# Asymptotic analysis of piezoelectric energy harvester

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## 1 Summary of the interested equations

Here we are interested in the classical model of a piezoelectric cantilever beam energy harvester, whose model is described using the following set of equations:

$$u'''' - \lambda^2 u = 0, \tag{1}$$

and the accompanying boundary conditions:

$$\begin{cases} u(0) = 0 \\ u'(0) = 0 \\ u''(1) + \dfrac{j\lambda\beta\alpha^2}{j\lambda\beta + 1} u'(1) = 0 \\ u'''(1) = 0 \end{cases}, \tag{2}$$

where $\lambda$ is the eigenvalues for the problem, $u$ denotes the displace function of the cantilever beam, $\beta$ is the dimensionless externally connected resistance, and $\alpha$ is the dimensionless piezoelectric coefficient. They can be expressed as follows

$$\lambda = \omega\sqrt{\frac{m_p l_p^4}{B_p}}, \quad \beta = R_l C_p \sqrt{\frac{B_p}{m_p l_p^4}}, \quad \alpha = e_p \sqrt{\frac{l_p}{C_p B_p}}, \tag{3}$$

where $\omega$ is angular frequency, $m_p$ is line mass density, $l_p$ is the length of the cantilever beam, $B_p$ is the bending stiffness, $C_p$ is the inherent capacitance of the piezoelectric layer, $e_p$ is the charge accumulation number, $R_l$ is the externally connected resistance. In practical applications, dielectric property of piezoelectric materials indicate that the parameter $\beta$ is changed from a very small value, which is close to a short-circuit condition to a very large value, which corresponds to an open-circuit condition. Thus we have that $0 \leq \beta \leq \infty$.

## 2 Asymptotic analysis when $\beta$ is small

Here we seek to find the behavior of the above system at a small value of connected resistance, i.e., $\beta \to 0$. In this case, we set $\beta$ to be the parameter for asymptotic expansion, and

$$\begin{aligned} \lambda^{(k)} &= \lambda_0^{(k)} + \beta\lambda_1^{(k)} + \beta^2\lambda_2^{(k)} + \cdots \\ u^{(k)} &= u_0^{(k)} + \beta u_1^{(k)} + \beta^2 u_2^{(k)} + \cdots \end{aligned} \tag{4}$$

where $\lambda^{(k)}$ and $u^{(k)}$ are the $k$th eigenvalue and eigenfunction respectively of the above mentioned system under perturbation. $\lambda_0^{(k)}$ and $u_0^{(k)}$ are the corresponding eigenvalue and eigenfunction of the unperturbed system at $\beta = 0$:

$$u'''' - \lambda_0^2 u = 0, \tag{5}$$

$$\begin{cases} u(0) = 0 \\ u'(0) = 0 \\ u''(1) = 0 \\ u'''(1) = 0 \end{cases}. \tag{6}$$

Obviously, the unperturbed system is a classical eigenvalue problem with the eigenvalues determined by

$$1+\cosh(\sqrt{\lambda_0})\cos(\sqrt{\lambda_0})=0 \tag{7}$$

whose first several values are

$$\frac{\sqrt{\lambda_0^{(1)}}}{\pi}=0.59686,\quad \frac{\sqrt{\lambda_0^{(2)}}}{\pi}=1.49418,\quad \frac{\sqrt{\lambda_0^{(3)}}}{\pi}=2.50025,\quad \frac{\sqrt{\lambda_0^{(4)}}}{\pi}=3.49999,\quad \cdots \tag{8}$$

Take the asymptotic expansions and substitute them into the previously derived system of equations, we have the following asymptotic expansions to different orders of $\beta$:
$O(\beta^0)$:

$$\begin{cases} u_0'''' - \lambda_0^2 u_0 = 0 \\ u_0(0)=0 \\ u_0'(0)=0 \\ u_0''(1)=0 \\ u_0'''(1)=0 \end{cases} \tag{9}$$

$O(\beta^1)$:

$$\begin{cases} u_1'''' - \left(\lambda_0^2 u_1 + 2\lambda_0 u_0 \lambda_1\right) = 0 \\ u_1(0)=0 \\ u_1'(0)=0 \\ u_1''(1)+j\alpha^2\lambda_0 u_0'(1)=0 \\ u_1'''(1)=0 \end{cases} \tag{10}$$

