the η resolution of jets for roughly 25% at a jet p_T of 20 GeV and > 5 times improvement for jet with p_T above 200 GeV, as measured from the difference between reconstructed jets and truth jets in MC simulation. Secondly, the pile-up correction is adopted to remove the excess energy due to in-time and out-of-time pile-up, which consists of two processes: an area-based p_T density subtraction applied on the top of each event; and a residual correction derived from the simulation. Thirdly, the absolute JES calibration corrects the jet four-momentum to the particle-level energy scale, using truth jets in di-jet MC events. Furthermore, the step of global sequential calibration uses calorimeter, track and MS-based variables to reduce the flavor dependence and energy leakage effects. Finally, the residual in situ calibration is adopted to correct jets in data by using well-measured objects eg. photons, Z bosons and calibrated jets.

B-jet tagging

Tagging of b-jets plays a important role in many physics analyses involving b- or t-quark. In the meantime, lots of analyses need to apply b-jet veto to suppress $t\bar{t}$ process. There are three major types of algorithms that have been developed to distinguish b-quark jets from light-quark (u,d,s) jets^[63]:

- **Impact parameter based algorithms (IP2D and IP3D):** b-hadrons usually have long lifetime (~ 1.5 ps, $c_\tau \sim 450$ μm), which leads to large impact parameter for tracks produced from b-hadron decay. The impact parameter taggers are developed based on these variables. The IP2D tagger makes use of the transverse impact parameter significance $d_0/\sigma(d_0)$ as discriminant, while IP3D tagger uses two-dimensional discriminant of both transverse and longitudinal impact parameter significances: $d_0/\sigma(d_0)$ and $z_0 \sin\theta/\sigma(z_0)$.
- **Secondary vertex finding algorithm (SV1)** makes use of the secondary vertex formed by decay products of b-hadron within the jet. All track pairs within a jet are tested for a two-track vertex hypothesis, and removed if they are likely to originate from a long-live particle decay (eg. K_s or Λ), hadronic interactions or photon conversions. After that, a new vertex is fitted with all tracks from remaining two-track vertices, and the outliers are removed from this set of tracks.
- **Decay chain multi-vertex algorithm (JetFitter)**^[64] exploits the topological structure of weak b- and c- hadron decays inside the jet and tries to reconstruct the full b-hadron decay chain. A Kalman filter is adopted to find a common line between primary vertex and b/c- vertices, as well as their position in this line, which gives

1697 a approximated flight path for the b-hadron. In this approach, the b- and c-hadron
 1698 vertices, whenever resolution allows, can be resolved, even when there is only a
 1699 single track associated to them.

1700 The final discrimination commonly used in many physics analyses is called **Multivariate**
 1701 **Algorithm (MV2)**, which is based on Boosted Decision Tree (BDT) implemented in the
 1702 TMVA package^[65] by combining the outputs from underlaying taggers mentioned above.
 1703 The MV2 was trained using jets in $t\bar{t}$ sample, where the b-jets are treated as signal while
 1704 the c- and light-flavor jets are treated as backgrounds. There are three kinds of MV2
 1705 depending on the fraction of c-jets in background for training: *MV2c00*, *MV2c10* and
MV2c20. Figure 4.13 presents the output score of MV2c10 for different flavor jets.

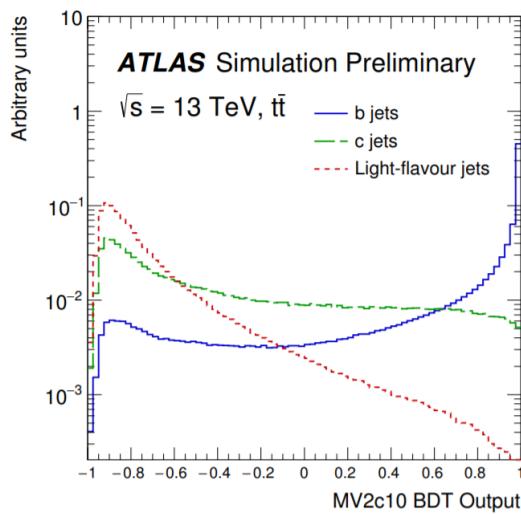


Fig. 4.13 MV2c10 BDT output for b- (solid blue), c- (dashed green) and light-flavour (dotted red) jets in $t\bar{t}$ events^[63].

1706

1707 4.2.6 Missing transverse energy

1708 Many interesting physics processes are with the involvement of neutrinos. Since they
 1709 do not interact with any materials in the detector, neutrinos cannot be detected directly;
 1710 but instead, they can result in imbalance in the plane transverse to the beam axis, where
 1711 momentum conservation is assumed. It is known as the missing transverse momentum
 1712 denoted as E_T^{miss} , which is obtained from the negative vector sum of the momenta of all
 1713 particles detected in a proton-proton collision event.

1714 The E_T^{miss} is measured using selected, reconstructed and calibrated hard objects in an
 1715 event. Its x- and y- components can be calculated as follow^[66]:

$$E_{x(y)}^{miss} = E_{x(y)}^{miss,e} + E_{x(y)}^{miss,\gamma} + E_{x(y)}^{miss,\tau} + E_{x(y)}^{miss,jets} + E_{x(y)}^{miss,\mu} + E_{x(y)}^{miss,soft} \quad (4.2)$$

where each object term is given by the negative vectorial sum of the momenta of the respective calibrated objects. The calorimeter signals are associated with the reconstructed objects in the following order: electrons, photons, hadronically decaying taus, jets, muons. The soft term is reconstructed from detected objects not match any hard object passing the selections, but associated with the primary vertex. Details of applied selections for each term are summarized in table 4.1.

Table 4.1 Overview of the contributions to E_T^{miss} .

Objects contributing to E_T^{miss}					
Priority	Type	Selections	Variables	Comments	
(1)	e	$ \eta < 1.37 \text{ or } 1.52 < \eta < 2.47$ $p_T > 10GeV$	$E_T^{miss,e}$	all e^\pm passing kinematic selections and medium reconstruction quality	
(2)	γ	$ \eta < 1.37 \text{ or } 1.52 < \eta < 2.47$ $p_T > 25GeV$	$E_T^{miss,\gamma}$	all γ passing kinematic selections and tight reconstruction quality, and without overlapping with (1)	
(3)	τ_{had}	$ \eta < 1.37 \text{ or } 1.52 < \eta < 2.47$ $p_T > 20GeV$	$E_T^{miss,\tau}$	all τ_{had} passing kinematic selections and medium reconstruction quality, and without overlapping with (1) and (2)	
(4)	μ	$ \eta < 2.7$ $p_T > 10GeV$	$E_T^{miss,\mu}$	all μ passing kinematic selections and medium reconstruction quality	
(5)	jet	$ \eta < 4.5$ $p_T > 60GeV$ --- or --- $2.4 < \eta < 4.5$ $20GeV < p_T < 60GeV$ --- or --- $ \eta < 2.4$ $20GeV < p_T < 60GeV$ $JVT > 0.59$	$E_T^{miss,jet}$	all jets passing kinematic selections and reconstruction quality (jet cleaning), and without overlap with (1)–(4)	
(6)	ID track	$p_T > 400MeV$ $ d_0 < 1.5mm$ $ z_0 \sin\theta < 1.5mm$ $\Delta R(track, e/\gamma cluster) > 0.05$ $\Delta R(track, \tau_{had}) > 0.2$	$E_T^{miss,soft}$	all ID tracks from the hard-scattering vertex passing kinematic selections and reconstruction quality, and not associated with any particle from (1), (3) or (4), or associated with a jet from (5)	

1721 Based on $E_{x(y)}^{miss}$, the magnitude of E_T^{miss} and the azimuthal angle ϕ^{miss} are computed:

$$E_T^{miss} = \sqrt{(E_x^{miss})^2 + (E_y^{miss})^2} \quad (4.3)$$

$$\phi^{miss} = \arctan(E_y^{miss}/E_x^{miss})$$

1722 In equation 4.2, each objects are required to pass certain reconstruction and calibrated
1723 criteria and selections mentioned above before taken as inputs.

1724 In figure 4.14, left plot shows the observed E_T^{miss} distribution for data and MC of
1725 $Z \rightarrow \mu\mu$ events that without genuine missing transverse momentum; and right plot shows
1726 the E_T^{miss} distribution for $W \rightarrow e\nu$ events that has genuine (true) missing transverse mo-
mentum due to real neutrino.

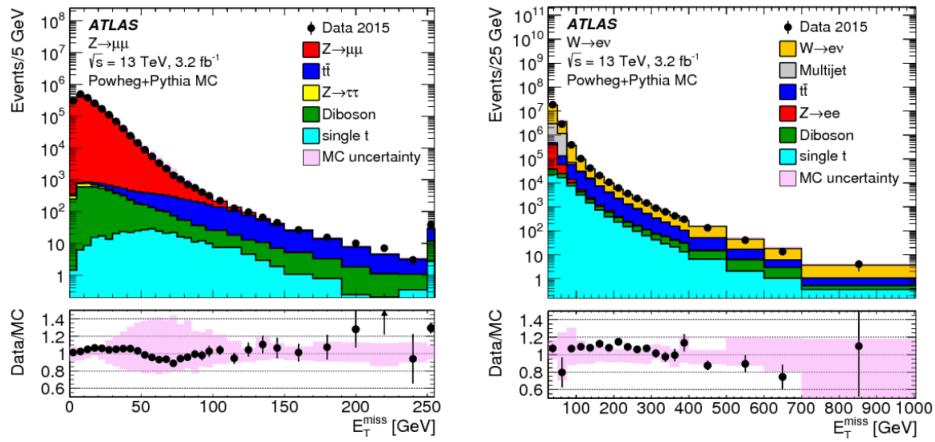


Fig. 4.14 Measured E_T^{miss} distribution for $Z \rightarrow \mu\mu$ events (left) and $W \rightarrow e\nu$ events (right).

1727

1728 **Chapter 5 Studies of SM ZZ production in $\ell\ell\ell'\ell'$ final** 1729 **state using pp collision data collected by ATLAS** 1730 **detector from 2015 to 2018**

1731 5.1 Introduction

1732 After the discovery of Higgs boson^[67-68], the examine of electroweak symmetry breaking (EWSB) becomes a main focus at the LHC. In addition to measuring the properties of 1733 Higgs boson directly, the vector boson scattering (VBS) process is another key avenue to 1734 probe EWSB^[69-71]. As introduced in section 2.1.3, in Standard Model (SM), the Higgs 1735 boson acts as “moderator” to unitarize high-energy longitudinal VBS amplitudes at the 1736 TeV scale. Therefore, studying high-energy behaviours of VBS is crucial to understand 1737 the mechanism of EWSB.

1739 Since there was no VBS process was observed prior to the LHC era, LHC provides 1740 an unexceptionable opportunity to study them due to its unprecedented high energy and 1741 luminosity. At LHC, the VBS process is typically studied through the measurements of 1742 electroweak (EW) production of two vector bosons radiated from quark-quark initial state, 1743 plus a pair hadronic jets with high energy in the back and forward regions (denoted as EW- 1744 $VVjj$). The quantum chromodynamics (QCD) production of $VVjj$ contains two QCD 1745 vertices at the lowest order (denoted as QCD- $VVjj$) is an irreducible background to the 1746 search of EW- $VVjj$ production. The features of EW- $VVjj$ production include a large 1747 invariant mass of jet pair (m_{jj}) and a significant separation of rapidity between two jets 1748 (Δy_{jj}). Figure 5.1 presents some typical Feynman diagrams of EW- and QCD- $ZZjj$ 1749 processes.

1750 The first evidence of the EW- $VVjj$ process was seen in same-sign WW channel (EW- 1751 $W^\pm W^\pm jj$) by ATLAS collaboration with 20.3 fb^{-1} 8 TeV data^[72], in which a 3.6σ excess 1752 was observed in data over the background-only prediction. In the LHC run-2, the obser- 1753 vation (with $> 5 \sigma$ statistical significance) of EW- $W^\pm W^\pm jj$ process has been reported 1754 in both ATLAS and CMS collaboration with 36 fb^{-1} 13 TeV data^[73-74]. In WZ channel 1755 (EW- $WZjj$), an observation with 5.3σ excess was also reported by the ATLAS col- 1756 laboration recently^[75]. As for the EW- $ZZjj$ production, it was searched by CMS using 1757 35.9 fb^{-1} 13 TeV data but no evidence was found^[76]. The EW production in ZZ final state 1758 (EW- $ZZjj$) is typically rare, whose fiducial cross section has an order of $O(0.1) \text{ fb}^{-1}$ in the 1759 final state where both Z bosons decay leptonically. But in the meantime, $ZZ \rightarrow \ell\ell\ell'\ell'$

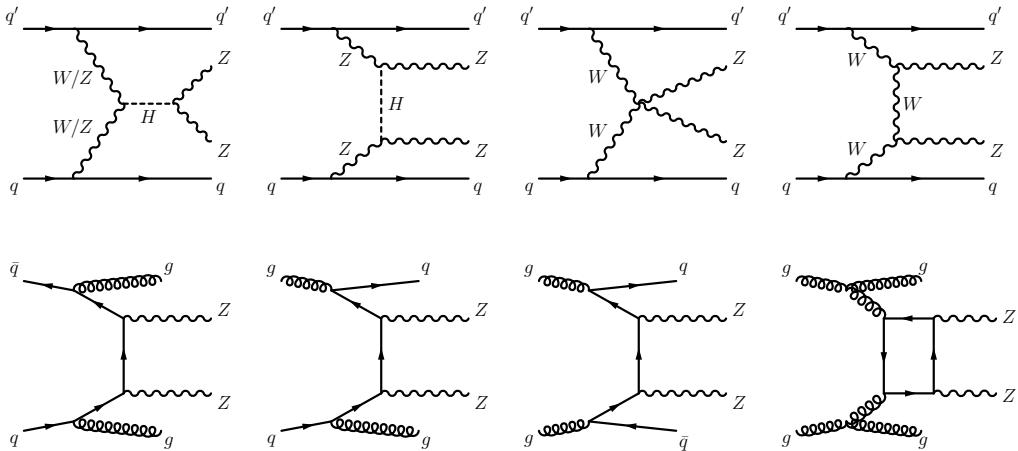


Fig. 5.1 Typical diagrams for the production of $ZZjj$, including the relevant EW VBS diagrams (first row) and QCD diagrams (second row).

process offers a extremely clean channel than all the others. So with more data collected in the LHC, the observation of EW- $ZZjj$ becomes possible.

This section presents the first observation of EW- $ZZjj$ production decaying to four charged leptons with two jets ($\ell\ell\ell'\ell' jj$) by ATLAS collaboration using the complete set of the LHC run-2 data with 139 fb^{-1} luminosity. It is a new milestone in the study of EWSB at the LHC, and completes the last missing part of observation of weak boson scattering for massive bosons. In the meantime, the measurement of fiducial cross-sections for SM ZZ production including both EW and QCD processes is also reported. The $ZZjj$ production involving intermediate τ -leptons from Z decays is considered as signal but has a negligible contribution to the selected events. Reducible backgrounds give minor contributions in the $\ell\ell\ell'\ell' jj$ channel are also studied. To further separate the EW signal and the QCD background, multivariate discriminant (MD) is trained using event kinematic information from simulated samples. The MD distribution is then used as discriminant in statistical fit to evaluate the signal strength of EW process.

5.2 Data and MC samples

5.2.1 Data samples

The data sets for this analysis include the full run-2 pp collision data collected by the ATLAS experiment during the years from 2015 to 2018. Data event is only used if it passed the latest Good Run List (GRL) released by the Data Quality group from ATLAS experiment as listed below, corresponding to an integrated luminosity of $139.0 \pm 2.4 \text{ fb}^{-1}$.
`data15_13TeV.periodAllYear_DetStatus-v89-pro21-02_Unknown_PHYS_StandardGRL_All_Good_25ns`

1781 data16_13TeV.periodAllYear_DetStatus-v89-pro21-01_DQDefects-00-02-04_PHYS_StandardGRL_All
1782 data17_13TeV.periodAllYear_DetStatus-v99-pro22-01_Unknown_PHYS_StandardGRL_All_Good_25ns
1783 data18_13TeV.periodAllYear_DetStatus-v102-pro22-04_Unknown_PHYS_StandardGRL_All_Good_25ns

1784 5.2.2 MC simulations

1785 The EW- $ZZjj$ production (signal) is modelled using MAD-
1786 GRAPH5_aMC@NLO 2.6.1^[77] with the matrix elements (ME) calculated in the leading-
1787 order (LO) approximation in perturbative QCD (pQCD) and with the NNPDF2.3LO^[78]
1788 parton distribution functions (PDF). The VBF Higgs process is also included.

1789 The QCD- $ZZjj$ production is modelled using SHERPA 2.2.2^[79] with the
1790 NNPDF3.0NNLO^[80] PDF, where events with up to one (three) outgoing partons are gen-
1791 erated at NLO (LO) in pQCD. The production of $ZZjj$ from the gluon-gluon initial state
1792 with a four-fermion loop or with an exchange of the Higgs boson has an order of α_S^4 in
1793 QCD, and is not included in the SHERPA simulation. A separate gg induced $ZZ + 2\text{jets}$
1794 sample is modelled using SHERPA 2.2.2 with the NNPDF3.0NNLO PDF and with an addi-
1795 tional 1.7 k-factor^[81] being applied. Then the interference between EW- and QCD- $ZZjj$
1796 is modelled with MADGRAPH5_aMC@NLO 2.6.1 calculated at LO.

1797 The diboson productions from QCD $WW \rightarrow \ell\nu qq$ as well as QCD and EW $WZ \rightarrow$
1798 $\ell\ell qq$ are modelled using SHERPA 2.2.2 with the NNPDF3.0NNLO PDF. The productions
1799 of semileptonic decays ($WW \rightarrow \ell\nu qq$ and $WZ \rightarrow qq\ell\ell$) are modelled using POWHEG-
1800 Box v2^[82] with the CT10 PDF^[83]. The triboson production is modelled using SHERPA
1801 2.2.2 with the NNPDF3.0NNLO PDF.

1802 For top-quark pair ($t\bar{t}$) production, the POWHEG-Box v2 is used with the CT10 PDF.
1803 The single top-quark production in t -channel, s -channel and Wt -channel are simulated
1804 using the POWHEG-Box v1 event generator^[84-86]. The productions of $t\bar{t}$ in association with
1805 vector boson(s) ($t\bar{t}V$) are modelled with MADGRAPH5_aMC@NLO 2.3.3 for $t\bar{t}W$ and $t\bar{t}Z$
1806 with $Z \rightarrow \nu\nu/qq$ decays, with SHERPA 2.2.1 for $t\bar{t}Z$ where the Z decays to dilepton, and
1807 with MADGRAPH5_aMC@NLO 2.2.2 for $t\bar{t}WW$ respectively.

1808 The $Z+jets$ processes are modelled using SHERPA 2.2.1 with the NNPDF3.0NNLO
1809 PDF, in which the ME is calculated for up to two partons with next-to-leading-order (NLO)
1810 accuracy in pQCD and up to four partons with LO accuracy.

1811 For all the samples except those from SHERPA, the parton showering is modelled with
1812 PYTHIA8^[41] using the NNPDF2.3^[78] PDF set, and the A14 set of tuned parameters^[87].
1813 While for SHERPA samples, the parton showering is simulated within the programme.

1814 All simulated events are processed with detector response simulated based on GEANT4

described in section 4.1. In addition, simulated inelastic pp collisions are overlaid to model additional pp collision in the same and neighbouring bunch crossings (pile-up), and reweighted to match the pile-up conditions in data. Moreover, all simulated events are processed using the same reconstruction algorithms as data. And the leptons and jets reconstruction, energy scale and resolution, and the leptons identification, isolation, trigger efficiencies for simulated events, as described in section 4.2, are all corrected to match the data measurements.

5.3 Objects and Event selection

5.3.1 Objects selection

The selection of analysis relies on the definition of multiple objects: *electrons*, *Muons*, and *jets*. Details of definition for each object are described as below:

Muon: To increase the acceptance range in reconstruction (reco) -level for $\ell\ell\ell'\ell'$ channel, all four types of muons (CB, ST, CT, ME muons, described in section 4.2.4) are used. The identified muons are then required to pass $p_T > 7$ GeV and $|\eta| < 2.7$, and satisfy the *Loose* identification criterion (see definition in sec 4.2.4). The impact parameter cuts are further applied to suppress the contribution from cosmic muons and non-prompt muons, with the value of: $|d_0/\sigma(d_0)| < 3.0$ and $|z_0 \sin\theta| < 0.5$ mm, where d_0 is the transverse impact parameter relative to the beam line, $\sigma(d_0)$ is its uncertainty, and z_0 is the longitudinal impact parameter relative to the primary vertex. In order to avoid muons associated with jets, all muons are required to be isolated and pass *FixedCutLoose* isolation criteria, which required $E_T^{\text{topocone}20}/p_T < 0.3$ and $p_T^{\text{varcone}30}/p_T < 0.15$.

Electron: As described in section 4.2.3, electrons are reconstructed from energy de- posits in the EM calorimeter matched to a track in the inner detector. The electron candi- dates must satisfy the *Loose* criterion valuing by the likelihood-based (LH) method. And electrons are required to have $p_T > 7$ GeV and $|\eta| < 2.47$. Moreover, the impact parame- ter requirements of $|d_0/\sigma(d_0)| < 5.0$ and $|z_0 \sin\theta| < 0.5$ mm are applied. Same as muon, all electrons are required to satisfy *FixedCutLoose* isolation criteria of $E_T^{\text{topocone}20}/p_T < 0.2$ and $p_T^{\text{varcone}20}/p_T < 0.15$.

Jets: Jet are key signatures for VBS processes. This analysis use the jets clustered using the anti- k_t algorithm with radius parameter $R = 0.4$, more details of jets' reconstruc- tion can be found in section 4.2.5. The jets are required to satisfy $p_T > 30$ (40) GeV in the $|\eta| < 2.4$ ($2.4 < |\eta| < 4.5$) region. To further reduce the effects of pile-up jets, a jet vertex tagger (JVT) is applied to jets with $p_T < 60$ GeV and $|\eta| < 2.4$ to select jets from

1848 hard-scattering vertex^[88].

1849 **Overlap removal:** An overlap-removal procedure is applied to selected leptons and
 1850 jets in this analysis. To enhance the selection efficiency, leptons are given higher priority
 1851 to be kept when overlapping with jets. With this lepton preferred method, the events of
 1852 EW signal after selection increases about 19% while background only increases about
 1853 14%. More details of the strategy is summarized in table 5.1.

	Reference objects	Criteria
Remove electrons	electrons	Share a track or have overlapping calorimeter cluster. Keep higher p_T electron
Remove muons	electrons	Share track and muon is calo-tagged
Remove electrons	muons	Share track
Remove jets	electrons	$\Delta R_{e-jet} < 0.2$
	muons	$\Delta R_{\mu-jet} < 0.2$ OR muon track is ghost-associated to jet $\text{AND } (N_{Trk}(jet) < 3 \text{ OR } (p_T^{jet}/p_T^\mu < 2 \text{ and } p_T^\mu/\Sigma_{Trk}P_t > 0.7))$

Table 5.1 Overlap removal criteria between pre-selection objects for the $\ell\ell\ell'\ell'$ channel. The overlap removal follows the order shown in this table. Once an object has been marked as removed, it does not participate in the subsequent stages of the overlap removal procedure.

1854 5.3.2 Event selection in reconstruction level

1855 The events are required to additionally be recorded by single or multi-lepton triggers,
 1856 with transverse momentum (p_T) thresholds varying from 8 to 26 GeV. The overall trigger
 1857 efficiency for selected inclusive $\ell\ell\ell'\ell' jj$ signal events in the analysis region are from 95
 1858 to 99%.

1859 The $\ell\ell\ell'\ell'$ quadruplets are formed by two opposite-sign, same-flavour (OSSF) lepton
 1860 pairs ($\ell^+\ell^-$), in which leptons are required to be separated by $\Delta R > 0.2$ in table 5.1.
 1861 At most one muon is allowed to be ME or CT muon. The p_T threshold of first three
 1862 leading leptons are 20, 20 and 10 GeV. If more than one quadruplets are found, the one
 1863 with minimum sum of difference between two dilepton pair masses and Z boson mass
 1864 ($|m_{l_1^+ l_1^-} - m_Z| + |m_{l_2^+ l_2^-} - m_Z|$) is selected. Both two dilepton pair masses are required to
 1865 be between 66 to 116 GeV. In addition, the invariant masses of all possible OSSF pairs
 1866 are required to be greater than 10 GeV to reject events from J/ψ or Y decay.

1867 For VBS topology, the two most energetic jets in different detector side ($y_{j1} \times y_{j2} < 0$)

are selected. Furthermore, the invariant mass of two jets (m_{jj}) is required to be greater than 300 GeV, while Δy_{jj} is required to be larger than 2. Table 5.2 summarizes the above selection requirements, which is defined as signal region (SR).

Electrons	$p_T > 7 \text{ GeV}, \eta < 2.47$ $ d_0/\sigma_{d_0} < 5 \text{ and } z_0 \times \sin \theta < 0.5 \text{ mm}$
Muons	$p_T > 7 \text{ GeV}, \eta < 2.7$ $ d_0/\sigma_{d_0} < 3 \text{ and } z_0 \times \sin \theta < 0.5 \text{ mm}$
Jets	$p_T > 30 (40) \text{ GeV for } \eta < 2.4 (2.4 < \eta < 4.5)$
ZZ selection	$p_T > 20, 20, 10 \text{ GeV for the leading, sub-leading and third leptons}$ Two OSSF lepton pairs with smallest $ m_{\ell^+\ell^-} - m_Z + m_{\ell'^+\ell'^-} - m_Z $ $m_{\ell^+\ell^-} > 10 \text{ GeV for lepton pairs}$ $\Delta R(\ell, \ell') > 0.2$ $66 < m_{\ell^+\ell^-} < 116 \text{ GeV}$
Dijet selection	Two most energetic jets with $y_{j_1} \times y_{j_2} < 0$ $m_{jj} > 300 \text{ GeV and } \Delta y_{jj} > 2$

Table 5.2 Summary of selection of physics objects and candidate events at detector level in the $\ell\ell\ell'\ell' jj$ signal region.

5.4 Background estimation

Table 5.3 summarizes the background yields for $ZZjj \rightarrow \ell\ell\ell'\ell' jj$ process in 139 fb^{-1} . Uncertainties on the predictions include both statistical and systematic components. "Others" includes minor contributions from non- ZZ processes including $Z+jets$, top-quark, triboson and $t\bar{t}V$ processes. Details of estimation for each source are described as below.

Process	$\ell\ell\ell'\ell' jj$
EW- $Z Z jj$	20.6 ± 2.5
QCD- $q\bar{q} \rightarrow ZZ$	77 ± 25
QCD- $gg \rightarrow ZZ$	13.1 ± 4.4
Others	3.2 ± 2.1
Total	114 ± 26
Data	127

Table 5.3 Observed data and expected signal and background yields in 139 fb^{-1} of luminosity. Minor backgrounds are summed together as 'Others'. Uncertainties on the predictions include both statistical and systematic components.

5.4.1 QCD backgrounds

The QCD- $Z Z jj$ production, which include both qq and gg initial processes, is the irreducible background in the search of EW- $Z Z jj$ production. A QCD-enriched control region, named as QCD CR, is defined to constrain the normalization of $Z Z$ background by reverting either the m_{jj} or Δy_{jj} requirements as:

$$m_{jj} < 300 \text{ GeV} \text{ or } \Delta y_{jj} < 2 \quad (5.1)$$

Then this normalization factor is included into statistical fit as a float parameter to properly treat the uncertainty correlations between SR and CR, while the shapes are taken from MC simulation. Table 5.4 shows the event yields of each background components in this CR. Uncertainties are statistical one only. The distributions of invariant mass of $\ell\ell\ell'\ell'$ and

Process	$\ell\ell\ell'\ell' jj$
EW- $Z Z jj$	3.9 ± 0
QCD- $Z Z jj$	136.9 ± 0.6
QCD- $gg Z Z jj$	16.8 ± 0.1
Diboson	0.3 ± 0.1
Triboson	1.6 ± 0.1
$Z+jets$	0
$t\bar{t}$	0
Total	159.5 ± 0.62
Data	152

Table 5.4 Observed data and expected signal and background yields in 139 fb^{-1} of luminosity. Diboson background in table includes all the other diboson processes discussed in section 5.2.2, except those with four-lepton final state. Uncertainties include only MC statistic. No events from $Z+jets$ and $t\bar{t}$ MC samples pass the selection, and are indicated as 0 in the table.

1885

1886 dijet in QCD CR are shown in figure 5.2.

5.4.2 Reducible backgrounds

Backgrounds from $Z+jets$, top-quark and WZ processes called reducible backgrounds can be estimated by data-driven method. These events usually contain two or three leptons from Z/W decays, together with heavy-flavor jets or misidentified components of jets reconstructed as leptons called "fake leptons". A *fake factor* method is used to estimate these backgrounds, where the lepton misidentification is measured in data with

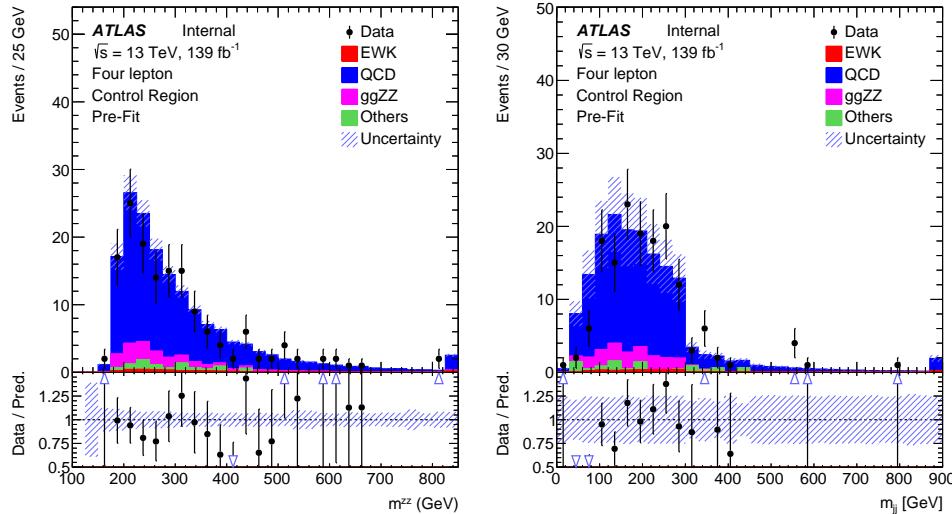


Fig. 5.2 Pre-fit m_{ZZ} and m_{ll} distribution in QCD-enriched CR.

the region enhances contributions from $Z+jets$ and top-quark processes. The method is described as below:

- Define a dedicated background dominant region to derive the fake factor for this background. The *fake factor* is defined as:

$$\mathcal{F} = \mathcal{N}_{good}/\mathcal{N}_{pool} \quad (5.2)$$

where \mathcal{N}_{good} refers to the number of good leptons passing all SR selection, while \mathcal{N}_{pool} denotes the number of poor leptons passing most SR selection but fail one certain requirement.

