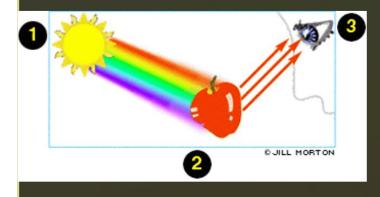
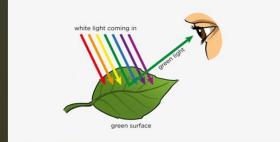
Table of Contents

- Background
- Method
- Results
- Conclusion & Significance
- Next steps

Background

How do we see the color of an object?

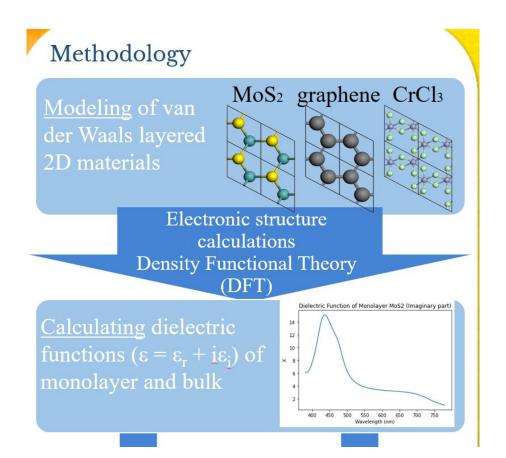




Three processes:

- Incident white light
- Certain color reflected by the object
- 3. Human <mark>eye</mark>s' perception

Brief Review of Method

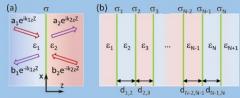


Brief Review of Method

multi-layers

single layer/bulk

Extrapolate for multi-layers' from transfer matrix method



$$\begin{bmatrix} b_1 \\ a_{N+1} \end{bmatrix} = \begin{bmatrix} M_{21}/M_{11} & (M_{11}M_{22} - M_{12}M_{21}) / M_{11} \\ 1/M_{11} & -M_{12}/M_{11} \end{bmatrix} \begin{bmatrix} a_1 \\ b_{N+1} \end{bmatrix}$$

$$\mathcal{M} = D_{1\to 2}P(d_{1,2})D_{2\to 3}P(d_{2,3})\cdots P(d_{N-1,N})D_{N\to N+1}$$

$$D_{1\to 2} = \frac{1}{2} \begin{bmatrix} 1 + \eta_s + \xi_s & 1 - \eta_s + \xi_s \\ 1 - \eta_s - \xi_s & 1 + \eta_s - \xi_s \end{bmatrix} P(\Delta z) = \begin{bmatrix} e^{-ik_z \Delta z} & 0 \\ 0 & e^{ik_z \Delta z} \end{bmatrix}$$

where **D** is the transmission matrix relating incident and reflected waves across the monolayer; **P** is the propagation matrix for the propagation of light in a homogenous medium; **M** is the transfer matrix relating incident and reflected waves across the multi-layer materials.

Calculate reflectivity spectra using S Gupta et al.

$$T = \frac{4n_1 n_2}{|n_1 + n_2 + \sigma_{2D} Z_{\text{vac}}|^2}$$

$$A = \frac{4n_1 Re\{\sigma_{2D}\} Z_{\text{vac}}}{|n_1 + n_2 + \sigma_{2D} Z_{\text{vac}}|^2}$$

$$R = \left| \frac{n_2 - n_1 + \sigma_{2D} Z_{\text{vac}}}{n_1 + n_2 + \sigma_{2D} Z_{\text{vac}}} \right|^2$$

