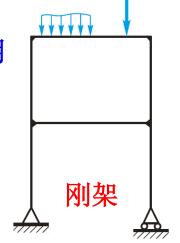
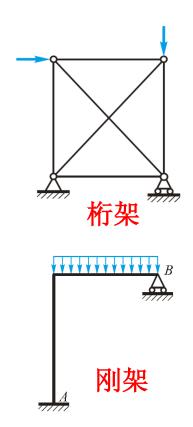


§ 14.1 超静定结构概述

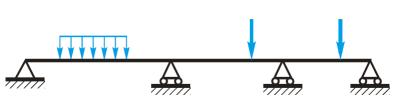
由直杆以铰节点相连接组成杆系, 若载 荷只作用于节点上,则每一杆件只承受拉伸或压缩,这种杆系称为桁架。

直杆以刚节点相连接组成杆系,在载荷作用下,各杆件可以承受拉、压、弯曲和扭转,这样的杆系称为刚架。





具有三个或三个以上支座的梁,通常称为连续梁。





时,外力也都作用于这一平面内,这种杆系称为平面杆系。

杆系各杆件的轴线在同一平面内,且它就是各杆件的形心主惯性平面,同

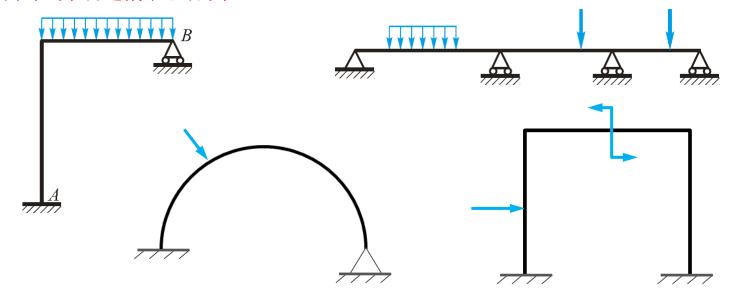




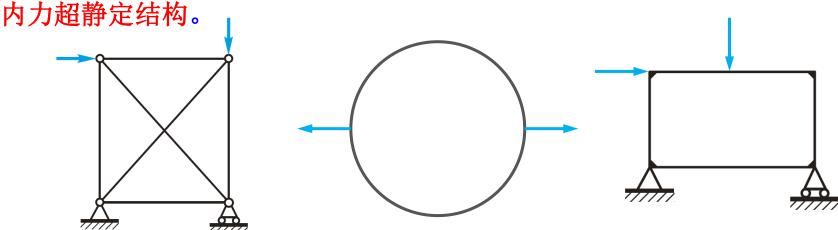
超静定结构:存在一些不是维持几何不变所必需的支座(约束)的结构 把这类约束称为多余约束,与多余约束对应的约束力称为多余约束力 超静定问题:平衡方程数*m*<未知数目*n*

超静定结构的形式:

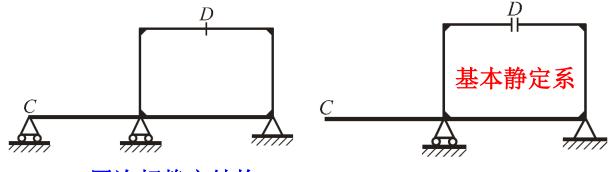
1. 超静定结构的支座约束力不能全由平衡方程求出,这种超静定结构, 称为外力超静定结构。



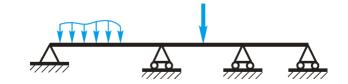
2. 超静定结构的内力不能全由平衡方程求出,这种超静定结构,称为



3. 既有外力超静定,又有内力超静定的结构。



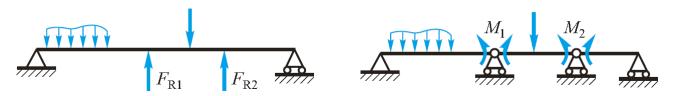
原超静定结构



解除超静定结构的某些约束后得到的静定结构, 称为原超静定结构的 基本静定系(◀: 基本静定系可以有不同的选择,不是唯一的)。



在基本静定系上,除原有载荷外,还有用相应的多余约束力代替被解除的多余约束。把载荷和多余约束力作用下的基本静定系称为相当系统。



求解静定结构位移的能量方法:

卡氏第一定理
$$F_i = \frac{\partial V_{\varepsilon}(\Delta_1, \Delta_2 \cdots \Delta_n)}{\partial \Delta_i}$$

卡氏第二定理
$$\Delta_i = \frac{\partial V_{\varepsilon}(F_1, F_2 \cdots F_n)}{\partial F_i}$$

余能定理
$$\Delta_i = \frac{\partial V_c(F_1, F_2 \cdots F_n)}{\partial F_i}$$

单位载荷法
$$1 \cdot \Delta = \int_{I} (\overline{M} d\theta^{*} + \overline{F}_{S} d\lambda^{*} + \overline{F}_{N} d(\Delta I)^{*} + \overline{T} d\varphi^{*})$$
$$1 \cdot \Delta = \int_{I} (\overline{M} \frac{M}{EI} dx + \overline{F}_{S} \cdot \frac{\alpha_{s} F_{S}}{GA} dx + \overline{F}_{N} \frac{F_{N}}{EA} dx + \overline{T} \frac{T}{GI_{P}} dx)$$

