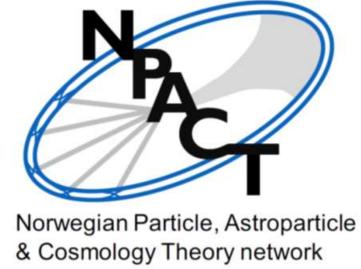
220 mode ringdown, Kerr deviation parameter from geometric deviations of final metric



Zaryab Ahmed



Introduction

- We are interested in the ringdown of black holes.
- We like to fit damped sinusoids (QNMs).
- Five parameters for each mode: Starting time, amplitude, phase, frequency, damping time.
- Perturb around a background BH (Kerr in GR), to find frequency and damping time (2 parameters -> 2 parameters in standard GR).
- Starting time, amplitude, phase determined by inspiral.

General Stationary, Axisymmetric Metrics have more than 2 parameters¹³

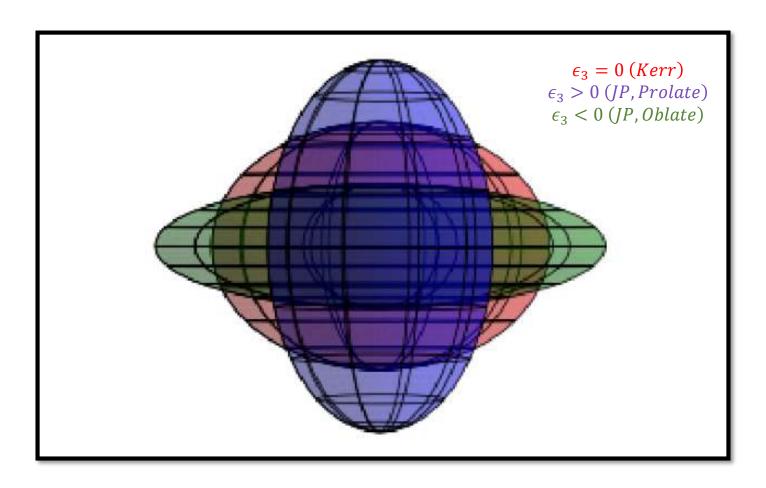
$$\begin{split} ds^2 &= -[1+h(r,\theta)](1-\frac{2Mr}{\Sigma})dt^2 - [1+h(r,\theta)]\frac{4aMr\sin^2\theta}{\Sigma}dtd\phi \\ &+ \frac{\Sigma[1+h(r,\theta)]}{\Delta+a^2\sin^2\theta h(r,\theta)}dr^2 + \Sigma d\theta^2 \\ &+ [\sin^2\theta(r^2+a^2+\frac{2a^2Mr\sin^2\theta}{\Sigma}) + h(r,\theta)\frac{a^2(\Sigma+2Mr)\sin^4\theta}{\Sigma}]d\phi^2 \\ \end{split} \\ &\text{where;} \quad h(r,\theta) &= \sum_{k=0}^{\infty} (\epsilon_{2k}+\epsilon_{2k+1}\frac{Mr}{\Sigma})(\frac{M^2}{\Sigma})^k \\ \end{split} \\ &g_{tr}^{JP} = -(1-\frac{2Mr}{\Sigma}) - \epsilon_3\frac{M^3(r-2M)}{r^4} \\ g_{rr}^{JP} &= \frac{\Sigma}{\Delta} + \epsilon_3\frac{M^3(r-2M)}{\Delta^2} \\ g_{\theta\theta}^{JP} &= \Sigma \\ g_{\theta\theta}^{JP} &= (r^2+a^2\frac{2Ma^2r\sin^2\theta}{\Sigma})\sin^2\theta + \epsilon_3\frac{a^2M^3(r+2M)}{r^3} \\ g_{t\phi}^{JP} &= -\frac{2Mar\sin^2\theta}{\Sigma} - \epsilon_3\frac{2aM^4}{r^4} \end{split}$$