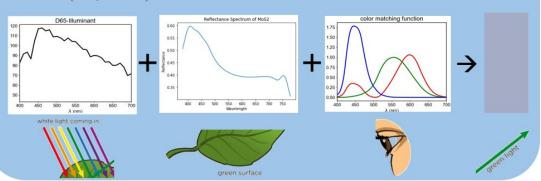
where **2D materials** are viewed as thin layers between two semi-infinite dielectric media slabs with normally incident and reflected waves across the interface; $\mathbf{n}_{1,2}$ are refractive indexes on the two sides of 2D materials (set to 1 as vacuum in our case); $\mathbf{\sigma}_{2D}$ is optical surface conductivity as a function of frequency and 2D wavevector; \mathbf{Z}_{vac} is impedance of vacuum.

$$\sigma_{2D}(q=0,\omega) = i\omega\varepsilon_0(1-\varepsilon_{3D}(q=0,\omega))L$$

Brief Review of Method

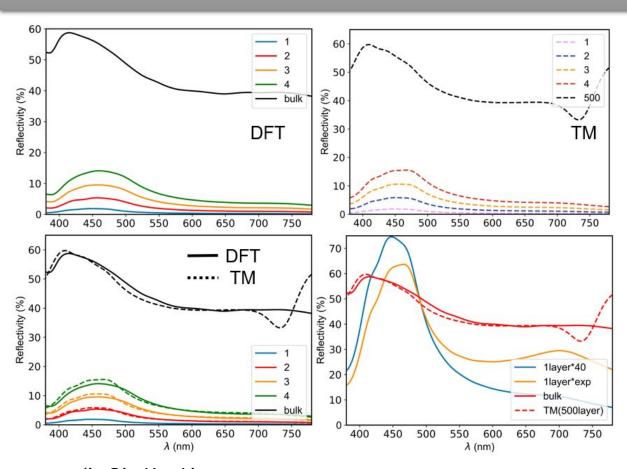
CIEXYZ Color Space → sRGB

where $\mathbf{x}(\lambda)$, $\mathbf{y}(\lambda)$ and $\mathbf{z}(\lambda)$ are the color matching function; $\mathbf{R}(\lambda)$ is the reflectivity obtained by DFT/transfer matrix method; $\mathbf{S}(\lambda)$ is the spectral power distribution of one of the standard CIE illuminant (D65, D50....).



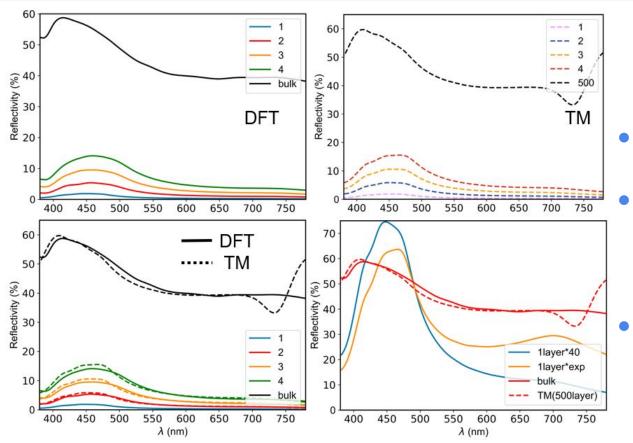
Results: Table of Selected 2D Materials

Types	Materials
TMD	MoS2, WSe2
Graphene	Graphene
Chromium Trihalides CrX3 (X=Cl, Br, I)	CrCl3



AA stacking Transfer matrix (TM) method: Interlayer separation 6.81Å

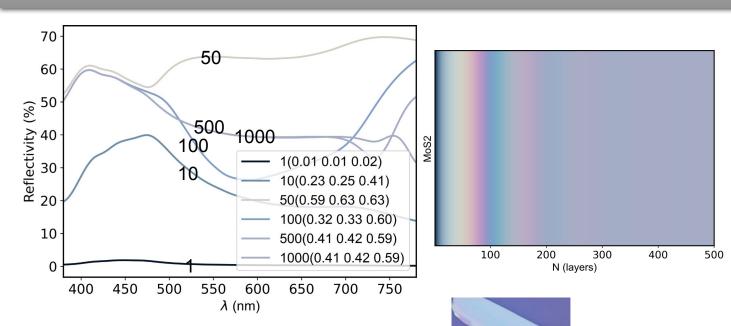
credit. Qin-Kun Li

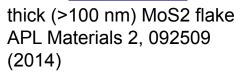


AA stacking Transfer matrix (TM) method: Interlayer separation 6.81Å

- TM produces good fit
 - o deviation after 700 nm
- Peaks between 400-550 nm
 - expect blue/ purple