- Define a fake control region, where one or two lepton(s) pass *poor* requirement while all the other leptons are required to have SR selection.
- The number of fake events are calculated as:

$$\mathcal{N}_{fake} = (N_{gggp} - N_{ggp}^{ZZ}) \times \mathcal{F} - (N_{gppp} - N_{gpp}^{ZZ}) \times \mathcal{F}^2 \quad (5.3)$$

with the subtraction of ZZ contribution, and the double counting between N_{gggp} and N_{gppp} .

For the definition of *poor* leptons: The poor electrons are defined as failing “Fixed-CutLoose” isolation requirement, or failing “LooseLH” electron ID requirement but satisfying “VeryLooseLH” WP. The poor muons are required to fail the “FixedCutLoose” isolation requirement or invert the impact parameter cut to be $3 < d_0/\sigma(d_0) < 10$. The dedicated $Z+jets$ and $t\bar{t}$ dominant regions are defined to calculate the fake factor respectively in the following subsections.

1. Fake factor for $Z+jets$

Fake factor for $Z+jets$ background is calculated in $Z+jets$ enriched region, where events with one SFOS lepton pair around Z mass associated with two jets are selected. The value of fake factor is driven from data, and as a function of p_T and η as shown in figure 5.3 for electrons and figure 5.4 for muons. During calculation, the contributions from non- $Z+jets$ backgrounds ($t\bar{t}$, ZZ , WZ) have been subtracted from data. The values calculated directly from $Z+jets$ MC are also shown in plots for comparison.

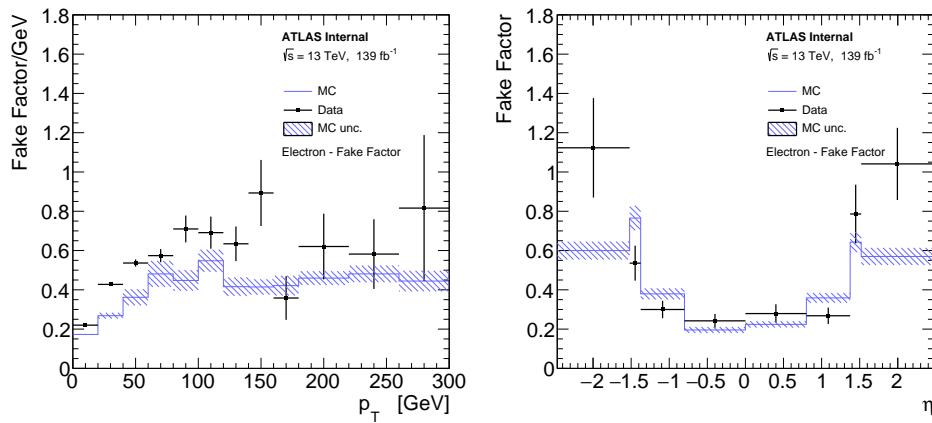


Fig. 5.3 Fake factor for $Z+jets$ background, constructed with additional electron, as a function of p_T (left) and η (right).

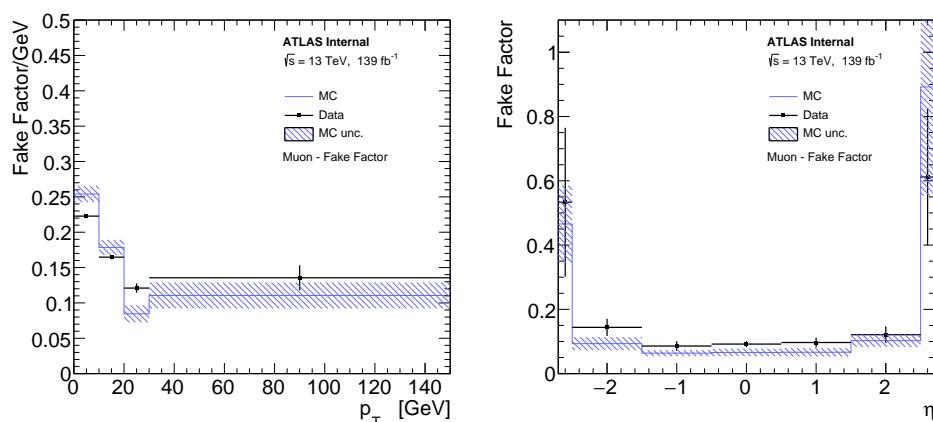


Fig. 5.4 Fake factor for $Z+jets$ background, constructed with additional muon, as a function of p_T (left) and η (right).

2. Fake factor for $t\bar{t}$

The fake factor for $t\bar{t}$ are calculated in $t\bar{t}$ dominant region by selecting one $e\mu$ -pair with additional two jets. For events with three leptons, $m_T^W < 60 \text{ GeV}$ cut is applied to reject

1921 the contribution from $t\bar{t} + W$ events. The m_T^W is defined as below:

$$m_T^W = \sqrt{2p_T^{l_3}E_T^{miss} \left[1 - \cos(\Delta\phi(p_T^{l_3}, E_T^{miss})) \right]} \quad (5.4)$$

1922 In addition, at least one b-jet is required to enhance the top component. The fake factors of
 1923 $t\bar{t}$ calculated from data as the function of p_T and η are shown in figure 5.5 for electrons and
 1924 5.6 for muons. The non- $t\bar{t}$ contributions including $Z+jets$, ZZ and WZ , are subtracted
 from data.

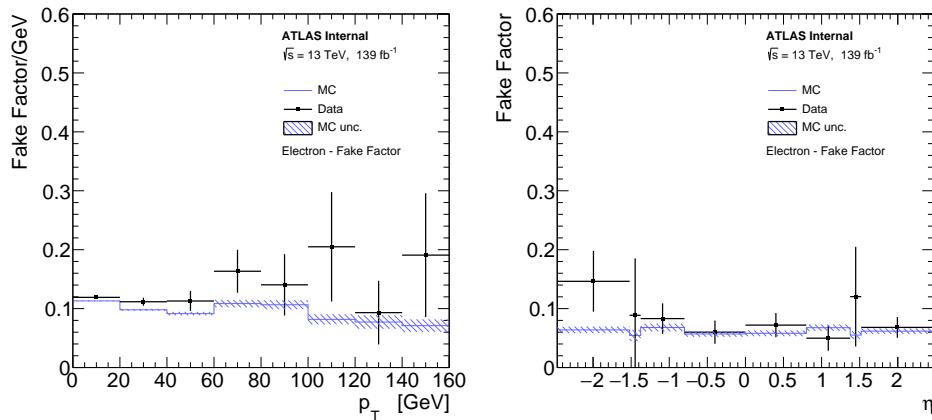


Fig. 5.5 Fake factor for $t\bar{t}$ background, constructed with additional electron, as a function of p_T (left) and η (right).

1925

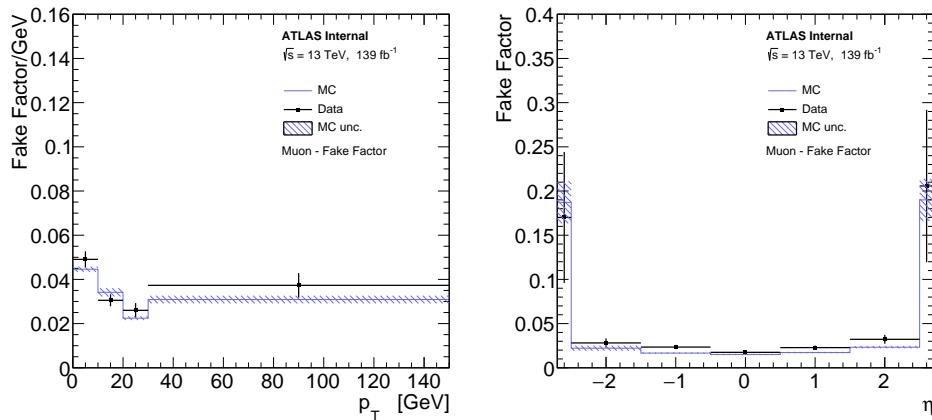


Fig. 5.6 Fake factor for $t\bar{t}$ background, constructed with additional muon, as a function of p_T (left) and η (right).

1926

3. Combination

1927 The fake factors calculated from each dedicated region are then combined together
 1928 according to their contributions in fake control region described previously. Figure 5.7
 1929 shows the m_{jj} distribution with data and major fake backgrounds in three different 41
 1930 channels.

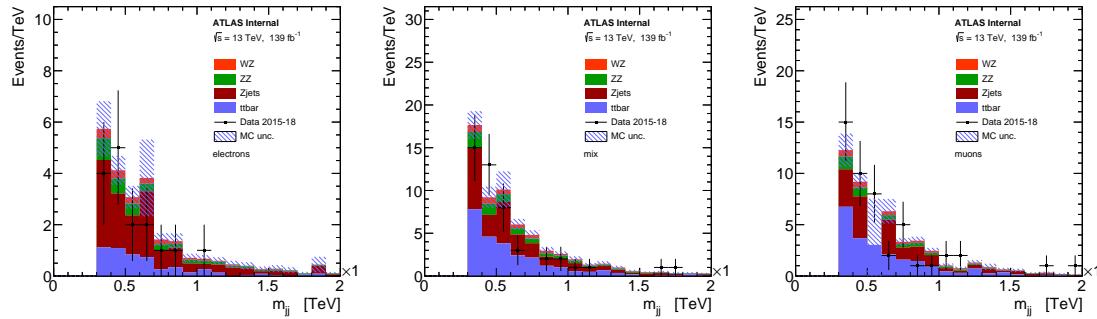


Fig. 5.7 m_{jj} distributions in fake control region in 4e (left), 2e2 μ (middle) and 4 μ (right) channel. The ratios between $Z+jets$ and $t\bar{t}$ ($Z+jets/t\bar{t}$) in each individual channel are: 2.59, 0.95, 0.74.

4. Systematics of fake estimation and results

The systematics of fake factor method can be measured by varying the parameters and selection requirements in fake factor calculation. In addition, due to the very limited data statistic in $\ell\ell\ell'\ell'$ channel, to be more conservative, the difference between data measurement and MC simulation are also considered as another systematics component. The sources of systematics that have been included are listed as below:

- Variations of isolation cut for the poor lepton definition up and down scaled by a factor of two.
- Variations of the yields of those subtracted MC in fake control region scaled by 30% up and down.
- The difference of fake factors between driven from data and from MC simulation.
- The difference of fake factors when changing to one bin measurement (instead of p_T or η dependent).
- The statistical uncertainties on fake factor in fake control region.

Table 5.5 summarizes the contribution of fake backgrounds in signal region under different systematic conditions mentioned above as well as the nominal one, together with their statistical uncertainties.

5.5 Systematics

The analysis includes both the statistical fit to MD distribution to search the EW- $Z Z jj$ process, as well as the cross section measurement of inclusive EW and QCD $Z Z jj$ process in fiducial volume. Therefore, theoretical and experimental uncertainties may affect the predictions background yields and shapes, correction factors from detector-level to particle-level measurement, as well as the $Z Z jj$ MD shapes and so on. Moreover, the

channel	4e	2e2 μ	4 μ	inclusive
Nominal estimate	0.678 ± 0.652	1.023 ± 0.740	0.566 ± 0.240	2.268 ± 1.015
F stat. uncertainty varied down	0.698 ± 0.622	0.872 ± 0.652	0.509 ± 0.214	2.079 ± 0.926
F stat. uncertainty varied up	0.657 ± 0.685	1.173 ± 0.840	0.622 ± 0.267	2.452 ± 1.116
One bin F	0.653 ± 0.590	0.594 ± 0.558	0.646 ± 0.313	1.892 ± 0.870
MC F	0.534 ± 0.471	1.415 ± 0.993	0.439 ± 0.184	2.389 ± 1.114
Isolation varied down	0.938 ± 0.686	0.552 ± 0.466	0.215 ± 0.107	1.704 ± 0.837
Isolation varied up	0.723 ± 0.646	1.104 ± 0.739	0.559 ± 0.237	2.386 ± 1.010
MC corr. varied down	0.697 ± 0.695	1.048 ± 0.811	0.832 ± 0.385	2.577 ± 1.136
MC corr. varied up	0.660 ± 0.614	0.984 ± 0.687	0.316 ± 0.159	1.961 ± 0.935

Table 5.5 Fake background estimations in the SR. For the nominal value the 2D fake factor together with the $Z+jets$ and $t\bar{t}$ combination applied. The other lines show the estimations with different uncertainty variations.

statistical uncertainties of simulated samples are also taken into account. Due to the extremely low cross section of $\ell\ell\ell'\ell'$ channel, the analysis is still data statistic dominant. This section describes the measurement of both theoretical and experimental systematics for $Z Z jj$ productions. The systematics for fake backgrounds have been elaborated in section 4.

5.5.1 Theoretical systematics

The theoretical systematics on EW- and QCD- $Z Z jj$ processes include the uncertainties from PDF, QCD scale, α_S and parton showering variations as summarized in table 5.6. The PDF uncertainty is estimated from the envelop of NNPDF internal variations and the difference between nominal and alternative PDF sets, following the PDF4LHC as introduced in section 2.2.1. The QCD scale uncertainty is estimated by varying the nominal renormalization scale (μ_R) and factorisation scale (μ_F) by a factor of 0.5 or 2.0. There are seven different configurations being considered, where the maximum of variations is chosen as final uncertainty. The parton showering uncertainty is estimated by comparing events with different parton showering setting between the nominal PYTHIA8 and the alternative HERWIG7^[89-90] algorithm. The α_S uncertainty is estimated by varying the value of α_S within ± 0.001 . Due to the lack of simulation sample for alternative parton showering on QCD- $Z Z jj$ process, the value of parton showering component is taken from the measurement of EW process.

Table 5.7 summarizes the uncertainties of each theoretical components in fiducial volume of SR, while table 5.8 shows the numbers in QCD-enriched CR region. For QCD process, the uncertainty is QCD scale dominant. Both of them are taken as inputs for

Process	EW- $Z Z jj$	QCD- $Z Z jj$	
PDFs	NNPDF30lo (nominal), CT14lo	NNPDF30nnlo (nominal), MMHT2014nnlo68cl, CT14nnlo	
α_S	0.118	0.117, 0.118 (nominal), 0.119	
QCD scale ($[\mu_R, \mu_F]$)	[0.5,0.5], [0.5,1], [1,0.5], [1,1], [1,2], [2,1], [2,2]	[0.5,0.5], [0.5,1], [1,0.5], [1,1], [1,2], [2,1], [2,2]	
Parton showering algorithm	PYTHIA8, HERWIG7	-	

Table 5.6 Summary of different variations for EW- and QCD- $Z Z jj$ theoretical uncertainties measurement.

statistical fit.

Process	PDF (%)	α_S (%)	QCD scale (%)	Parton shower (%)
EW	+5.9 -5.9		+6.1 -5.6	+3.3 -3.3
qqQCD	+2.0 -1.0	+2.6 -2.6	+34.2 -22.8	

Table 5.7 Summary of theoretical uncertainties for the fiducial volume (SR) for both EW and QCD $q q$ -initial processes.

1976

Process	PDF (%)	α_S (%)	QCD scale (%)	Parton shower (%)
EW $\ell\ell\ell'\ell'$	+6.1 -6.1		+0.8 -1.1	+10.1 -10.1
qqQCD $\ell\ell\ell'\ell'$	+2.0 -1.0	+2.6 -2.6	+31.5 -22.0	

Table 5.8 Summary of theoretical uncertainties for the control region for EW and qqQCD processes.

1977

The uncertainties of QCD gg -induced process ($gg \rightarrow ZZ$) as the function of MD discriminant is shown in figure 5.8 for both fiducial volume (SR) and QCD CR.

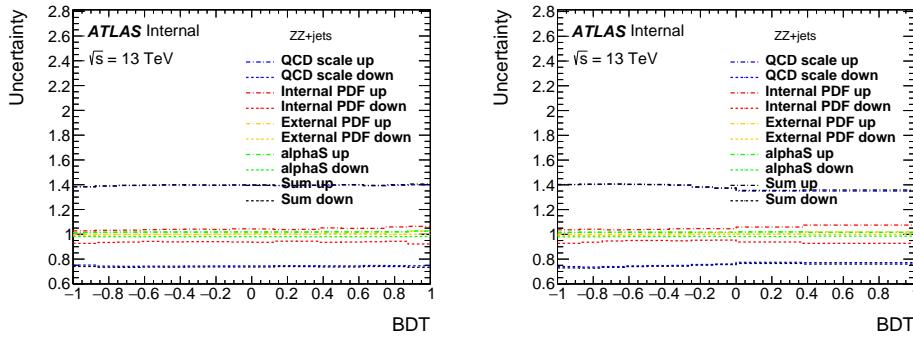


Fig. 5.8 The theoretical uncertainties for $gg \rightarrow ZZ$ background in particle-level SR (left) and CR (right).

1978

5.5.2 Experimental systematics

The dominant experimental uncertainties are from the luminosity uncertainty, the momentum scale and resolution of leptons and jets, as well as the lepton reconstruction and selection efficiency. Some smaller uncertainties, such as trigger efficiency and pile-up correction, are also considered. Table 5.9 lists the major systematic components from leptons and jets for signal and major background processes in $\ell\ell\ell'\ell'$ channel. The total uncertainties for sources from electron, muon and jet respectively, as well as the sum (quadratic sum) of them are also summarized in this table.

name	EW- $Z Z jj$	QCD qq -initial	QCD gg
nominal yield	20.61	76.69	13.10
EG_RESOLUTION_ALL	$\pm^{0.00\%}_{0.03\%}$	$\pm^{0.02\%}_{0.04\%}$	$\pm^{0.01\%}_{1.41\%}$
EG_SCALE_ALL	$\pm^{0.03\%}_{0.05\%}$	-0.04%	$\pm^{0.01\%}_{0.06\%}$
EL_EFF_ID_TOTAL_1NPCOR_PLUS_UNCOR	$\pm^{2.66\%}_{2.58\%}$	$\pm^{2.60\%}_{2.53\%}$	$\pm^{2.65\%}_{2.57\%}$
EL_EFF_Iso_TOTAL_1NPCOR_PLUS_UNCOR	$\pm 0.70\%$	$\pm 0.47\%$	$\pm 0.42\%$
EL_EFF_Reco_TOTAL_1NPCOR_PLUS_UNCOR	$\pm 0.55\%$	$\pm 0.55\%$	$\pm 0.63\%$
JET_EtaIntercalibration_NonClosure	-0.01%	-0.03%	0%
JET_GroupedNP_1	$\pm 1.97\%$	$\pm^{11.82\%}_{10.14\%}$	$\pm^{16.21\%}_{12.92\%}$
JET_GroupedNP_2	$\pm 0.23\%$	$\pm 1.26\%$	+5.3%
JET_GroupedNP_3	$\pm 0.55\%$	$\pm 2.94\%$	$\pm^{3.14\%}_{0.12\%}$
JET_JER_SINGLE_NP	0.11%	+5.47%	+6.31%
JET_JvtEfficiency	$\pm 0.04\%$	$\pm 0.12\%$	$\pm 0.15\%$
MUON_EFF_ISO_STAT	$\pm 0.09\%$	$\pm 0.08\%$	$\pm 0.07\%$
MUON_EFF_ISO_SYS	$\pm 0.54\%$	$\pm 0.55\%$	$\pm 0.56\%$
MUON_EFF_RECO_STAT	$\pm 0.15\%$	$\pm 0.19\%$	$\pm 0.15\%$
MUON_EFF_RECO_STAT_LOWPT	$\pm 0.06\%$	$\pm 0.02\%$	$\pm 0.03\%$
MUON_EFF_TTVA_STAT	$\pm 0.06\%$	$\pm 0.07\%$	$\pm 0.06\%$
MUON_EFF_TTVA_SYS	$\pm 0.03\%$	$\pm 0.4\%$	$\pm 0.03\%$
MUON_ID	$\pm 0.03\%$	$\pm 0.02\%$	<0.001%
MUON_MS	-0.05%	$\pm^{0.04\%}_{0.01\%}$	<0.001%
MUON_SAGITTA_RESBIAS	$\pm 0.01\%$	$\pm 0.02\%$	<0.001%
MUON_SAGITTA_RHO	+1.13%	-0.73%	$\pm 1.00\%$
MUON_SCALE	$\pm 0.02\%$	$\pm^{0.03\%}_{0.02\%}$	<0.001%
PRW_DATASF	$\pm 0.5\%$	$\pm^{0.42\%}_{1.02\%}$	$\pm^{2.17\%}_{1.46\%}$
Electron Exp	$\pm^{2.8\%}_{2.7\%}$	$\pm^{2.70\%}_{2.62\%}$	$\pm^{2.75\%}_{2.64\%}$
Muon Exp	$\pm 1.3\%$	$\pm 1.3\%$	$\pm 1.04\%$
Jet Exp	$\pm 2.0\%$	$\pm^{13.39\%}_{10.64\%}$	$\pm^{18.54\%}_{13.57\%}$
Total experimental uncertainties	$\pm^{3.7\%}_{4.0\%}$	$\pm^{13.72\%}_{11.11\%}$	$\pm^{18.90\%}_{13.57\%}$

Table 5.9 Experimental uncertainties in $\ell\ell\ell'\ell'$ channel with the luminosity of 139 fb^{-1} . The "Electron Exp", "Muon Exp" and "Jet Exp" represent the quadrature of the respective sources from electron, muon, and jets.

1987 In addition, the uncertainty of the combined 2015 to 2018 integrated luminosity is
1988 1.7%^[91] in ATLAS experiment, obtained using the LUCID-2 detector^[92] for the primary
1989 luminosity measurements.

1990 On top of them, a systematic uncertainty for MD distribution with different pile-up
1991 ($\langle\mu\rangle$) is also considered for QCD- $Z Z jj$ background by comparing the distributions be-

tween events with low and high pile-up conditions. A boundary of $\langle\mu\rangle = 33$ is used to defined low/high pile-up according to the average $\langle\mu\rangle$ for signal (about 34.5) and QCD background (about 33). Figure 5.9 shows the MD distribution in SR (left) and QCD CR (right) in two different PU conditions, the difference as function of MD is then taken into account as additional shape uncertainty for statistical fit.

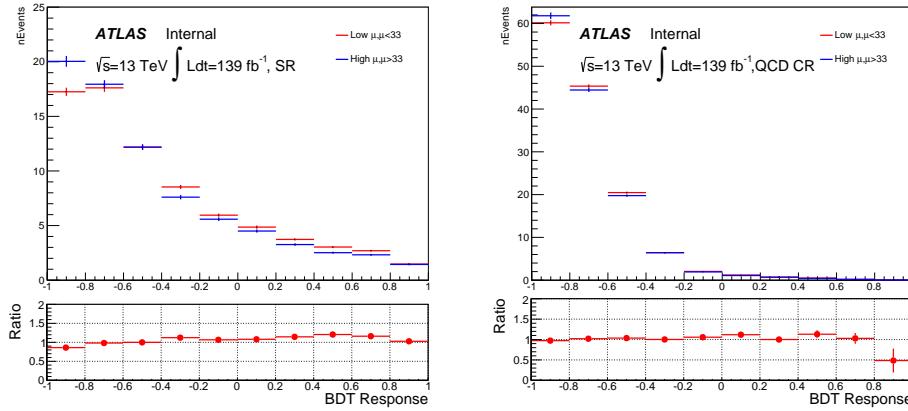


Fig. 5.9 MD distribution for QCD- $Z Z jj$ process in low and high pile-up events for SR (left) and CR (right).

Moreover, a conservative uncertainty is signed to QCD- $Z Z jj$ process by comparing the sample modelled by SHERPA generator (nominal) with MADGRAPH5_aMC@NLO. The MD shape difference for both SR (left) and QCD CR (right) are shown in figure 5.10. The modelling uncertainty is then calculated from the envelop between nominal and alternative samples as function of MD as one additional shape uncertainty.

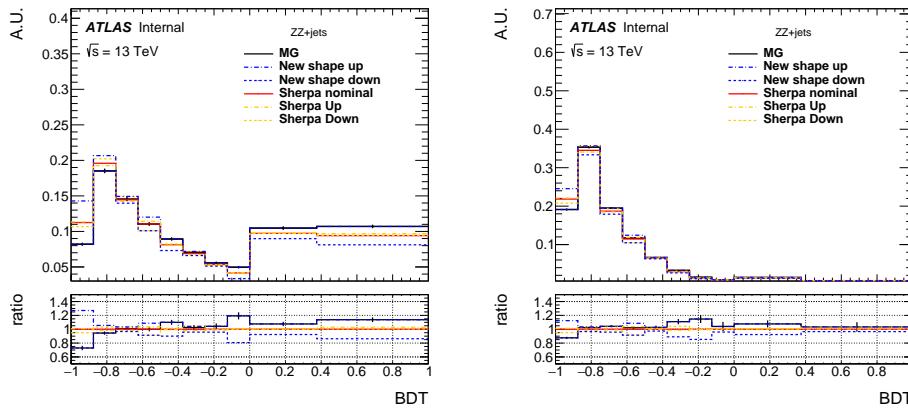


Fig. 5.10 MD shape difference for QCD $q\bar{q} \rightarrow ZZ$ background between different SHERPA theoretical uncertainties and sample from MADGRAPH5_aMC@NLO on SR (left) and CR (right).

5.6 Measurement of fiducial cross section

The fiducial cross section for inclusive $Z Z j j$ production, which includes both EW and QCD components, is then measured. The definition of fiducial volume, which is used for cross section measurement, follows closely to the detector-level selection but use physics objects in particle-level, which are reconstructed in simulation from stable final-state particles, prior to their interactions with the detector.

For electrons and muons, QED final-state radiation is for the most part recovered by adding the four-momenta of surrounding photons that are not originating from hadrons and within an angular distance $\Delta R < 0.1$ to the lepton four-momentum, called lepton “dressing” in truth level. Particle-level jets are built with anti- k_T algorithm with radius parameter $R = 0.4$ using all final-state particles except leptons and neutrinos as inputs. Comparing to the events selection in detector-level in section 5.3, in particle-level, the selected dilepton pair mass required is relaxed to be within 60 to 120 GeV for the reasons of reducing the migration effect, as well as being more compatibility with CMS publication^[76]. All other kinematic selections are the same as the definition in detector-level.

5.6.1 Calculation of C-factor

C-factor is defined as the ratio between the number of selected events in detector-level and the number of particle-level events in fiducial volume (FV):

$$C = \frac{N_{detector-level}}{N_{FV.}} \quad (5.5)$$

The value of C-factor for each $Z Z j j$ process are calculated from each individual simulation samples as listed in table 5.10 together with their systematics.

Process	C	ΔC (stats.)	ΔC (sys.)	ΔC (theo.)
EWK $Z Z j j$	0.663	± 0.002	$\pm^{0.032}_{0.031}$	NA
QCD $q\bar{q} \rightarrow Z Z$	0.702	± 0.003	$\pm^{0.061}_{0.051}$	$\pm^{0.015}_{0.018}$
QCD $gg \rightarrow Z Z$	0.741	± 0.021	$\pm^{0.143}_{0.072}$	± 0.002

Table 5.10 C Factor of different $Z Z j j$ processes.

Then the C from different processes are combined together to be used as inputs for cross section calculation:

$$C = \sum_i \frac{N_{FV.}^i}{\sum_j N_{FV.}^j} \times C_i = 0.699 \pm 0.003(stat.) \pm^{0.011}_{0.013}(theo.) \pm 0.028(exp.) \quad (5.6)$$

The stats. refers to the statistical uncertainty from MC simulation statistics. The theo. and

2025 exp. denote the theoretical and experimental uncertainties described in section 5.5.

2026 **5.6.2 Result of fiducial cross section**

2027 The cross section in fiducial volume is computed as:

$$\sigma^{FV} = \frac{N_{data} - N_{bkg}}{C \times Lumi} \quad (5.7)$$

2028 where N_{data} and N_{bkg} denote the number of events selected from detector-level selection
2029 from data and sum of backgrounds, and C is the C-factor calculated above, Lumi represents
2030 the integrated luminosity of data from 2015 to 2018 of 139 fb^{-1} . Table 5.11 shows the
2031 fiducial cross section for $\ell\ell\ell'\ell'$ final state measured from equation 5.7, as well as the
2032 predicted cross section measured from MC simulation directly.

Measured fiducial σ [fb]	Predicted fiducial σ [fb]
$1.27 \pm 0.12(\text{stat}) \pm 0.02(\text{theo}) \pm 0.07(\text{exp}) \pm 0.01(\text{bkg}) \pm 0.03(\text{lumi})$	$1.14 \pm 0.04(\text{stat}) \pm 0.20(\text{theo})$

Table 5.11 Measured and predicted fiducial cross-sections in $\ell\ell\ell'\ell' jj$ final-state. Uncertainties due to different sources are presented.

2033 The measured cross section has a total uncertainty of 11%, and is found to be compat-
2034 ible with SM prediction. This measurement is still dominant by data statistic.

2035 **5.7 Search for EW- $ZZjj$**

2036 **5.7.1 MD discriminant**

2037 To further separate the EW- $ZZjj$ component from QCD- $ZZjj$, a MD based on Gra-
2038 dient Boosted Decision Tree (BDTG) algorithm^[93] is trained with simulated events via
2039 TMVA framework^[65]. For $\ell\ell\ell'\ell'$ channel, training is performed between EW (signal)
2040 and QCD (background) processes. Twelve event kinematic variables sensitive to the char-
2041 acteristics of the EW signal is used as input features in training. Table 5.12 lists those
2042 input variables with the order of their ranking provided by TMVA tool. The jet-related
2043 information provides larger sensitivity. Then the MD distributions in both SR and QCD
2044 CR region are used for statistical fit.

2045 **5.7.2 Profile likelihood ratio method**

2046 To examine the compatibility between data and the signal-plus-background hypoth-
2047 esis, a test statistic is driven by using the profile likelihood ratio method. The binned

Rank	Variables	Description
1	m_{jj}	Dijet invariant mass
2	p_T^{j1}	p_T of the leading jet
3	p_T^{j2}	p_T of the sub-leading jet
4	$\frac{p_T(ZZjj)}{H_T(ZZjj)}$	p_T of the $Z Z jj$ system divided by the scalar p_T sum of Z bosons and two jets
5	$y_{j1} \times y_{j2}$	Product of jet rapidities
6	Δy_{jj}	Rapidity difference between two jets
7	Y_{Z2}^*	Rapidity of the second Z boson
8	Y_{Z1}^*	Rapidity of the Z boson reconstructed from the lepton pair with the mass closer to the Z boson mass
9	p_T^{ZZ}	p_T of 4l system
10	m_{ZZ}	Invariant mass of 4l system
11	p_T^{Z1}	p_T of the Z boson reconstructed from the lepton pair with the mass closer to the Z boson mass
12	$p_T^{\ell3}$	p_T of the third lepton

Table 5.12 Input features for the training of MD.

2048 likelihood function is given as

$$\mathcal{L}(\mu, \sigma) = \prod_i^{\text{bins}} \mathcal{L}_{\text{poiss}}(N_{\text{data}} | \mu s(\theta) + b(\theta))_i \times \mathcal{L}_{\text{gauss}}(\theta)_i \quad (5.8)$$

2049 where the Poisson term presents the statistical fluctuations of the data and a Gaussian
 2050 term models the pdf of auxiliary measurement to constrain the systematics. μ denotes the
 2051 signal strength of EW- $Z Z jj$ process, computed as the ratio between measured (expected)
 2052 cross section to the SM prediction. θ presents the nuisance parameter, which is the set of
 2053 parameters that parameterize the effect of systematic uncertainties described in section 5.5
 2054 following the Gaussian distribution. N_{data} is the number of selected data events, while the
 2055 $s(\theta)$ is the expected signal yield and $b(\theta)$ is the expected background yield as the function

2056 of nuisance parameters.