- Non-GR (not a vacuum solution), parametric deviation, reduces to Kerr
- $\epsilon_0 = \epsilon_1 = 0$, for the metric to be asymptotically flat
- $\epsilon_2 = 4.6 \times 10^{-4}$, Lunar Laser Ranging experiment
- ϵ_3 , first unconstrained parameter $\rightarrow h(r,\theta) = \epsilon_3 \frac{M^3 r}{\Sigma^2}$
- Real and imaginary parts of the QNM spectrum for equatorial orbits in Johanssen-Psaltis geometry up to linear order in ϵ_3 :

$$\omega_{R}^{JP} = \omega_{R}^{K} + \epsilon_{3} \left(\frac{1}{81\sqrt{3}M} + \frac{10}{729M} \chi + \frac{47}{1458\sqrt{3}M} \chi^{2} \right)$$

$$\omega_{I}^{JP} = \omega_{I}^{K} - \epsilon_{3} \left(\frac{1}{486M} \chi + \frac{16}{2187\sqrt{3}M} \chi^{2} \right)$$
(2 params -> 3 params)

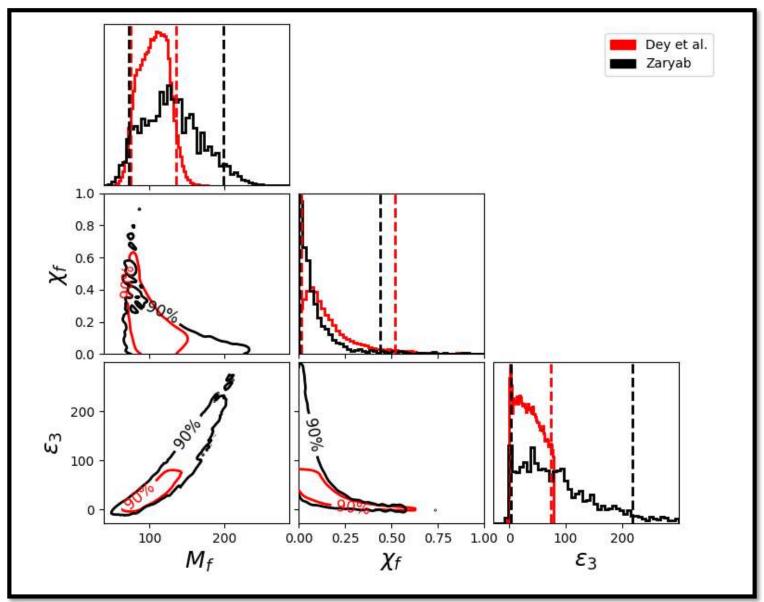
{1} A metric for rapidly spinning black holes suitable for strong-field test of the no-hair theorem https://link.aps.org/doi/10.1103/PhysRevD.83.124015



A visual representation of how ϵ_3 affects the ergospheres of Kerr and JP spacetimes.