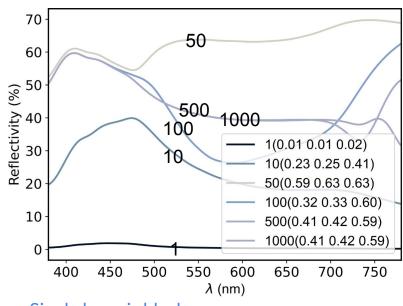
 Single layer and bulk: no linear or exponential relationship

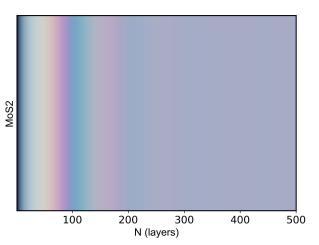






crystalline bulk MoS2 single crystal grown by Sn flux method CrystEngComm, 2015,17, 4026





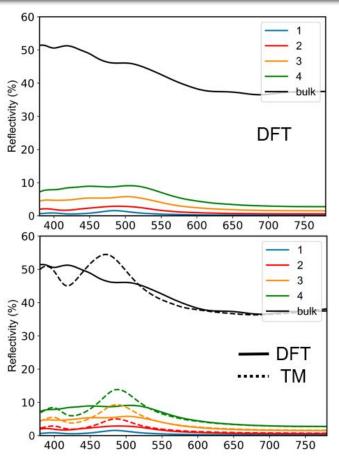
- Single layer is black
- 50 layers is creamy white
- Color is stable after ~400 layers
- Pattern:
 - pink/blue stripes fading periodically APL Materials 2, 092509
 - "blending yogurt"

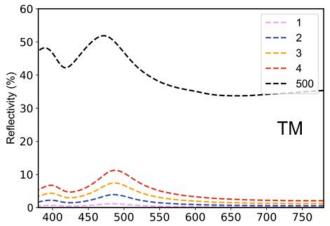


thick (>100 nm) MoS2 flake APL Materials 2, 092509 (2014)



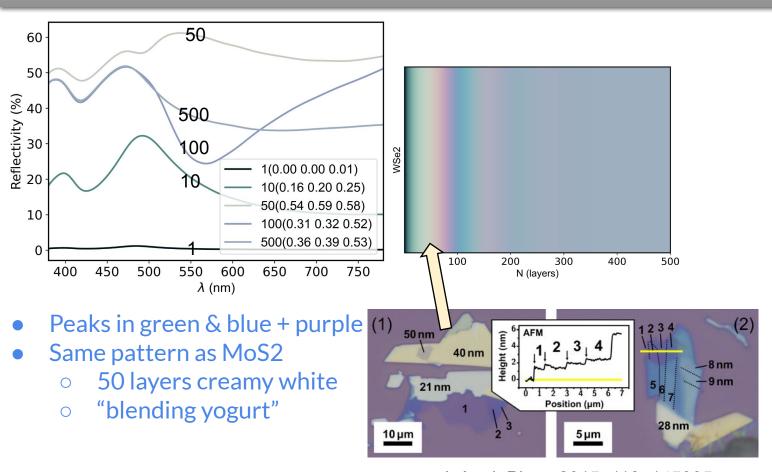
crystalline bulk MoS2 single crystal grown by Sn flux method CrystEngComm, 2015,17, 4026