§ 14.2 用能量法解超静定系统

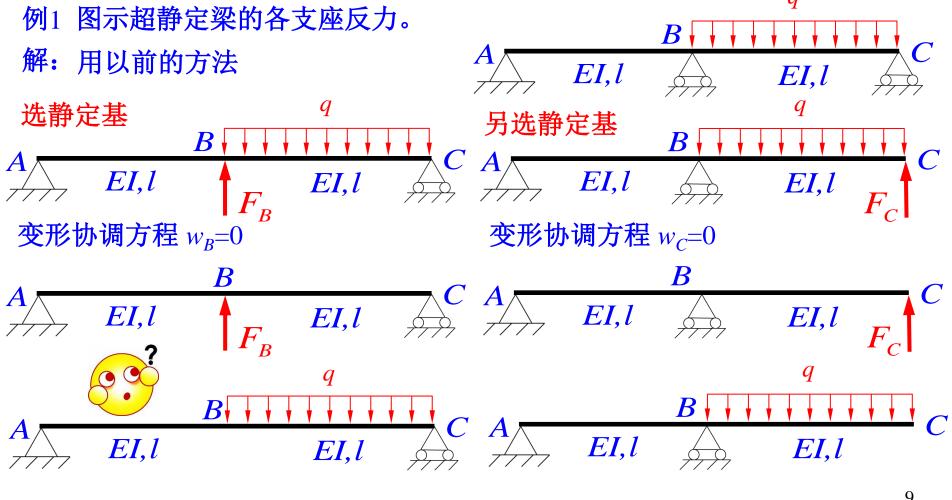
求解超静定问题的基本方法 (平衡方程数m<未知数目n)

- (1) 选取基本静定系
- (2) 变形几何相容方程 (compatibility equations)
- (3) 本构方程 (constitutive equations)

用能量法来建立! (不同之处) 建立p个补充方程

平衡方程(m个)

(4)原问题变成静定问题,求解出所有的约束反力,进而完成 内力、应力、位移的计算。



解:用能量方法

选静定基

建立补充方程

卡氏第二定理

$$\begin{split} \frac{\partial V_{\varepsilon}}{\partial F_{C}} &= 0 & M_{AB}(x) = F_{A}x = (F_{c} - \frac{1}{2}ql)x & M_{BC}(x) = F_{c}x - \frac{1}{2}qx^{2} \\ V_{\varepsilon} &= \int_{0}^{l} \frac{M_{BC}^{2}(x)}{2El} \mathrm{d}x + \int_{0}^{l} \frac{M_{AB}^{2}(x)}{2El} \mathrm{d}x \\ \frac{\partial V_{\varepsilon}}{\partial F_{C}} &= \int_{0}^{l} \frac{M_{BC}(x)}{El} \cdot \frac{\partial M_{BC}(x)}{\partial F_{C}} \mathrm{d}x + \int_{0}^{l} \frac{M_{AB}(x)}{El} \cdot \frac{\partial M_{AB}(x)}{\partial F_{C}} \mathrm{d}x \\ &= \int_{0}^{l} \frac{(F_{C}x - \frac{1}{2}qx^{2})}{El} \cdot x \mathrm{d}x + \int_{0}^{l} \frac{(F_{C} - \frac{1}{2}ql)x}{El} \cdot x \mathrm{d}x \end{split} \qquad F_{B} = \frac{5}{8}ql \end{split}$$

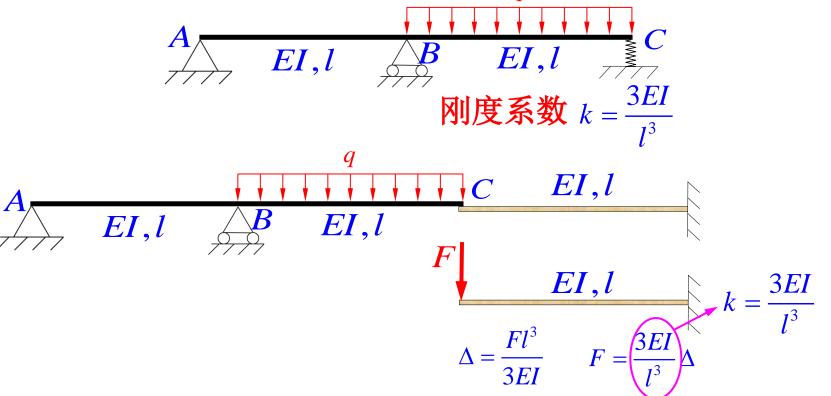
 $F_A = \frac{1}{2}(F_C l - \frac{1}{2}ql^2) = F_C - \frac{1}{2}ql \quad (\Sigma M_P = 0)$

去除B支座或 A支座均可!

$$= \int_{0}^{\frac{(4CN-24N)}{EI}} \cdot x dx + \int_{0}^{\frac{(4C-24N)N}{EI}} \cdot x dx \qquad F_{B} = \frac{5}{8} q l$$

$$= \frac{1}{EI} \left(\frac{1}{3} F_{C} l^{3} - \frac{1}{8} q l^{4} + \frac{1}{3} F_{C} l^{3} - \frac{1}{6} q l^{4} \right) = 0 \qquad F_{C} = \frac{7}{16} q l \qquad F_{A} = -\frac{1}{16} q l$$

若将超静定梁的C 支座改为刚度系数 $k = \frac{3EI}{l^3}$ 的弹簧,试求各 支座的反力。



解: 选静定基同前

利用能量法建立补充方程

卡氏第二定理

$$\frac{\partial V_{\varepsilon}}{\partial F_{C}} = \frac{F_{C}}{\Delta} \frac{F_{C}}{k}$$

$$V_{\varepsilon} = \int_{0}^{l} \frac{M_{BC}^{2}(x)}{2EI} dx + \int_{0}^{l} \frac{M_{AB}^{2}(x)}{2EI} dx \qquad M_{BC}(x)$$

$$= \int_{0}^{l} \frac{(F_{c}x - \frac{1}{2}qx^{2})^{2}}{2EI} dx + \int_{0}^{l} \frac{[(F_{c} - \frac{1}{2}ql)x]^{2}}{2EI} dx$$