2057 The test statistic q_μ is defined as:

$$q_\mu = -2 \ln \left(\frac{\mathcal{L}(\mu, \hat{\theta}_\mu)}{\mathcal{L}(\hat{\mu}, \hat{\theta})} \right) \quad (5.9)$$

2058 in which $\mathcal{L}(\hat{\mu}, \hat{\theta})$ is the unconditional likelihood with respect to both μ and θ , and $\mathcal{L}(\mu, \hat{\theta}_\mu)$
 2059 is the conditional likelihood for a constant μ . Signal-like data distributions are more likely
 2060 to have a low test-statistic (q_μ close to 0) while the contributions of background-like data
 2061 have a larger q_μ . Under the background-only hypothesis, the compatibility of the observed
 2062 (Asimov) data with the prediction is calculated to obtain the observed (expected) signifi-
 2063 cance respectively.

2064 **5.7.3 Fitting procedure**

2065 A profile likelihood fit is performed on MD discriminant to extract the EW- $Z Z jj$
 2066 signal from backgrounds. The binning of MD distributions in SR is optimized to maximize
 2067 the sensitivity for detecting EW signal. The normalization of QCD- $Z Z jj$ production
 2068 (μ_{QCD}^{III}) in $\ell\ell\ell'\ell'$ channel is determined by data from simultaneously fit in SR and QCD
 2069 CR as described in section 5.4. The signal strength of EW- $Z Z jj$ production (μ_{EW}) is
 2070 taken as parameter of interest and floated in the fit. The effects of the uncertainties related
 2071 to normalizations and shapes described previously in section 5.5 of background processes
 2072 in the MD distribution are all taken into account.

2073 In most case, a common nuisance parameter is used for each source of systematic in
 2074 all bins and all channels. The statistical uncertainties for simulated samples are uncorre-
 2075 lated among all bins, and the background uncertainties only applied to their corresponding
 2076 backgrounds. Furthermore, to be more conservative, the generator modelling uncertainty
 2077 for QCD- $Z Z jj$ production mentioned in section 5.5 is separated to be two nuisance pa-
 2078 rameters in low and high MD region.

2079 **5.7.4 Result of statistical fit**

2080 The results of statistical fit for $\ell\ell\ell'\ell'$ final state are presented in table 5.13. To drive
 2081 expected results, the observed data is used for QCD CR to extract normalization factor of
 2082 QCD component (μ_{QCD}^{III}), while in SR, asimov data built from background prediction and
 2083 signal model with SM assumed cross section is used.

2084 As a conclusion, the background-only hypothesis is rejected at 5.5σ (3.9σ) for data
 2085 (expectation), which leads to the observation of EW- $Z Z jj$ production.

	μ_{EW}	$\mu_{\text{QCD}}^{\ell\ell\ell'\ell' jj}$	Significance Obs. (Exp.)
$\ell\ell\ell'\ell' jj$	1.54 ± 0.42	0.95 ± 0.22	$5.48 (3.90) \sigma$

Table 5.13 Observed μ_{EW} and $\mu_{\text{QCD}}^{\ell\ell\ell' jj}$, as well as the observed and expected significance from the individual $\ell\ell\ell'\ell'$ channel. The full set of systematic uncertainties are included.

Figure 5.11 shows the post-fit MD distributions for $\ell\ell\ell'\ell'$ channel in SR (left) and QCD CR (right). The EW- $ZZjj$ cross section measured in $\ell\ell\ell'\ell'$ channel is extracted to be $0.94 \pm 0.26 \text{ fb}$.

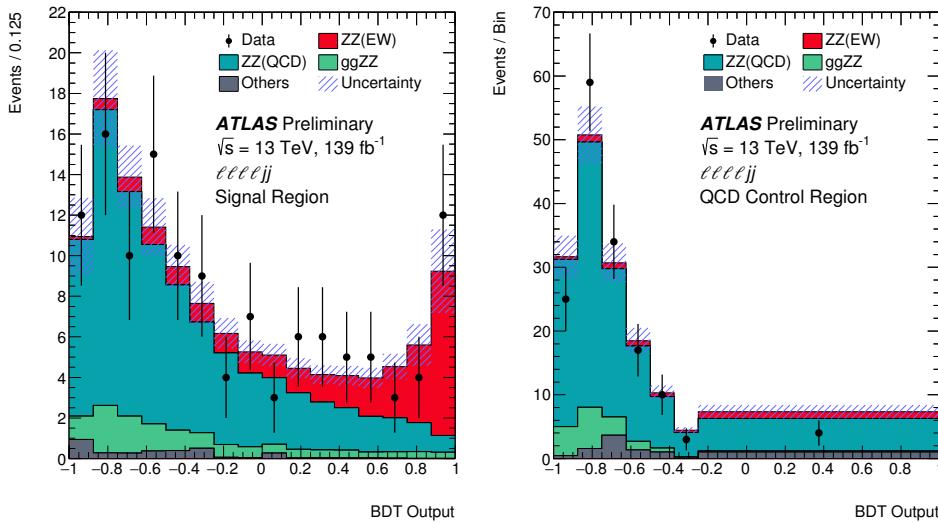


Fig. 5.11 Observed and expected multivariate discriminant distributions after the statistical fit in the $\ell\ell\ell'\ell'$ SR (left) and QCD CR (right). The error bands include the experimental and theoretical uncertainties, as well as the uncertainties in μ_{EW} and $\mu_{\text{QCD}}^{\ell\ell\ell' jj}$. The error bars on the data points show the statistical uncertainty on data.

Figure 5.12 shows the m_{jj} distribution in SR (left) and QCD CR (right), where the normalization of EW and QCD processes are scaled according to their observed value in table 5.13. High m_{jj} region is more sensitive for EW- $ZZjj$ events detection from this figure. Figure 5.13 shows the spectrum of invariant mass of $\ell\ell\ell'\ell'$ system (m_{ZZ}) in SR also with the normalization of EW and QCD processes scaled.

Figure 5.14 is the display of one event candidate of EW- $ZZjj$ production in $2e2\mu$ final state with two jets in forward and backward region.

5.8 Prospect study of EW- $ZZjj$ production in HL-LHC

The High-Luminosity Large Hadron Collider (HL-LHC) project aims to increase the luminosity by a factor of 10 beyond the LHC's design value to increase the potential for discoveries after 2025. The designed luminosity will reach 3000 fb^{-1} with the centre-of-

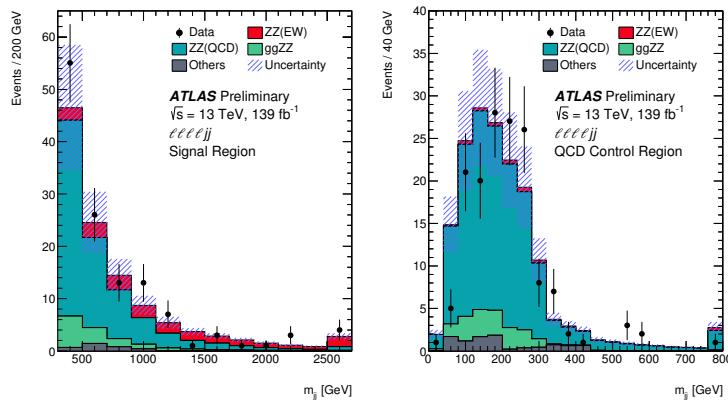


Fig. 5.12 Observed and expected m_{jj} distributions in SR (left) and QCD CR (right). The error bands include the expected experimental and theoretical uncertainties. The error bars on the data points show the statistical uncertainty. The contributions from the QCD and EW production of $Z Z j j$ events are scaled by 0.96 and 1.35, respectively, corresponding to the observed normalization factors in the statistical fit. The last bin includes the overflow events.

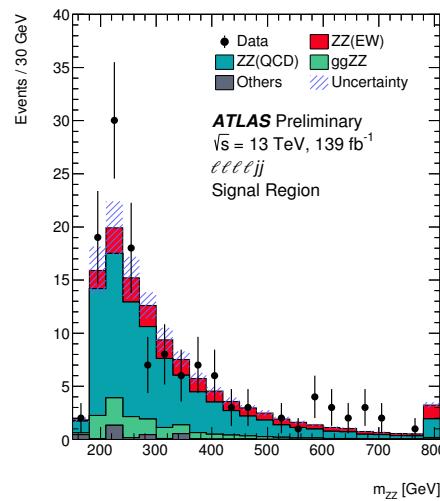


Fig. 5.13 Observed and expected m_{ZZ} spectrum in SR. The error bands include the expected experimental and theoretical uncertainties. The error bars on the data points show the statistical uncertainty. The contributions from the QCD and EW production of $Z Z j j$ events are scaled by 0.96 and 1.35, respectively, corresponding to the observed normalization factors in the statistical fit. The last bin includes the overflow events.

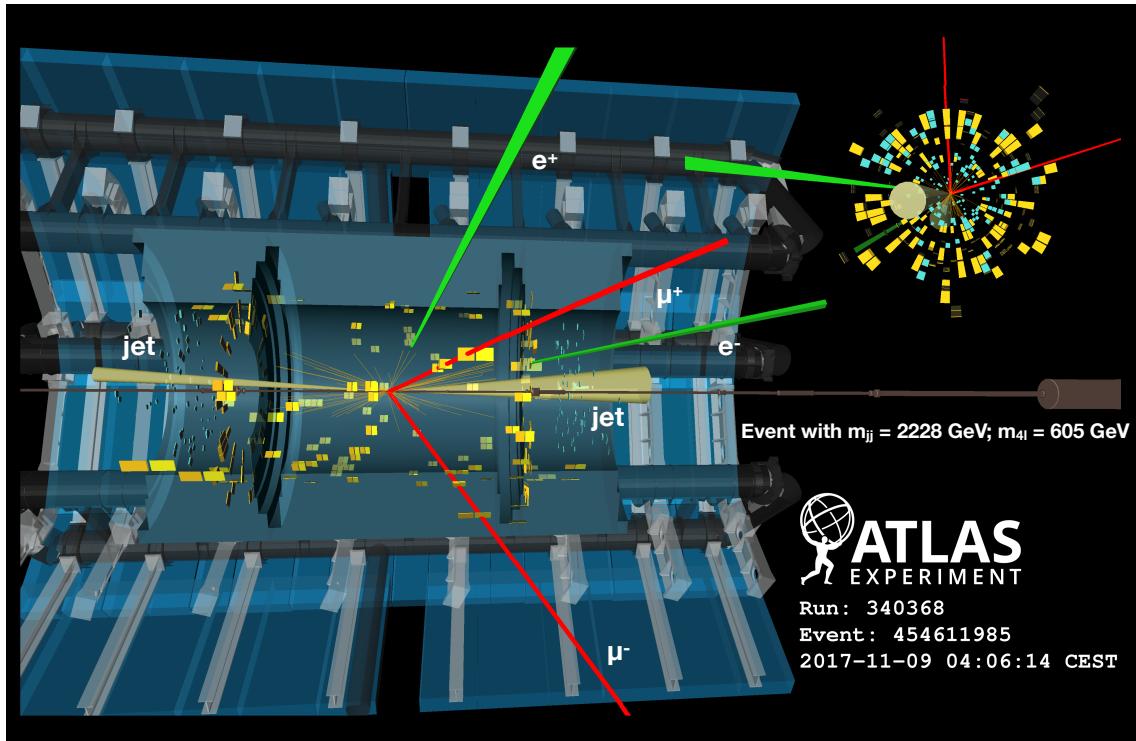


Fig. 5.14 Display of an event candidate of EW- $Z Z jj$ production in $2e2\mu$ channel in last MD bin ($0.875 < \text{MD} < 1.0$). The invariant mass of the di-jet (four-lepton) system is 2228 (605) GeV.

2100 mass energy of 14 TeV.

2101 As introduced in previous sections, with full run-2 data of 139 fb^{-1} collected by AT-
 2102 LAS detector at the LHC, the EW- $Z Z jj$ production is the last channel of observation for
 2103 VBS processes with massive boson due to its very low cross section in $Z Z$ decay. So
 2104 we expect that this channel will benefit significantly from the increased luminosity at the
 2105 HL-LHC, and can be studied in great details for this known mechanism.

2106 In this section, a prospect study is performed for EW- $Z Z jj$ production at the HL-
 2107 LHC in the $\ell\ell'\ell'\ell'$ channel. The study uses 3000 fb^{-1} of simulated pp collision data at a
 2108 centre-of-mass energy of 14 TeV as expected to be recorded by the ATLAS detector at
 2109 HL-LHC. All simulated events are produced at particle-level, and the detector effects of
 2110 leptons and jets reconstruction and identification are estimated by corrections assuming
 2111 the mean number of interactions per bunch crossing ($\langle\mu\rangle$) of 200.

2112 5.8.1 The ATLAS detector at HL-LHC

2113 As the expectation of HL-LHC, the new Inner Tracker (ITk)^[94] will extend the track-
 2114 ing acceptance capability of ATLAS detector to pseudorapidity ($|\eta|$) up to 4.0. By in-
 2115 cluding a forward muon trigger, the upgraded Muon Spectrometer^[95] is also expected to
 2116 provide muon identification capabilities to $|\eta|$ up to 4.0. In addition, the new high gran-

ularity timing detector (HGTD)^[96] that designed to mitigate the pile-up (PU) effects is also expected to be installed in the forward region of $2.4 < |\eta| < 4.0$. More details of expected performance of the upgraded ATLAS detector at the HL-LHC has been reported in Ref.^[97].

5.8.2 Simulation

The analysis is performed using particle-level events. The samples are generated at $\sqrt{s} = 14$ TeV and with a fast simulation based on the setting for ATLAS detector at the HL-LHC. The signal in this analysis is EW- $ZZjj$ process, while only the dominant irreducible background of QCD- $ZZjj$ is considered. Both signal and background are generated using SHERPA with the NNPDF3.0NNLO PDF set. The signal sample is modelled with two jets at Matrix Element (ME) level. The background is generated with up to one (three) outgoing partons at NLO (LO) in pQCD. As a quick study, other minor backgrounds such as fake backgrounds from $Z+jets$ and top-quark processes, as well as Diboson without 4l final-state and Triboson processes are not considered into this analysis. Furthermore, for hard scattering events, the pile-up collisions are set with a mean value of 200 interactions per bunch crossing. Signal and background yields are then scaled to an integrated luminosity of 3000 fb^{-1} as expected at the HL-LHC.

5.8.3 Event selection

The analysis selection follows closely to the one in ATLAS run-2 analysis as described in section 5.3. Here are some changes according to the expectation of the HL-LHC scenario for ATLAS detector:

- Extend the lepton (both electron and muon) identification in forward with $|\eta| < 4.0$
- Pile-up (PU) jet suppression is applied with a PU rejection factor of 50 for all PU jets in the region of $|\eta| < 3.8$, based on the expected ATLAS detector performance at the HL-LHC.
- The jets are required to have $p_T > 30$ (70) GeV in the $|\eta| < 3.8$ ($3.8 < |\eta| < 4.5$) region.
- For two selected jets, tight the m_{jj} requirement to > 600 GeV, and require $\Delta\eta_{jj} > 2$.

In addition, a fiducial volume, which is used to study the expected precision of the cross-section measurements, is defined at particle-level with the same kinematic requirements listed above.

Table 5.14 summarized the number of selected signal and background events normalized to 3000 fb^{-1} . In addition to the *baseline* selection listed above, to compare the

2150 different detector scenarios at the HL-LHC, two alternative selections are also studied:

- 2151 • Reduce the lepton η region to 2.7, to understand the effect due to forward lepton
2152 reconstruction and identification with the upgraded ATLAS detector.
- 2153 • Only apply the PU jet suppression with region $|\eta| < 2.4$, to measure the improve-
2154 ment of *baseline* by extending the rejection range of PU jets at the HL-LHC with
2155 the installation of HGTD.

Selection	$N_{\text{EW-ZZjj}}$	$N_{\text{QCD-ZZjj}}$	$N_{\text{EW-ZZjj}} / \sqrt{N_{\text{QCD-ZZjj}}}$
Baseline	432 ± 21	1402 ± 37	11.54 ± 0.58
Leptons with $ \eta < 2.7$	373 ± 19	1058 ± 33	11.46 ± 0.62
PU jet suppression only in $ \eta < 2.4$	536 ± 23	15470 ± 120	4.31 ± 0.19

**Table 5.14 Comparison of event yields for signal ($N_{\text{EW-ZZjj}}$) and background ($N_{\text{QCD-ZZjj}}$) pro-
cesses, and expected significance of EW- $ZZjj$ processes, normalized to 3000 fb^{-1}
data at 14 TeV, with baseline and alternative selections. Uncertainties in the table
refer to expected data statistical uncertainty at 14 TeV with 3000 fb^{-1} .**

2156 From this table, one can see the extended track coverage increases the $\ell\ell\ell'\ell'jj$ events
2157 by 15 to 30%, via improving the lepton efficiency. But the significance of searching for
2158 EW- $ZZjj$ process does not improve so much due to the large increment of background
2159 events.

2160 Figure 5.15 shows the kinematic distributions of di-jet invariant mass (m_{jj}), the ZZ
2161 invariant mass (m_{ZZ}) and the ϕ separation of two Z bosons ($|\Delta\phi(ZZ)|$) as well as the
2162 centrality of the ZZ system. The ZZ centrality is defined as:

$$ZZ \text{ centrality} = \frac{|y_{ZZ} - (y_{j1} + y_{j2})/2|}{|y_{j1} - y_{j2}|} \quad (5.10)$$

2163 To measure the event yield, the top panel shows the stack distribution for EW- and QCD-
2164 $ZZjj$ processes, while bottom panel is the ratio between two processes.

2165 5.8.4 Systematics

2166 According to studies in section 5.5, the dominant systematic in $\ell\ell\ell'\ell'$ channel is from
2167 theoretical systematic for QCD- $ZZjj$ background process. Different sizes of systematics
2168 have been studied, at a factor of 5, 10 and 30% on background modelling. The 5% un-
2169 certainty is an optimal estimation when there is enough data events from QCD-enriched
2170 control region at the HL-LHC that can be used to constrain the theoretical normaliza-
2171 tion on QCD- $ZZjj$ process. The 30% one is a conservative estimation, in which the
2172 uncertainties are directly calculated from different PDF sets and QCD renormalization

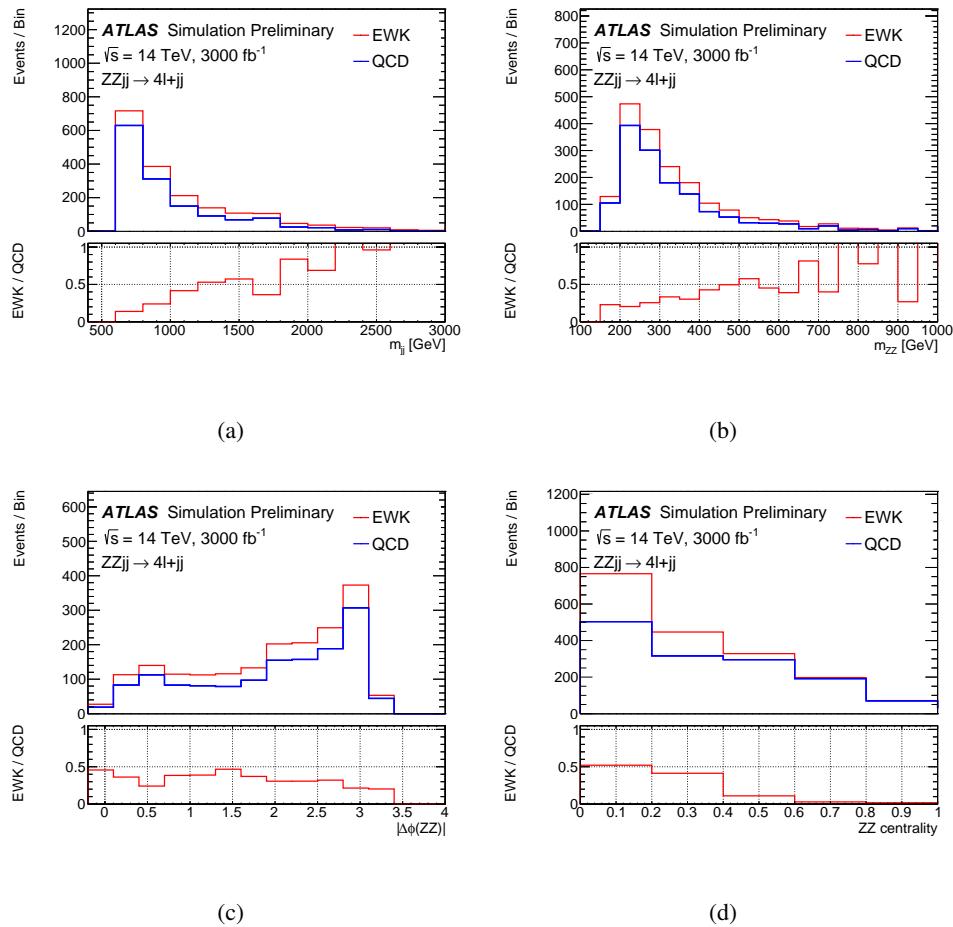


Fig. 5.15 Detector-level distributions of EW- and QCD- $Z Z j j$ processes with selected events in defined phase space at 14 TeV of (a) $m_{j j}$, (b) $m_{Z Z}$, (c) $|\Delta\phi(Z Z)|$, (d) $Z Z$ centrality, normalized to 3000 fb^{-1} .

and factorization scales, following recommendation from the PDF4LHC mentioned in section 5.5.

For experimental sources, the jet systematics have been checked following the setting provided by the HL-LHC in Ref. [97], and the uncertainties are within 5% level, which is smaller than run-2 measurement at 10%. Figure 5.16 depicts the up and down variations for jet uncertainty provided by the HL-LHC performance tool as function of dijet invariant mass ($m_{j j}$). Therefore, a conservative 5% uncertainty is used as experimental uncertainty.

Since the final result relies greatly on the uncertainties, especially the theoretical uncertainties on QCD- $Z Z j j$ production. So results with different uncertainty conditions are shown as below:

- The case with statistical uncertainty of simulated samples only.
- The case with statistical and experimental uncertainties (5%)
- The case with statistical, experimental and additional theoretical uncertainties at

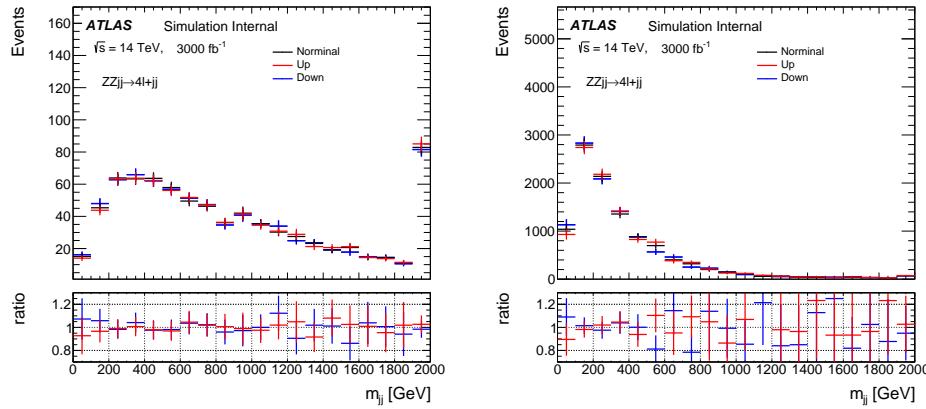


Fig. 5.16 Jet variations on m_{jj} distribution for EW- $Z Z jj$ (left) and QCD- $Z Z jj$ (right) processes with luminosity of 3000 fb^{-1} at 14 TeV. *Upgrade Performance Function* is used to extract the uncertainties with *baseline* setting.

2186 5%, 10% and 30% levels respectively.

2187 Three different sources of uncertainties are treated as uncorrelated and summed up
2188 quadratically.

2189 5.8.5 Results

2190 In this analysis, instead of a statistical fit, the expected significance of EW- $Z Z jj$
2191 production is calculated as:

$$\text{Significance} = \frac{S}{\sqrt{\sigma(B)_{\text{stat.}}^2 + \sigma(B)_{\text{syst.}}^2}}, \quad (5.11)$$

2192 where S presents the number of selected signal events, and $\sigma(B)_{\text{stat.}}$ and $\sigma(B)_{\text{syst.}}$ denote
2193 the statistical and systematic (exp. + theo.) uncertainties from background processes. The
2194 statistical uncertainty is computed from expected data yield with an integrated luminosity
2195 of 3000 fb^{-1} .

2196 Base on baseline selection of $m_{jj} > 600 \text{ GeV}$, a additional scan over different m_{jj} cuts
2197 are also performed with a step of 50 GeV under different systematic conditions, as shown
2198 in figure 5.17.

2199 In addition, the expected differential cross section of EW- $Z Z jj$ process is measured
2200 in the defined phase space at 14 TeV, as a function of m_{ZZ} and m_{jj} , shown in figure 5.18.
2201 The expected differential cross sections are calculated as:

$$\sigma = \frac{N_{\text{pseudo-data}} - N_{\text{QCD-ZZjj}}}{L * C_{\text{EW-ZZjj}}} \quad (5.12)$$

$$C_{\text{EW-ZZjj}} = \frac{N_{\text{EW-ZZjj}}^{\text{det.}}}{N_{\text{EW-ZZjj}}^{\text{part.}}}$$

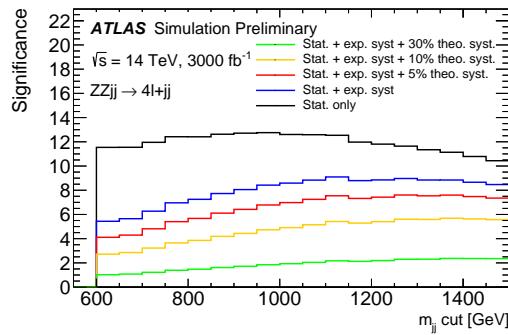


Fig. 5.17 The expected significance of EW- $Z Z_{jj}$ processes as a function of different m_{jj} cut with 3000 fb^{-1} , under conditions of different sizes of theoretical uncertainties on the QCD- $Z Z_{jj}$ background modelling. The statistical uncertainty is estimated from expected data yield at 14 TeV with 3000 fb^{-1} . Different uncertainties are summed up quadratically.

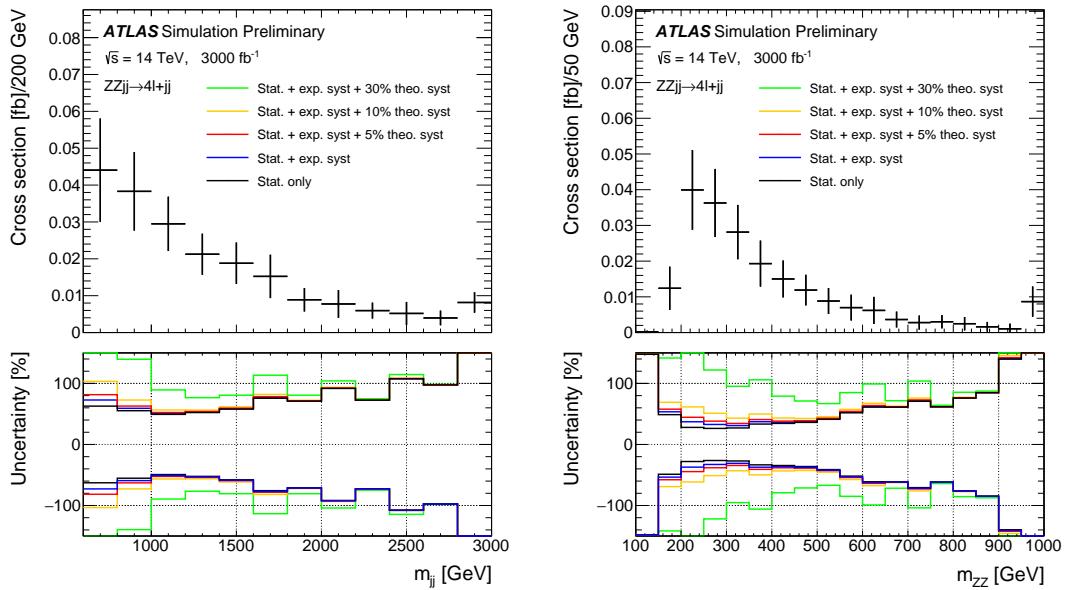


Fig. 5.18 The projected differential cross-sections at 14 TeV for the EW- $Z Z_{jj}$ processes as a function of m_{jj} (left) and m_{ZZ} (right). The top panel shows measurement with statistical only case, where statistical uncertainty is estimated from expected data yield at 14 TeV with 3000 fb^{-1} . The bottom panel shows impact of different sizes of systematic uncertainties.

where $N_{pseudo-data}$ denotes the expected number of data events with 3000 fb^{-1} luminosity at 14 TeV, and $N_{QCD-ZZjj}$ and $N_{EW-ZZjj}$ are the number of predicted events of QCD- $Z Z jj$ and EW- $Z Z jj$ processes in particle-level. The $C_{EW-ZZjj}$ factor represents the detector efficiency for EW- $Z Z jj$ processes introduced in section 5.6.1. The interference between EW- and QCD- $Z Z jj$ processes is ignored due to its minor contribution.

The number of expected integrated cross section as well as its uncertainty under different systematic conditions are shown in table 5.15 with 3000 fb^{-1} luminosity at 14 TeV. The statistical uncertainty is at 10% level when with such large luminosity. The result is dominated by systematics and can reach 100% level when theoretical modelling uncertainty is 30% for QCD- $Z Z jj$ processes.

	Cross section [fb]	Stat. only	Plus exp.	Plus 5% theo.	Plus 10% theo.	Plus 30% theo.
EW- $Z Z jj$	0.21	± 0.02	± 0.04	± 0.05	± 0.08	± 0.21

Table 5.15 Summary of expected cross-section measured with different theoretical uncertainties. The statistical uncertainty is computed from expected data yield with 3000 fb^{-1} at 14 TeV. Different uncertainties are treated as uncorrelated and summed quadratically.

2212 5.9 Conclusion

2213 The fiducial cross section for inclusive $ZZjj$ production is measured in this section,
2214 with a total relative uncertainty of 11% for the $\ell\ell\ell'\ell'$ final state, and found to be com-
2215 patible with the SM prediction. The observation of electroweak production of two jets in
2216 association with a Z -boson pair decay to $\ell\ell\ell'\ell'$ final state using 139 fb^{-1} of 13 TeV pp
2217 collision data collected by ATLAS experiment at the LHC is presented in this section.
2218 The search for electroweak production of two jets in association with a Z -boson pair is
2219 based on multivariate discriminants (MD) to enhance the separation between the signal
2220 and backgrounds. In $\ell\ell\ell'\ell'$ final state, the background-only hypothesis is rejected with
2221 an observed (expected) significance of 5.5 (3.9) σ , which gives the first observation of
2222 electroweak production in $ZZjj$ channel.