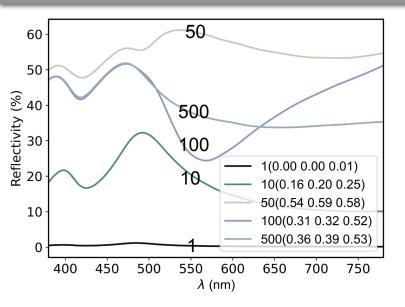


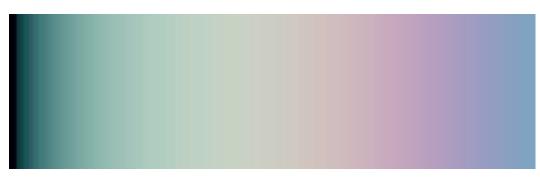
AA stacking Transfer matrix (TM) method: Interlayer separation 7.03Å

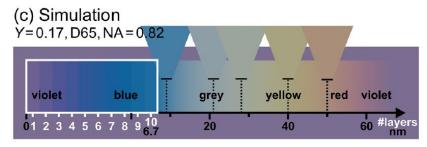
- DFT has plateau instead of peaks
- TMM produces wave-like curves
 - based on monolayer



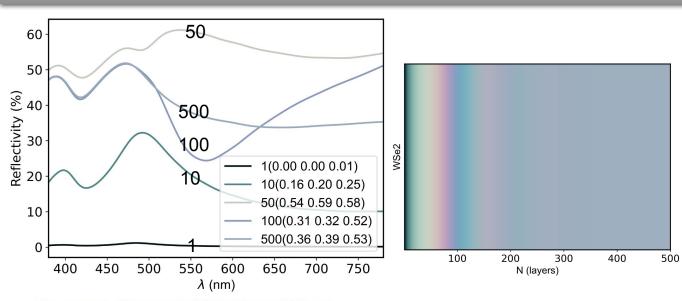
J. Appl. Phys. 2015, 118, 145305

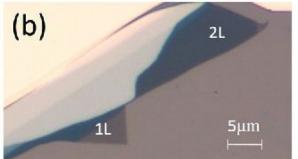






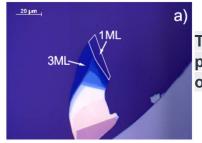
J. Appl. Phys. 2015, 118, 145305





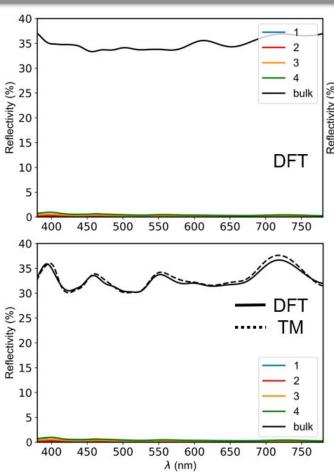
Optical microscopic image of pre-marked WSe2 monolayer, bilayer, and thin bulk flakes;

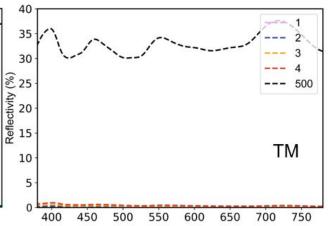
101901-3 Yan et al. Appl. Phys. Lett. 105, 101901 (2014)



Temperature dependence of photoluminescence lifetime of atomically thin WSe2 layer

graphene

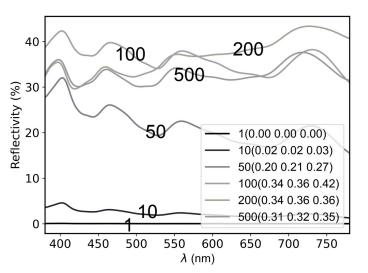


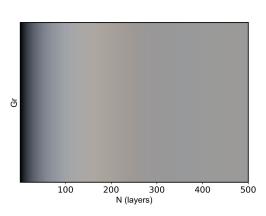


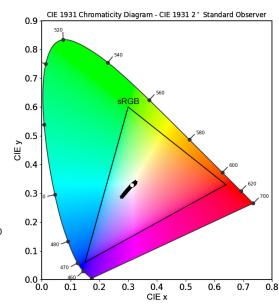
AA stacking Transfer matrix (TM) method: Interlayer separation 3.53Å