$$\frac{\partial V_{\varepsilon}}{\partial F_{C}} = \int_{0}^{l} \frac{(F_{C}x - \frac{1}{2}qx^{2})}{EI} \cdot x dx + \int_{0}^{l} \frac{(F_{C} - \frac{1}{2}ql)x}{EI} \cdot x dx = -\frac{F_{C}}{k}$$

$$EI, l$$

$$EI, l$$

$$EI, l$$

$$EI, l$$

$$k = \frac{3EI}{l^3} F_C$$

$$M_{AB}(x) = F_A x = (F_c - \frac{1}{2}ql)x$$

$$M_{BC}(x) = F_c x - \frac{1}{2} q x^2$$

$$k = \frac{3EI}{l^3}$$

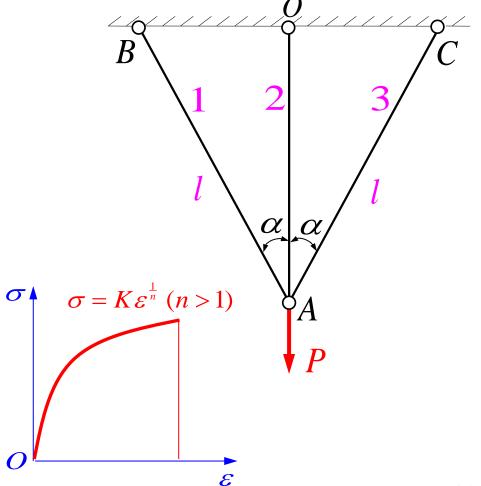
$$EI, l$$
 B EI, l C ∂V_{ε} $= \frac{F_C}{\Delta k}$ ∂F_C ∂

更一般化讨论: 设弹簧的刚度系数为 k $\frac{1}{EI} \left(\frac{1}{3} F_C l^3 - \frac{1}{8} q l^4 + \frac{1}{3} F_C l^3 - \frac{1}{6} q l^4 \right) = -\frac{F_C}{k} A$ $\frac{l^3}{FL}(\frac{2}{3}F_C - \frac{7}{24}ql) = -\frac{F_C}{k}$ $\tilde{F}_{C}(\frac{F_{C}}{al})$ $(1 + \frac{2}{3} \frac{kl^3}{EI}) F_C = \frac{7}{24} \frac{kl^3}{EI} q l$ $0.5_{\rm f}$ $F_C = \frac{\frac{7}{24} \frac{k l^3}{EI}}{1 + \frac{2}{3} \frac{k l^3}{EI}} q l$ 无量纲化可使讨论 k=0 时,F=0 静定问题 更广泛和深入! $k \to \infty$ 时, $F \to \frac{7}{16}ql$ 给出变化曲线? $\tilde{F}_C = \frac{\frac{7}{24}k}{1 + \frac{2}{3}\tilde{k}}$ $\tilde{F}_C = \frac{F_C}{al}$ $\tilde{k} = \frac{k}{FI/I^3}$ 无量纲化!

例2 求图示结构各杆的内力。 各杆的横截面面积为A。

考虑以下两种情形:

- 1. 三杆均为线弹性材料, 弹性模量为E;
- 2. 三杆均为非线性弹性材料, 有 $\sigma = K\varepsilon^{\frac{1}{n}}$ (n > 1) ;



解: (1) 线弹性问题

选静定基如图

$$2F_1 \cos \alpha + F_2 = P$$
 $F_1 = (P - F_2)/(2\cos \alpha)$

利用能量法建立补充方程

卡氏第二定理
$$\dfrac{\partial V_{\varepsilon}}{\partial F_2}=0$$

$$V_{\varepsilon} = \frac{F_2^2 l_2}{2EA} + \frac{F_1^2 l_1^2}{2EA} \times 2$$
 $l_2 = l \cos \alpha$

$$\frac{\partial V_{\varepsilon}}{\partial F_2} = \frac{F_2 l_2}{EA} + \frac{F_1 l_1}{EA} \times 2 \times \left(-\frac{1}{2\cos\alpha}\right) = 0$$

$$\frac{F_2 l \cos \alpha}{EA} - \frac{(P - F_2)l}{2EA \cos^2 \alpha} = 0 \qquad F_2 = \frac{P}{1 + 2 \cos^3 \alpha}$$

解: (2) 非线性弹性问题

选静定基同前

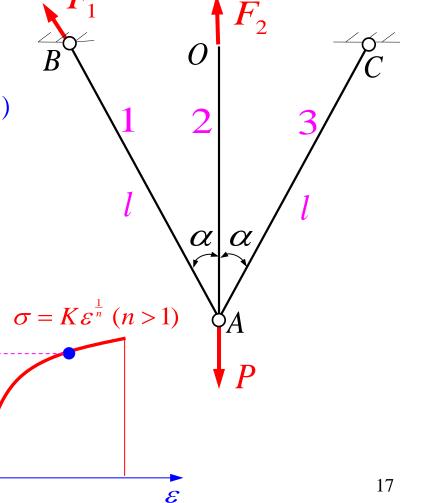
$$2F_1 \cos \alpha + F_2 = P$$
 $F_1 = (P - F_2)/(2\cos \alpha)$

利用能量法建立补充方程

余能定理
$$\frac{\partial V_C}{\partial F_2} = 0$$

 $V_C = \int_V v_C dV = v_{C2} \cdot Al_2 + v_{C1} \cdot Al_1 \times 2$ $v_C = \int_V c d\sigma = V_C \cdot Al_2 + v_{C1} \cdot Al_1 \times 2$