2223 In addition, the prospect study for the EW- $ZZjj$ production at the HL-LHC in the
2224 $\ell\ell\ell'\ell'$ channel, using 3000 fb^{-1} simulated pp collision data at a centre-of-mass energy
2225 of 14 TeV has been presented. The precision of the expected measurements of the inte-
2226 grated and differential cross sections as a function of dijet or $\ell\ell\ell'\ell'$ invariant mass are
2227 shown. Under the assumption of theoretical uncertainty for the QCD- $ZZjj$ processes and
2228 experimental uncertainty for jets being constraint at 5% level respectively, with statistical
2229 uncertainty in 3000 fb^{-1} being considered, the observation of the EW- $ZZjj$ process can
2230 reach a significance of 7σ .

2231 **Chapter 6 Search for heavy resonances decaying into**
2232 **a pair of Z bosons in $\ell\ell\ell'\ell'$ final state using pp**
2233 **collision data collected by ATLAS detector from 2015 to**
2234 **2018**

2235 **6.1 Introduction**

2236 A new particle was discovered by the ATLAS and CMS Collaborations at the
2237 LHC^[67-68] in 2012. Both two experiments have confirmed that the properties including
2238 spin, couplings and parity of this new particle are consistent with Higgs boson predicted in
2239 the Standard Model (SM), which is an important milestone in understanding of the mecha-
2240 nism of EWSB. Nevertheless, the possibility that this newly discoved particle is just a part
2241 of the extended Higgs sector as predicted by various extensions in the SM cannot be ruled
2242 out. There are many models predicted the existence of new heavy resonances decaying
2243 into dibosons, such as a heavy spin-0 neutral Higgs boson^[98] and the two-Higgs-doublet
2244 models (2HDM)^[99], as well as the spin-2 Kaluza–Klein (KK) excitations of the graviton
2245 (G_{KK})^[100].

2246 Though with smaller branching ratio comparing to semileptonic or fully hadronic de-
2247 cay channels, the $\ell\ell\ell'\ell'$ final state has its unique sensitivity in mass range smaller than
2248 1 TeV region due to its good mass resolution and relative smaller experimental and the-
2249 oretical systematics. This section presents the search for heavy resonance decaying into
2250 a pair of Z bosons to the $\ell\ell\ell'\ell'$ final state, in which ℓ denotes to either an electron
2251 or a muon. Several signal hypotheses are considered. The first hypothesis is a heavy
2252 Higgs boson (spin-0 resonance) under the narrow-width approximation (NWA). Then as
2253 several theoretical models prefer non-negligible natural widths, the models under large-
2254 width approximation (LWA), assuming widths of 1%, 5%, 10% and 15% of the resonance
2255 mass, are also studied. In addition, the graviton excitations (spin-2 resonance) under the
2256 Randall–Sundrum model are also searched. It is assumed that the heavy resonance is
2257 produced predominantly via the gluon-gluon Fusion (ggF) and the Vector Boson Fusion
2258 (VBF) productions, but with the unknown ratio of two production rates. So the results are
2259 separated for ggF and VBF production modes. To gain more sensitivity, the $\ell\ell\ell'\ell'$ events
2260 are classified into ggF- and VBF-enriched categories. Moreover, for the NWA model, the
2261 categorizations are studied under both cut-based and multivariate (MVA) -based methods,
2262 the details of categorization are shown in following sections.

2263 The search uses the four-lepton invariant mass in the range of 200 GeV to 2000 GeV for
2264 signal hypothesis of spin-0 resonance under the NWA model, and from 400 GeV to
2265 2000 GeV for the one under the LWA models. And the spin-2 graviton signals are searched
2266 in the mass range from 600 GeV to 2000 GeV. The data collected by ATLAS detector at
2267 the LHC from 2015 to 2018 at the centre-of-mass energy of 13 TeV is used. In case of no
2268 excess, upper limits on the production rate of different signal hypotheses are computed
2269 from statistical fits to m_{4l} distribution.

2270

6.2 Data and MC samples

2271

6.2.1 Data samples

2272 The data used in this analysis are collected by ATLAS detector at the centre-of-mass
2273 energy of 13 TeV during the years of 2015 to 2018. Only events passing the latest Good
2274 Run List (GRL) released by the Data Quality group from ATLAS experiment as listed in
2275 section 5.2.1 corresponding to an integrated luminosity of $139.0 \pm 2.4 \text{ fb}^{-1}$ are used.
2276 Table 6.1 listed the recorded integrated luminosity, average and peak pile-up of each year's
data.

Table 6.1 Summary of the recorded integrated luminosity (lumi), average and peak pile-up (PU) of data from 2015 to 2018.

Year	recorded integrated lumi	lumi after GRL	average PU	peak PU
2015	3.86 fb^{-1}		13.4	28.1
2016	35.6 fb^{-1}	36.2 fb^{-1}	25.1	52.2
2017	46.9 fb^{-1}	44.3 fb^{-1}	37.8	79.8
2018	60.6 fb^{-1}	58.5 fb^{-1}	36.1	88.6

2278

6.2.2 Background MC simulations

2279 Background processes considered in this analysis include ZZ ($q\bar{q} \rightarrow ZZ$, $gg \rightarrow ZZ$), triboson (WWZ , WZZ , ZZZ), $Z+jets$ and top-quark ($t\bar{t}$, ttV) processes.

2280 The QCD $q\bar{q} \rightarrow ZZ$ process is modelled using SHERPA 2.2.2^[79] with the
2281 NNPDF3.0NNLO^[80] PDF, where events with up to one (three) outgoing partons are gen-
2282 erated at NLO (LO) in pQCD. The production of ZZ from the gluon-gluon initial state
2283 with a four-fermion loop or with an exchange of the Higgs boson, which has an order of α_S^4
2284 in QCD, is not included in this SHERPA simulation. So a separate gg induced ZZ sample
2285 including the continuum background, the SM Higgs boson, and the interference contri-
2286 bution is modelled using SHERPA 2.2.2 with the NNPDF3.0NNLO PDF set, and with an
2287

additional 1.7 k-factor^[81] being applied. The EW- $ZZjj$ production is simulated using SHERPA 2.2.2 with the NNPDF3.0NNLO PDF, and the $ZZZ \rightarrow \ell\ell\ell'\ell'qq$ process is also taken into account in this sample.

The $Z+jets$ events are generated using SHERPA 2.2.2 with the NNPDF3.0NNLO PDF, in which the ME is calculated for up to two partons with next-to-leading-order (NLO) accuracy in pQCD and up to four partons with LO accuracy. The $Z+jets$ events are normalized using the next-to-next-to-leading-order (NNLO) cross section. The triboson processes with full leptonic decays and at least four prompt charged leptons are generated using SHERPA 2.1.1. For top-quark pair ($t\bar{t}$) production and the single top-quark productions in t -channel, s -channel and Wt -channel, the PowHEG-Box v2 is used with the CT10 PDF. The productions of $t\bar{t}$ in association with Z boson(s) ($t\bar{t}Z$) is modelled with MADGRAPH5_aMC@NLO.

6.2.3 Signal MC simulations

One model considered in this analysis is heavy spin-0 resonance under the Narrow Width Approximation (NWA) simulated using PowHEG-Box v2 MC event generator with the CT10 PDF. The gluon-gluon fusion (ggF) production mode and vector-boson fusion (VBF) production mode are calculated separately with matrix elements up to NLO in QCD. The PowHEG-Box is interfaced to PYTHIA8 for parton showering, and for decaying the Higgs boson into the $H \rightarrow ZZ \rightarrow \ell\ell\ell'\ell'$ final states. Events of NWA signal are generated at mass points between 200 GeV to 2000 GeV using the step of 100 (200) GeV up to (above) 1 TeV in both ggF and VBF production modes.

In addition, heavy Higgs boson events under the Large Width Approximation (LWA) with widths of 1%, 5%, 10% and 15% of the boson mass are generated using MADGRAPH5_aMC@NLO 2.3.2 interfaced to PYTHIA8. Only ggF production is considered. Mass points between 400 GeV to 2000 GeV are simulated with the step of 100 (200) GeV up to (above) 1 TeV. To describe jet multiplicity, MADGRAPH5_aMC@NLO is used to simulate process of $pp \rightarrow H + \geq 2\text{jets}$ at NLO in QCD with the FxFx merging scheme^[101].

Spin-2 Kaluza–Klein (KK) gravitons (G_{KK}) from the Bulk Randall–Sundrum model^[102] are also studied in this analysis. Events are generated by MADGRAPH5_aMC@NLO at LO in QCD, which is then interfaced to PYTHIA8 for parton showering. The G_{KK} -gluon coupling $k/\overline{M}_{\text{Planck}}$, where k is the curvature scale of the extra dimension and $\overline{M}_{\text{Planck}}$ is the reduced Planck mass, is set to 1. The width of the resonance is correlated with the coupling $k/\overline{M}_{\text{Planck}}$ and in this configuration is around

2322 6% of its mass. The mass of the G_{KK} is the only free parameter in this simplified model.
2323 Mass points between 600 GeV to 2 TeV with 200 GeV spacing were generated.

2324

6.3 Analysis selections

2325

6.3.1 Objects selection

2326 Similar to VBSZZ analysis in section 5.3, the selection of this analysis relies on the
2327 definition of multiple objects: *electrons*, *Muons*, and *jets*. Details of definitions for each
2328 object are described as below:

2329 **Electron:** As described in section 4.2.3, electrons are reconstructed from energy
2330 deposits in the EM calorimeter matched to a track in the inner detector. The electron
2331 candidates satisfying the *Loose* criterion valuing by the likelihood-based (LH) method
2332 are selected, with a selection efficiency ranging from 90% for transverse momentum
2333 $p_T = 20$ GeV to 96% for $p_T > 60$ GeV. In addition, the electrons are required to have
2334 $p_T > 7$ GeV, $|\eta| < 2.47$ and $|z_0 \sin\theta| < 0.5$ mm.

2335 **Muon:** To increase the acceptance range in reco-level for $\ell\ell\ell'\ell'$ channel, all four
2336 types of muons (CB, ST, CT, ME muons, described in section 4.2.4) are used. The CT
2337 muons are required to pass $p_T > 15$ GeV and $|\eta| < 0.1$, while the ST muons are also
2338 limited in $|\eta| < 0.1$ region. The ME muons are only used in the region of $2.5 < |\eta| < 2.7$.
2339 And at most one CT, ST or ME muon is allowed in one $\ell\ell\ell'\ell'$ quadruplet. The Muon
2340 candidates are required to pass $p_T > 5$ GeV and $|\eta| < 2.7$, and satisfy the *Loose* identifi-
2341 cation criterion with an efficiency of at least 98.5%. The impact parameter requirements
2342 of $|d_0| < 1$ mm and $|z_0 \sin\theta| < 0.5$ mm are further applied.

2343 **Jets:** Jets are clustered using the anti- k_t algorithm with radius parameter $R =$
2344 0.4 implemented in the FASTJET package as described in section 4.2.5. The ‘particle
2345 flow’ (PFlow) objects^[103], which combines measurements from both the tracker and the
2346 calorimeter, are used as inputs to the FASTJET package. The energy deposited in the
2347 calorimeter by all charged particles is removed, and the jet reconstruction is performed on
2348 an ensemble of PFlow objects consisting of the remaining calorimeter energy and tracks
2349 which are matched to the hard interaction. This improves the accuracy of the charged-
2350 hadron measurement, while retaining the calorimeter measurements of neutral-particle
2351 energies. Compared to only using topological clusters, jets reconstructed with the particle
2352 flow algorithm with $p_T > 30$ GeV have approximately 10% better transverse momentum
2353 resolution. The jets used in this analysis are then required to have $p_T > 30$ GeV and
2354 $|\eta| < 4.5$. To further reduce the effects of pile-up jets, a jet vertex tagger (JVT) is applied

2355 to jets with $p_T < 60$ GeV and $|\eta| < 2.4$.

2356 **Overlap removal:** As the selected jet and lepton candidates can be reconstructed
 2357 from same detector information, an overlap-removal procedure is applied. For electron
 2358 and muon sharing the same ID track, the electron is selected in the case that the muon is
 2359 calorimeter-tagged and does not have a MS track, or is a segment-tagged muon, otherwise
 2360 the muon is selected. The jet overlapping with electron (muon) within a cone of size of
 2361 $\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.2(0.1)$ are removed.

2362 **6.3.2 Event selection**

2363 First of all, the four-lepton events are required to pass single or multi-lepton triggers.
 2364 Due to the increasing of peak luminosity and pile-up, the p_T and E_T thresholds of triggers
 2365 increase slightly during the data-taking periods from 2015 to 2018. Table 6.2 summarizes
 2366 the triggers used for $\ell\ell\ell'\ell'$ channel. The overall trigger efficiency for selected signal
 2367 events passing final selection is around 98%.

**Table 6.2 Summary of the p_T (E_T) trigger thresholds (in GeV) employed for the muon (electron)
 trigger selection in the year of 2015, 2016, 2017, and 2018.**

Trigger item	Trigger threshold			
	2015	2016	2017	2018
single muon	$\mu 20; \mu 50; \mu 60$	$\mu 24; \mu 26; \mu 40; \mu 50$	$\mu 26; \mu 50; \mu 60$	$\mu 26; \mu 50; \mu 60$
single electron	$e24; e60; e120$	$e26; e60; e140; e300$	$e26; e60; e140; e300$	$e26; e60; e140; e300$
dimuon	$2\mu 10; \mu 18_{-\mu 8}$	$2\mu 10; 2\mu 14; \mu 22_{-\mu 8}$	$2\mu 14; \mu 22_{-\mu 8}$	$2\mu 14; \mu 22_{-\mu 8}$
dielectron	$2e12$	$2e15; 2e17$	$2e17; 2e24$	$2e17; 2e24$
electron-muon	$e24_{-\mu 8}$	$e24_{-\mu 8}; e26_{-\mu 8}$ $e17_{-\mu 14}; e7_{-\mu 24}; 2e12_{-\mu 10}; e12_{-\mu 10}$	$e26_{-\mu 8}$	$e26_{-\mu 8}$
trimuon	$\mu 18_{-2\mu 4}$	$\mu 11_{-2\mu 4}; \mu 6_{-2\mu 4}; \mu 20_{-2\mu 4}; 3\mu 4$ $3\mu 6$	$4\mu 4; \mu 20_{-2\mu 4}; 3\mu 4$	$\mu 20_{-2\mu 4}$
trilepton	$e17_{-2e9}$	$e17_{-2e9}; e17_{-2e10}$	$e24_{-2e12}$	$e24_{-2e12}$

2368 The $\ell\ell\ell'\ell'$ quadruplets are formed by two opposite-sign, same-flavour (OSSF) lepton
 2369 pairs ($\ell^+\ell^-$). The p_T threshold of first three leading leptons are required to be 20, 15 and
 2370 10 GeV. If there are more than one combination of lepton pairing in quadruplet, the pairing
 2371 is selected by keeping it with the mass of lepton pairs closest (leading pair, refers as m_{12})
 2372 and second closest (sub-leading pair, refers as m_{34}) to Z boson mass. The mass of leading
 2373 pair is required to satisfy $50 < m_{12} < 106$ GeV, while the sub-leading pair is required to
 2374 be less than 115 GeV and larger than 50 GeV.

2375 The two lepton pairs in quadruplet are required to have angular separation with $\Delta R >$
 2376 0.1. To suppress the contribution from $J/\psi \rightarrow \ell\ell$ decays, for 4 μ and 4 e quadruplets,
 2377 the events are rejected if any opposite-sign same-flavour lepton pair is found with mass

2378 below 5 GeV. If there are more than one quadruplets from different channels in event at
 2379 this point, the one with highest expected signal rate is selected in the order of 4μ , $2e2\mu$,
 2380 $4e$. The transverse impact-parameter significance ($|d_0|/\sigma_{d_0}$) for muons (electrons) is than
 2381 required to be smaller than 3 (5) to suppress the backgrounds from heavy-flavour hadrons.

2382 In addition, the track- and calorimeter- based isolation criteria is required for all elec-
 2383 trons and muons to further suppress the reducible backgrounds of $Z+jets$ and $t\bar{t}$. For
 2384 lepton isolation selection, the two track- and calorimeter- based variables, $E_T^{topocone}$ and
 2385 $p_T^{varcone}$ as described in section 4.2.4 (section 4.2.3) for muons (electrons), are vulnerable
 2386 to pileup. For track-based variable, this is because of additional tracks in the event. The
 2387 definition of $p_T^{varcone}$ attempts to limit the tracks used in the calculation to those from the
 2388 vertex via a loose cut of $|z_0 \sin(\theta)| < 3$, which proved to be too loose in new pile-up
 2389 regime 2017 and 2018 datasets. So new track-based variable is used, by adding a require-
 2390 ment that the track be used in determining the vertex, or that, if not, it both pass the cut on
 2391 $|z_0 \sin(\theta)|$ and not be used in determining any other vertex, which makes the track-based
 2392 variable to be more isolation-robust in the high pile-up regime. The new variable is named
 2393 as `ptvarcone[cone]_TightTTVA_pt[pT cut]`, where [cone] is the cone size and [p_T cut] is
 2394 the cutoff for including tracks in the calculation.

2395 For calorimeter-based variable, the calculation of $E_T^{topocone}$ corrects the pile-up effects
 2396 by subtracting an average pileup contribution computed over the whole detector. But with
 2397 the increasing of energy density of pile-up events, the root mean square (RMS) of $E_T^{topocone}$
 2398 variable increases, which leads to the increment of possibility that the pile-up fluctuations
 2399 are not be accounted for correctly. One possible solution is that use particle-flow (PFlow)
 2400 method to calculate the calorimeter isolation. As part of PFlow reconstruction process,
 2401 it assigns the clusters to tracks which improves the track-cluster association for better de-
 2402 termination of the raw value of the E_T in the cone. And using PFlow jets to calculate the
 2403 pileup correction provides a further improvement. So a resulting variable named `neflow-`
 2404 `isol[cone]` is used. Finally, a requirement of isolation, called *FixedCutPFlowLoose*, which
 2405 gives better performance in high piup-up condition is applied to electrons and muons as:
 2406
$$(\max(\text{ptcone20_TightTTVA_pt500}, \text{ptvarcone30_TightTTVA_pt500}) + 0.4 \times \text{neflow-} \\ 2407 \text{isol20}) / p_T < 0.16$$

2408 On the top of impact parameter cut and lepton isolation cut, the four-lepton candidates
 2409 are also required to originate from a common vertex to reduce $Z+jets$ and $t\bar{t}$ backgrounds.
 2410 This is ensured by applying a vertex fit χ^2 cut of 4 ID tracks of lepton candidates satisfying
 2411 $\chi^2/N_{dof} < 6$ (9) for events in 4μ ($4e$ and $2e2\mu$) channel(s).

2412 To improve the mass resolution, the QED process of final state radiation (FSR) pho-

2413 tons in Z boson decays are taken into account in the reconstruction of Z bosons. The
2414 four-momentum of any reconstructed photon that is consistent with having been radiated
2415 from lepton(s) in leading pair are added into final state. Moreover, the four-momenta of
2416 leptons in both (leading and sub-leading) pairs are recomputed by performing a Z -mass-
2417 constrained kinematic fit, which uses a Breit–Wigner Z boson line-shape and Gaussian
2418 function with width set to the expected lepton resolution per lepton to model the momen-
2419 tum response function. The Z -mass-constrained mass improves the $m_{4\ell}$ resolution by up
2420 to 15% depending on m_H .

2421 In summary, table 6.3 lists a comprehensive object and event level selection as de-
2422 scribed above. Table 6.4 to 6.7 shows the cutflow of NWA ggF and VBF signal at the
2423 mass points of 600 and 1000 GeV as examples.

2424

6.3.3 Event categorizations

2425 To improve the sensitivity of search in both VBF and ggF production mode in NWA
2426 model, events are classified into the VBF- and ggF- enriched categories. With the statistic
2427 increasing in full run-2 data, a multivariate (MVA) based classifier has been studied for
2428 NWA signal, while in the meantime the traditional cut-based classifier is also used as a
2429 model-independent result for all three (NWA, LWA, graviton) models.

2430

1. Cut-based categorization

2431 There are four categories in total: one VBF-enriched category and three ggF-enriched
2432 categories. The categorization is defined based on kinematic cuts:

- 2433 • VBF-CBA-enriched category: Events have at least two selected jets as defined in
2434 section 6.3.1, with the two leading jets being separated by $|\Delta\eta_{jj}| > 3.3$ and invariant
2435 mass satisfying $m_{jj} > 400$ GeV;
- 2436 • ggF-CBA-enriched categories: The remaining events that are not classified into
2437 VBF-enriched category. Then events are categorized into three channels based on
2438 lepton-flavor, namely ggF_2e2μ, ggF_4e and ggF_4μ.

2439

2. MVA-based categorization

2440 In order to target different production modes, two types of classifiers, one dedicate to
2441 VBF production while the other one for ggF, have been trained using deep neural network
2442 technique. Details of two classifiers are described as below:

2443

DNN models

2444 Figure 6.1 shows the architecture of VBF (left) and ggF (right) network. The VBF
2445 network includes three parts: two recurrent neural networks (RNNs) and one multilayer
2446 perceptron (MLP). The ggF network consists of one RNN and one MLP.

Table 6.3 Summary of the object and event selection requirements.

Physics Objects	
ELECTRONS	
Loose Likelihood quality electrons with hit in innermost layer, $E_T > 7$ GeV and $ \eta < 2.47$	
Interaction point constraint: $ z_0 \cdot \sin \theta < 0.5$ mm (if ID track is available)	
MUONS	
Loose identification with $p_T > 5$ GeV and $ \eta < 2.7$	
Calo-tagged muons with $p_T > 15$ GeV and $ \eta < 0.1$, segment-tagged muons with $ \eta < 0.1$	
Stand-alone and silicon-associated forward restricted to the $2.5 < \eta < 2.7$ region	
Combined, stand-alone (with ID hits if available) and segment-tagged muons with $p_T > 5$ GeV	
Interaction point constraint: $ d_0 < 1$ mm and $ z_0 \cdot \sin \theta < 0.5$ mm (if ID track is available)	
JETS	
anti- k_T jets with <i>bad-lose</i> identification, $p_T > 30$ GeV and $ \eta < 4.5$	
OVERLAP REMOVAL	
Jets within $\Delta R < 0.2$ of an electron or $\Delta R < 0.1$ of a muon are removed	
VERTEX	
At least one collision vertex with at least two associated track	
PRIMARY VERTEX	
Vertex with the largest p_T^2 sum	
Event Selection	
QUADRUPLET	- Require at least one quadruplet of leptons consisting of two pairs of same-flavour opposite-charge leptons fulfilling the following requirements:
SELECTION	<ul style="list-style-type: none"> - p_T thresholds for three leading leptons in the quadruplet: 20, 15 and 10 GeV - Maximum one calo-tagged or stand-alone muon or silicon-associated forward per quadruplet - Leading di-lepton mass requirement: $50 < m_{12} < 106$ GeV - Sub-leading di-lepton mass requirement: $50 < m_{34} < 115$ GeV - $\Delta R(\ell, \ell') > 0.10$ for all leptons in the quadruplet - Remove quadruplet if alternative same-flavour opposite-charge di-lepton gives $m_{\ell\ell} < 5$ GeV - Keep all quadruplets passing the above selection
ISOLATION	<ul style="list-style-type: none"> - Contribution from the other leptons of the quadruplet is subtracted - FixedCutPFlowLoose WP for all leptons
IMPACT	- Apply impact parameter significance cut to all leptons of the quadruplet
PARAMETER	- For electrons: $d_0/\sigma_{d_0} < 5$
SIGNIFICANCE	- For muons: $d_0/\sigma_{d_0} < 3$
BEST	- If more than one quadruplet has been selected, choose the quadruplet
QUADRUPLET	with highest Higgs decay ME according to channel: 4μ , $2e2\mu$, $2\mu2e$ and $4e$
VERTEX	- Require a common vertex for the leptons:
SELECTION	- $\chi^2/\text{ndof} < 5$ for 4μ and < 9 for others decay channels

Table 6.4 Cutflow table for a narrow-width ggF signal sample at $m_H = 600$ GeV. N_{event} denotes the number of events selected after each cut is applied, normalized to 139 fb^{-1} , according to the expected upper limit on the cross section. The acceptances (the proportion of events selected relative to the initial number of events) are also included.

	N_{event}	$N_{\text{event}}/\text{BR}(ZZ \rightarrow \ell\ell\ell'\ell')$	Acc. [%]	Acc. · $\text{BR}(ZZ \rightarrow \ell\ell\ell'\ell') \cdot 1000$
Initial	17.902	3 964.3	100.00	4.516
Lepton selection	6.247	1 383.4	34.90	1.576
SFOS	5.758	1 275.1	32.16	1.453
Kinematic cuts	5.754	1 274.2	32.14	1.452
Z_1 Mass	5.726	1 267.9	31.98	1.444
Z_2 Mass	5.112	1 132.0	28.56	1.290
J/ψ Veto	5.111	1 131.9	28.55	1.289
ΔR	5.111	1 131.7	28.55	1.289
Isolation	4.864	1 077.0	27.17	1.227
Impact parameters	4.796	1 062.1	26.79	1.210
Vertex requirement	4.786	1 059.8	26.73	1.207
Trigger	4.783	1 059.1	26.72	1.207
“Badjet” veto	4.763	1 054.7	26.61	1.201

Table 6.5 Cutflow table for a narrow-width ggF signal sample at $m_H = 1000$ GeV. N_{event} denotes the number of events selected after each cut is applied, normalized to 139 fb^{-1} , according to the expected upper limit on the cross section. The acceptances (the proportion of events selected relative to the initial number of events) are also included.

	N_{event}	$N_{\text{event}}/\text{BR}(ZZ \rightarrow \ell\ell\ell'\ell')$	Acc. [%]	Acc. · $\text{BR}(ZZ \rightarrow \ell\ell\ell'\ell') \cdot 1000$
Initial	5.603	1 240.8	100.00	4.516
Lepton selection	2.141	474.1	38.21	1.725
SFOS	1.944	430.5	34.70	1.567
Kinematic cuts	1.943	430.3	34.68	1.566
Z_1 Mass	1.932	427.8	34.48	1.557
Z_2 Mass	1.715	379.7	30.61	1.382
J/ψ Veto	1.715	379.7	30.60	1.382
ΔR	1.714	379.6	30.60	1.382
Isolation	1.640	363.2	29.27	1.322
Impact parameters	1.620	358.6	28.90	1.305
Vertex requirement	1.616	357.8	28.84	1.302
Trigger	1.615	357.7	28.83	1.302
“Badjet” veto	1.609	356.2	28.71	1.297

Table 6.6 Cutflow table for a narrow-width VBF signal sample at $m_H = 600$ GeV. N_{event} denotes the number of events selected after each cut is applied, normalized to 139 fb^{-1} , according to the expected upper limit on the cross section. The acceptances (the proportion of events selected relative to the initial number of events) are also included.

	N_{event}	$N_{\text{event}}/\text{BR}(ZZ \rightarrow \ell\ell\ell'\ell')$	Acc. [%]	Acc. · $\text{BR}(ZZ \rightarrow \ell\ell\ell'\ell') \cdot 1000$
Initial	12.143	2 688.9	100.00	4.516
Lepton selection	4.307	953.7	35.47	1.602
SFOS	3.975	880.2	32.74	1.478
Kinematic cuts	3.972	879.6	32.71	1.477
Z_1 Mass	3.953	875.4	32.56	1.470
Z_2 Mass	3.545	785.0	29.19	1.318
J/ψ Veto	3.545	785.0	29.19	1.318
ΔR	3.544	784.9	29.19	1.318
Isolation	3.418	756.9	28.15	1.271
Impact parameters	3.368	745.9	27.74	1.253
Vertex requirement	3.362	744.5	27.69	1.250
Trigger	3.360	744.0	27.67	1.250
“Badjet” veto	3.340	739.7	27.51	1.242

Table 6.7 Cutflow table for a narrow-width VBF signal sample at $m_H = 1000$ GeV. N_{event} denotes the number of events selected after each cut is applied, normalized to 139 fb^{-1} , according to the expected upper limit on the cross section. The acceptances (the proportion of events selected relative to the initial number of events) are also included.

	N_{event}	$N_{\text{event}}/\text{BR}(ZZ \rightarrow \ell\ell\ell'\ell')$	Acc. [%]	Acc. · $\text{BR}(ZZ \rightarrow \ell\ell\ell'\ell') \cdot 1000$
Initial	3.827	847.4	100.00	4.516
Lepton selection	1.474	326.5	38.53	1.740
SFOS	1.351	299.1	35.30	1.594
Kinematic cuts	1.350	299.0	35.28	1.593
Z_1 Mass	1.341	297.0	35.04	1.583
Z_2 Mass	1.195	264.6	31.23	1.410
J/ψ Veto	1.195	264.6	31.23	1.410
ΔR	1.195	264.6	31.22	1.410
Isolation	1.161	257.1	30.34	1.370
Impact parameters	1.148	254.1	29.99	1.354
Vertex requirement	1.146	253.8	29.95	1.352
Trigger	1.145	253.6	29.93	1.352
“Badjet” veto	1.139	252.2	29.77	1.344

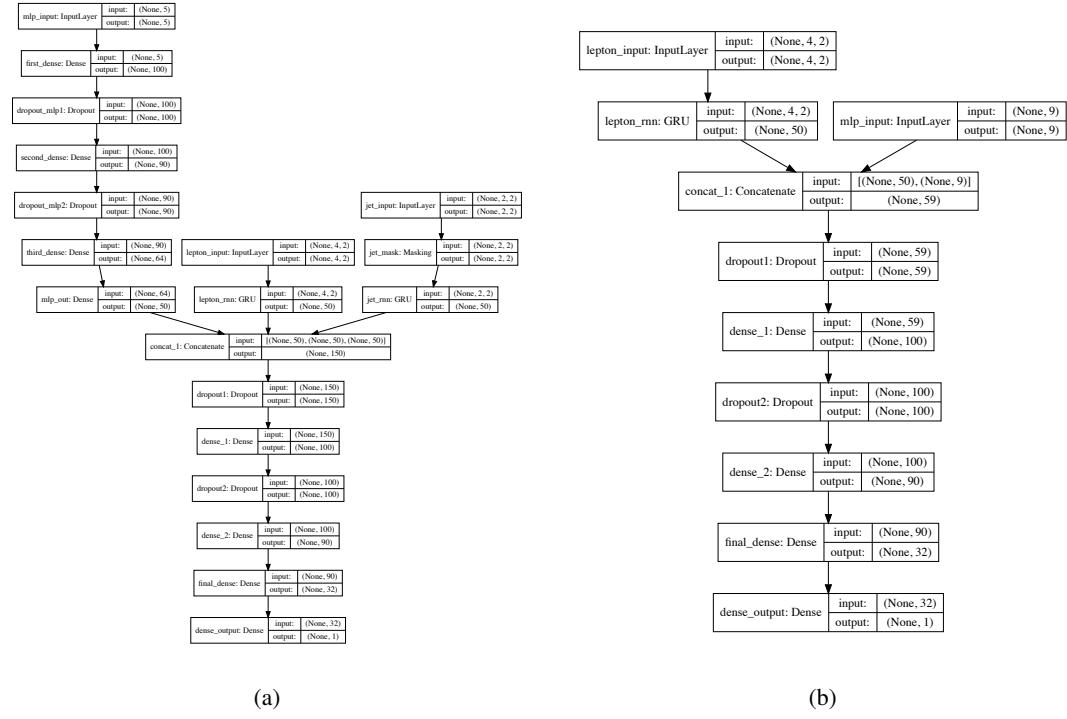


Fig. 6.1 (a) VBF DNN architecture diagram. (b) ggF DNN architecture.

