# [X/Fe] Marks the Spot: Mapping Chemical Azimuthal Variations in the Galactic Disk with APOGEE

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ABSTRACT

Chemical cartography of the Galactic disk provides insights to its structure and assembly history over cosmic time. In this work, we use chemical cartography to explore chemical gradients and azimuthal substructure in the Milky Way disk with giant stars from APOGEE DR17. We confirm the existence of a radial metallicity gradient in the disk of  $\Delta [Fe/H]/\Delta R \sim -0.066 \pm 0.0004 \text{ dex/kpc}$  and a vertical metallicity gradient of  $\Delta [Fe/H]/\Delta Z \sim -0.164 \pm 0.001$ . We find azimuthal variations ( $\pm 0.1$  dex) on top of the radial metallicity gradient that have been previously established with other surveys. The APOGEE giants show strong correlations with stellar age and the intensity of azimuthal variations in [Fe/H]; older stellar populations show the largest deviations from the radial metallicity gradient. Beyond iron, we show that other elements (e.g., Mg, O) display azimuthal variations at the  $\pm 0.05$ dex-level across the Galactic disk. We illustrate that moving into the orbit-space could help constrain the mechanisms producing these azimuthal metallicity variations. These results suggest that the spiral arms of the Galaxy are not solely responsible for azimuthal metallicity variations and other Galactic processes are at play.

Keywords: Galactic archaeology (2178), Milky Way disk (1050), Stellar abundances (1577)

# 1. INTRODUCTION

The field of Galactic archaeology exists to answer 26 long-standing questions about the processes that drive <sup>27</sup> Galactic formation and evolution (Eggen et al. 1962; 28 Searle & Zinn 1978). We can use the Milky Way and its 29 resolved Galactic components as a laboratory to answer 30 questions about Galactic evolution and characterize the 31 hierarchical formation (Davis et al. 1985) of the Milky 32 Way.

The current investigation of the Galactic processes 34 that drive the Milky Way's evolution has exploded 35 within the last few decades due to wide-field missions 36 aimed at mapping the stellar content of the Milky Way 37 such as APOGEE (Majewski et al. 2017a) and Gaia 38 (Gaia Collaboration et al. 2022b). Employing stars as 39 the key witnesses to Galactic evolution, we are able to

40 take stellar information (stellar parameters, kinematics, 41 chemical abundances, etc.) and apply them with spa-42 tial information to create information-dense maps of the 43 Milky Way in a process known as chemical cartography 44 (Hayden et al. 2015). Measuring the distribution of el-45 ements throughout the Milky Way disk can inform us 46 about global and secular processes that dominate over 47 space and time (e.g., Hawkins et al. 2015). To postulate which Galactic phenomena are the most

49 influential, observations of the Galactic disk (both thick 50 and thin) are needed to identify which signatures and 51 patterns prevail. One of the most prominent trends in 52 the Milky Way is the existence of the negative radial 53 and vertical metallicity gradients (an incomprehensive 54 list includes: Mayor 1976; Andrievsky et al. 2002; Ma-55 grini et al. 2009; Luck & Lambert 2011; Bergemann et al. 56 2014; Xiang et al. 2015; Yan et al. 2019; Hawkins 2022). 57 These gradients could provide supporting evidence for 58 certain formation theories of the Milky Way, such as 59 the 'inside-out' formation theory (Larson 1976). This

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theory suggests that the negative radial ( $\Delta[Fe/H]/\Delta R$ ) metallicity gradient in the Milky Way could be caused by the inner Galaxy forming first and fast, leading to a higher metallicity concentration towards the Galactic center. As time proceeds, the outer Galaxy starts to form causing it to be metal-poor compared to its centralized counterpart (Matteucci & Francois 1989; Chiappini et al. 1997).

While the metallicity gradients are some of the most prominent features observed in the Galaxy, there are more subtle chemical characteristics that appear to be washed out by these strong trends. Chemical azimuthal substructure in the Milky Way has been previously identified using a variety of different tracers such as HII regions (Balser et al. 2011, 2015) and Cepheids (Lemasle et al. 2008; Pedicelli et al. 2009). Recently, using large-scale stellar surveys, azimuthal variations in [Fe/H] throughout the disk are quantifiable on the level of  $\sim 0.1$  dex (Hawkins 2022; Poggio et al. 2022; Imig et al. 2023).