 $v_C = \int \mathcal{E} d\sigma \quad \sigma_2 = F_2/A$ $v_{c2} = \int_0^{\sigma_2} \left(\frac{\sigma}{K}\right)^n d\sigma = \frac{1}{n+1} \cdot \frac{(\sigma_2)^{n+1}}{K^n}$ $= \frac{1}{n+1} \cdot \frac{(F_2)^{n+1}}{K^n A^{n+1}}$



$$\begin{aligned} v_{c2} &= \int_{0}^{\sigma_{2}} (\frac{\sigma}{K})^{n} d\sigma = \frac{1}{n+1} \cdot \frac{(\sigma_{2})^{n+1}}{K^{n}} = \frac{1}{n+1} \cdot \frac{(F_{2})^{n+1}}{K^{n}A^{n+1}} \\ v_{c1} &= \int_{0}^{\sigma_{1}} (\frac{\sigma}{K})^{n} d\sigma = \frac{1}{n+1} \cdot \frac{(\sigma_{1})^{n+1}}{K^{n}} = \frac{1}{n+1} \cdot \frac{(F_{1})^{n+1}}{K^{n}A^{n+1}} \\ &= \frac{1}{n+1} \cdot \frac{(P-F_{2})^{n+1}}{K^{n}A^{n+1}(2\cos\alpha)^{n+1}} \qquad F_{1} = (P-F_{2})/(2\cos\alpha) \end{aligned}$$

$$V_{C} &= \int_{V} v_{C} dV = v_{C2} \cdot Al_{2} + v_{C1} \cdot Al_{1} \times 2$$

$$&= \frac{1}{n+1} \cdot \frac{(F_{2})^{n+1}}{K^{n}A^{n+1}} \cdot Al_{2} + \frac{1}{n+1} \cdot \frac{(P-F_{2})^{n+1}}{K^{n}A^{n+1}(2\cos\alpha)^{n+1}} \cdot Al_{1} \times 2 \qquad l_{2} = l\cos\alpha$$

$$\frac{\partial V_{C}}{\partial F_{2}} &= \frac{(F_{2})^{n}}{K^{n}A^{n+1}} \cdot Al\cos\alpha + \frac{(P-F_{2})^{n}}{K^{n}A^{n+1}(2\cos\alpha)^{n+1}} \times (-1) \cdot Al \times 2 = 0$$

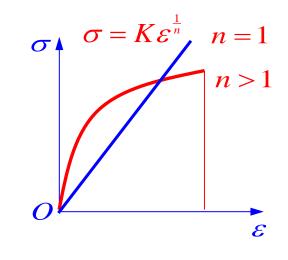
$$\frac{\partial V_C}{\partial F_2} = \frac{(F_2)^n}{K^n A^{n+1}} \cdot Al \cos \alpha + \frac{(P - F_2)^n}{K^n A^{n+1} (2\cos \alpha)^{n+1}} \times (-1) \cdot Al \times 2 = 0$$

$$(F_2)^n \cos \alpha = \frac{(P - F_2)^n}{2^n \cos^{n+1} \alpha}$$

$$(F_2)^n 2^n \cos^{n+2} \alpha = (P - F_2)^n$$

$$F_2 \cdot 2(\cos \alpha)^{\frac{n+2}{n}} = P - F_2$$

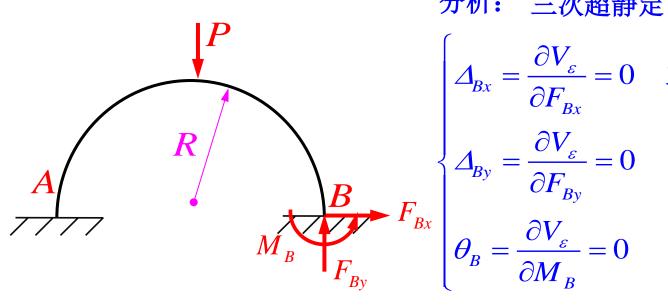
$$F_2 = \frac{P}{1 + 2(\cos\alpha)^{\frac{n+2}{n}}}$$



当
$$n=1$$
 (线弹性情形): 结果可退化到情形 (1)

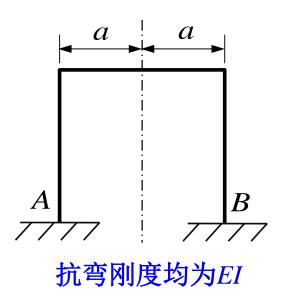
$$F_2 = \frac{P}{1 + 2\cos^3\alpha}$$

两端固定的半圆环在对称截面处受一集中力P作用,圆环轴线半径为 R,弯曲刚度为EI。不计剪力和轴力对圆环变形的影响。试求对称截 面上的内力。

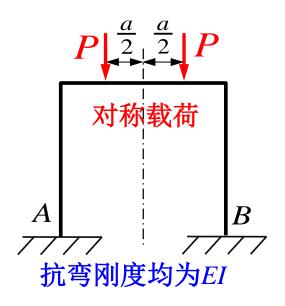


三次超静定问题

§ 14.3 对称和反对称性质的利用

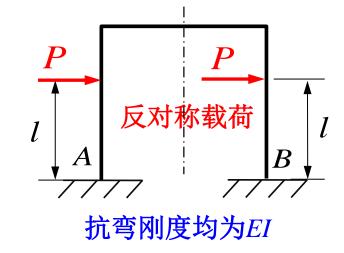


结构的几何形状、支承条件和杆件的刚度都对称于某一对称轴,称为对称结构!