- Multiple peaks
- TM produces good fitting

graphene

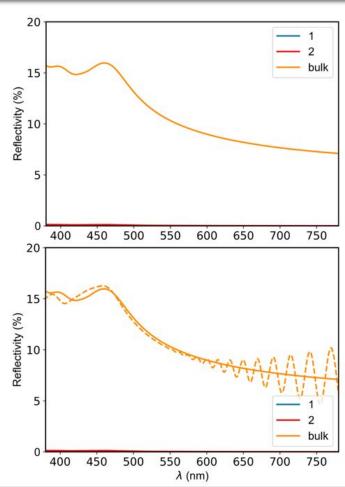


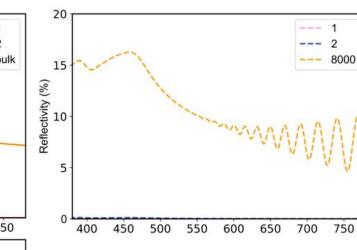




- Less "colorful" than other materials
 - low reflectivities & multiple peaks
 - o due to simple structure and composition?

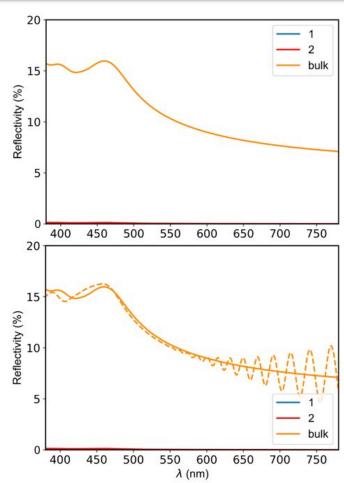
CrCl3

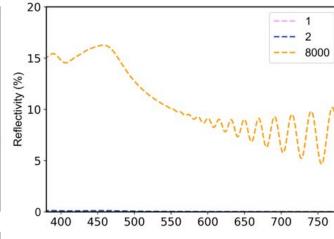




AA stacking Transfer matrix (TM) method: Interlayer separation 6.01Å

CrCl3

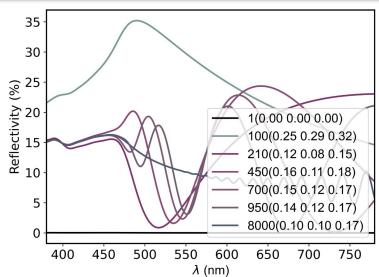


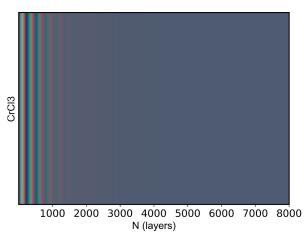


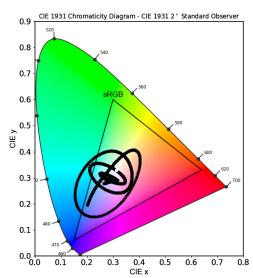
AA stacking Transfer matrix (TM) method: Interlayer separation 6.01Å

- Bulk has big drop
- expect mixture of blue and purple
- TM produces wave-like fluctuations near the end