One motivator for looking for angle-dependent chemi-81 cal variations in our Galaxy is the confirmed existence of 82 azimuthal variations in galaxies beyond the Milky Way. 83 Ho et al. (2017) found [O/H] azimuthal substructure in 84 NGC 1365, proposing that the spiral arms in this galaxy 85 plays a large role in the dispersion of metals throughout 86 the galaxy. Hwang et al. (2019) show clear azimuthal 87 metallicity variations in SDSS IV MaNGA galaxies in 88 Figure 13. Kreckel et al. (2019) observed subtle az-89 imuthal variations in 4 out of 8 in their sample of nearby 90 galaxies from the PHANGS-MUSE survey, with ranging 91 associations between the metallicity variations and the 92 spiral arms. This range in correlations with the spiral 93 arms calls into question if spiral arms are the most ade-94 quate explanation for the causes of azimuthal chemical 95 variations.

There are two distinct mechanisms that spiral arms can generate to induce azimuthal metallicity variations. The first process is radial migration/churning (Sellwood & Binney 2002) which, in this context, is the morphing of stellar orbits that alters the angular momentum of an orbit without the addition of excess energy. The second process is kinematic heating/blurring which would induce changes (heat) in the stellar orbital parameters without increasing the angular momentum. An example of this is the spiral arms dynamically evoking changes in the motions of stars along the leading and trailing edges of the spiral arms (Grand et al. 2012, 2016).

In simulations, Khoperskov et al. (2018) found that azimuthal variations in metallicity may arise from the dynamics of stellar disks alone, without the need of radial migration to reshape the stellar population. This is

due to dynamically cooler populations in the disk showing a larger contribution to spiral arms than dynamically hotter populations, leading to azimuthal variations in metallicity. Khoperskov & Gerhard (2021) claimed that in various phase-space coordinates, stars in an angular momentum overdensity caused by the spiral arms also exhibit a mean metallicity that is systematically higher than stars not in an angular momentum overdensity.

Using Gaia DR3 stars, Poggio et al. (2022) found that the azimuthal variations present in the disk correlate with where the spiral arms are predicted to be using different samples of bright stars (using effective temperature as a proxy for age) within  $\sim 4$  kpc of the Sun. Interestingly, they concluded the inverse of the simulated results from Di Matteo et al. (2013), in that the youngest population exhibits the most pronounced inhomogenities.

Debattista et al. (2024) quantified azimuthal metallicity variations ( $\delta$ [Fe/H]) in a Milky Way-like galaxy simulation and found variations that are comprable with the magnitude of the azimuthal variations seen in the Galaxy. Their azimuthal variations in metallicity were coincident with the spiral density waves and present in both young and old populations of stars. When looking at the pattern speeds of the  $\delta$ [Fe/H] variations, they found that these pattern speeds matched those of the spirals, pointing to the spiral arms as the root cause of these variations.

There are a handful of explanations, aside from spi-141 ral arms, for these azimuthal variations that have been 142 introduced in the literature. One possibility for these 143 angle-dependent trends could be secular processes, such 144 as influence and migration due to the Galactic bar (Di 145 Matteo et al. 2013). Specifically, it has been speculated 146 that the presence of azimuthal variations in an old stel-147 lar population is a probe of bar activity. Using N-body 148 simulations, Filion et al. (2023) found that the bar in-149 duces azimuth-dependent trends in stellar radii, which 150 then coincides with angle-dependent variations in metal-151 licity using younger to intermediate-age stars ( $\sim 1-4$ 152 Gyr).

A recent study done suggested that the azimuthal variations could arise from interactions with a satellite galaxy, such as Sagittarius (Carr et al. 2022). Torques from the gravitational interaction between an external satellite and the disk of the host galaxy can cause raise dial migration of the stellar population where inward-migrating stars would be more metal-poor on average compared to in-situ populations and outward-migrating populations would be more metal-rich than the in-situ outer Galactic population (with an assumed negative raidal metallicity gradient). This migration response will

be induced with respect to the location of the perturber, which then stimulate the azimuthal variations between inward and outward migrating populations. With the authors' simulations, they posited that the influence of Sagittarius will be the most prominent in the outer disk, a conclusion also drawn from Laporte et al. (2018). Observationally, Hwang et al. (2019) found gas-phase azimuthal metallicity variations in close or interacting galactic pairs, providing evidence for merger-induced azimuthal metallicity variations.