1. 若载荷的作用位置、大小和方向都对称于结构的对称轴,称为对称载荷! (即绕对称轴对折后,左右两部分的载荷重合(作用点重合,大小相等,方向相同)

2. 若载荷的作用位置、大小对称于结构的对称 轴,但方向是反对称的,称为反对称载荷! (即绕对称轴对折后,左右两部分的载荷作 用点重合,大小相等,方向相反)



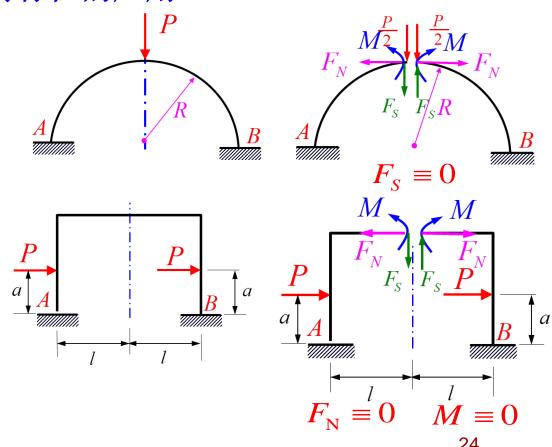
对称截面上的内力分类

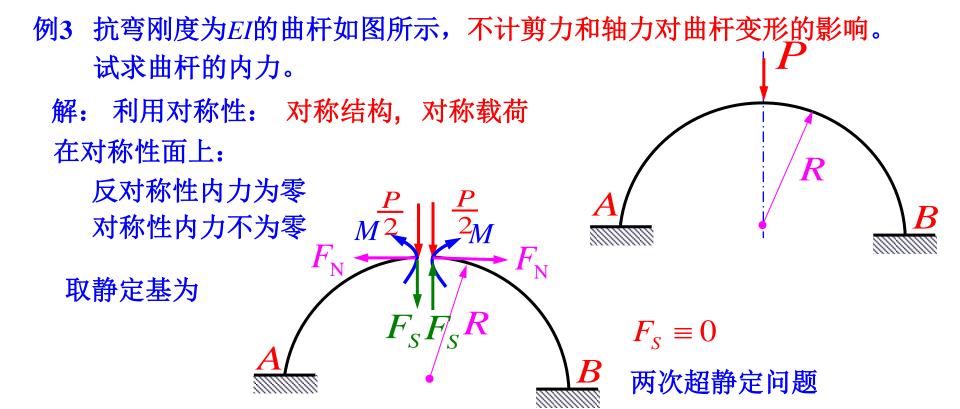
对称内力: 轴力 $F_{\rm N}$, 弯矩M 反对称内力: 剪力 F_s , 扭矩T轴力 F_{N} 弯矩M 对称 剪力 F_s 反对称 扭矩T 反对称

前提条件:结构是对称的 在对称性面上内力的特征:

- 1. 受对称载荷作用 反对称性内力为零 对称性内力不为零
- 2. 受反对称载荷作用 对称性内力为零 反对称性内力不为零

对称性的应用





利用卡氏第二定理可写出

$$\frac{\partial V_{\varepsilon}}{\partial M} = 0 \qquad \frac{\partial V_{\varepsilon}}{\partial F_{N}} = 0$$

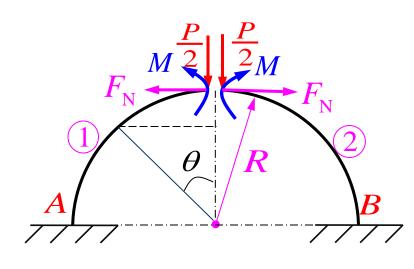
上面两式的物理意义: 在对称面处, 轴向相 对位移为零;两侧截面的相对转角为零。

不计剪力和轴力对圆环变形的影响

$$V_{\varepsilon} = \int \frac{M^{2}(\theta)}{2EI} R d\theta = V_{\varepsilon 1} + V_{\varepsilon 2} \qquad M_{1}(\theta) = M + F_{N}(R - R\cos\theta) - \frac{P}{2}R\sin\theta$$

$$V_{\varepsilon 1} = V_{\varepsilon 2} \qquad V_{\varepsilon}(\theta) = \int_{0}^{\frac{\pi}{2}} \frac{[M_{1}(\theta)]^{2}}{2EI} R d\theta \times 2 \qquad V_{\varepsilon 1}(\theta) = \int_{0}^{\frac{\pi}{2}} \frac{[M_{1}(\theta)]^{2}}{2EI} R d\theta$$

$$\frac{\partial V_{\varepsilon}(\theta)}{\partial M} = 2 \times \int_{0}^{\frac{\pi}{2}} \frac{M_{1}(\theta)}{EI} \cdot \frac{\partial M_{1}(\theta)}{\partial M} R d\theta = 0 \qquad \frac{\partial V_{\varepsilon}(\theta)}{\partial F_{N}} = 2 \times \int_{0}^{\frac{\pi}{2}} \frac{M_{1}(\theta)}{EI} \cdot \frac{\partial M_{1}(\theta)}{\partial F_{N}} R d\theta = 0$$



$$V_{\varepsilon 1}(\theta) = \int_0^{\frac{\pi}{2}} \frac{[M_1(\theta)]^2}{2EI} R d\theta$$

$$\frac{\partial V_{\varepsilon}(\theta)}{\partial F_{N}} = 2 \times \int_{0}^{\frac{\pi}{2}} \frac{M_{1}(\theta)}{EI} \cdot \frac{\partial M_{1}(\theta)}{\partial F_{N}} R d\theta = 0$$