For training, the VBF and ggF signal samples at the masses of 200, 300, 400, 500, 600, 700, 800, 900, 1000, 1200, 1400 GeV are used with positive label. The VBF (ggF) signals are only used for VBF (ggF) classifier. The background including simulated samples of QCD and EW $q\bar{q} \rightarrow ZZ$ processes as well as $gg \rightarrow ZZ$ process summed according to their cross section are signed with negative labels. In addition to the selections described in section 6.3.2, the events used for VBF network are required to have $N_{\text{jets}} \geq 2$, while $N_{\text{jets}} < 2$ is required for events in ggF network.

In order to assign equivalent importance to signals with different mass assumptions, during the training, signal events are reweighted to follow the $m_{4\ell}$ distribution from background, as shown in figure 6.2 (figure 6.3) before (left) and after(right) reweighting for VBF (ggF) samples.

Input features

Table 6.8 (table 6.9) lists the input features used for VBF (ggF) network during the training. For VBF network, one RNN (the other one) takes the p_T and η of p_T -ordered four leptons (two leading jets) as input features, which intends to study the time relationship from particle decay between leptons (jets). For ggF network, the only one RNN model takes the p_T and η of p_T -ordered four leptons as inputs.

Evaluation of models

Figure 6.4 shows the output of “ggF-classifier” and “VBF-classifier” for data, SM

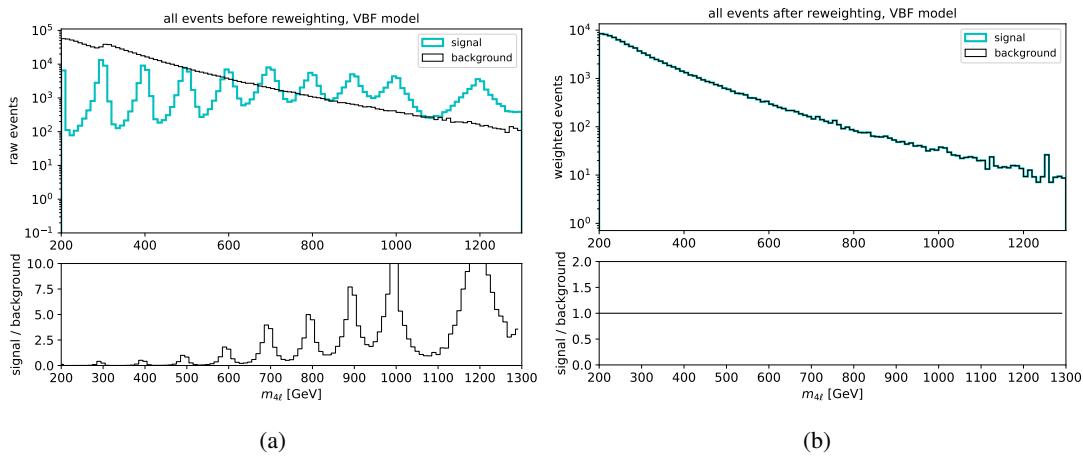


Fig. 6.2 (a) $m_{4\ell}$ distribution of raw (unweighted) training events for VBF signal (blue) and background (black); (b) $m_{4\ell}$ distribution of weighted VBF signal (blue) and background (black) used at training time.

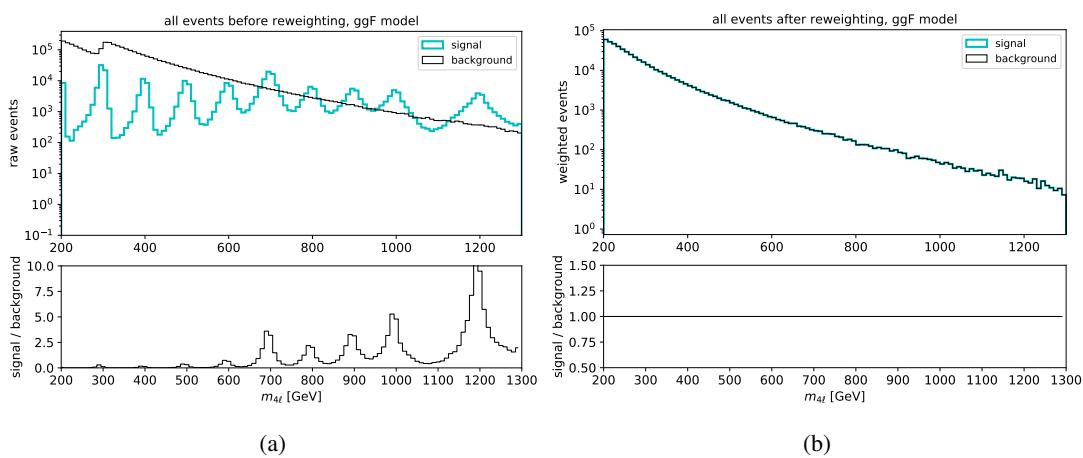


Fig. 6.3 (a) $m_{4\ell}$ distribution of raw (unweighted) training events for ggF signal (blue) and background (black); (b) $m_{4\ell}$ distribution of weighted ggF signal (blue) and background (black) used at training time.

Table 6.8 Input features used in the “VBF-classifier” for the $\ell\ell\ell'\ell'$ analysis. The RNN stands for the recurrent neural network and MLP for the multilayer perceptron.

Model	Inputs	Description
RNN	$p_T^{j0,j1}$	transverse momenta of the two leading jets
	$\eta^{j0,j1}$	pseudorapidity of the two leading jets
	$p_T^{\ell0,\ell1,\ell2,\ell3}$	transverse momenta of the four leptons
	$\eta^{\ell0,\ell1,\ell2,\ell3}$	pseudorapidity of the four leptons
MLP	$m_{4\ell}$	invariant mass of the four lepton system
	m_{jj}	invariant mass of the two leading jet system
	p_T^{jj}	transverse momentum of the two leading jet system
	$\Delta\eta_{H,j}$	difference in pseudorapidity between the four lepton system and the leading jet
	$\min\Delta R_{jZ}$	minimum distance between one of the two lepton pairs and a jet

backgrounds and an example signal at 600 GeV. The ggF and VBF signals cross section are set to be one hundred times of their observed upper limit described in section 6.7.3 for ggF output and fifty times of the observed upper limit for VBF output for best visibility.

Then the optimal cut at output score from each classifier is chosen based on an overall good performance of classifier to have a large significance improvement while retaining a high signal efficiency. Figure 6.5 shows the significance improvements of MVA-based cuts when comparing with cut-based one at different VBF (left) and ggF (right) mass samples, where the significance is calculated as:

$$Z = \sqrt{2 \left(n \ln \left[\frac{nb + \sigma^2}{b^2 + n\sigma^2} \right] - \frac{b^2}{\sigma^2} \ln \left[1 + \frac{\sigma^2(n-b)}{b(b+\sigma^2)} \right] \right)} \quad (6.1)$$

Cut at 0.5 (0.8) for VBF (ggF) classifier is chosen as shown in solid lines.

Then the events passing VBF classifier are categorized into VBF-MVA-enriched category. Otherwise, the events failing VBF classifier but passing ggF classifier are categorized into ggF-MVA-high category, which is further split into 3 channels. All remaining events are sorted into one additional ggF-MVA-low category. Thus there are five categories defined in MVA-based categorization. In summary, cuts applied in categorization are defined as follow, and these different phase spaces are also illustrated in figure 6.6.

- VBF-MVA-enriched category: Events have at least two selected jets ($N_{\text{jets}} \geq 2$), and with $DNN_{\text{VBF}} > 0.8$;

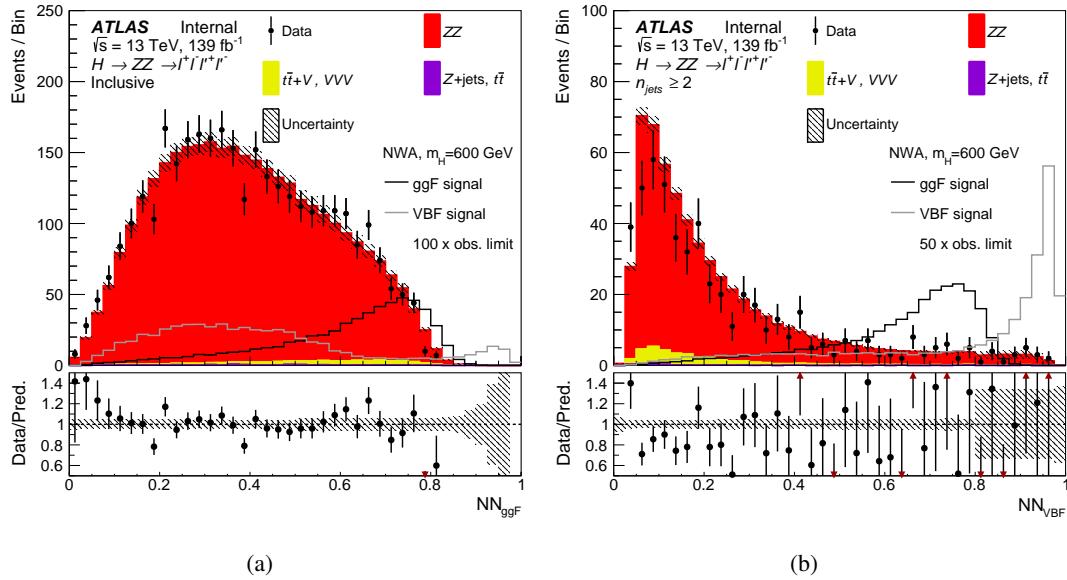


Fig. 6.4 The output score of “ggF-classifier” (a) and “VBF-classifier” (b) with the events passing the common event selections for the data, the SM backgrounds and an example of a NWA signal with a mass of 600 GeV. For the “VBF-classifier”, an additional requirement of at least two jets in the event is applied. The signals cross section are set to one hundred times of the observed limit for the “ggF-classifier” and fifty times of the observed limit for the “VBF -classifier”. The ZZ backgrounds are scaled by the normalisation factors shown in Table 6.15. The lower panels show the ratio of data to prediction. Only statistical and experimental systematic uncertainties are included.

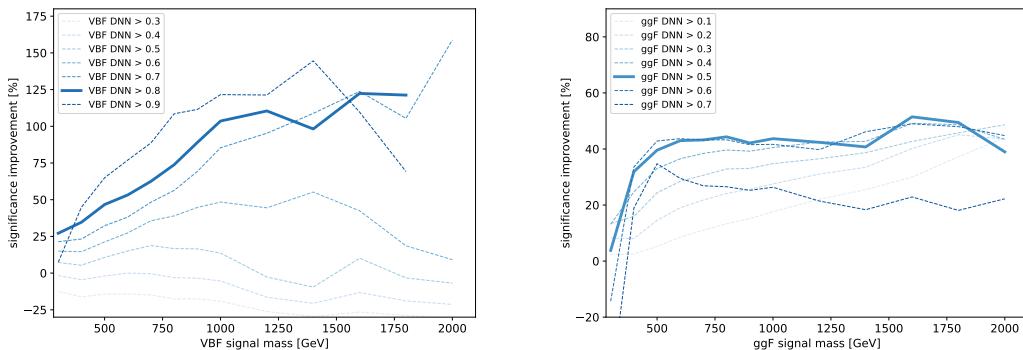


Fig. 6.5 Significance improvements of the MVA-based over the cut-based categorization of the VBF (ggF) category for VBF (ggF) signal samples from 300 to 2000 GeV for seven different cuts on the VBF (ggF) output score. The optimal cut of 0.8 (0.5) for VBF (ggF) score is chosen as the solid line, while other alternative cuts are plotted with dashed lines. For VBF category, results at 2000 GeV for cuts of 0.8 and 0.9 are missing due to a lack of background events passing this tight selection.

Table 6.9 Input features used in the “ggF-classifier” for the $\ell\ell\ell'\ell'$ analysis. The RNN stands for the recurrent neural network and MLP for the multilayer perceptron.

Model	Inputs	Description
RNN	$p_T^{\ell0,\ell1,\ell2,\ell3}$	transverse momenta of the four leptons
	$\eta^{\ell0,\ell1,\ell2,\ell3}$	pseudorapidity of the four leptons
MLP	$m_{4\ell}$	invariant mass of the four lepton system
	$p_T^{4\ell}$	transverse momentum of the four lepton system
	$\eta^{4\ell}$	pseudorapidity of the four lepton system
	$\cos \theta^*$	production angle of the leading Z defined in the four lepton rest frame
	$\cos \theta_1$	angle between the negative final state lepton and the direction of flight of leading Z in the Z rest frame
	$\cos \theta_2$	angle between the negative final state lepton and the direction of flight of sub-leading Z in the Z rest frame
	Φ	angle between the decay planes of the four final state leptons expressed in the four lepton rest frame
	p_T^{j0}	transverse momentum of the leading jet
	η^{j0}	pseudorapidity of the leading jet

- 2483 • ggF-MVA-high categories: ($N_{\text{jets}} \geq 2 \ \&\& \ DNN_{\text{VBF}} \leq 0.8 \ \&\& \ DNN_{\text{ggF}} >$
 2484 $0.5) || (N_{\text{jets}} < 2 \ \&\& \ DNN_{\text{ggF}} > 0.5);$
 2485 • ggF-MVA-low category: All remaining events that fail VBF and ggF cuts men-
 2486 tioned above.

2487 6.3.4 Signal acceptance

2488 The signal acceptance is defined as the ratio of events passing all analysis selection in
 2489 each category to the total number of simulated events in whole phase space. In denom-
 2490 inator, the events with τ final states are not taken into account. And the contribution of
 2491 τ -lepton decay to electrons and muons final states is found to be negligible.

2492 Figure 6.7 and 6.8 show the acceptance of NWA signals in DNN- and Cut- based
 2493 categorization, estimated by merging the three signal MC campaigns, mc16a, mc16d and
 2494 mc16e. A 3-rd order polynomial fit is applied for each category.

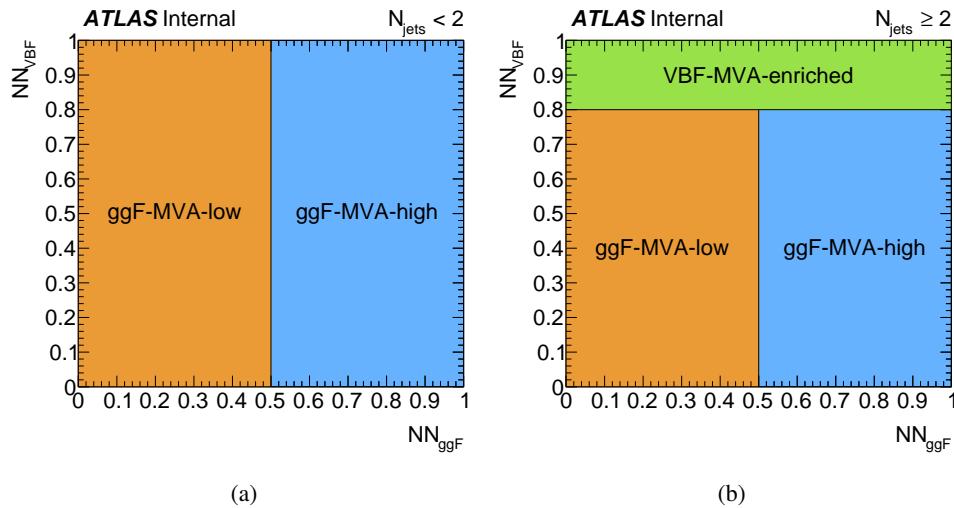


Fig. 6.6 Illustration of the MVA-based VBF and ggF event classification for events with (a) $N_{\text{jets}} < 2$ and (b) $N_{\text{jets}} \geq 2$.

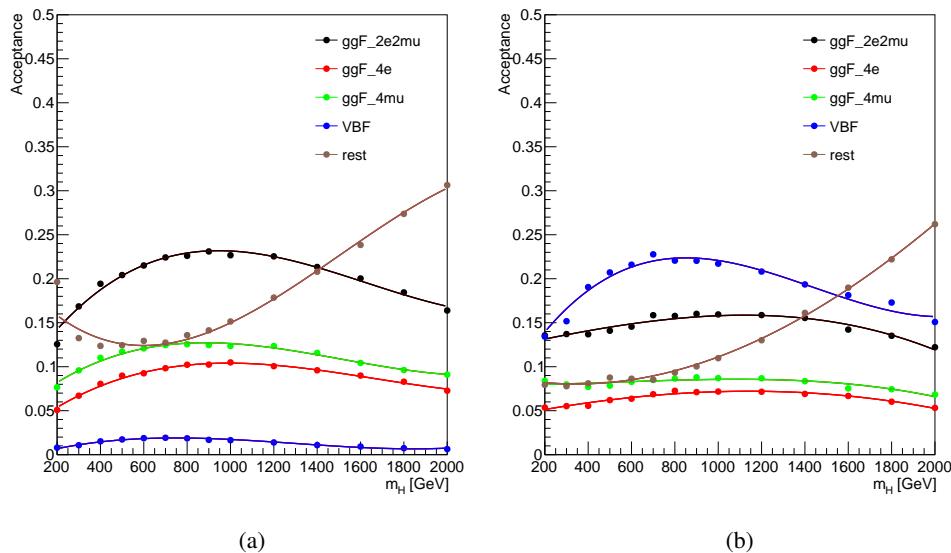


Fig. 6.7 NWA acceptance as a function of m_H for the MVA-based categorization for the samples of (a) ggF production; (b) VBF production.

2495 6.4 Background estimation

2496 In this analysis, 97% of total expected background events are from irreducible ZZ
 2497 backgrounds, which includes about 86% quark-antiquark annihilation ($q\bar{q} \rightarrow ZZ$), 10%
 2498 of gluon-induced production ($gg \rightarrow ZZ$) and around 1% of EW vector boson scattering
 2499 ($q\bar{q} \rightarrow ZZ$ EW) contribution. For $q\bar{q} \rightarrow ZZ$ EW, although it has small contribution
 2500 in total background events after analysis selection, it's important for VBF category with
 2501 about 16% contribution.

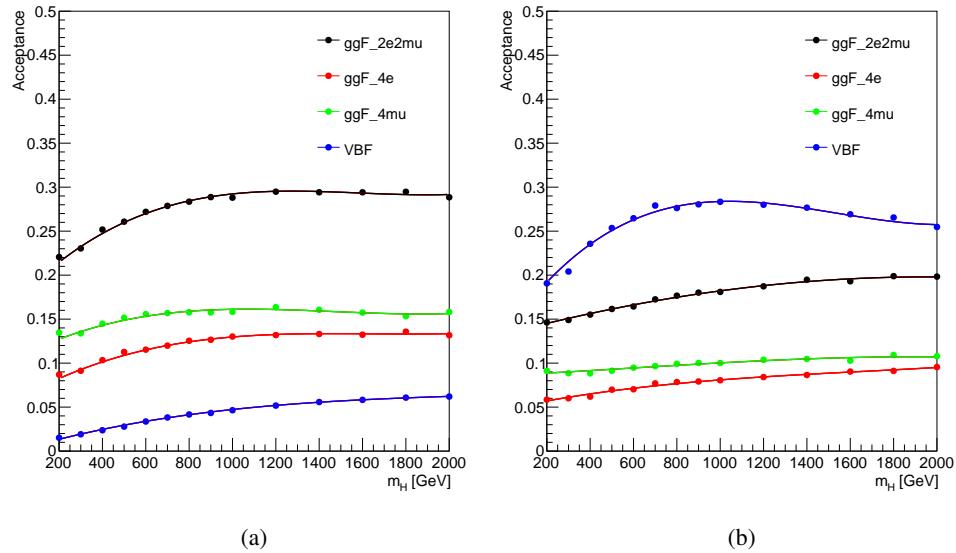


Fig. 6.8 NWA acceptance as a function of m_H for the Cut-based categorization for the samples of (a) ggF production mode; (b) VBF production mode.

2502 In addition to irreducible backgrounds, events from $Z+jets$ and $t\bar{t}$ processes, represent
 2503 as reducible backgrounds, contribute at a few percent level and can be measured using data
 2504 driven method that will be described briefly later. Additional background called ‘Others’,
 2505 including ttV and triple-V (VVV) processes, has tiny contribution and is estimated from
 2506 MC simulation directly.

2507 6.4.1 Irreducible backgrounds

2508 The Irreducible backgrounds have events with four prompt leptons. The normalization
 2509 of two dominant backgrounds $q\bar{q} \rightarrow ZZ$ and $gg \rightarrow ZZ$ are taken from data by statistical
 2510 fit, and the normalization of small $q\bar{q} \rightarrow ZZ$ EW background is measured directly from
 2511 MC simulation.

2512 The $m_{4\ell}$ shapes of all three background components are taken from MC samples and
 2513 then parameterized by an empirical function for each of them in each category respectively.
 2514 Details of background modellings are illustrated as below:

2515 The empirical function used for background parameterization is:

$$f(m_{4\ell}) = C_0 H(m_0 - m_{4\ell}) f_1(m_{4\ell}) + H(m_{4\ell} - m_0) f_2(m_{4\ell}), \quad (6.2)$$

where,

$$f_1(x) = \left(\frac{x - a_4}{a_3} \right)^{a_1-1} \left(1 + \frac{x - a_4}{a_3} \right)^{-a_1-a_2},$$

$$f_2(x) = \exp \left[b_0 \left(\frac{x - b_4}{b_3} \right)^{b_1-1} \left(1 + \frac{x - b_4}{b_3} \right)^{-b_1-b_2} \right],$$

$$C_0 = \frac{f_2(m_0)}{f_1(m_0)}.$$

The function consists of two parts, the first part f_1 describes the $m_{4\ell}$ spectrum in low mass region where both Z bosons decay on-shell, while the second one f_2 covers distribution at high mass tail. The transition between the low- and high- mass parts is presented in function 6.2 by the Heaviside step function $H(x)$ at the transition point m_0 . The m_0 is chosen to optimize the smoothness of the function, and practically $m_0 = 260$ (350) GeV is used for $q\bar{q} \rightarrow ZZ$ ($gg \rightarrow ZZ$ and $q\bar{q} \rightarrow ZZ$ EW). Besides, the continuity of two functions at m_0 is ensured by the factor C_0 applied to f_1 . The coefficients a_i in f_1 and b_i in f_2 are shape parameters obtained by fitting to $m_{4\ell}$ distribution from each MC simulated sample.

Figure 6.9 to 6.11 shows the fitting results of $q\bar{q} \rightarrow ZZ$, $gg \rightarrow ZZ$, $q\bar{q} \rightarrow ZZ$ EW backgrounds in four cut-based categories (ggF-CBA-enriched- $2e2\mu$, ggF-CBA-enriched- $4e$, ggF-CBA-enriched- 4μ and VBF-CBA-enriched). Figure 6.12 to 6.14 shows the fitting results of those backgrounds in five MVA-based categories (ggF-MVA-high- $2e2\mu$, ggF-MVA-high- $4e$, ggF-MVA-high- 4μ , ggF-MVA-low and VBF-MVA-enriched).

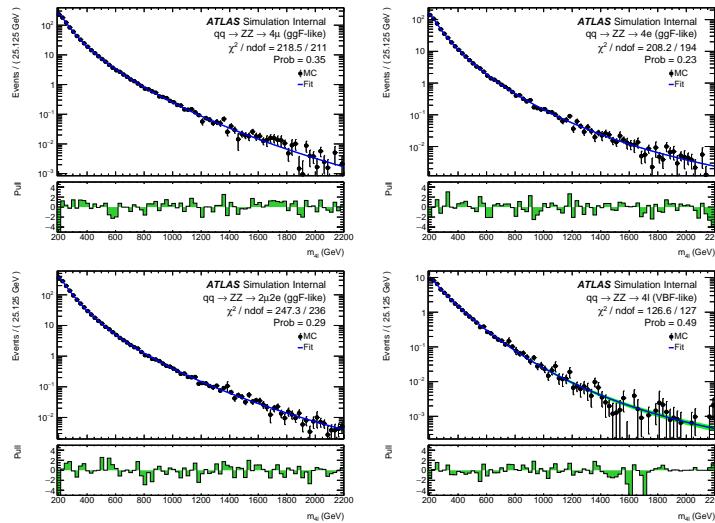


Fig. 6.9 Distributions of the $m_{4\ell}$ invariant mass fit projections of the $q\bar{q} \rightarrow ZZ$ background samples for the 4μ , $4e$ and $2\mu2e$ final states in the ggF-CBA-enriched category, and the 4ℓ inclusive VBF-CBA-enriched category. Cut-based categorization is used.

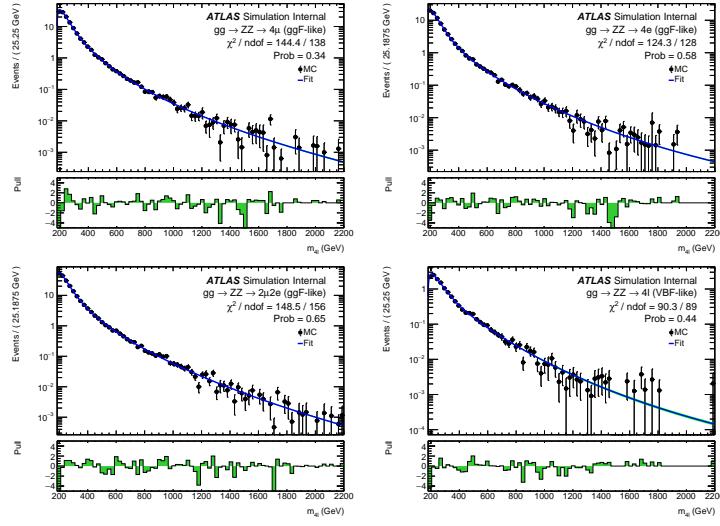


Fig. 6.10 Distributions of the $m_{4\ell}$ invariant mass fit projections of the $gg \rightarrow ZZ$ background samples for the 4μ , $4e$ and $2\mu2e$ final states in the ggF-CBA-enriched category, and the 4ℓ inclusive VBF-CBA-enriched category. Cut-based categorization is used.

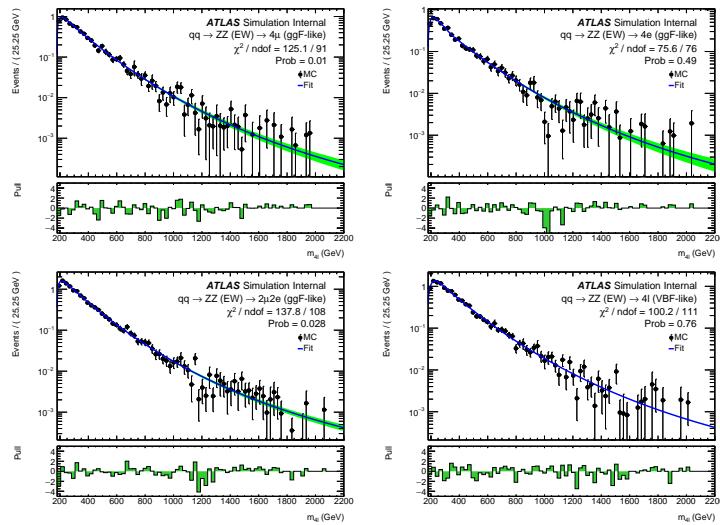


Fig. 6.11 Distributions of the $m_{4\ell}$ invariant mass fit projections of the $q\bar{q} \rightarrow ZZ$ (EW) background samples for the 4μ , $4e$ and $2\mu2e$ final states in the ggF-CBA-enriched category, and the 4ℓ inclusive VBF-CBA-enriched category. Cut-based categorization is used.

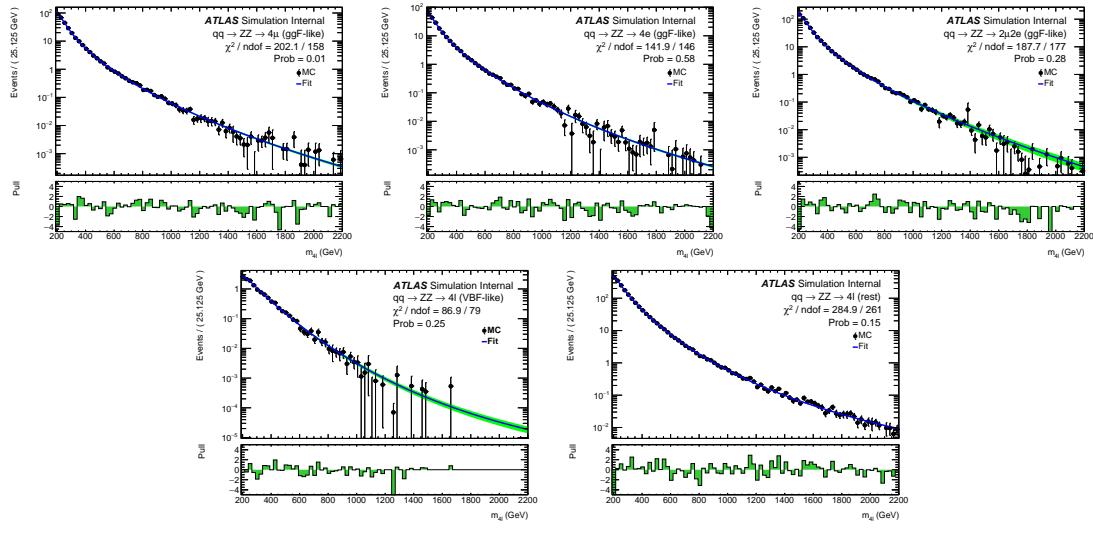


Fig. 6.12 Distributions of the $m_{4\ell}$ invariant mass fit projections of the $q\bar{q} \rightarrow ZZ$ background samples for the 4μ , $4e$ and $2\mu 2e$ final states in the ggF-MVA-high category, the 4ℓ inclusive ggF-MVA-low category and VBF-MVA-enriched category. DNN-based categorization is used.

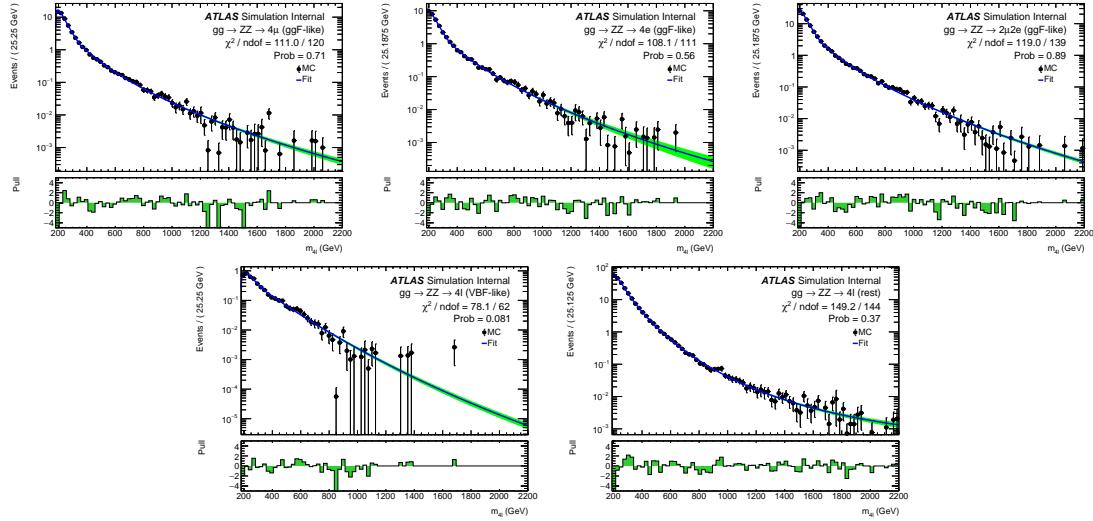


Fig. 6.13 Distributions of the $m_{4\ell}$ invariant mass fit projections of the $gg \rightarrow ZZ$ background samples for the 4μ , $4e$ and $2\mu 2e$ final states in the ggF-MVA-high category, the 4ℓ inclusive ggF-MVA-low category and VBF-MVA-enriched category. DNN-based categorization is used.