Most likely, there will be a contribution from all of the aforementioned dynamical processes that could cause azimuthal variations. Empirical observations are necessary to try to disentangle the processes that are responsible for the observed azimuthal metallicity variations in the Galactic disk. While the azimuthal variations of [Fe/H] in the Galactic disk have been characterized by different tracer populations such as *APOGEE* red giants (Eilers et al. 2022) and LAMOST OBAF-type stars as well as *Gaia* (Hawkins 2022; Poggio et al. 2022), little work has been done when looking beyond iron towards other elements, investigating the effect that stellar population age has on the intensity of azimuthal variations, and examining the dynamical perspective.

Throughout this work, we aim to confirm the az-189 imuthal variations in the Galactic disk and investigate 190 any correlation with the spiral arms of the Galaxy. We 191 additionally split our sample by age to quantify the ef-192 fect that stellar age has on the intensity of the azimuthal variations, examine azimuthal variations in elements be-194 yound Fe, and probe the dynamical origins of these az-195 imuthal metallicity variations. In Section 2, we delin-196 eate the dataset adopted for this project. In Section 3, we explain the methods taken to achieve our goals. In Section 4.1 we present our radial and vertical [Fe/H] gra-199 dients, and highlight the azimuthal deviations from the 200 radial gradient in Section 4.2. We quantify azimuthal variations in multiple other elements in Section 4.3. We 202 separate our sample into distinct age bins and charac-203 terize how stellar populations of different ages have disparities in their azimuthal variations in Section 4.4. In Section 4.5, we link the stellar chemistry to kinematics. 206 We summarize our results in Section 5.

## 2. DATA

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The initial sample of stars came from the seventeenth data release (DR17, Abdurro'uf et al. 2022) of the Apache Point Observatory Galactic Evolution Experiment (APOGEE) (Majewski et al. 2017b). APOGEE is a large scale near-infrared (15140Å  $< \lambda < 16960$ Å) stellar spectroscopic survey. The survey spans both hemi-spheres, consisting of a spectrograph on the 2.5-meter

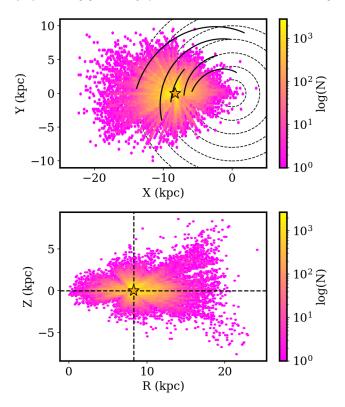


Figure 1. The face-on (top) and edge-on (bottom) distribution of our thin disk sample of 202,510 stars in which the the hexagonal bins are colored by the logarithmic number of stars. The black contours in the top panel are the spiral arms of the Milky Way determined by Reid et al. (2019). The orange star in both panels represents the Sun.

<sup>215</sup> Sloan Foundation Telescope at the Apache Point Ob-<sup>216</sup> servatory in New Mexico, USA as well as the 2.5-meter <sup>217</sup> Irénée du Pont Telescope at the Las Campanas Obser-<sup>218</sup> vatory in Chile.

We made use of the astroNN value-added catalog (Leung & Bovy 2019a)<sup>1</sup> of abundances, distances, and ages
for APOGEE DR17 sources. astroNN is an open source
python package developed for the neural network trained
on the APOGEE data and is designed to be a general
package for deep learning in astronomy. The APOGEEastroNN catalog contains results from applying astroNN
neural nets on APOGEE DR17 spectra to infer stellar
parameters, abundances trained with ASPCAP DR17
(Holtzman et al., in prep.), distances retrained with
Gaia eDR3 (Gaia Collaboration et al. 2021) from Leung & Bovy (2019b), and ages trained with APOKASC(Mackereth et al. 2019) in combination with lowmetallicity asteroseismic ages (Montalbán et al. 2021).

<sup>&</sup>lt;sup>1</sup> The astroNN package is available here: https://github.com/ henrysky/astroNN

The neural network from Leung & Bovy (2019a) mim- ics a 'standard' spectroscopic analysis by using the full wavelength range to deduce stellar parameters and specific sections of the spectrum to derive individual elemental abundances. astroNN takes into account incomplete and noisy training data while applying dropout variational inference to find uncertainties on the measurements. This catalog contains stellar parameters [This catalog contains stellar parameters of the measurement abundances with precisions at the  $\sim 0.03$  dex level, agreeing quite well with the traditional ASPCAP pipeline and producing a smaller scatter with tighter uncertainties.