$$\begin{split} & M_{1}(\theta) = \mathbf{M} + \mathbf{F_{N}}(R - R\cos\theta) - \frac{P}{2}R\sin\theta \\ & \frac{\partial V_{\varepsilon}(\theta)}{\partial M} = \mathbf{2} \times \int_{0}^{\frac{\pi}{2}} \frac{M_{1}(\theta)}{EI} \cdot \frac{\partial M_{1}(\theta)}{\partial M} R d\theta = 0 \\ & \frac{2}{EI} \times \int_{0}^{\frac{\pi}{2}} [M + F_{N}(R - R\cos\theta) - \frac{P}{2}R\sin\theta] \cdot 1 \cdot R d\theta = 0 \\ & (\frac{1}{2}\pi - 1)R \cdot \mathbf{F_{N}} + \frac{1}{2}\pi \cdot \mathbf{M} - \frac{P}{2}R = 0 \\ & (\frac{3}{4}\pi - 2)R \cdot \mathbf{F_{N}} + (\frac{1}{2}\pi - 1) \cdot \mathbf{M} - \frac{P}{4}R = 0 \end{split} \qquad \begin{aligned} F_{N} &= \frac{4 - \pi}{\pi^{2} - 8} P \\ M &= \frac{2(\pi - 3)}{\pi^{2} - 8} P R \end{aligned} \\ & \frac{2}{EI} \times \int_{0}^{\frac{\pi}{2}} [M + F_{N}(R - R\cos\theta) - \frac{P}{2}R\sin\theta] \cdot R(1 - \cos\theta) \cdot R d\theta = 0 \end{aligned}$$

$$\frac{\partial V_{\varepsilon}(\theta)}{\partial F_{N}} &= \mathbf{2} \times \int_{0}^{\frac{\pi}{2}} \frac{M_{1}(\theta)}{EI} \cdot \frac{\partial M_{1}(\theta)}{\partial F_{N}} R d\theta = 0 \end{aligned}$$

变形协调方程:
$$\frac{\partial V_{\varepsilon}}{\partial M} = 0$$
 $\frac{\partial V_{\varepsilon}}{\partial F_{N}} = 0$

物理意义: 在对称面处, 轴向相对位移为零; 两侧截面的相对转角为零。

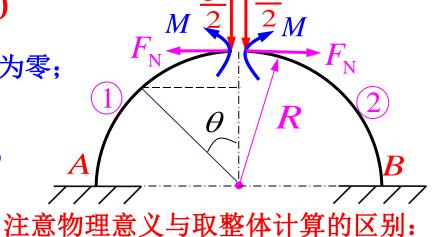
$$V_{\varepsilon 1} = V_{\varepsilon 2}$$
 $V_{\varepsilon 1}(\theta) = \int_0^{\frac{\pi}{2}} \frac{[M_1(\theta)]^2}{2EI} R d\theta$

$$V_{\varepsilon}(\theta) = \int_0^{\frac{\pi}{2}} \frac{[M_1(\theta)]^2}{2EI} R d\theta \times 2$$

$$\frac{\partial V_{\varepsilon}(\theta)}{\partial M} = 2 \times \int_{0}^{\frac{\pi}{2}} \frac{M_{1}(\theta)}{EI} \cdot \frac{\partial M_{1}(\theta)}{\partial M} R d\theta = 0$$

$$\frac{\partial V_{\varepsilon}(\theta)}{\partial F_{N}} = 2 \times \int_{0}^{\frac{\pi}{2}} \frac{M_{1}(\theta)}{EI} \cdot \frac{\partial M_{1}(\theta)}{\partial F_{N}} R d\theta = 0$$

可以取结构的一半来计算!



在对称面处,轴向位移为零;截面的 转角为零。

$$\frac{\partial V_{\varepsilon}(\theta)}{\partial M} = \int_{0}^{\frac{\pi}{2}} \frac{M_{1}(\theta)}{EI} \cdot \frac{\partial M_{1}(\theta)}{\partial M} Rd\theta = 0$$

$$\frac{\partial V_{\varepsilon}(\theta)}{\partial F_{N}} = \int_{0}^{\frac{\pi}{2}} \frac{M_{1}(\theta)}{EI} \cdot \frac{\partial M_{1}(\theta)}{\partial F_{N}} Rd\theta = 0$$

例4 材料为线弹性,抗弯刚度为*EI*的刚架如图所示,不计剪力和轴力对 刚架变形的影响。试求刚架的支座约束反力。

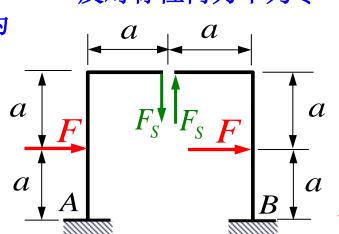
解: 利用对称性:

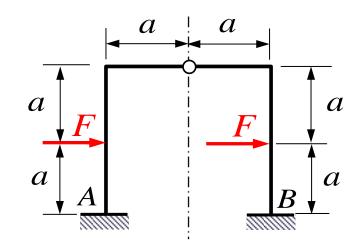
对称结构, 反对称载荷

在对称性面上: 对称性内力为零

反对称性内力不为零

取静定基为





 $QF_s \neq 0$ 一次超静定问题

利用卡氏第二定理,有
$$\frac{\partial V_{\varepsilon}}{\partial F_{S}} = 0$$

$$V_{\varepsilon} = 2 \times \left[\int_{0}^{a} \frac{M_{DE}^{2}(x)}{2EI} dx + \int_{0}^{a} \frac{M_{CD}^{2}(x)}{2EI} dx + \int_{0}^{a} \frac{M_{AC}^{2}(x)}{2EI} dx \right] = 2 \times \left[\int_{0}^{a} \frac{(F_{S}x)^{2}}{2EI} dx + \int_{0}^{a} \frac{(F_{S}a)^{2}}{2EI} dx + \int_{0}^{a} \frac{(F_{S}a+F_{X})^{2}}{2EI} dx \right] = 2 \times \left[\int_{0}^{a} \frac{(F_{S}x)^{2}}{2EI} dx + \int_{0}^{a} \frac{(F_{S}a)^{2}}{2EI} dx + \int_{0}^{a} \frac{(F_{S}a+F_{X})^{2}}{2EI} dx \right] = 2 \times \left[\int_{0}^{a} \frac{(F_{S}x)^{2}x}{EI} dx + \int_{0}^{a} \frac{(F_{S}a)^{2}x}{EI} dx + \int_{0}^{a} \frac{(F_{S}a+F_{X})^{2}x}{EI} dx \right] = 0$$