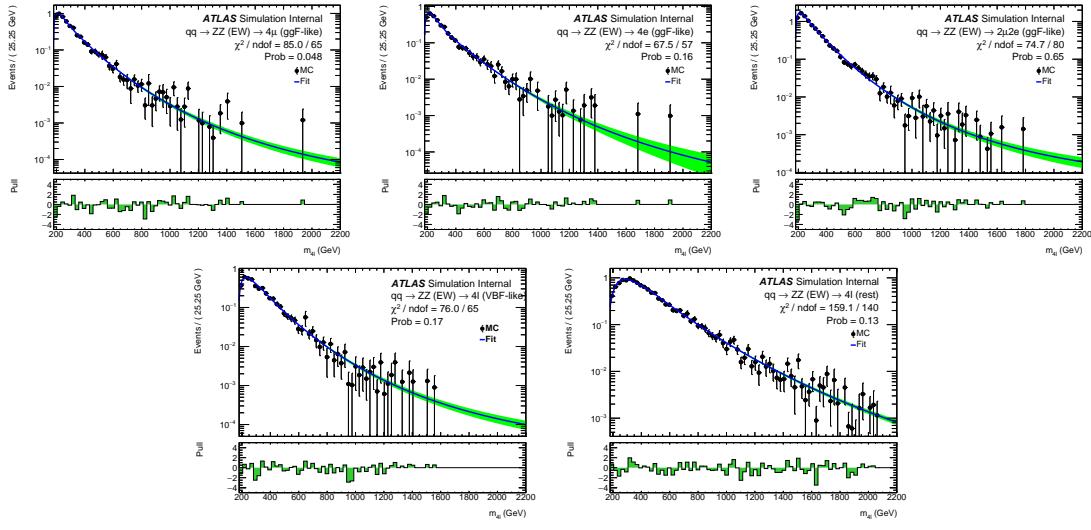


Fig. 6.14 Distributions of the $m_{4\ell}$ invariant mass fit projections of the $q\bar{q} \rightarrow ZZ$ (EW) background samples for the 4μ , $4e$ and $2\mu2e$ final states in the ggF-MVA-high category, the 4ℓ inclusive ggF-MVA-low category and VBF-MVA-enriched category. DNN-based categorization is used.

6.4.2 Reducible backgrounds

Similar as section 5.4, the reducible backgrounds include $Z+jets$ (consists of both heavy- and light-flavour jets), top quark pair, and WZ production, which contain fake and non-isolated leptons. The simulations are not very robust in terms of the selection efficiencies. Thus, the data-driven method is applied to estimate the normalization of those processes in different control regions (CRs). The estimations in this analysis are performed separately for $\ell\ell + \mu\mu$ and $\ell\ell + ee$ final states, with slightly different approaches for “muon” and “electron” backgrounds.

The “electron” backgrounds mostly come from process of a Z boson with light-flavour jets ($Z+LF$) misidentified as electrons. The large contribution of “muon” backgrounds come from heavy-flavour jets produced in association with a Z boson ($Z+HF$) or in the decays of top quark. The estimations are done following the common H4l studies without a specific $m_{4\ell}$ range requirement^[104], and then the corresponding fraction of event yield in $m_{4\ell} > 200$ GeV is calculated from MC simulation.

$\ell\ell + \mu\mu$ final states

The normalizations of “muon” backgrounds are extracted from simultaneous fits of the leading lepton pair’s invariant mass (m_{12}) in four orthogonal CRs:

- **Inverted d_0 CR:** this CR is formed by inverting the d_0 selection for at least one lepton in subleading lepton pair while the leptons in leading pair are required to pass all standard selection. This CR enhances $Z+HF$ and $t\bar{t}$ as leptons from heavy-

flavour hadronic decays are characterised by large d_0 .

- $e\mu + \mu\mu$ **CR**: this CR is formed using an opposite-charge different-flavour dilepton in leading pair. It aims to enhance $t\bar{t}$ background as the leading lepton pair cannot come from Z boson decay.
- **Inverted isolation CR**: in this CR, leptons in leading pair are required to satisfy all standard analysis selection, while for leptons in subleading pair, they are required to pass d_0 selection but have at least one of them failing isolation selection. This CR enhances the events from $Z+LF$ processes while suppress $Z+HF$ by d_0 cut.
- **Same-sign CR**: in this CR, the leptons in subleading pair are required to have same-charge, while the leading pair still passes standard selection. This CR is not dominant by any specific background since all reducible backgrounds could have sizable contribution in it.

The fit results of normalizations are then propagated to signal region (SR) by applying transfer factors to account the difference of selection efficiencies between SR and CRs.

The transfer factors are computed using $Z + \mu$ MC samples.

$\ell\ell + ee$ final states

The “electron” backgrounds are estimated in $3\ell + X$ CR, where X denotes the lower p_T electron in the subleading pair. The selection and identification criterias for X are relaxed , while other three leptons must satisfy the standard selection. In this case, X could be a light-flavour jet, a photon conversion or an electron from heavy-flavour hadron decay. Moreover, the subleading pair is required to have same charge dilepton to ensure the orthogonality to the signal region. The normalization of backgrounds are obtained based on a fit to the number of hits in the innermost ID layer in CR, and the transfer factors are computed from $Z + e$ simulated sample.

The $m_{4\ell}$ shapes of reducible backgrounds are obtained from MC simulation in signal region, and then smoothed by an one-dimensional kernel estimation, which models the input data as a superposition of Gaussian kernels, one for each data point with contributing $1/N$ to total integral N ^[105]. The difference from using different smoothing strength (ρ) in kernel estimation is taken into account as additional shape uncertainties for these reducible backgrounds.

6.5 Signal modelling

The parameterization of $m_{4\ell}$ distributions based on simulated samples for signals are described in this section. Several signal models are studied, including heavy Higgs like

2583 narrow-width signal (NWA) and large-width signal (LWA), as well as the modelling of
2584 Randall-Sundrum graviton (RSG) signal.

2585 6.5.1 Modelling of narrow-width signal

2586 For narrow-width (NWA) signal, the $m_{4\ell}$ width is totally determined by detector reso-
2587 lution, which is modelled by the sum of a Crystal Ball (C) function^[106-107] and a Gaussian
2588 (G) function:

$$P_s(m_{4\ell}) = f_C \cdot C(m_{4\ell}; \mu, \sigma_C, \alpha_C, n_C) + (1 - f_C) \cdot G(m_{4\ell}; \mu, \sigma_G) \quad (6.3)$$

2589 The two functions share the same central value μ , while the resolution parameters, σ_C
2590 and σ_G , are different. In the Crystal Ball function, the parameters α_C and n_C model the
2591 shape of non-Gaussian tail, and the fraction parameter f_C is used to ensure the relative
2592 normalization between two functions.

2593 The parameters are obtained by fitting to signal MC simulations combining the mc16a,
2594 mc16d and mc16e campaigns for each category at each mass points from 200 GeV to
2595 2000 GeV respectively, and the shape of ggF and VBF signals are found to be similar.
2596 Figure 6.15 shows the $m_{4\ell}$ distribution and fitted curves for ggF production at mass from
2597 200 GeV to 2000 GeV in $2e2\mu$ channel as examples.

2598 Then the $C + G$ parameters are fitted with a polynomial function as the function of
2599 generated mass points (m_H), as an example shown in figure 6.16 for $2e2\mu$ channel. The
2600 fitting quality can be measured by the Pearson's χ^2 , which is within 3 (2) for $2e2\mu$ ($4e$ and
2601 4μ) channel.

2602 In addition, possible difference on the signal yield extracted from parameterization and
2603 MC simulation is studied. Figure 6.17 shows this difference by computing $\frac{N_{\text{reco}} - N_{\text{fit}}}{N_{\text{fit}}}$, where
2604 N_{reco} denotes the total number of reconstructed events observed from MC simulation at
2605 that mass point and N_{fit} depicts the number of events obtained from the fitted PDF. The
2606 differences are treated as an additional systematic uncertainty with the value of 2% (1%)
2607 for $2e2\mu$ ($4e$ and 4μ) channel in statistical fit.

2608 In summary, the final interpolated signal shapes for the ggF production mode are
2609 shown together in figure 6.18 for mass points with step of 100 GeV from 200 GeV to
2610 3000 GeV.

2611 6.5.2 Modelling of large-width signal

2612 The $m_{4\ell}$ shape of heavy Higgs model in large-width (LWA) hypothesis can be de-
2613 scribed by a convolution of a truth distribution and a resolution from detector effect. The

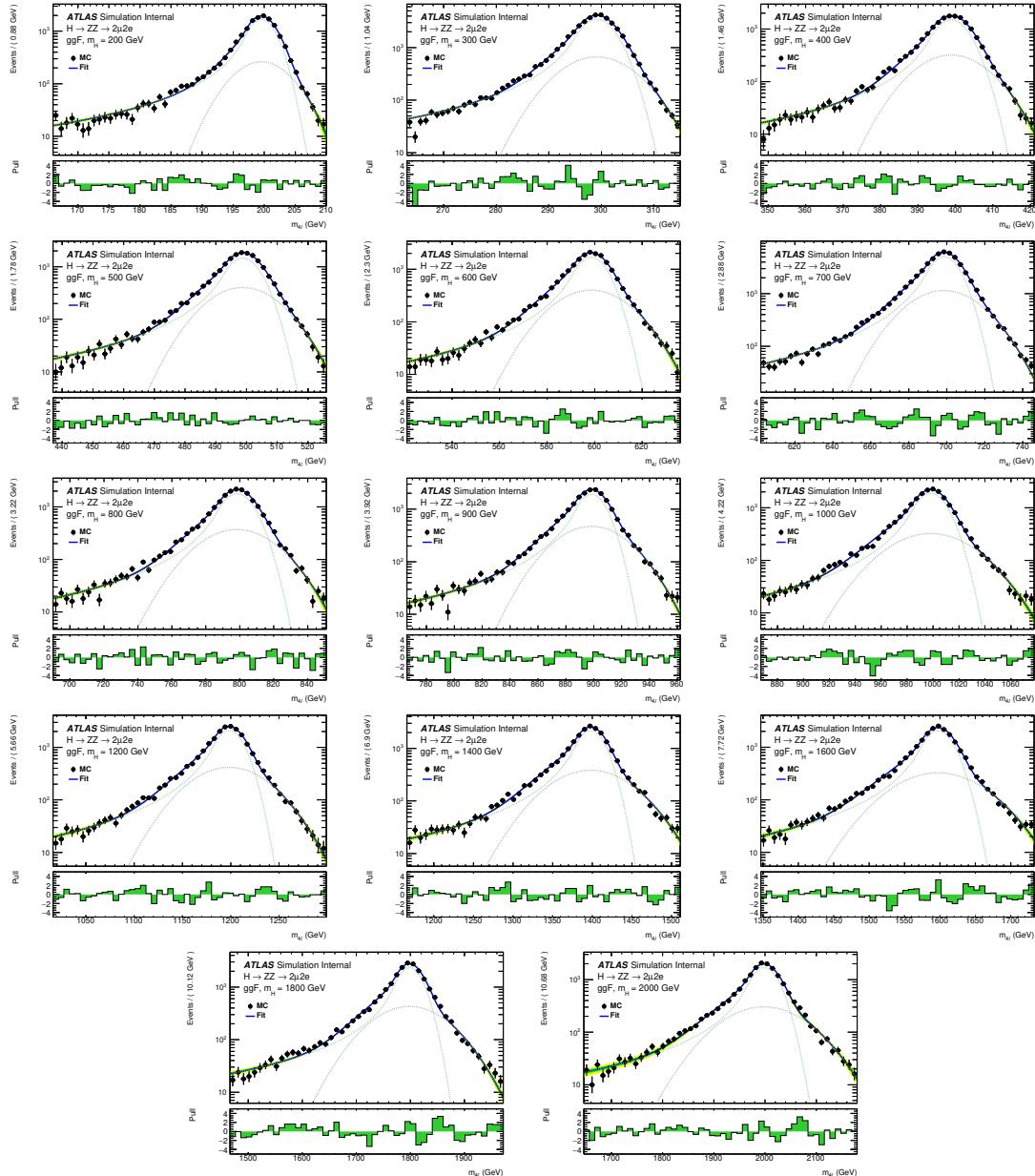


Fig. 6.15 Distributions of the $m_{2\mu 2e}$ and fit projection for signal samples between 200 to 3000 GeV for ggF production mode. Three MC campaigns, mc16a, mc16d and mc16e, are combined. The lower panel in each plot shows the pull distribution.

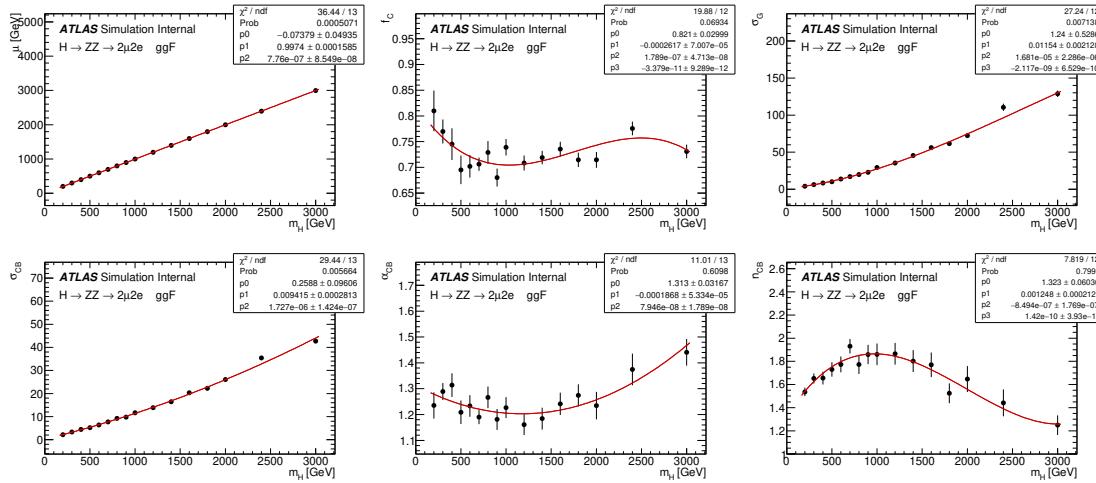


Fig. 6.16 Polynomial fits of the parameters μ , f_C , σ_G , σ_C , n_c and α_c for the signal $\mathcal{C} + \mathcal{G}$ model in the $2\mu 2e$ channel as a function of m_H for the ggF production mode. The combination of the mc16a, mc16d and mc16e MC campaigns is used.

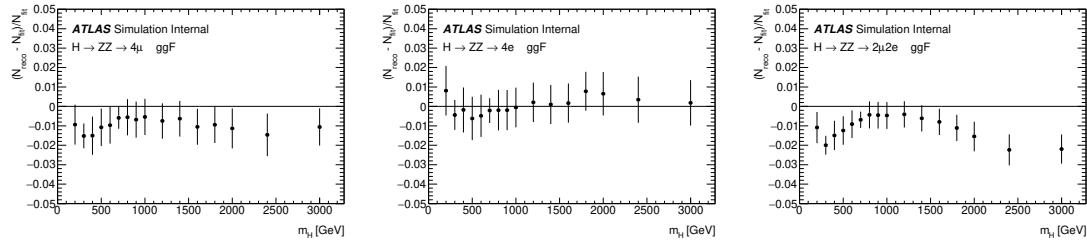


Fig. 6.17 The difference between MC simulation and parameterization of 4μ (left), $4e$ (middle) and $2\mu 2e$ (right) for the ggF production mode. The combination of the mc16a, mc16d and mc16e MC campaigns is used.

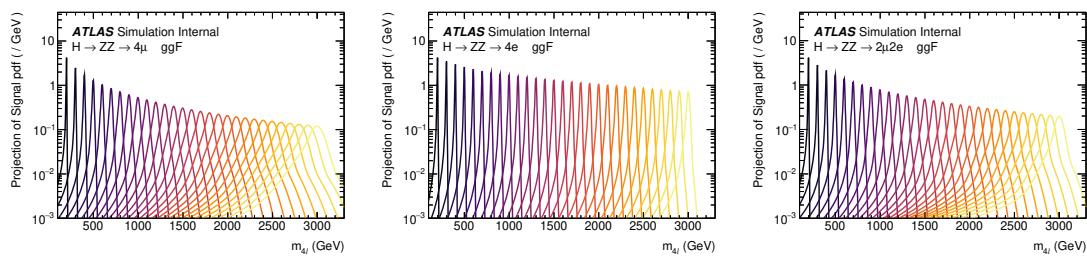


Fig. 6.18 The final signal shapes for the ggF production mode, interpolated from the polynomial fit parameters.

detector resolution effect is the one modelled by the function described in NWA parameterization, as in NWA model the truth level width is negligible.

The differential parton cross section for the heavy Higgs model can be written as^[108]:

$$\sigma_{gg \rightarrow H \rightarrow ZZ}(s) = \frac{1}{2s} \int d\Omega |A_{gg \rightarrow H}(s, \Omega)|^2 \frac{1}{|s - s_H|^2} |A_{H \rightarrow ZZ}(s, \Omega)|^2 \quad (6.4)$$

where $A_{gg \rightarrow H}(s, \Omega)$ and $A_{H \rightarrow ZZ}(s, \Omega)$ are corresponding Higgs production and decay amplitudes, and $\frac{1}{|s - s_H|}$ denotes the Higgs propagator and Ω represents the phase space of the process.

Using the definition of a partial width,

$$\Gamma_{H \rightarrow F}(s) = \frac{1}{2\sqrt{s}} \int d\Omega |A_{H \rightarrow F}(s, \Omega)|^2 \quad (6.5)$$

the parton cross section can be rewritten as,

$$\sigma_{gg \rightarrow H \rightarrow ZZ}(s) = 2 \frac{1}{|s - s_H|^2} \times \Gamma_{H \rightarrow gg}(s) \times \Gamma_{H \rightarrow ZZ}(s) \quad (6.6)$$

with the components computed in Ref^[108-109]:

$$\begin{aligned} \frac{1}{s - s_H} &= \frac{1 + i \cdot \bar{\Gamma}_H / \bar{m}_H}{s - \bar{m}_H^2 + i \cdot s \cdot \bar{\Gamma}_H / \bar{m}_H} \\ \bar{m}_H &= \sqrt{\bar{\Gamma}_H^2 + m_H^2} \\ \bar{\Gamma}_H &= \bar{m}_H \cdot \frac{\Gamma_H}{m_H} \end{aligned} \quad (6.7)$$

$$\Gamma_{H \rightarrow ZZ}(s) = C \cdot s^{\frac{3}{2}} \cdot \left[1 - \frac{4m_Z^2}{s} + \frac{3}{4} \left(\frac{4m_Z^2}{s} \right)^2 \right] \cdot \left[1 - \frac{4m_Z^2}{s} \right]^{\frac{1}{2}} \quad (6.8)$$

$$\begin{aligned} \Gamma_{H \rightarrow gg}(s) &= C \cdot s^{\frac{3}{2}} \cdot |A_t(\tau_t)|^2 \\ A_t(\tau) &= 2 \frac{\tau + (\tau - 1)f(\tau)}{\tau^2} \\ \tau_t &= \frac{s}{4m_t^2} \\ f(\tau) &= \begin{cases} \arcsin^2(\sqrt{\tau}), & \tau \leq 1 \\ -\frac{1}{4} \left[\log \frac{1+\sqrt{1-\tau^{-1}}}{1-\sqrt{1-\tau^{-1}}} - i\pi \right]^2, & \tau > 1 \end{cases} \end{aligned} \quad (6.9)$$

where m_f stands for the mass of a fermion f , and Γ_H denotes an assumed total width of the heavy Higgs boson.

At the LHC, the $m_{4\ell}$ line shape can be defined by a hadron cross section that is derived from equation 6.6 by multiplication with gluon-gluon luminosity \mathcal{L}_{gg} described in^[110]. Meanwhile, the cross section is rewritten as a function of $m_{4\ell}$ instead of s , which will give an extra power of mass dependence in the formula:

$$\sigma_{pp \rightarrow H \rightarrow ZZ}(m_{4\ell}) = 2 \cdot m_{4\ell} \cdot \mathcal{L}_{gg} \cdot \frac{1}{|s - s_H|^2} \cdot \Gamma_{H \rightarrow gg}(m_{4\ell}^2) \cdot \Gamma_{H \rightarrow ZZ}(m_{4\ell}^2) \quad (6.10)$$

The analytical shapes of truth level $m_{4\ell}$ distribution of gg2VV MC samples is shown on figure 6.19.

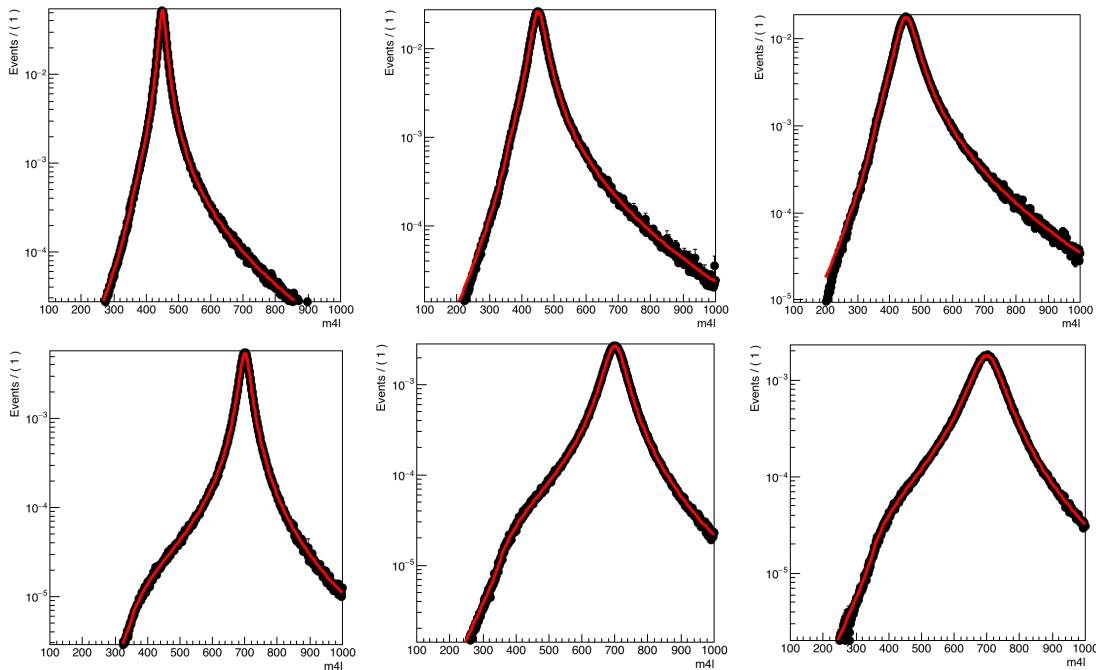


Fig. 6.19 Comparison of the analytical shape to a truth $m_{4\ell}$ distribution of gg2VV MC samples for $m_H = 450$ GeV (top), 700 GeV (bottom) and width equal to 5% (left), 10% (middle), 15% (right) of the mass.

The reconstruction level signal shape can then be modelled by the analytical truth shape convoluted with detector effects modelled in section 6.5.1. A comparison between the modelled shape and reconstruction level MC simulation for signal mass above 400 GeV (for ggF production in $2e2\mu$ channel as an example) are shown in figure 6.20, the shapes are well compatible between each other. This modelling is not valid for lower masses due to the rapid change of detector resolution.

6.5.3 Modelling of interference

There are three processes sharing the same gg initial state and ZZ final state:

- The SM $gg \rightarrow ZZ$ process with an amplitude A_B

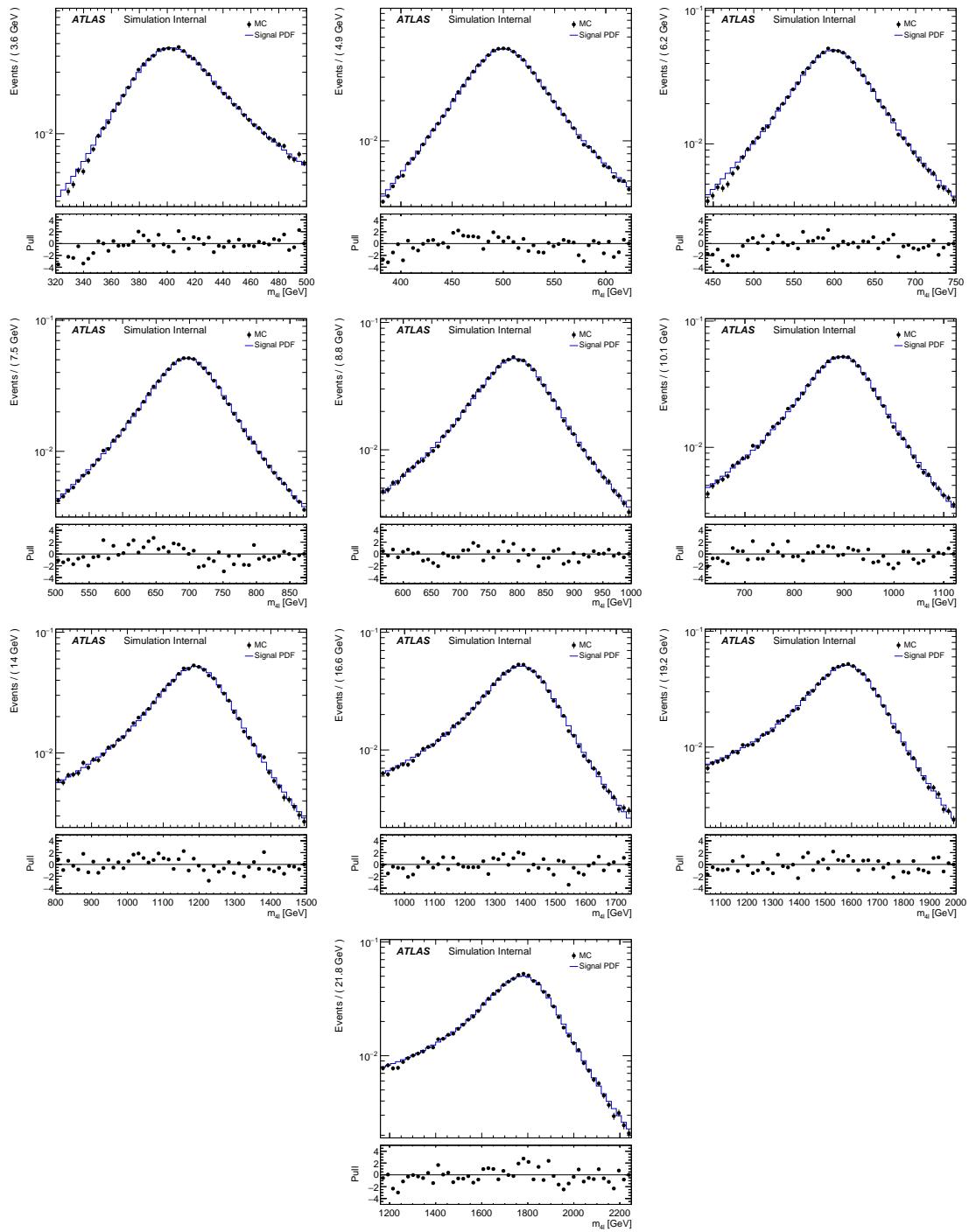


Fig. 6.20 Comparison between the analytical shape convoluted with detector effects and the reconstructed $m_{2\mu 2e}$ MC distribution for mass points ranging from 400 to 1800 GeV and width equal to 15% of the mass.

- The SM (light) Higgs at mass of around 125 GeV with an amplitude A_h

- The BSM heavy Higgs we are searching in this analysis with an amplitude A_H

The three processes can interfere with each other due to the same initial and final states.

The parton cross section for these processes can be written as:

$$\begin{aligned}
 \sigma_{gg \rightarrow (X) \rightarrow ZZ}(s) &= \frac{1}{2s} \int d\Omega |A_h(s, \Omega) + A_H(s, \Omega) + A_B(s, \Omega)|^2 \\
 &= \frac{1}{2s} \int d\Omega \left(|A_h(s, \Omega)|^2 + |A_H(s, \Omega)|^2 + |A_B(s, \Omega)|^2 \right) + \\
 &\quad + \frac{1}{s} \int d\Omega \left(\text{Re}[A_h(s, \Omega) \cdot A_B^*(s, \Omega)] \right. \\
 &\quad \left. + \text{Re}[A_H(s, \Omega) \cdot A_B^*(s, \Omega)] + \text{Re}[A_H(s, \Omega) \cdot A_h^*(s, \Omega)] \right) \\
 &\quad + \frac{1}{s} \text{Re} \left[\frac{1}{s - s_H} \int d\Omega \cdot A_H^P(s, \Omega) \cdot A_H^D(s, \Omega) \cdot A_B^*(s, \Omega) \right] \\
 &\quad + \frac{1}{s} \int d\Omega \cdot \text{Re} \left[A_H^P(s, \Omega) \cdot \frac{1}{s - s_H} \cdot A_H^D(s, \Omega) \cdot A_h^{P*}(s, \Omega) \cdot \frac{1}{(s - s_h)^*} \cdot A_h^{D*}(s, \Omega) \right]
 \end{aligned} \tag{6.11}$$

The first term in equation 6.11 denotes the on-shell SM Higgs contribution, which is negligible in this analysis. The second term corresponds to the heavy Higgs contribution, whose line shape has been described in previous section. The third term is the $gg \rightarrow ZZ$ continuum process, while the forth term is the interference between SM Higgs and $gg \rightarrow ZZ$ continuum. The fifth and sixth terms are the interferences between heavy Higgs and $gg \rightarrow ZZ$ continuum (H-B), and between heavy Higgs and SM Higgs (H-h) that we are interested in. More details about the parameterization of these two interferences are described as below.

1. Interference between heavy Higgs and $gg \rightarrow ZZ$ continuum

The parton cross section of this interference term has been written down in equation 6.11. By assuming that this function has a smooth behaviour, it can be replaced with complex polynomial:

$$\int d\Omega \cdot A_H^P(s, \Omega) \cdot A_H^D(s, \Omega) \cdot A_B^*(s, \Omega) \approx (a_0 + a_1 \cdot \sqrt{s} + \dots) + i \cdot (b_0 + b_1 \cdot \sqrt{s} + \dots) \tag{6.12}$$

The parameters a_i and b_i can be extracted by fitting to the $m_{4\ell}$ distribution from truth level MC simulation after analysis selection. Since the signal mass and width does not enter into this function, the parameters should be independent for every tested signal hypothesis.