CrCI3



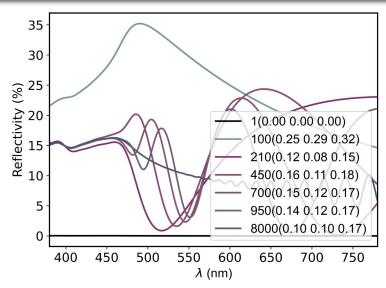


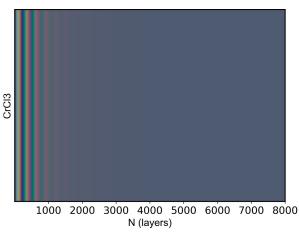


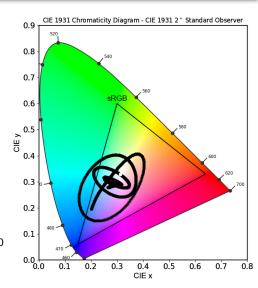


Bulk single crystal of CrCl3 Nat. Phys. 2019 15, 1255

CrCI3







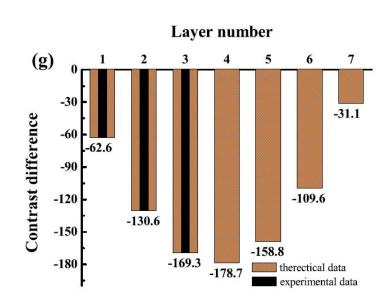
- Pink curves have two peaks
 - violet + red → magenta
- Constant color at ~3000 layers
 - takes longer than other materials



Bulk single crystal of CrCl3 Nat. Phys. 2019 15, 1255

Applications

Thickness identification for material characterization

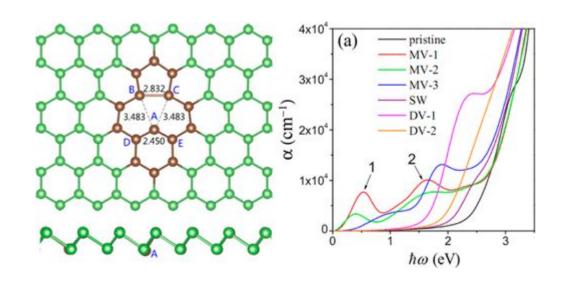


Optical thickness identification of few-layer MoS₂ deposited by chemical vapor deposition

Zusong Zhu *et al* 2019 *Mater. Res. Express* 6 045025

Applications

2. Defect identification



Optical Identification of Topological Defect Types in Monolayer Arsenene by First-Principles Calculation

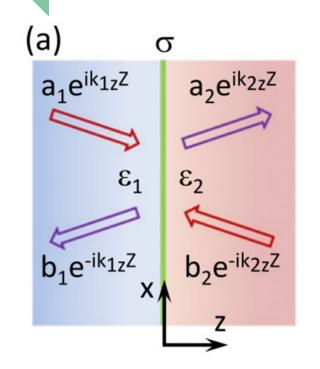
J. Phys. Chem. C 2016, 120, 43, 24917-24924

Feedback

- color not resulted from defects
- gold



Yang Yiyi Zhu Zien Wu Xinan



$$\left[egin{array}{c} a_1 \ b_1 \end{array}
ight] = D_{1 o 2} \left[egin{array}{c} a_2 \ b_2 \end{array}
ight]$$

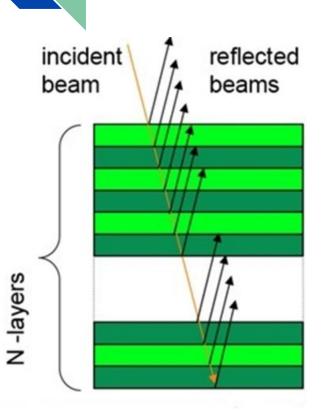
$$egin{align} D_{1 o\,2,m} &= rac{1}{2} egin{bmatrix} 1 + \eta_m + \xi_m & 1 - \eta_m - \zeta_m \xi_m \ 1 - \eta_m + \zeta_m \xi_m & 1 + \eta_m - \xi_m \end{bmatrix} \ & ilde{r} &= rac{b_1}{a_1} = rac{D_{21}}{D_{22}}, ilde{t} &= rac{a_2}{a_2} = rac{1}{D_{22}} \end{split}$$

for p polarization:

$$ilde{r}_p=rac{1-\eta_p+\xi_p}{1+\eta_p+\xi_p}=rac{ ilde{\sigma}_p}{2+ ilde{\sigma}_p Z_{vac}},\; ilde{t}_p=rac{2}{1+\eta_p+\xi_p}=rac{2}{2+ ilde{\sigma}_p Z_{vac}}.$$

for p polarization:

$$ilde{r}_s = rac{1 - \eta_s - \xi_s}{1 + \eta_s + \xi_s} = rac{- ilde{\sigma}_s}{2 + ilde{\sigma}_s Z_{vac}}, \; ilde{t}_s = rac{2}{1 + \eta_s + \xi_s} = rac{2}{2 + ilde{\sigma}_s Z_{vac}}$$