Posterior Distribution GW150914 220 mode,

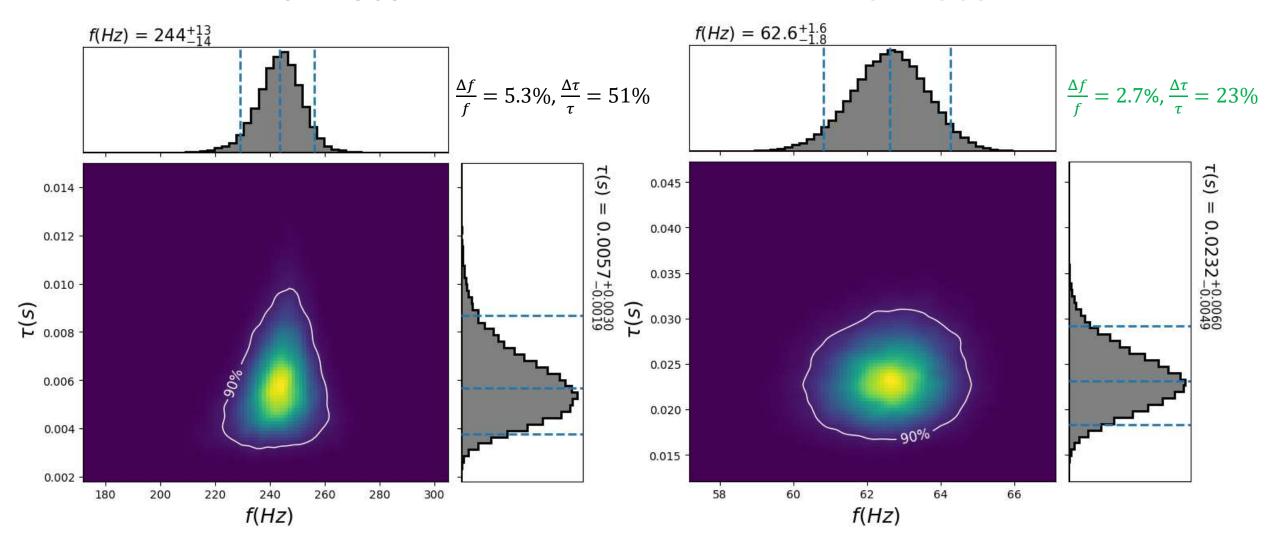
 $(\epsilon_3 = [-30,300], M=[20,300])$

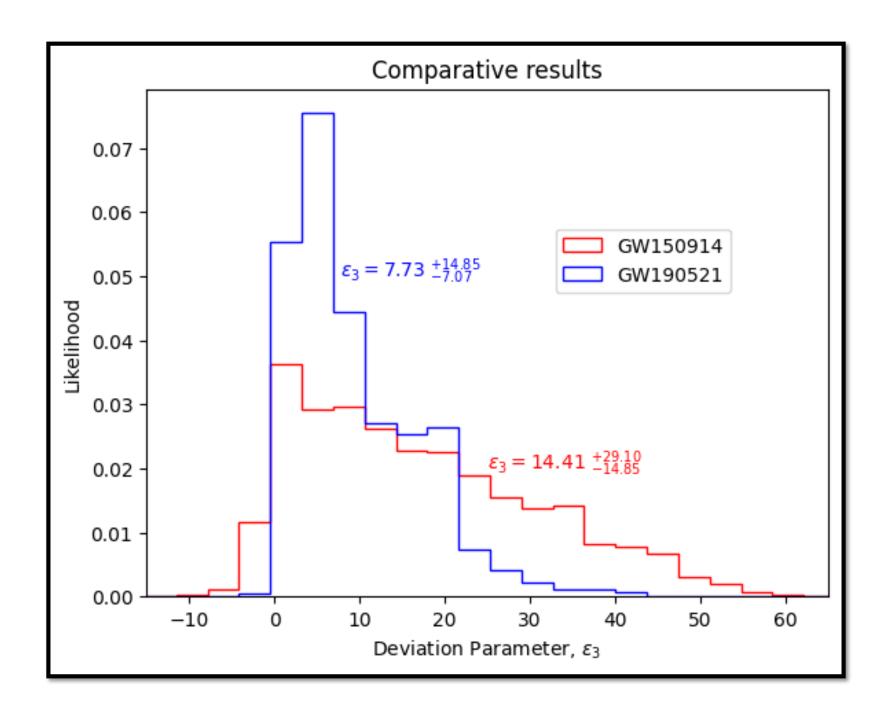


Frequency-Tau Distribution

• GW150914

• GW190521

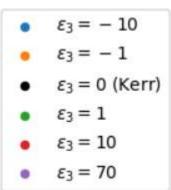




Conclusions

- Ringdown is sensitive to post-merger metric.
- Spin is anti-correlated with ϵ_3 in the ringdown.
- GW190521 is a good candidate for ringdown analysis.
- A single ringdown mode is unlikely to be competitive with the inspiral.
- $\omega_R^{330} = \frac{3}{2}\omega_R^{220}$, in the eikonal-limit.
- Use multimode observation in GW190521 ©

THANK YOU!