$$\frac{\partial V_{\varepsilon}}{\partial F_{S}} = \frac{2}{EI} \times \left(F_{S} \frac{a^{3}}{3} + F_{S}a^{3} + F_{S}a^{3} + F_{B}a^{2} \right) = 0$$

$$F_{Ay} - F_{S} = 0 \Longrightarrow F_{Ax} = -F (\leftarrow)$$

$$F_{Ay} - F_{S} = 0 \Longrightarrow F_{Ay} = F_{S} = -\frac{3}{14} F (\downarrow)$$

$$M_{A} - Fa - F_{S}a = 0 \Longrightarrow M_{A} = Fa + F_{S}a = Fa - \frac{3}{14} Fa = \frac{11}{14} Fa$$

同样,B支座的约束反力也可求出!

若要求C点的水平位移?

$$w_{Cx} = \frac{\partial V_{\varepsilon}}{\partial F} \Big|_{F_s = -\frac{3}{14}F} V_{\varepsilon}$$
取结构的一半来计算!

计算V。时,把前述已计算出的剪力代入

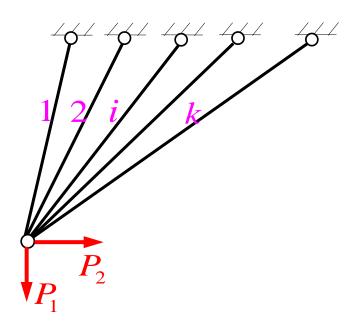
$$V_{\varepsilon} = \int_{0}^{a} \frac{(F_{S}x)^{2}}{2EI} dx + \int_{0}^{a} \frac{(F_{S}a)^{2}}{2EI} dx + \int_{0}^{a} \frac{(F_{S}a + Fx)^{2}}{2EI} dx$$

$$F_{S} = -\frac{3}{14}F$$

$$V_{\varepsilon} = \int_{0}^{a} \frac{\left(-\frac{3}{14}Fx\right)^{2}}{2EI} dx + \int_{0}^{a} \frac{\left(-\frac{3}{14}Fa\right)^{2}}{2EI} dx + \int_{0}^{a} \frac{\left(-\frac{3}{14}Fa + Fx\right)^{2}}{2EI} dx$$

$$w_{Cx} = \frac{\partial V_{\varepsilon}}{\partial F} = \int_0^a \frac{(-\frac{3}{14}Fx)\cdot(-\frac{3}{14}x)}{EI} dx + \int_0^a \frac{(-\frac{3}{14}Fa)\cdot(-\frac{3}{14}a)}{EI} dx + \int_0^a \frac{(-\frac{3}{14}Fa+Fx)\cdot(-\frac{3}{14}a+x)}{EI} dx$$
$$= \frac{F}{EI} \left[\left(\frac{3}{14} \right)^2 \times \frac{a^3}{3} + \left(\frac{3}{14} \right)^2 \times a^3 + \left(\frac{3}{14} \right)^2 \times a^3 - 2 \times \frac{3}{14} a \times \frac{a^2}{2} + \frac{a^3}{3} \right] = \frac{19Fa^3}{84EI}$$

例5 由 $k(k\geq 3)$ 根直杆组成的杆系,在节点A处用铰连接在一起,并受到水平荷载 P_1 和铅垂荷载 P_2 作用,如图所示。已知各杆的材料相同,拉压弹性模量均为E,横截面面积分别为 A_1 、 $A_2...A_k$,试求杆系中各杆的内力。



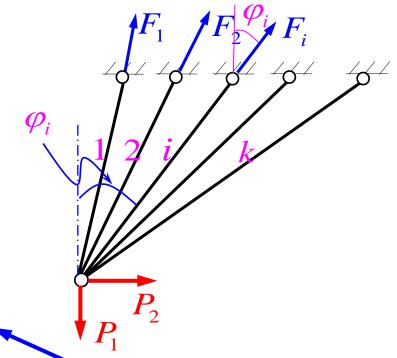
解(1): (k-2) 次超静定问题

利用卡氏第二定理可写出

静力平衡方程

$$\sum_{i=1}^{k} F_i \cos \varphi_i - P_1 = 0$$

$$\sum_{i=1}^{k} F_i \sin \varphi_i + P_2 = 0$$



$$F_{k-1} = F_{k-1}(F_1, F_2 \cdots F_{k-2})$$

$$F_k = F_k(F_1, F_2 \cdots F_{k-2})$$

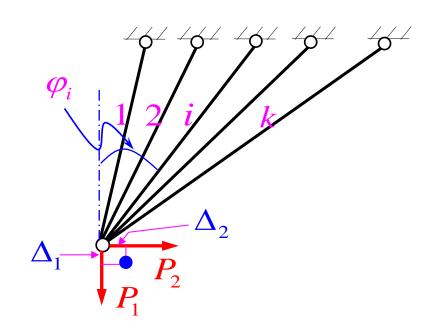
解(2):利用卡氏第一定理

$$V_{\varepsilon} = V_{\varepsilon}(\Delta_1, \Delta_2)$$

$$\frac{\partial V_{\varepsilon}(\Delta_1, \Delta_2)}{\partial \Delta_1} = P_1$$

$$\frac{\partial V_{\varepsilon}(\Delta_1, \Delta_2)}{\partial \Delta_2} = P_2$$

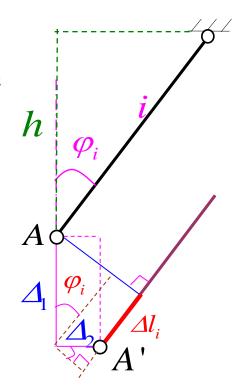
$$V_{\varepsilon}(\Delta_1, \Delta_2) = ?$$