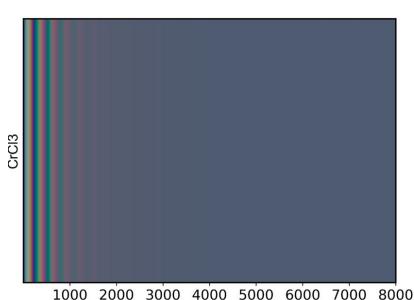


$$egin{aligned} ilde{r}_{total} &= ilde{r} + ilde{r} ilde{t}^2 \Delta^2 + ilde{r} ilde{t}^4 \Delta^4 + \cdots + ilde{r} ilde{t}^{2(N-1)} \Delta^{2(N-1)} \ & ext{Let} \qquad ilde{x} = ilde{t}^2 \Delta^2 = ilde{t}^2 e^{i2kd} \ & ilde{r}_{total} = ilde{r} (1 + ilde{x} + ilde{x}^2 + \cdots + ilde{x}^{N-1}) = ilde{r} rac{1 - ilde{x}^N}{1 - ilde{x}} \ & R_N = ilde{r}_{total}^* ilde{r}_{total} = \left| rac{ ilde{r}}{1 - ilde{x}}
ight|^2 \cdot |1 - ilde{x}^N|^2 \ & ilde{x} = ilde{t}^2 e^{i2kd} = T e^{2i(kd + heta)} \qquad T = ilde{t}^* ilde{t} pprox 1, ilde{t} = \sqrt{T} \, e^{i heta} \end{aligned}$$

$$R_{\scriptscriptstyle N}\!=\!\left|rac{ ilde{r}}{1- ilde{x}}
ight|^2\cdot\left(1-2{
m Re}(ilde{x}^{\scriptscriptstyle N})+| ilde{x}^{\scriptscriptstyle N}|^2
ight)\!=\!\left|rac{ ilde{r}}{1- ilde{x}}
ight|^2\cdot\left[1-2T^{\scriptscriptstyle N}{
m Re}ig(e^{2iN(kd+ heta)}ig)\!+T^{2N}ig]$$

$$R_{N} = ilde{r}_{total} ^{st} ilde{r}_{total} = \left| rac{ ilde{r}}{1 - ilde{x}}
ight|^{2} \cdot |1 - ilde{x}^{N}|^{2}$$

$$R_N = \left|rac{ ilde{r}}{1- ilde{x}}
ight|^2 \cdot \left(1-2 ext{Re}\left(ilde{x}^N
ight) + \left| ilde{x}^N
ight|^2
ight) = \left|rac{ ilde{r}}{1- ilde{x}}
ight|^2 \cdot \left[1-2T^N ext{Re}\left(e^{2iN(kd+ heta)}
ight) + T^{2N}
ight]$$



N (layers)

when the number of layers increases by δN , and satisfy the following relation:

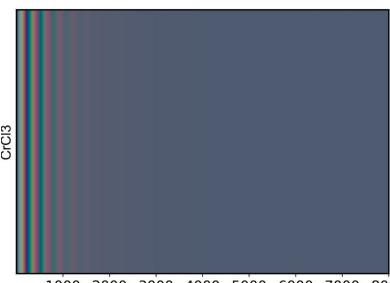
$$\delta N \cdot (kd + \theta) = m\pi$$

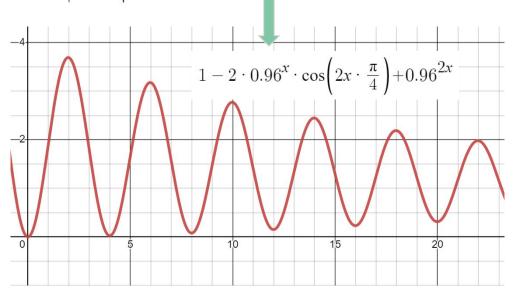
Then the periodicity factor will not change, which happens to be the most important part of the change. so we will see the same color as for the previous layers, but the overall brightness is reduced due to more absorption.