-1.00 -0.75 -0.50 -0.25 0.00

dim.less spin, a/M

0.000 -

-0.025 -

-0.050

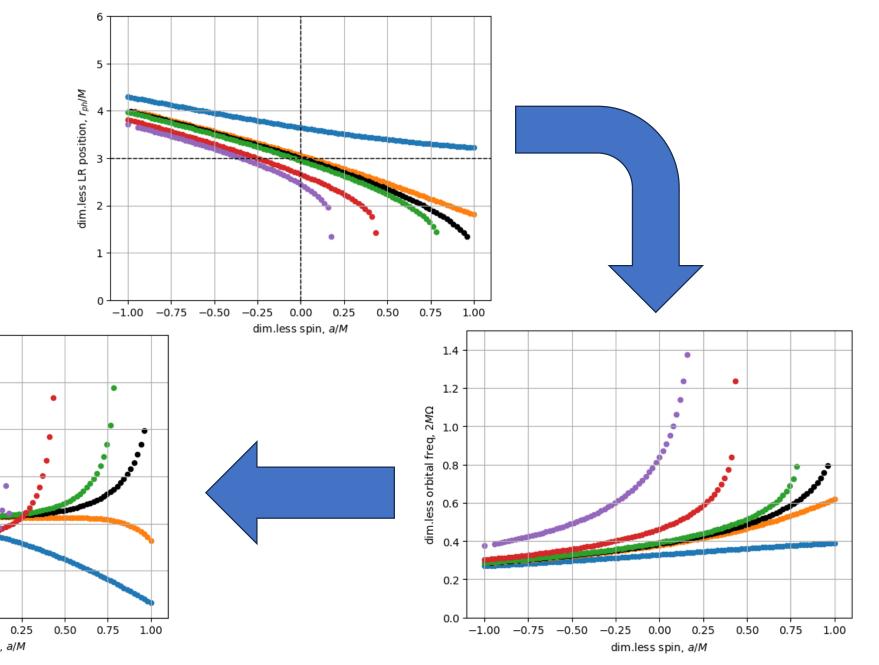
-0.075

-0.100

-0.125 -

-0.150 ⊥

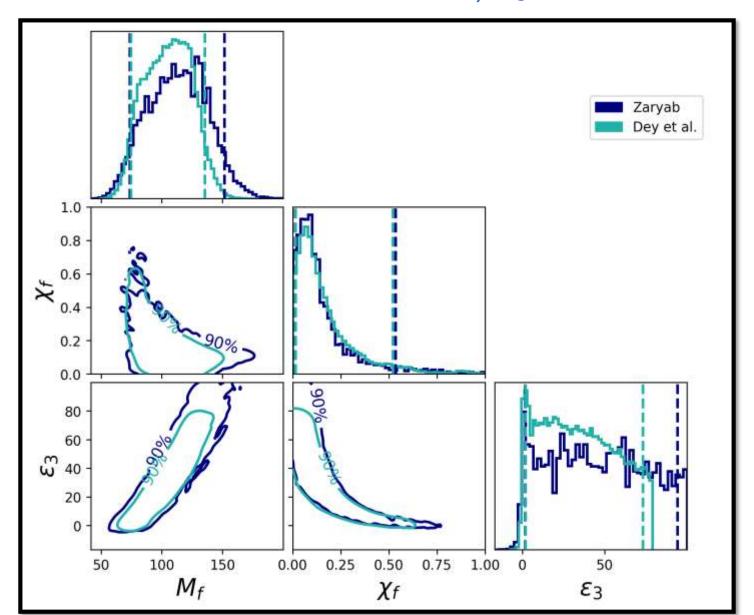
Lyapunov exp,



Posterior Distribution GW150914, $(\epsilon_3 = [-30,100], M=[20,200])$

Dey et al. priors

- $\epsilon_3 = [-30,100]$
- $\chi = [0,0.99]$
- M=[20,200]



