## Asymptotic analysis of piezoelectric energy harvester

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### 1 Summary of the interested equations

Here we are interested in the classical model of a piezoelectric cantilever beam energy harvester, whose model is described using the following set of equations:

$$u'''' - \lambda^2 u = 0, (1)$$

and the accompanying boundary conditions:

$$\begin{cases} u(0) = 0 \\ u'(0) = 0 \end{cases}$$

$$u''(1) + \frac{j\lambda\beta\alpha^2}{j\lambda\beta + 1}u'(1) = 0$$

$$u'''(1) = 0$$
(2)

where  $\lambda$  is the eigenvalues for the problem, u denotes the displace function of the cantilever beam,  $\beta$  is the dimensionless externally connected resistance, and  $\alpha$  is the dimensionless piezoelectric coefficient. They can be expressed as follows

$$\lambda = \omega \sqrt{\frac{m_p l_p^4}{B_p}}, \quad \beta = R_l C_p \sqrt{\frac{B_p}{m_p l_p^4}}, \quad \alpha = e_p \sqrt{\frac{l_p}{C_p B_p}}, \tag{3}$$

where  $\omega$  is angular frequency,  $m_p$  is line mass density,  $l_p$  is the length of the cantilever beam,  $B_p$  is the bending stiffness,  $C_p$  is the inherent capacitance of the piezoelectric layer,  $e_p$  is the charge accumulation number,  $R_l$  is the externally connected resistance. In practical applications, dielectric property of piezoelectric materials indicate that the parameter  $\beta$  is changed from a very small value, which is close to a short-circuit condition to a very large value, which corresponds to an open-circuit condition. Thus we have that  $0 \le \beta \le \infty$ .

## 2 Asymptotic analysis when $\beta$ is small

Here we seek to find the behavior of the above system at a small value of connected resistance, i.e.,  $\beta \to 0$ . In this case, we set  $\beta$  to be the parameter for asymptotic expansion, and

$$\lambda^{(k)} = \lambda_0^{(k)} + \beta \lambda_1^{(k)} + \beta^2 \lambda_2^{(k)} + \cdots$$

$$u^{(k)} = u_0^{(k)} + \beta u_1^{(k)} + \beta^2 u_2^{(k)} + \cdots$$
(4)

where  $\lambda^{(k)}$  and  $u^{(k)}$  are the kth eigenvalue and eigenfunction respectively of the above mentioned system under perturbation.  $\lambda_0^{(k)}$  and  $u_0^{(k)}$  are the corresponding eigenvalue and eigenfunction of the unperturbed system at  $\beta=0$ :

$$u'''' - \lambda_0^2 u = 0, (5)$$

$$\begin{cases} u(0) = 0 \\ u'(0) = 0 \\ u''(1) = 0 \\ u'''(1) = 0 \end{cases}$$
 (6)

Obviously, the unperturbed system is a classical eigenvalue problem with the eigenvalues determined by

$$1 + \cosh(\sqrt{\lambda_0})\cos(\sqrt{\lambda_0}) = 0 \tag{7}$$

whose first several values are

$$\frac{\sqrt{\lambda_0^{(1)}}}{\pi} = 0.59686, \quad \frac{\sqrt{\lambda_0^{(2)}}}{\pi} = 1.49418, \quad \frac{\sqrt{\lambda_0^{(3)}}}{\pi} = 2.50025, \quad \frac{\sqrt{\lambda_0^{(4)}}}{\pi} = 3.49999, \quad \cdots$$
 (8)

Take the asymptotic expansions and substitute them into the previously derived system of equations, we have the following asymptotic expansions to different orders of  $\beta$ :  $O(\beta^0)$ :

$$\begin{cases}
 u_0'''' - \lambda_0^2 u_0 = 0 \\
 u_0(0) = 0 \\
 u_0'(0) = 0 \\
 u_0''(1) = 0 \\
 u_0'''(1) = 0
\end{cases} \tag{9}$$

 $O(\beta^1)$ :

$$\begin{cases}
 u_1'''' - \left(\lambda_0^2 u_1 + 2\lambda_0 u_0 \lambda_1\right) = 0 \\
 u_1(0) = 0 \\
 u_1'(0) = 0 \\
 u_1''(1) + j\alpha^2 \lambda_0 u_0'(1) = 0 \\
 u_1'''(1) = 0
\end{cases} \tag{10}$$

 $O(\beta^2)$ :

$$\begin{cases}
 u_2'''' - (\lambda_0^2 u_2 + 2\lambda_0 u_1 \lambda_1 + \lambda_1^2 u_0 + 2\lambda_0 u_0 \lambda_2) = 0 \\
 u_2(0) = 0 \\
 u_2'(0) = 0 \\
 u_2''(1) + \alpha^2 \lambda_0 u_0'(1) + j\alpha^2 \left[\lambda_0 u_1'(1) + \lambda_1 u_0'(1)\right] = 0 \\
 u_2'''(1) = 0
\end{cases} \tag{11}$$