$$\begin{split} V_{\varepsilon}(\Delta_{1}, \Delta_{2}) &= ? & \Delta l_{i} = \Delta_{1} \cos \varphi_{i} - \Delta_{2} \sin \varphi_{i} \\ V_{\varepsilon}(\Delta_{1}, \Delta_{2}) &= \sum_{i=1}^{k} \frac{EA_{i}}{2l_{i}} \left(\Delta l_{i}\right)^{2} = \sum_{i=1}^{k} \frac{EA_{i}}{2l_{i}} \left(\Delta_{1} \cos \varphi_{i} - \Delta_{2} \sin \varphi_{i}\right)^{2} \\ V_{\varepsilon}(\Delta_{1}, \Delta_{2}) &= \sum_{i=1}^{k} \frac{EA_{i}}{2} \cdot \frac{\cos \varphi_{i}}{h} \cdot \left(\Delta_{1} \cos \varphi_{i} - \Delta_{2} \sin \varphi_{i}\right)^{2} \\ \frac{\partial V_{\varepsilon}(\Delta_{1}, \Delta_{2})}{\partial \Delta_{1}} &= P_{1} & \frac{\partial V_{\varepsilon}(\Delta_{1}, \Delta_{2})}{\partial \Delta_{2}} = P_{2} \end{split}$$

$$\sum_{i=1}^{\frac{EA_i}{h}} \cos \varphi_i \cdot (\Delta_1 \cos \varphi_i - \Delta_2 \sin \varphi_i) \cdot \cos \varphi_i = P_1$$

$$\sum_{i=1}^{k} \frac{EA_i}{h} \cos \varphi_i \cdot (\Delta_1 \cos \varphi_i - \Delta_2 \sin \varphi_i) \cdot (-\sin \varphi_i) = P_2$$



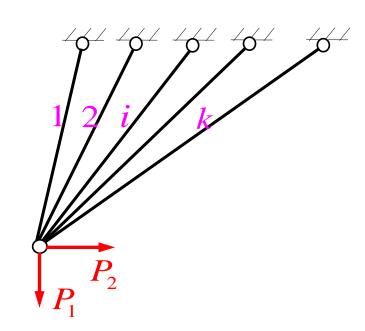
$$\begin{split} \sum_{i=1}^{k} \frac{EA_{i}}{h} \cos \varphi_{i} \cdot \left(\Delta_{1} \cos \varphi_{i} - \Delta_{2} \sin \varphi_{i}\right) \cdot \cos \varphi_{i} &= P_{1} \\ \sum_{i=1}^{k} \frac{EA_{i}}{h} \cos \varphi_{i} \cdot \left(\Delta_{1} \cos \varphi_{i} - \Delta_{2} \sin \varphi_{i}\right) \cdot \left(-\sin \varphi_{i}\right) &= P_{2} \\ \left(\sum_{i=1}^{k} \frac{EA_{i}}{h} \cos \varphi_{i} \cdot \left(\Delta_{1} \cos \varphi_{i} - \Delta_{2} \sin \varphi_{i}\right) \cdot \left(-\sin \varphi_{i}\right) &= P_{2} \\ \left(\sum_{i=1}^{k} \frac{EA_{i}}{h} \cos \varphi_{i} \cdot \left(\Delta_{1} \cos \varphi_{i} - \Delta_{2} \sin \varphi_{i}\right) \Delta_{1} - \left(\sum_{i=1}^{k} \frac{EA_{i}}{h} \cos \varphi_{i} \cdot \sin \varphi_{i}\right) \Delta_{2} &= P_{1} \\ -\left(\sum_{i=1}^{k} \frac{EA_{i}}{h} \cos \varphi_{i} \cdot \sin \varphi_{i}\right) \Delta_{1} + \left(\sum_{i=1}^{k} \frac{EA_{i}}{h} \cos \varphi_{i} \cdot \sin^{2} \varphi_{i}\right) \Delta_{2} &= P_{2} \\ C_{1}\Delta_{1} - C_{2}\Delta_{2} &= P_{1} \\ -C_{2}\Delta_{1} + C_{3}\Delta_{2} &= P_{2} \\ \Delta_{1} &= \frac{P_{1}C_{3} + P_{2}C_{2}}{C_{1}C_{3} - C_{2}^{2}} \quad \Delta_{2} &= \frac{P_{1}C_{2} + P_{2}C_{1}}{C_{1}C_{3} - C_{2}^{2}} \\ F_{Ni} &= EA \frac{\Delta I_{i}}{I_{i}} &= EA \frac{\cos \varphi_{i}}{h} \Delta I_{i} &= EA \frac{\cos \varphi_{i}}{h} \left(\Delta_{1} \cos \varphi_{i} - \Delta_{2} \sin \varphi_{i}\right) \end{split}$$

小 结

方法(1)中,以多余未知力作为基本未知量, 这种以力为基本未知量求解超静定问题的方法, 统称为力法。

方法(2)中,以未知的节点位移作为基本未知量,这种以位移为基本未知量求解超静定问题的方法,统称为位移法。

力法和位移法是求解超静定系统的两种基本方法。通常力法应用较为广泛。



工程实际的结构,超静定次数都是比较高的,对于超静定次数比较高的结构,教材中介绍了以下几种方法:

力法解超静定结构:

正则方程(§14.2)

三弯矩方程(§14.4多跨连续梁)

矩阵位移法: 教材第十七章

适用于采用编程计算!

关于力法和位移法的进一步讨论将在结构力学课程中详细介绍。





谢谢大家!

P159: 14.4(c)
P161: 14.12
P162: 14.15

对应第6版题号 P152: 14.4(c); P154-155: 14.12, 14.15

下次课介绍 动载荷和交变应力