Post-Kerr QNM Validity

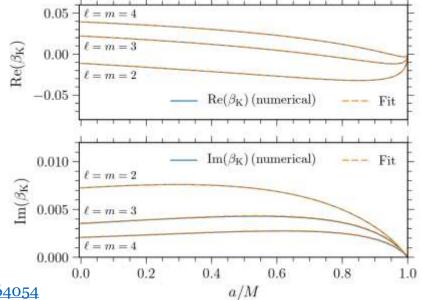
- QNM spectrum used for GW analysis based on Kerr metric
- Ensure that it is valid for a non-Kerr geometry
- Introduce an offset function, β_K

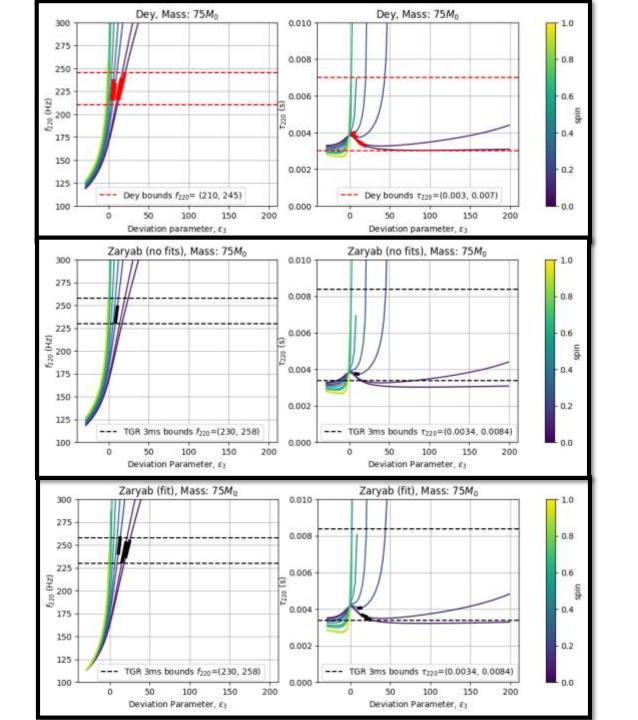
$$\omega_{\rm obs} = \sigma + \beta_{\rm K}$$

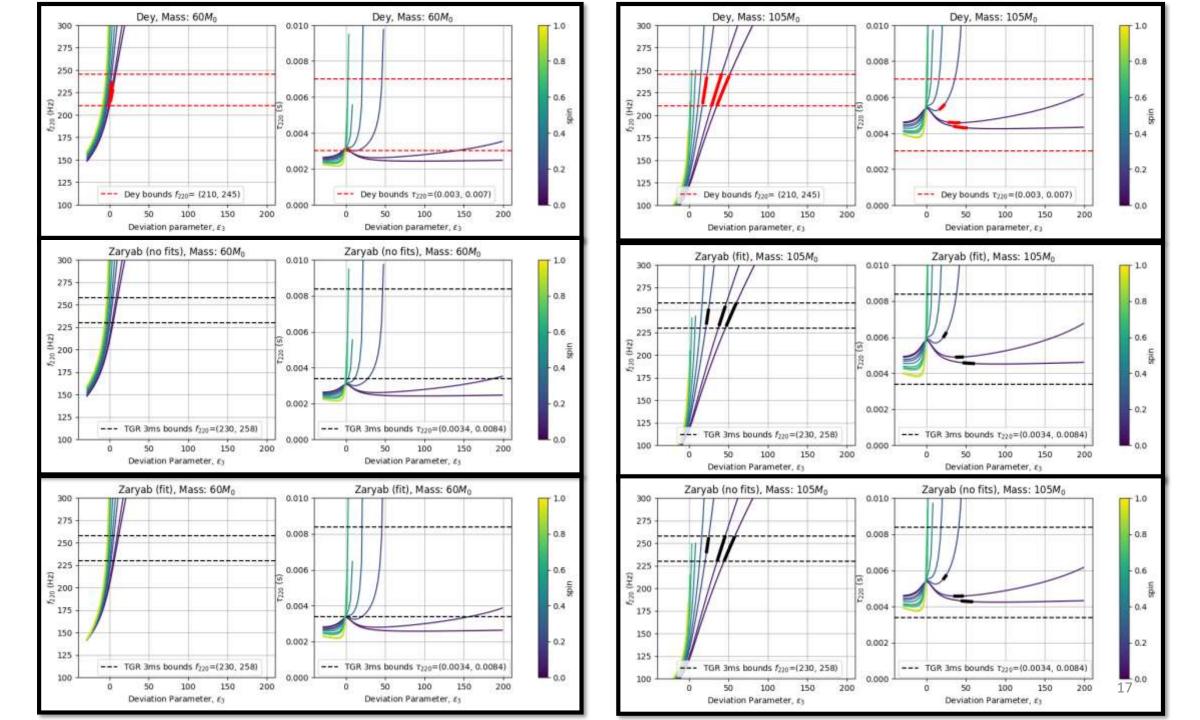
$$\omega_{\rm K} = \sigma_{\rm K} + \beta_{\rm K}$$