$$N o \infty, \; \tilde{x}^N = T^N e^{2iN(kd+\theta)} o \infty$$
 $R_\infty(\lambda) = \left| \frac{\tilde{r}}{1-\tilde{x}} \right|^2$

$$R_N = ilde{r}_{total} * ilde{r}_{total} = \left| rac{ ilde{r}}{1 - ilde{x}}
ight|^2 \cdot |1 - ilde{x}^N|^2$$

$$R_N = \left|rac{ ilde{r}}{1- ilde{x}}
ight|^2 \cdot \left(1-2 ext{Re}\left(ilde{x}^N
ight) + \left| ilde{x}^N
ight|^2
ight) = \left|rac{ ilde{r}}{1- ilde{x}}
ight|^2 \cdot \left[1-2T^N ext{Re}\left(e^{2iN(kd+ heta)}
ight) + T^{2N}
ight]$$





1000 2000 3000 4000 5000 6000 7000 8000 N (layers)

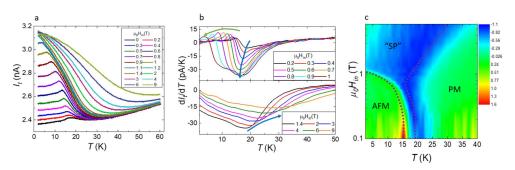
CrCl3 Color Mismatch

CrX3 has complex magnetic properties



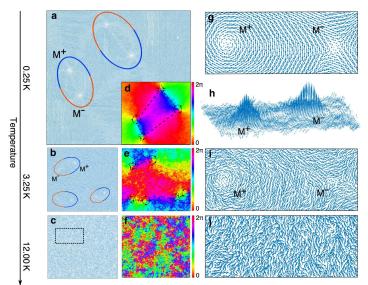
$$\sigma_{2D}(q=0,\omega) = i\omega\varepsilon_0(1-\varepsilon_{3D}(q=0,\omega))L$$

"The van der Waals magnetic insulator CrCl3 shows layered antiferromagnetism down to the bilayer limit with the magnetic moments in the plane of the layers and little or no anisotropy within the plane."



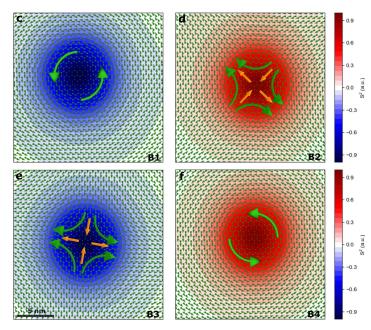
Nano Letters 2019 19 (6), 3993-3998

CrCl3 Color Mismatch



"Our studies show that the overall isotropic in-plane magnetized 2D material such as ML CrCl3 could provide the opportunity for generating meron-type topological defects without involving any other interactions."

Nat Commun 11, 4724 (2020)



"Their dynamics is determined by the interplay between the strong in-plane dipolar interactions and the weak out-of-plane magnetic anisotropy stabilising a vortex core within a radius of 8–10 nm."

Nat Commun 12, 185 (2021)

Outlook for future work

1) Comparison with other methods of measuring the number of layers by optical properties (Li, et al. (2017). Advanced Functional Materials.)



- 2) Consider the effect of substrate reflection
- Comparison with results from membrane calculation software (like TFCalc)
- 4) Consider different types of light sources and matching functions(different races; microscope)
- 5) Different stacking methods for 2D materials