$O(\beta^2)$:

$$\begin{cases} u_2'''' - \left(\lambda_0^2 u_2 + 2\lambda_0 u_1 \lambda_1 + \lambda_1^2 u_0 + 2\lambda_0 u_0 \lambda_2\right) = 0 \\ u_2(0)=0 \\ u_2'(0)=0 \\ u_2''(1)+\alpha^2\lambda_0 u_0'(1)+j\alpha^2\left[\lambda_0 u_1'(1)+\lambda_1 u_0'(1)\right]=0 \\ u_2'''(1)=0 \end{cases} \tag{11}$$

# 3 Asymptotic analysis when $\beta$ is large

Here we seek to find the behavior of the above system at a large value of connected resistance, i.e., $\beta \to \infty$. In this case, we set $\frac{1}{\beta}$ to be the parameter for asymptotic expansion and

$$\begin{aligned} \lambda^{(k)} &= \tilde{\lambda}_0^{(k)} + \left(\frac{1}{\beta}\right)\tilde{\lambda}_1^{(k)} + \left(\frac{1}{\beta}\right)^2\tilde{\lambda}_2^{(k)} + \cdots \\ u^{(k)} &= \tilde{u}_0^{(k)} + \left(\frac{1}{\beta}\right)\tilde{u}_1^{(k)} + \left(\frac{1}{\beta}\right)^2\tilde{u}_2^{(k)} + \cdots \end{aligned} \tag{12}$$

where $\tilde{\lambda}^{(k)}$ and $\tilde{u}^{(k)}$ are the $k$th eigenvalue and eigenfunction respectively of the above mentioned system under perturbation. $\tilde{\lambda}_0^{(k)}$ and $\tilde{u}_0^{(k)}$ are the corresponding eigenvalue and eigenfunction of the unperturbed system at $\beta=\infty$:
$O(\frac{1}{\beta^0})$:

$$\begin{cases} \tilde{u}_0'''' - \tilde{\lambda}_0^2\tilde{u}_0 = 0 \\ \tilde{u}_0(0)=0 \\ \tilde{u}_0'(0)=0 \\ \tilde{u}_0''(1)+\alpha^2\tilde{u}_0'(1)=0 \\ \tilde{u}_0'''(1)=0 \end{cases} \tag{13}$$

$O(\frac{1}{\beta^1})$:

$$\begin{cases} \tilde{u}_1'''' - \left(\tilde{\lambda}_0^2 u_1 + 2\tilde{\lambda}_0 \tilde{u}_0 \tilde{\lambda}_1\right) = 0 \\ \tilde{u}_1(0) = 0 \\ \tilde{u}_1'(0) = 0 \\ \tilde{u}_1''(1) + \alpha^2 \tilde{u}_1'(1) + \dfrac{j\alpha^2}{\tilde{\lambda}_0} \tilde{u}_0'(1) = 0 \\ \tilde{u}_1'''(1) = 0 \end{cases} \tag{14}$$

$O(\frac{1}{\beta^2})$:

$$\begin{cases} \tilde{u}_2'''' - \left(\tilde{\lambda}_0^2 \tilde{u}_2 + 2\tilde{\lambda}_0 \tilde{u}_1 \tilde{\lambda}_1 + \tilde{\lambda}_1^2 \tilde{u}_0 + 2\tilde{\lambda}_0 \tilde{u}_0 \tilde{\lambda}_2\right) = 0 \\ \tilde{u}_2(0) = 0 \\ \tilde{u}_2'(0) = 0 \\ \tilde{u}_2''(1) + \left[\alpha^2 \tilde{u}_2'(1) - \dfrac{\alpha^2}{\tilde{\lambda}_0^2} \tilde{u}_0'(1)\right] + j \left[\dfrac{\alpha^2}{\tilde{\lambda}_0} \tilde{u}_1'(1) - \dfrac{\alpha^2 \tilde{\lambda}_1}{\tilde{\lambda}_0^2} \tilde{u}_0'(1)\right] = 0 \\ \tilde{u}_2'''(1) = 0 \end{cases} \tag{15}$$