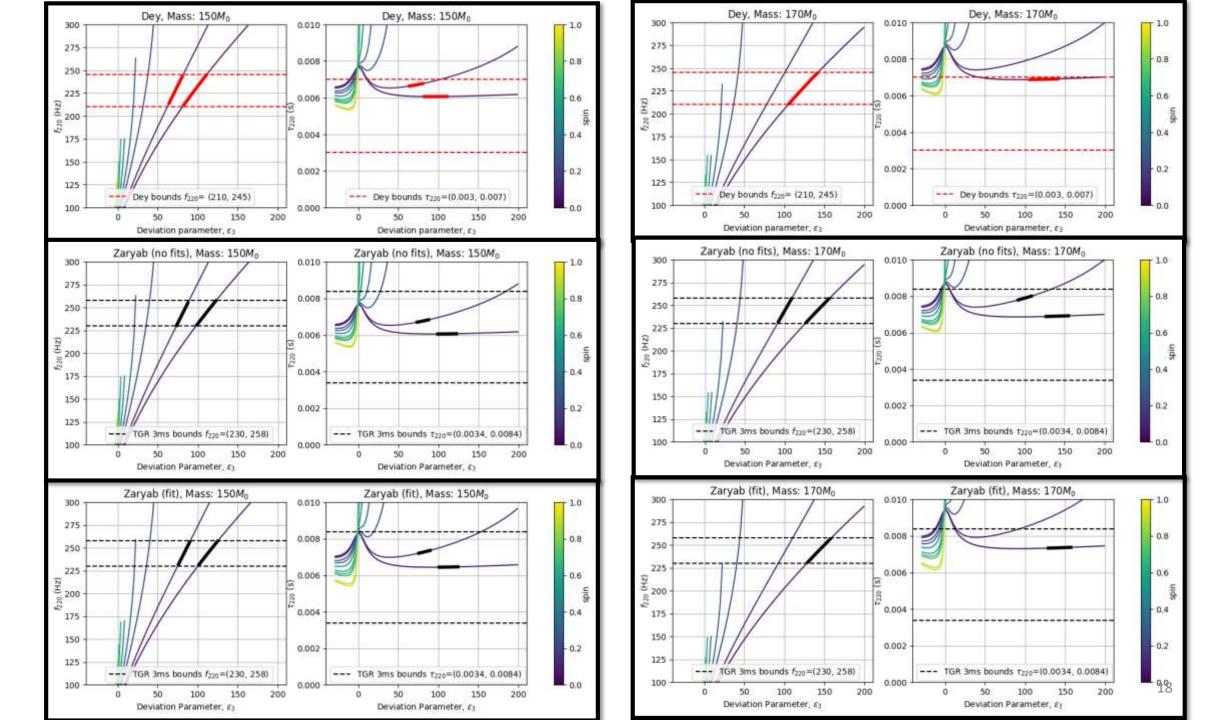
$$\omega_{\rm obs} - \omega_{\rm K} = \sigma - \sigma_{\rm K} \neq 0$$

$$\sigma = \sigma_{\rm K} + \delta \sigma$$









Extras