Same as description for equation 6.10, the parton cross section can be transformed

2660 into a hadron cross section as a function of $m_{4\ell}$:

$$\sigma_{pp}(m_{4\ell}) = \mathcal{L}_{gg} \cdot \frac{1}{m_{4\ell}} \cdot \text{Re} \left[\frac{1}{s - s_H} \cdot ((a_0 + a_1 \cdot m_{4\ell} + \dots) + i \cdot (b_0 + b_1 \cdot m_{4\ell} + \dots)) \right] \quad (6.13)$$

2661 where the propagators are shown in equation 6.7.

2662 Figure 6.21 shows the distributions of interference function obtained by simultaneous
2663 fitting to $m_{4\ell}$ shape from truth level H-B interference simulation at different mass in $2e2\mu$
2664 channel as an example.

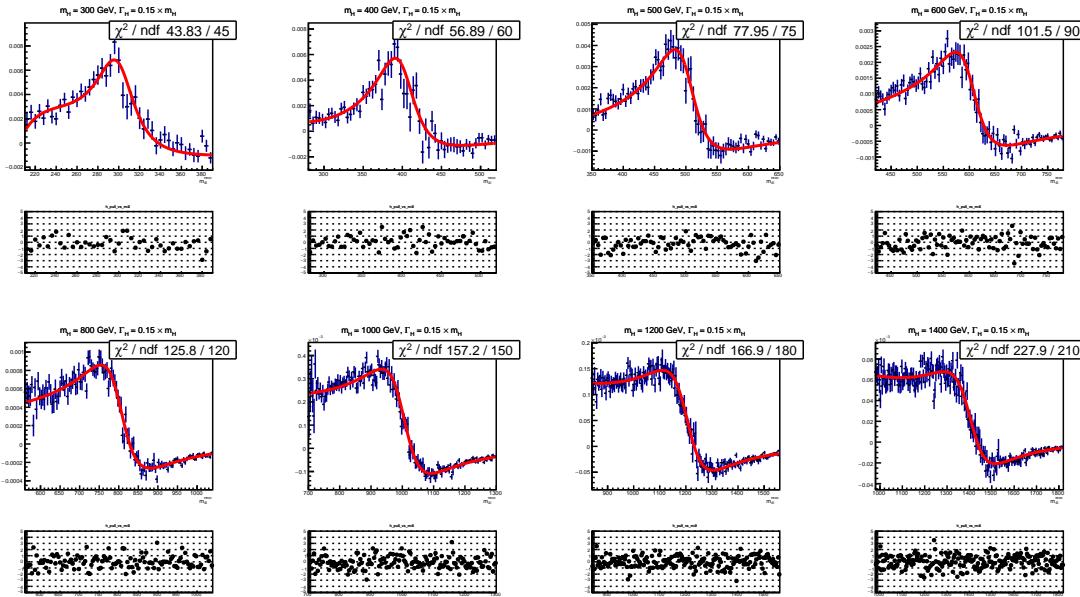


Fig. 6.21 The interference (H-B) model fitted to the truth $m_{4\ell}$ MC distribution after signal region selection for $2\mu 2e$ channel.

2665 **2. Interference between heavy Higgs and SM Higgs**

2666 The parton cross section of this interference term has been written down in equa-
2667 tion 6.11. By assuming the production and decay amplitudes are the same for heavy Higgs
2668 boson and SM Higgs boson, the cross section function can be simplified to:

$$\sigma_{gg}(s) = \frac{1}{s} \int d\Omega \cdot \text{Re} \left[\frac{1}{s - s_H} \cdot \frac{1}{(s - s_h)^*} \right] \cdot |A_{gg \rightarrow H}(s, \Omega)|^2 |A_{H \rightarrow ZZ}(s, \Omega)|^2 \quad (6.14)$$

2669 Taking into account Equation 6.5:

$$\sigma_{gg}(s) = 4 \cdot \text{Re} \left[\frac{1}{s - s_H} \cdot \frac{1}{(s - s_h)^*} \right] \cdot \Gamma_{H \rightarrow gg}(s) \cdot \Gamma_{H \rightarrow ZZ}(s) \quad (6.15)$$

2670 where the propagators are described in equation 6.7, and the partial widths are de-
2671 scribed in equations 6.8 and 6.9.

2672 Same as previous procedure, the parton cross section can be transformed to a hadron
 2673 cross section as a function of $m_{4\ell}$:

$$\sigma_{pp}(m_{4\ell}) = 4 \cdot m_{4\ell} \cdot \mathcal{L}_{gg} \cdot \text{Re} \left[\frac{1}{s - s_H} \cdot \frac{1}{(s - s_h)^*} \right] \cdot \Gamma_{H \rightarrow gg}(m_{4\ell}) \cdot \Gamma_{H \rightarrow ZZ}(m_{4\ell}) \quad (6.16)$$

2674 The modelling procedure of interference is the same as the way for large-width signal
 2675 described in section 6.5.2. The truth line shape is measured as analytical function from
 2676 equation 6.16, and then convolute with detector effect from NWA parameterization to get
 2677 the reconstruction level shape.

2678 For LWA signal model, these two interferences are carefully token into account, and
 2679 the integration of the pure LWA signal with the interferences is used for further studies.
 2680 Figure 6.22 shows the signal model for large-width scenario at mass points of 400 GeV,
 2681 600 GeV, 800 GeV, for three different signal widths: 5%, 10%, 15%, with and without
 2682 interference. Additionally, the contribution of the interference between heavy Higgs and
 2683 SM Higgs (H-h) is shown together with the one between heavy Higgs and SM $gg \rightarrow$
 2684 ZZ background (H-b). One can see the interference effect on signal shape becomes less
 2685 important when going to higher mass.

2686 6.5.4 Modelling of spin-2 RS Graviton signal

2687 The search for Randall-Sundrum (RS) graviton is performed in mass region between
 2688 600 to 2000 GeV. The width of resonance is determined by the $k/\overline{M}_{\text{Planck}}$, which, as
 2689 mentioned in section 6.2.3, is set to be 1. In this configuration, the width of signal is
 2690 expected to be about 6% of its mass.

The reconstructed $m_{4\ell}$ lineshape of graviton is also built by convolving the truth-level lineshape with a detector resolution function, where the detector resolution effect is modelled by a Gaussian + Crystal Ball function, whose parameters are taken from the NWA signal parameterization in section 6.5.1. And for truth-level shape, for graviton, it's modelled as the product of a relativistic Breit-Wigner (RBW) term, a term corresponding to the squared matrix element of the production process and a parton luminosity term \mathcal{L} as given in [111]. So the truth lineshape of $m_{4\ell}$ is token from:

$$m_{4\ell}^{\text{Truth}} \sim \mathcal{L}_{gg} \cdot s^2 \cdot \frac{s(1+s)(1+2s+2s^2)}{(s^2 - m_G^2)^2 + m_G^2 \Gamma^2}$$

2691 The truth-level signal model is extracted by fitting to MC simulation at truth-level with
 2692 the mass m_G and width Γ parameters floating at each mass points respectively. And then

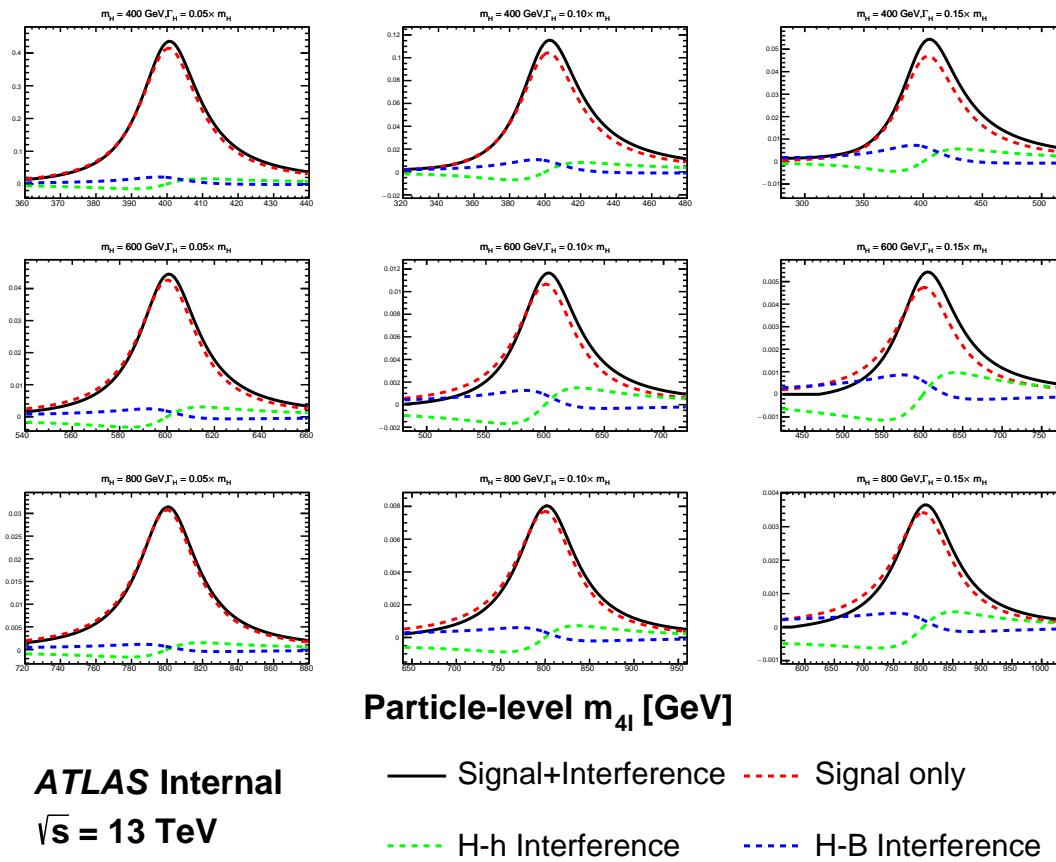


Fig. 6.22 The signal modelling for the large-width scenario at m_H of 400 GeV (top), 600 GeV (middle) and 800 GeV (bottom), as well as three different signal width: 5% (left), 10% (middle) and 15% (right). The contribution of the interference between heavy Higgs and SM Higgs (H - h) is shown together with the one between heavy Higgs and SM $gg \rightarrow ZZ$ background (H - b).

the two parameters are parameterized as the function of m_H by a linear fit as shown in figure 6.23.

The final signal model is obtained by convolving the truth-level lineshape with the detector resolution function. To verify the result, figure 6.24 compares the $m_{4\ell}$ lineshape from parameterization with the one observed from reconstructed-level MC simulation in $2e2\mu$ channel at masses of 600 GeV, 1600 GeV and 2000 GeV as examples.

6.6 Systematic uncertainties

This section describes the sources and value of theoretical and experimental systematic uncertainties considered in this analysis. In addition, as mentioned in previous sections, the uncertainties of irreducible background modelling, reducible background shape smoothing procedure and signal yield difference between simulation and parameterization

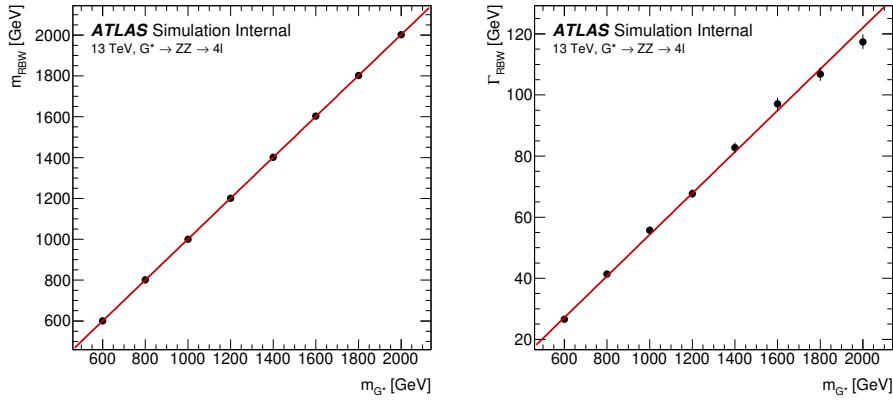


Fig. 6.23 Fitted parameters of the graviton RBW, m_{RBW} and Γ_{RBW} , as a function of the graviton resonance mass, m_G .

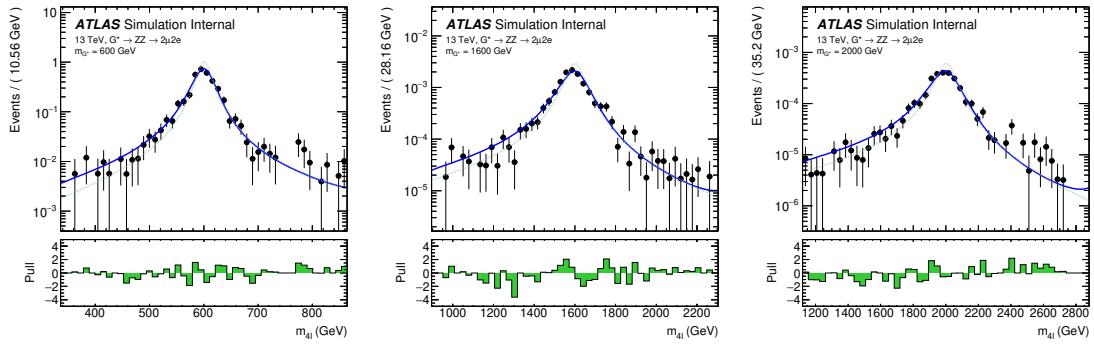


Fig. 6.24 Reconstructed m_{4l} distributions in the $2\mu 2e$ channel with the final signal model superimposed for each RS graviton signal sample at masses of 600 GeV, 1600 GeV and 2000 GeV. The lower panel in each plot shows the pull distribution. The dashed green lines show the truth-level graviton signal models for reference.

2704 are all taken into account.

2705 6.6.1 Theoretical uncertainties

2706 The theoretical modelling uncertainties include the PDF variations, missing QCD
 2707 higher-order corrections via the variations of factorisation and renormalization scales, and
 2708 the parton showering uncertainties.

2709 1. Theoretical uncertainties for signal

2710 The PDF, QCD scale and parton showering uncertainties affecting the acceptance dif-
 2711 ference originating from analysis selection for signal are taken into account in different
 2712 categories. The acceptance uncertainties are calculated on the acceptance factor which
 2713 extrapolates from the fiducial space to the full phase space by a simple ratio:

$$A = \frac{N_{fiducial}}{N_{total}} \quad (6.17)$$

2714 For PDF uncertainties, the standard derivations of 100 PDF replicas of NNPDF3.0
 2715 NNLO, as well as comparison to two external PDF sets: MMHT2014 NNLO, CT14
 2716 NNLO are considered. For missing QCD higher-order corrections, the effects are studied
 2717 with truth events by comparing weights corresponding to variations of the renormalization
 2718 and factorization scale factors, up and down by a factor of two, and the envelop of different
 2719 variations is used. The parton showering uncertainties are estimated by comparing events
 2720 with different setting via PYTHIA8.

2721 Systematic uncertainties are studied for both cut- and MVA- based event categorizations, in two different categories: the inclusive ggF-CBA-enriched and VBF-CBA-enriched category for cut-based analysis and in three different categories: inclusive ggF-MVA-high, ggF-MVA-low and VBF-MVA-enriched category for MVA-based one. This section shows the MVA-based results as an example.

2726 Table 6.10 and 6.11 show the theoretical uncertainties mentioned above for ggF and
 2727 VBF signal respectively in MVA-based categorization.

Table 6.10 Summary of acceptance uncertainties of PDF, QCD scale and parton shower variations for ggF production. The MVA-based categorization is used.

Categories	PDF	QCD Scale	Parton Shower
ggF-MVA-high	0.40%	0.06%	2.03%
ggF-MVA-low	0.56%	0.07%	4.86%
VBF-MVA-enriched	0.53%	0.09%	3.43%

Table 6.11 Summary of acceptance uncertainties of PDF, QCD scale and parton shower variations for VBF production. The MVA-based categorization is used.

Categories	PDF	QCD Scale	Parton Shower
ggF-MVA-high	0.18%	1.20%	0.41%
ggF-MVA-low	0.43%	0.26%	0.36%
VBF-MVA-enriched	0.23%	3.19%	0.85%

2728 **2. Theoretical uncertainties for SM background processes**

2729 The theoretical uncertainties of irreducible ZZ backgrounds are considered in terms
 2730 of both the variations of shape of $m_{4\ell}$ distributions and the acceptance originating from
 2731 the event selection.

2732 The PDF and QCD scale uncertainties are considered by using the same method as
 2733 described for signal. The parton showering uncertainties for those SHERPA samples are
 2734 evaluated by varying the resummation scale by a factor of 2, changing the CKKW setting
 2735 and using different showering option, following the PMG recommendation in ref.^[112],
 2736 and the quadratic sum between the uncertainties in different kinds of showering option

2737 is taken as final result of uncertainties. Moreover, the shape uncertainty associated with
2738 electroweak higher-order correction for $q\bar{q} \rightarrow ZZ$ process is also taken into account.

2739 Same as for signals, these theoretical uncertainties for irreducible backgrounds are
2740 studied for both cut- and MVA- based event categorizations. The value of shape uncer-
2741 tainties vary from less than 1% at low mass region to 50% at high mass tail due to large
2742 statistic fluctuation. As for the acceptance uncertainties, the values vary from about 1%
2743 for PDF variations to 40% for parton showering variations. The VBF category has relative
2744 larger uncertainties.

2745 Table 6.12 summarizes the acceptance uncertainties of PDF, QCD scale, and parton
2746 showering variations for the dominant background: $q\bar{q} \rightarrow ZZ$.

**Table 6.12 Summary of acceptance uncertainties of PDF, scale, and parton showering varia-
tions for QCD $q\bar{q} \rightarrow ZZ$ background. The MVA-based categorization is used.**

Categories	PDF	QCD Scale	Parton showering
ggF-MVA-high	1.15%	10.16 %	3.71%
ggF-MVA-low	1.04%	3.26 %	3.80%
VBF-MVA-enriched	2.91%	27.90 %	23.82%

2747

6.6.2 Experimental systematics

2748 The signal and background predictions used in this analysis are also affected by vari-
2749 ous sources of experimental systematic uncertainties. Similar as described in section 5.5.2,
2750 the dominant experimental uncertainties in this analysis come from the energy/momentum
2751 scales and reconstruction and identification efficiencies of the leptons and jets, as well as
2752 the luminosity uncertainty. The systematic uncertainties are calculated using the recom-
2753 mendations from the Combined Performance (CP) groups of ATLAS experiment. Ta-
2754 ble 6.13 summarizes the experimental systematics considered in this analysis that affect
2755 either the normalization of total event yield or the shape of $m_{4\ell}$ distribution. The impact
2756 of those systematics in statistical fit are studied in section 6.7.

2757

6.7 Results in $\ell\ell\ell'\ell'$ channel

2758 The statistical treatment in searching for heavy resonances in $ZZ \rightarrow \ell\ell\ell'\ell'$ final state
2759 is described in this section. Results are presented in both cut- and MVA- based analysis.

2760 6.7.1 Statistical procedure

2761 The upper limits on heavy resonances are obtained using the unbinned profile likelihood fits. $m_{4\ell}$ is the discriminant. The likelihood function is a product of a Poisson term
 2762 representing the probability of observing n events and a weighted sum of both signal and
 2763 background probability distribution functions (PDFs) evaluated at all observed events.

$$L(x_1..x_n|\sigma_{ggF}, \sigma_{VBF}) = \text{Pois}(n|S_{ggF} + S_{VBF} + B) \left[\prod_{i=1}^n \frac{S_{ggF}f_{ggF}(x_i) + S_{VBF}f_{VBF}(x_i) + Bf_B(x_i)}{S_{ggF} + S_{VBF} + B} \right] \quad (6.18)$$

2765 where f_X s are the probability distribution functions of signal and backgrounds modelled
 2766 in section 6.5 and 6.4, S_X and B are the normalizations of signal and sum of backgrounds.

2767 The parameters of interest (POI) in the search is σ_{ggF} (and σ_{VBF} only for NWA signal),
 2768 which is the cross section of signal model in ggF (and VBF) production mode. In the case
 2769 of there are two POIs, when testing one POI, the other one is profiled along with other
 2770 nuisance parameters (except left unconstrained) during the minimization. These POIs
 2771 enter the likelihood inside the expected signal yields S_{ggF} and S_{VBF} as:

$$S_{ggF(VBF)} = \sigma_{ggF(VBF)} \times BR(S \rightarrow ZZ) \times A \times C \times \int \mathcal{L} \quad (6.19)$$

2772 where $A \times C$ is the signal acceptance as parameterized in 6.3.4, and $\int \mathcal{L} = 139 \text{ fb}^{-1}$ is the
 2773 integrated luminosity of the dataset.

2774 The dependence of the expected number of signal and background events (normaliza-
 2775 tions) and the shape of the PDFs on the systematic uncertainties measured in section 6.6
 2776 is described by a set of nuisance parameters (NPs) θ_i . The Gaussian constraints are ap-
 2777 plied to those NPs. The constraints are implemented as additional ‘penalty’ terms added
 2778 to the likelihood which increase the negative log-likelihood when any nuisance parameter
 2779 is shifted from its nominal value. The final likelihood function $L(\sigma_{ggF}, \sigma_{VBF}, m_H, \theta_i)$ is
 2780 therefore a function of $\sigma_{ggF}, \sigma_{VBF}, m_H$, and θ_i .

2781 Furthermore, the normalization of SM background $pp \rightarrow ZZ$, including both $q\bar{q} \rightarrow$
 2782 ZZ and $gg \rightarrow ZZ$, is a free parameter (μ_{ZZ}) and profiled during the minimization.
 2783 Floating ZZ normalization in fit takes the advantage of reducing the dependence on the-
 2784 ory predictions and their associated uncertainties, especially given that the increased data
 2785 luminosity would provide precise determination of the SM ZZ background rate.

2786 At the end, the upper limit on production cross-section $\sigma_{ggF(VBF)}$ at a given heavy
 2787 resonance model is obtained by setting the mass of signal m_H parameter as constant at the
 2788 desired value, and maximising the likelihood function with respect to nuisance parameters.

2789 The CL_s ^[113] method is used to obtain exclusion limits.

2790 6.7.2 Likelihood fit under background-only hypothesis for MVA-based
2791 analysis

2792 Both MVA- and cut-based analysis are studied by performing likelihood fit to the
2793 (pseudo-) data under the background-only hypothesis and under different signal models.
2794 Due to the same background estimation and modelling procedures, as well as the same
2795 method of systematic measurements, this section only shows the results of background-
2796 only fits for MVA-based analysis under the model of heavy Higgs resonance with narrow-
2797 width as an example. The final results of interpretation in both MVA- and cut- based
2798 analysis in all signal models described in section 6.5 will be measured in next section.

2799 First of all, table 6.14 summarized the expected and observed number of events for
2800 region of $m_{4\ell} > 200$ GeV together with their systematic uncertainties after background-
2801 only fit. The post-fit $m_{4\ell}$ spectrum in each category is shown in figure 6.25.

2802 To inspect the likelihood model, pulls and constraints as well as the correlation matrix
2803 of NPs are studied by performing a background only fit. Figure 6.26 shows the pulls and
2804 constraints when fitting to pseudo-data (top) and observed data (bottom). Figure 6.27
2805 shows the correlation matrix, only for NPs with correlation between each others greater
2806 than 0.1 when fitting to pseudo-data. The normalization of ZZ background is taken from
2807 data for one category each, as shown in table 6.15.

2808 The impact of a systematic uncertainty on the result depends on the production mode
2809 and the mass hypothesis. To check the impact of systematic uncertainties on expected sig-
2810 nal sensitivity, a NP ranking study is performed using signal injected Asimov data with
2811 the injected cross section close to 95% CLs upper limit at the masses of 400 GeV and
2812 1000 GeV. The results are shown in table 6.16. For ggF production, at lower masses, the
2813 systematic uncertainties of parton showering variation for signal, the luminosity uncer-
2814 tainty, and the parametrization of signal acceptance dominate, while at higher masses, the
2815 shape uncertainties from PDF variation for ZZ ($q\bar{q} \rightarrow ZZ$ and $gg \rightarrow ZZ$) background
2816 become important, as also seen in VBF production mode. In addition for VBF, jet related
2817 uncertainties become more important comparing to ggF production. Moreover, the domi-
2818 nate uncertainties include the acceptance uncertainty from QCD scale variation for signal
2819 and the luminosity uncertainty.

2820 6.7.3 Interpretations

2821 1. Spin-0 resonance with NWA

2822 In the absence of a specific model, the ratio of ggF and VBF production mode is un-
2823 known for this additional heavy scalar. For this reason, the fits for ggF and VBF processes

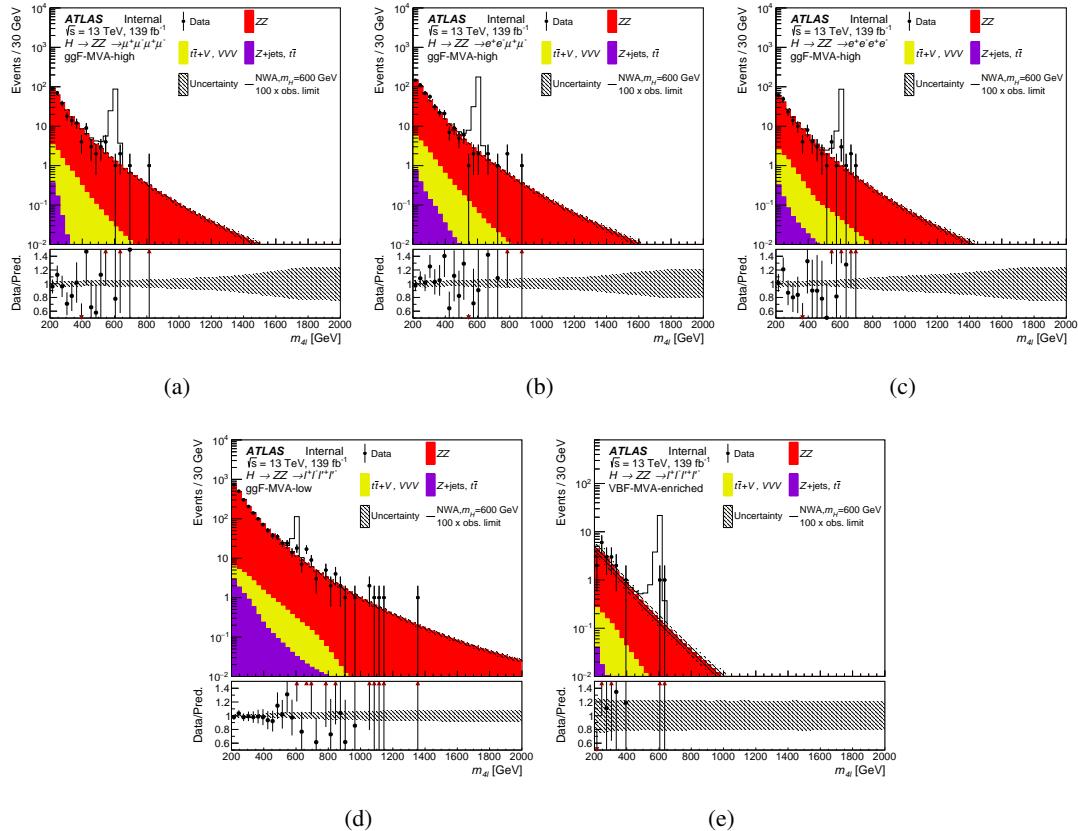


Fig. 6.25 Distribution of the four-lepton invariant mass $m_{4\ell}$ in the $\ell\ell\ell'\ell'$ search for (a), (b), (c) the ggF-MVA-high categories, (d) the ggF-MVA-low category and (e) the VBF-MVA-enriched category. The backgrounds are determined from a combined likelihood fit to the data under the background-only hypothesis. The simulated signal at 600 GeV is normalized to a cross section corresponding to one hundred times the observed upper limit given in section 6.7.3. The error bars on the data points indicate the statistical uncertainty, while the systematic uncertainty in the prediction is shown by the hatched band. The lower panels show the ratio of data to prediction.

are done separately, and in each case the cross section of the untested process is allowed to be a free parameter in the statistical fit. The observed and expected upper limit at 95% confidence level (CL) on the $\sigma \times BR(H \rightarrow ZZ)$ of a narrow scalar resonance for both ggF (left) and VBF (right) production mode with the integrated luminosity of 139 fb^{-1} is shown in figure 6.28 (6.29) for MVA- (cut-) based analysis. No excess over 2σ is found.

2. Spin-0 resonance with LWA

In the case of LWA model, only ggF production mode is studied. The interference between the heavy scalar and SM Higgs boson (H-h), as well as the heavy scalar and SM $gg \rightarrow ZZ$ continuum background (H-B) as modelled in section 6.5.3 are taken into account. The upper limit at 95% confidence level (CL) on ggF cross section times branch ratio ($\sigma_{ggF} \times BR(H \rightarrow ZZ)$) is shown in figure 6.30 for a width of 1, 5, 10 and 15% of

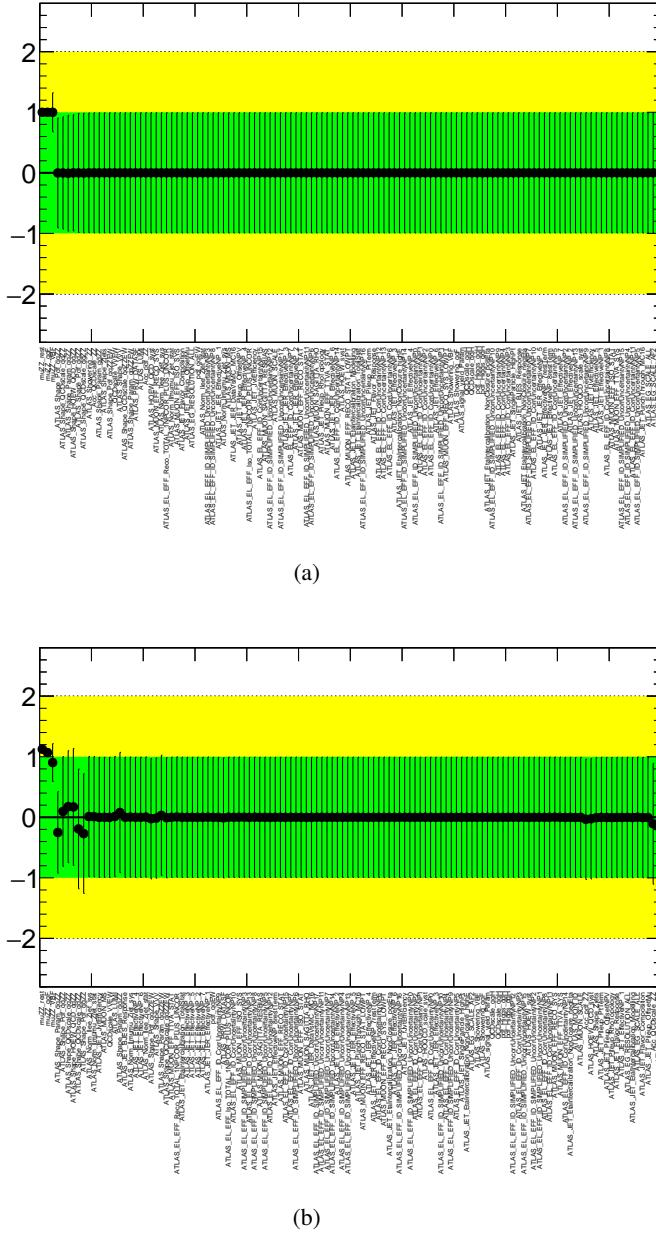
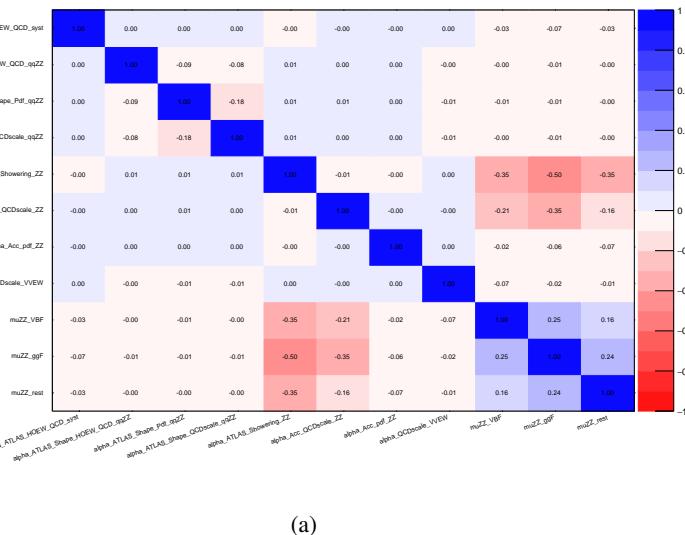


Fig. 6.26 Pulls and constraints of nuisance parameters after a background only fit to (a) Asimov data and (b) observed data in the $\ell\ell\ell'\ell'$ channel. The Asimov data is generated with background data only, and the observed data includes datasets from 2015 to 2018.