## 3 Asymptotic analysis when $\beta$ is large

Here we seek to find the behavior of the above system at a large value of connected resistance, i.e.,  $\beta \to \infty$ . In this case, we set  $\frac{1}{\beta}$  to be the parameter for asymptotic expansion and

$$\lambda^{(k)} = \tilde{\lambda}_0^{(k)} + \left(\frac{1}{\beta}\right) \tilde{\lambda}_1^{(k)} + \left(\frac{1}{\beta}\right)^2 \tilde{\lambda}_2^{(k)} + \cdots$$

$$u^{(k)} = \tilde{u}_0^{(k)} + \left(\frac{1}{\beta}\right) \tilde{u}_1^{(k)} + \left(\frac{1}{\beta}\right)^2 \tilde{u}_2^{(k)} + \cdots$$
(12)

where  $\tilde{\lambda}^{(k)}$  and  $\tilde{u}^{(k)}$  are the kth eigenvalue and eigenfunction respectively of the above mentioned system under perturbation.  $\tilde{\lambda}_0^{(k)}$  and  $\tilde{u}_0^{(k)}$  are the corresponding eigenvalue and eigenfunction of the unperturbed system at  $\beta = \infty$ :  $O(\frac{1}{30})$ :

$$\begin{cases}
\tilde{u}_0'''' - \lambda_0^2 \tilde{u}_0 = 0 \\
\tilde{u}_0(0) = 0 \\
\tilde{u}_0'(0) = 0 \\
\tilde{u}_0''(1) + \alpha^2 \tilde{u}_0'(1) = 0 \\
\tilde{u}_0'''(1) = 0
\end{cases} \tag{13}$$

$$O(\frac{1}{\beta^1})$$
:

$$\begin{cases}
\tilde{u}_{1}^{""} - \left(\tilde{\lambda}_{0}^{2}u_{1} + 2\tilde{\lambda}_{0}\tilde{u}_{0}\tilde{\lambda}_{1}\right) = 0 \\
\tilde{u}_{1}(0) = 0 \\
\tilde{u}_{1}^{"}(0) = 0 \\
\tilde{u}_{1}^{"}(1) + \alpha^{2}\tilde{u}_{1}^{"}(1) + \frac{j\alpha^{2}}{\tilde{\lambda}_{0}}\tilde{u}_{0}^{"}(1) = 0 \\
\tilde{u}_{1}^{"'}(1) = 0
\end{cases}$$
(14)

$$O(\frac{1}{\beta^2})$$
:

$$\begin{cases}
\tilde{u}_{2}^{""} - \left(\tilde{\lambda}_{0}^{2}\tilde{u}_{2} + 2\tilde{\lambda}_{0}\tilde{u}_{1}\tilde{\lambda}_{1} + \tilde{\lambda}_{1}^{2}\tilde{u}_{0} + 2\tilde{\lambda}_{0}\tilde{u}_{0}\tilde{\lambda}_{2}\right) = 0 \\
\tilde{u}_{2}(0) = 0 \\
\tilde{u}_{2}'(0) = 0
\end{cases}$$

$$\tilde{u}_{2}'(1) + \left[\alpha^{2}\tilde{u}_{2}'(1) - \frac{\alpha^{2}}{\tilde{\lambda}_{0}^{2}}\tilde{u}_{0}'(1)\right] + j\left[\frac{\alpha^{2}}{\tilde{\lambda}_{0}}\tilde{u}_{1}'(1) - \frac{\alpha^{2}\tilde{\lambda}_{1}}{\tilde{\lambda}_{0}^{2}}\tilde{u}_{0}'(1)\right] = 0$$

$$\tilde{u}_{2}^{"}(1) = 0$$

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# 4 Asymptotic analysis in terms of small $\alpha^2$

Directly using the eigenvalue analysis method for linear boundary value problem, we arrive at the equation for the eigenvalue  $\lambda$ :

$$\sqrt{\lambda} \left[ 1 + \left( \frac{e^{\sqrt{\lambda}} + e^{-\sqrt{\lambda}}}{2} \right) \cos \sqrt{\lambda} \right] + \frac{j\beta\lambda\alpha^2}{1 + j\beta\lambda} \left[ \left( \frac{e^{\sqrt{\lambda}} - e^{-\sqrt{\lambda}}}{2} \right) \cos \sqrt{\lambda} + \left( \frac{e^{\sqrt{\lambda}} + e^{-\sqrt{\lambda}}}{2} \right) \sin \sqrt{\lambda} \right] = 0$$
(16)