2835 m_H .

2836 3. Spin-2 RS Graviton resonance

2837 The observed and expected 95% upper limit on the cross section times branching ratio
 2838 for RS Graviton (RSG) scenario is shown in figure 6.31. Same as LWA case, only $4e$,
 2839 4μ and $2e2\mu$ channel of ggF production mode are used. On top of the expected and ob-
 2840 served upper limits in this model, a predicted cross section as function of mass provided
 2841 by theorist is also shown in the figure. Comparing with the observed result provided by
 2842 $ZZ \rightarrow \ell\ell\ell'\ell'$ decay, this spin-2 graviton is excluded up to a mass of 1500 GeV.



(a)

Fig. 6.27 Correlation of nuisance parameters after a background only fit to Asimov data in the $\ell\ell\ell'\ell'$ channel. The Asimov data is generated with background data only.

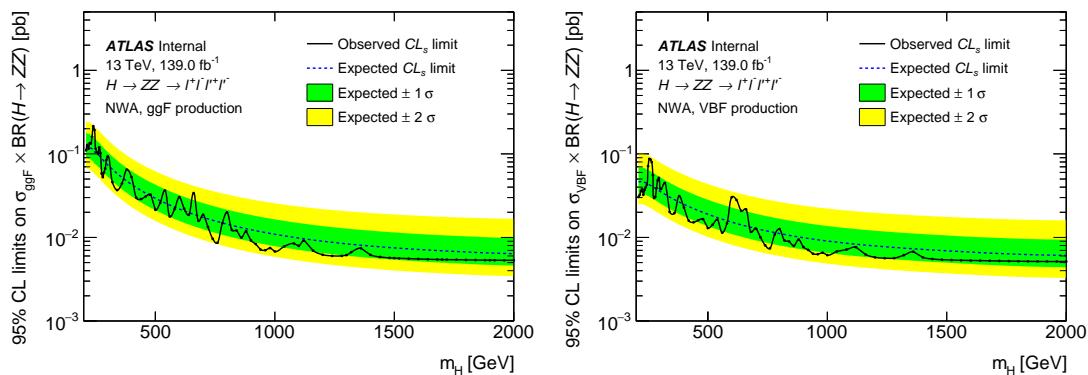


Fig. 6.28 The expected and observed upper limits at 95% CL on $\sigma \times BR(H \rightarrow ZZ)$ using the MVA-based analysis for ggF (left) and VBF (right) production. The green and yellow bands represent the $\pm 1\sigma$ and $\pm 2\sigma$ uncertainties in the expected limits.

4. Summary of interpretation

As a summary, figure 6.32 shows the comparison of expected and observed 95% CL upper limits between different models described above.

Figure 6.33 compares the expected 95% CL upper limits as a function of the NWA resonance mass in this analysis with full run-2 data and the one in previous publication^[114] with the integrated luminosity of 36.1 fb^{-1} . With a significant increase of integrated luminosity and an improved analysis strategy, comparing to the previous publication, the expected sensitivities of searching for narrow-width heavy resonance reduce by up to 70% in MVA-based analysis, where 50% of reduction is due to luminosity increase while other improvement mainly comes from inviting multivariate method.

Figure 6.34 shows the display of one candidate event passing analysis selection in

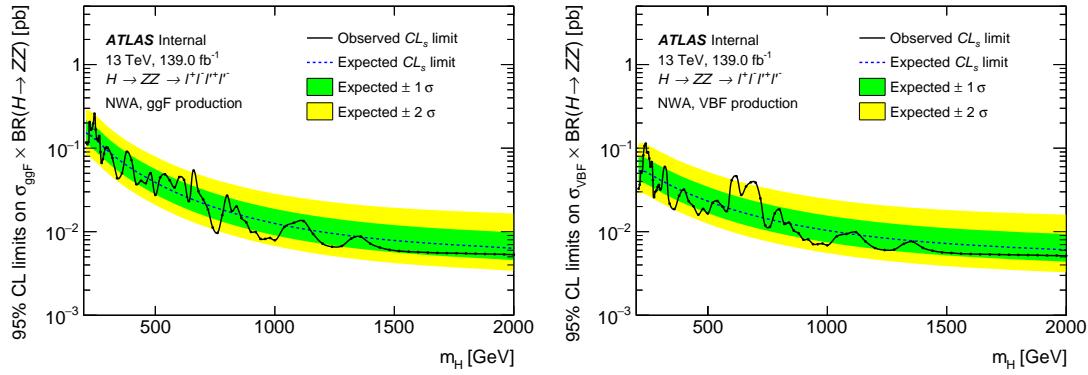


Fig. 6.29 The expected and observed upper limits at 95% CL on $\sigma \times BR(H \rightarrow ZZ)$ using the cut-based analysis for ggF (left) and VBF (right) production. The green and yellow bands represent the $\pm 1\sigma$ and $\pm 2\sigma$ uncertainties in the expected limits.

2854 four-muon final state with four-muon invariant mass of 1.34 TeV.

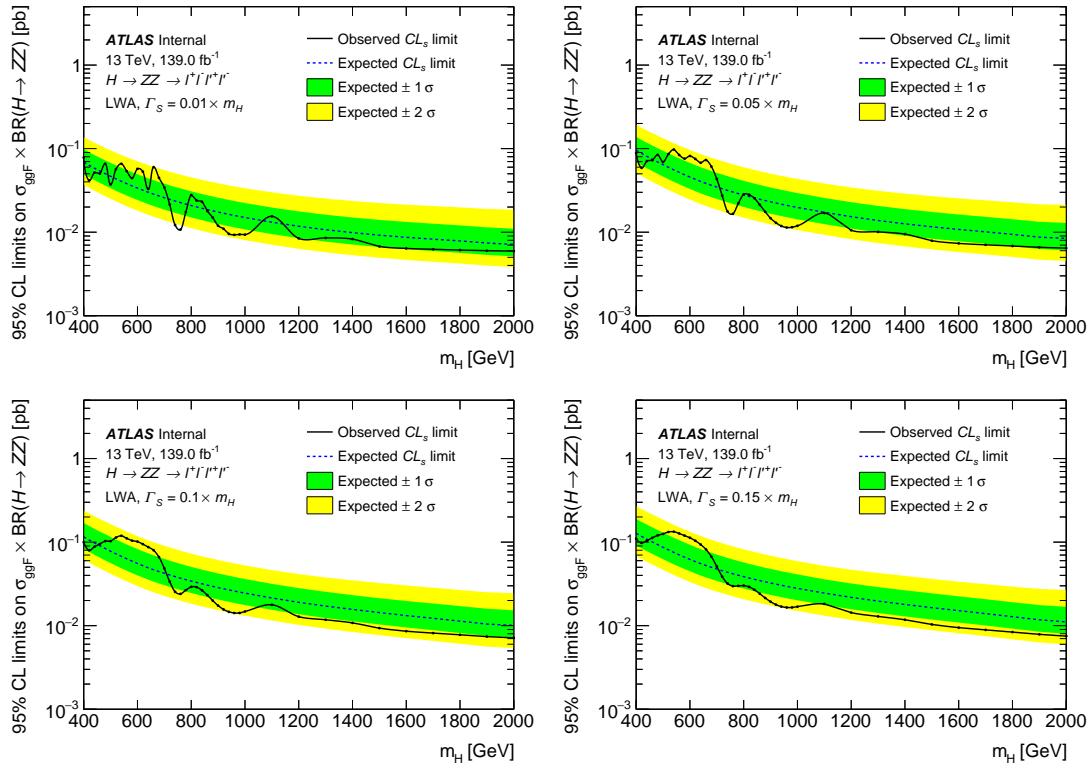


Fig. 6.30 The upper limits at 95% confidence level on $\sigma_{ggF} \times BR(H \rightarrow ZZ)$ as a function of the heavy resonance mass m_H for the ggF production mode with an intrinsic width of 1% (top left), 5% (top right), 10% (bottom left) and 15% (bottom right) for both the case where interference with Standard Model processes is considered. The green and yellow bands represent the $\pm 1\sigma$ and $\pm 2\sigma$ uncertainties in the expected limits.

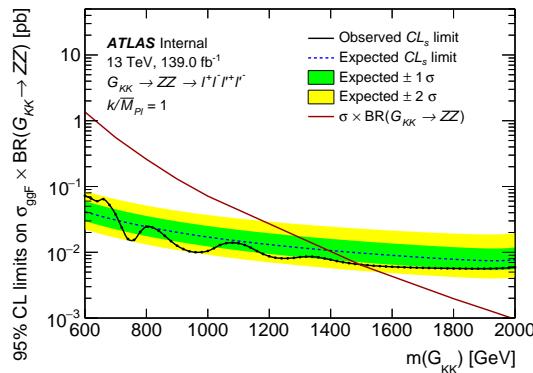


Fig. 6.31 The upper limits at 95% confidence level on $\sigma_{ggF} \times BR(G_{KK} \rightarrow ZZ)$ as a function of the heavy resonance mass $m(G_{KK})$ for the ggF production mode in RS Graviton model. The green and yellow bands represent the $\pm 1\sigma$ and $\pm 2\sigma$ uncertainties in the expected limits.

Table 6.13 A list of the experimental systematics considered in this analysis. The NPs have been separated by whether they only affect the normalisation (left column) or if they affect the shape (right column) of the $m_{4\ell}$ distribution. They are further subdivided into the primary objects that they affect.

Normalisation NPs	Shape NPs
Electrons	
EL_EFF_ID_CorrUncertaintyNP [0–15]	EG_RESOLUTION_ALL
EL_EFF_ID_SIMPLIFIED_UncorrUncertaintyNP [0–17]	EG_SCALE_ALLCORR
EL_EFF_Iso_TOTAL_1NPCOR_PLUS_UNCOR	EG_SCALE_E4SCINTILLATOR
EL_EFF_Reco_TOTAL_1NPCOR_PLUS_UNCOR	EG_SCALE_LARCALIB_EXTRA2015
	EG_SCALE_LARTEMPERATURE_EXT
	EG_SCALE_LARTEMPERATURE_EX
Muons	
MUON_EFF_ISO_STAT	MUON_ID
MUON_EFF_ISO_SYS	MUON_MS
MUON_EFF_RECO_STAT	MUON_SAGITTA_RESBIAS
MUON_EFF_RECO_STAT_LOWPT	MUON_SAGITTA_RHO
MUON_EFF_RECO_SYS	MUON_SCALE
MUON_EFF_RECO_SYS_LOWPT	
MUON_EFF_TTVA_STAT	
MUON_EFF_TTVA_SYS	
Jets	
	JET_BJES_Response
	JET_EffectiveNP_[1–7]
	JET_EffectiveNP_8restTerm
	JET_EtaIntercalibration_Mod
	JET_EtaIntercalibration_Non
	JET_EtaIntercalibration_Non
	JET_EtaIntercalibration_Non
	JET_EtaIntercalibration_Tot
	JET_Flavor_Composition
	JET_Flavor_Response
	JET_JER_DataVsMC
	JET_JER_EffectiveNP_[1–6]
	JET_JER_EffectiveNP_7restTer
	JET_Pileup_OffsetMu
	JET_Pileup_OffsetNPV
	JET_Pileup_PtTerm
	JET_Pileup_RhoTopology
	JET_PunchThrough_MC16
	JET_SingleParticle_HighPt
Other	
HOEW_QCD_syst	
HOEW_syst	
HOQCD_scale_syst	
PRW_DATASF	

Table 6.14 Expected and observed numbers of events for $m_{4\ell} > 200$ GeV, together with their systematic uncertainties, for three MVA-based categories. The expected number of events, as well as their uncertainties, are obtained from a likelihood fit to the data under the background-only hypothesis. The uncertainties of the ZZ normalisation factors, presented in table 6.15, are also taken into account.

Process	VBF-enriched category		ggF-enriched categories			the “rest” category
			4 μ channel	2e2 μ channel	4e channel	
$q\bar{q} \rightarrow ZZ$	11 ± 4		232 ± 10	389 ± 17	154 ± 7	2008 ± 47
$gg \rightarrow ZZ$	3 ± 2		37 ± 6	64 ± 10	26 ± 4	247 ± 19
ZZ (EW)	4.1 ± 0.4		4.5 ± 0.2	7.5 ± 0.4	3 ± 0.2	14.3 ± 0.7
$Z+jets, t\bar{t}$	0.08 ± 0.02		0.6 ± 0.1	1.7 ± 0.4	0.8 ± 0.1	8.8 ± 2.1
$t\bar{t}V, VVV$	0.97 ± 0.1		9.8 ± 0.2	17.5 ± 0.4	7.8 ± 0.2	21.9 ± 0.5
Total background	19 ± 4.5		285 ± 11.7	479 ± 19.7	192 ± 8.1	2301 ± 50.7
Observed	19		271	493	191	2301

Table 6.15 ZZ normalization factor in each category, obtained from a likelihood fit to the data under the background-only hypothesis.

Normalization factor	Fitted value
$\mu_{ZZ}^{ggF-MVA-high}$	1.07 ± 0.047
$\mu_{ZZ}^{ggF-MVA-low}$	1.12 ± 0.026
$\mu_{ZZ}^{VBF-MVA-enriched}$	0.91 ± 0.314

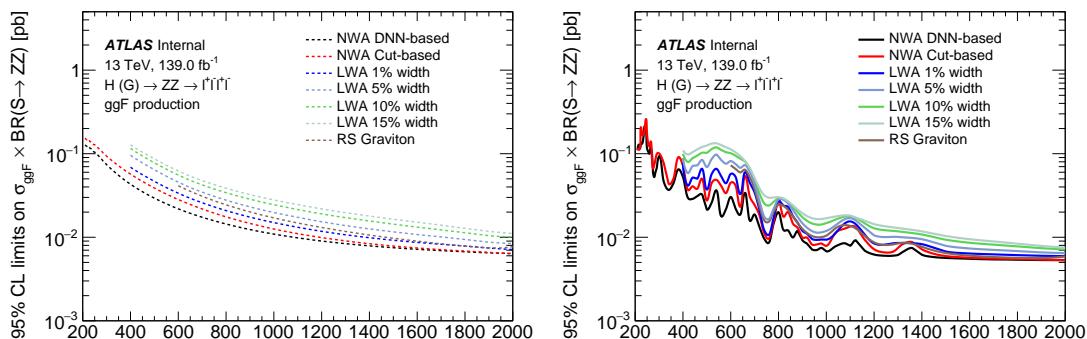


Fig. 6.32 The expected (left) and observed (right) upper limits at 95% confidence level on $\sigma \times BR(S \rightarrow ZZ)$ for ggF production mode at different assumptions.

Table 6.16 Impact of the leading systematic uncertainties, the data statistic uncertainties, as well as the total uncertainties on the predicted signal event yield with the cross section times branching ratio being set to the expected upper limit, expressed as a percentage of the signal yield for the ggF (left) and VBF (right) production modes at $m_H = 400$ and 1000 GeV.

ggF production		VBF production	
Systematic source	Impact [%]	Systematic source	Impact [%]
$m_H = 400$ GeV			
Parton showering of ggF	2.3	QCD scale of VBF	2.7
Luminosity	1.8	Jet flavor composition	2.5
PDF of $q\bar{q} \rightarrow ZZ$	1.6	Luminosity	1.8
Signal yield parameterization	1.4	Jet energy scale (in-su calibration)	1.6
Data stat. uncertainty	48	Data stat. uncertainty	57
Total Uncertainty	49	Total Uncertainty	58
$m_H = 1000$ GeV			
PDF of $q\bar{q} \rightarrow ZZ$	2.5	QCD scale of VBF	2.3
Parton showering of ggF	2.4	PDF of $q\bar{q} \rightarrow ZZ$	2.2
PDF of $gg \rightarrow ZZ$	1.9	Luminosity	1.8
Luminosity	1.8	PDF of $gg \rightarrow ZZ$	1.6
Data stat. uncertainty	84	Data stat. uncertainty	92
Total Uncertainty	86	Total Uncertainty	93

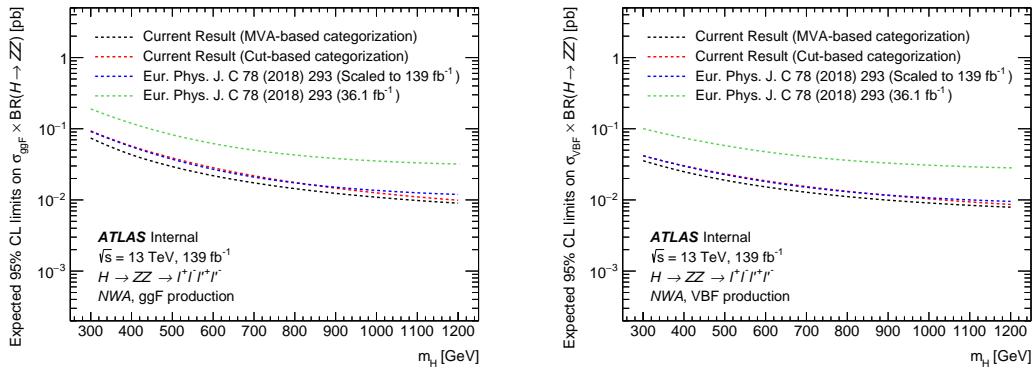


Fig. 6.33 Comparisons of the expected upper limits at 95% CL on the cross section times branching ratio as a function of the heavy resonance mass m_H for the ggF production mode (left) and for the VBF production mode (right) in the case of the NWA. The expected limits from the previous publication are shown in the green dashed line and are projected to the 139 fb^{-1} as shown in the blue dashed line. In addition, the current results based on either cut-based categorisation or the multivariate-based categorisation are shown in red and black lines.

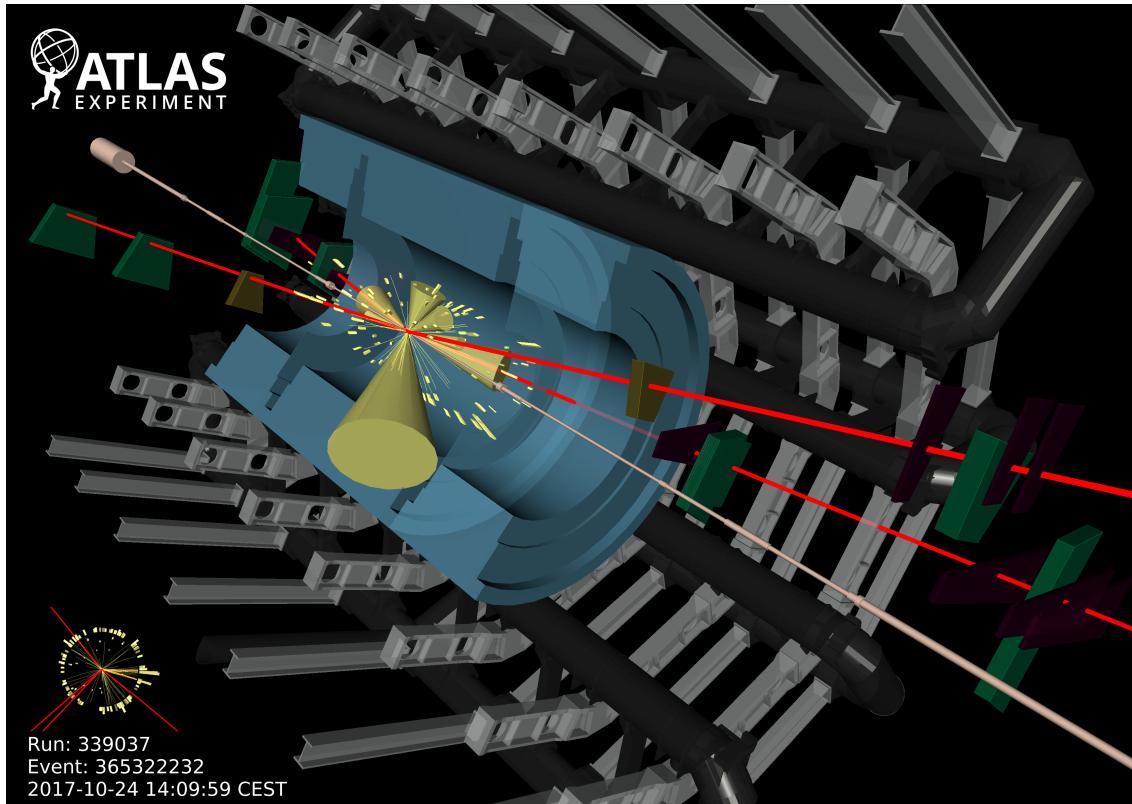


Fig. 6.34 Display of one candidate event in 4μ final state with the mass of 1.35 TeV.

2855 6.8 Conclusion

2856 Searches of heavy resonances decaying into a pair of Z boson to $\ell\ell\ell'\ell'$ final state
2857 are performed using 139 fb^{-1} of 13 TeV pp collision data collected by ATLAS experiment
2858 at the LHC. The results are interpreted as 95% CL upper limits on the production cross
2859 section of a spin-0 and spin-2 resonances under different theoretical models. The search
2860 range of the hypothetical resonances is between 200 GeV to 2000 GeV depending on the
2861 signal model.

2862 The spin-0 resonance is assumed to be a heavy Higgs like scalar produced predom-
2863 inantly from gluon–gluon fusion (ggF) and vector-boson fusion (VBF) decays, and it is
2864 studied under both the narrow-width approximation and with the large-width assumption.
2865 For narrow-width approximation, limits on cross section of heavy scalar decaying into
2866 two Z bosons are set separately for ggF and VBF production modes, under DNN- and
2867 cut- based analysis. In DNN-based analysis, the 95% CL upper limit range is from 215 fb
2868 at $m_H = 240 \text{ GeV}$ to 5.3 fb at $m_H = 2000 \text{ GeV}$ for ggF production mode, and from 87 fb
2869 at $m_H = 255 \text{ GeV}$ to 5.1 fb at $m_H = 1960 \text{ GeV}$ for VBF production mode. In cut-based
2870 analysis, the 95% CL upper limit range is from 259 fb at $m_H = 245 \text{ GeV}$ to 5.3 fb at
2871 $m_H = 2000 \text{ GeV}$ for ggF production mode, and from 113 fb at $m_H = 240 \text{ GeV}$ to 5.1 fb
2872 at $m_H = 2000 \text{ GeV}$ for VBF production mode. DNN-based analysis gains about 20% im-
2873 provement on upper limits at lower mass region comparing to the cut-based analysis, while
2874 for mass above 1500 GeV, both analyses perform closely. For large-width approximation,
2875 limits are studied on ggF production rate at four different widths assumptions: 1%, 5%,
2876 10% and 15% of resonance’s mass, with the interference between the heavy scalar and the
2877 SM Higgs boson as well as the heavy scalar and the SM $gg \rightarrow ZZ$ continuum background
2878 taken into account. The maximum and minimum of upper limits are obtained as 78 fb at
2879 $m_H = 400 \text{ GeV}$ to 5.9 fb at $m_H = 2000 \text{ GeV}$ for 1% width; 98 fb at $m_H = 540 \text{ GeV}$ to 6.4
2880 fb at $m_H = 2000 \text{ GeV}$ for 5% width; 119 fb at $m_H = 540 \text{ GeV}$ to 7.1 fb at $m_H = 2000 \text{ GeV}$
2881 for 10% width; 133 fb at $m_H = 540 \text{ GeV}$ to 7.5 fb at $m_H = 2000 \text{ GeV}$ for 15% width. Last
2882 but not least, the framework of the Randall–Sundrum model with a graviton excitation
2883 spin-2 resonance with $m(G_{KK}) < 1500 \text{ GeV}$ is excluded at 95% CL.

2884

Chapter 7 Summary

2885 On December 3rd, 2018, the LHC finished its second run (run-2) after three fantastic
2886 years. Thanks to run-2 with largely increased statistic, we now know the masses of the
2887 Higgs boson, top quark and W boson to considerably greater precision. And also confirm
2888 the Standard Model as a stable theory.

2889 In this dissertation, various physics processes in $ZZ \rightarrow \ell\ell\ell'\ell'$ final state are studied,
2890 taking the advantage of full run-2 pp collision data in the LHC. Using this signature,
2891 we measured the fiducial cross section of ZZ production to $\ell\ell\ell'\ell' jj$ channel in SM,
2892 which is an important physics process and major background in many analysis with ZZ
2893 production, eg. Higgs analysis (HZZ). In addition, we searched the electroweak ZZ
2894 production via vector boson scattering in associated with 2-jet process in $\ell\ell\ell'\ell'$ final
2895 state. In the meantime, the searches of heavy resonances decaying into a pair of Z bosons
2896 to $\ell\ell\ell'\ell'$ final state for several different hypothetical resonances are conducted in this
2897 dissertation. The results of several analyses are summarized as below:

2898 **Measurement of fiducial cross section of ZZ production in $\ell\ell\ell'\ell' jj$ final state**

2899 The fiducial cross section of inclusive SM $ZZ \rightarrow \ell\ell\ell'\ell'$ production is measured to
2900 be:

$$\sigma_{ZZ \rightarrow \ell\ell\ell'\ell'}^{fid} = 1.27 \pm 0.12(stat) \pm 0.02(theo) \pm 0.07(exp) \pm 0.01(bkg) \pm 0.03(lumi) \quad (7.1)$$

2901 which is found to be compatible with the SM prediction. The ZZ cross section is calcu-
2902 lated with up to one (three) outgoing partons at NLO (LO) using **SHERPA** 2.2.2 for QCD
2903 production, and in LO using **MADGRAPH5_aMC@NLO** 2.6.1 for EW production. The
2904 total uncertainty is 11%, the analysis is still data static dominant (data statistic uncertainty
2905 is about 9.5%).

2906 **Observation of electroweak ZZ production in $\ell\ell\ell'\ell' jj$ final state**

2907 Thanks to the largely increased data statistic collected by ATLAS experiment in the
2908 LHC run-2, the electroweak ZZ production (EW- $ZZjj$) to $\ell\ell\ell'\ell'$ channel in associa-
2909 tion with two jets is observed with a significant deviation from the background-only
2910 hypothesis. The signal strength of EW- $ZZjj$ production, the normalization of QCD-
2911 $ZZjj$ production, as well as the observed and expected statistical significance measured

2912 in $\ell\ell\ell'\ell' jj$ channel are found to be:

$$\begin{aligned}\mu_{\text{EW}} &= 1.54 \pm 0.42 \\ \mu_{\text{QCD}} &= 0.95 \pm 0.22\end{aligned}\quad (7.2)$$

Obs. (Exp.) Significance = 5.48 (3.90) σ

2913 Then in this dissertation, the differential cross section and expected significance of
2914 EW- $Z Z jj$ production, using 3000 fb^{-1} simulated pp collision data at a centre-of-mass
2915 energy of 14 TeV to be recorded by ATLAS experiment at the HL-LHC, are studied via
2916 simulations. The HL-LHC will for sure give us more opportunity to probe rare process
2917 like $Z Z \rightarrow \ell\ell\ell'\ell'$ in the future.

2918 **Searches of heavy $Z Z$ resonances in $\ell\ell\ell'\ell'$ final state**

2919 Searches of heavy $Z Z$ resonances are performed in four-lepton invariant mass $m_{4\ell}$
2920 range from 200 GeV to 2000 GeV. Data are found to agree with the background-only
2921 hypothesis, and 95% CL upper limits are set on the production rate under the models of:

- Spin-0 heavy Higgs under narrow-width approximation (NWA).

2923 Search range is from 200 GeV to 2000 GeV.

2924 In DNN-based analysis, the limits are range from 215 fb at $m_H = 240$ GeV to 5.3
2925 fb at $m_H = 2000$ GeV for ggF production mode, and from 87 fb at $m_H = 255$ GeV
2926 to 5.1 fb at $m_H = 1960$ GeV for VBF production mode.

2927 The DNN-based analysis is found to be at most 20% better than cut-based results.

- 2928 • Spin-0 heavy Higgs under large-width approximation (LWA) with the width of 1,
2929 5, 10, 15% of its mass.

2930 Search range is from 400 GeV to 2000 GeV, and only ggF production is studied.

2931 The maximum and minimum of upper limits are obtained as 78 fb at $m_H = 400$ GeV
2932 to 5.9 fb at $m_H = 2000$ GeV for 1% width; 98 fb at $m_H = 540$ GeV to 6.4 fb at $m_H =$
2933 2000 GeV for 5% width; 119 fb at $m_H = 540$ GeV to 7.1 fb at $m_H = 2000$ GeV for
2934 10% width; 133 fb at $m_H = 540$ GeV to 7.5 fb at $m_H = 2000$ GeV for 15% width.

- 2935 • Spin-2 graviton excitation under the Randall–Sundrum model.

2936 Search range is from 600 GeV to 2000 GeV, and only ggF production is studied.

2937 The maximum and minimum of limits are 73 fb at $m_H = 600$ GeV and 5.6 fb
2938 at $m_H = 1880$ GeV for ggF production mode. And the mass of graviton below
2939 1500 GeV is excluded comparing the observed results with theoretical prediction.

2940
2941 In summary, the $Z Z \rightarrow \ell\ell\ell'\ell'$ production presented in this dissertation are consis-
2942 tent with SM prediction. This result completes the observation of weak boson scattering

2943 for massive bosons, which is a new milestone reached in the study of electroweak symme-
2944 try breaking. In the meantime, no indication of new physics is observed. We are looking
2945 forward the HL-LHC, with greatly increased luminosity and higher centre-of-mass energy,
2946 which should enhance the sensitivity for new physics search and precise measurement for
2947 rare process like $\ell\ell\ell'\ell'$ final state.

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