For the distances in this work, we adopted the weighted distance in APOGEE-astronn. This parameters ter is a weighted combination of the astronn distance (spectro-photometric calibrated distances from Leung & Bovy 2019b) and Gaia parallax (Gaia Collaboration et al. 2022a). The Galactocentric X, Y, and Z coordinates were found as transformations from the Galactocentric cylindrical radius, azimuth, and vertical height given by astronn. The Galactocentric positions and velocities were computed assuming the Sun is located at 8.125 kpc from the Galactic center (GRAVITY Collaboration et al. 2018), 20.8 pc above the Galactic midplane (Bennett & Bovy 2019) and has radial, rotational, and vertical velocities of -11.1, 242, and 7.25 km/s, respectively (Schönrich et al. 2010; Bovy et al. 2012).

To explore the effect of stellar age on the chemical azimuthal variations, we sorted our sample into
3 distinct groups with the ages in *APOGEE*-astronn.
Due to the tendency to under-predict old ages, Leung
5 & Bovy (2019a) and Mackereth et al. (2019) used a
5 non-parametric Locally Weighted Scatterplot Smooth5 ing (LOWESS) to correct the ages. While this incon5 sistency mainly applies to stars older than 8+ Gyr, we
5 followed the recommendation and applied the LOWESS
5 corrected ages for the analysis in this work.

To explore the intersection of chemistry and dynamics, we aim to quantify how metallicity excess ( $\delta[\text{Fe/H}]$ ) changes with respect to the dynamical parameters radial action  $(J_r)$ , angular momentum  $(L_z)$  which is proportional to the azimuthal action  $(J_\phi)$ , vertical action  $(J_z)$ , orbital eccentricity (e), maximum height above the Galactic plane  $(Z_{\text{max}})$ , and total energy. These parameters were found by integrating the stellar orbits with the Gala code designed primarily for Galactic dynamics evaluations (Price-Whelan 2017). The mass-model of the Milky Way used for the gala.dynamics.orbit function was the MilkyWayPotential2022 that has been fit to the rotation curve in Eilers et al. (2019) and incorpo-

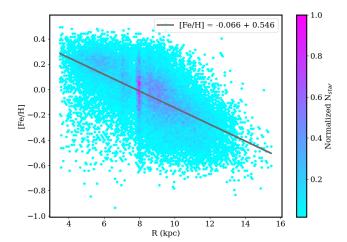


Figure 2. The [Fe/H] abundances with respect to Galactocentric radius of our planar thin disk sample. The artifact at R  $\sim$  8 kpc is an observational effect of the over-representation of stars in the solar neighborhood. We fit a linear model to the colored data points (column-normalized planar thin disk sample) which is represented by the grey line. From this, we obtain a metallicity gradient of  $\Delta [{\rm Fe/H}]/\Delta R$  -0.066±0.0004 dex/kpc and a y-intercept of 0.546 dex for the stars in our planar thin disk sample.

<sup>284</sup> rates the phase-space spiral in the solar neighborhood <sup>285</sup> set by Darragh-Ford et al. (2023).

With the *APOGEE*-astroNN catalog, we employed the following cuts to obtain a set of stars that is well-sampled and a reliable representation of the population we aim to characterize:

- The Galactocentric distances are an integral part
  of this work to characterize the metallicity gradient in the Galactic disk. To ensure we have precise positions we place an error cut on the distance
  measurements obtained from the astroNN dataset,
  selecting stars with errors < 30%.</li>
- 2. To minimize systematic effects and dwarf contamination in our sample, we removed stars outside of the effective temperature range  $3500 < T_{\rm eff} < 5000$  K as well as any star with log g > 3.6 dex. This selection criterion establishes that we are sampling the true red giant section of the Color-Magnitude Diagram (CMD).
- 3. We removed any stars with an error larger than 0.08 dex on [Fe/H], [O/H], or [Mg/H] due to these elements being the most relevant in our chemical cartography. The limit of 0.08 dex was chosen based on previous studies that have found the dispersion of line-to-line abundances in APOGEE is  $\sim 0.08$  dex for [Fe/H] (Chen et al. 2015; Hawkins

et al. 2016) We placed no selection cut on any of the other elements to maintain a balance between the robustness and size of our sample. For completeness we include all of the elements in Section 11 but we refrain from drawing any conclusions about Galactic evolution from the elements with no error cuts (elements other than [Fe/H], [O/H], and [Mg/H]).

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4. Chemical composition and evolution varies differently among the thin disk, thick disk, and halo. We select only the kinematic thin disk stars based on the processes outlined in Ramírez et al. (2013) (full derivation shown in Appendix A). In this work, the authors outlined a probabilistic approach to assign membership probabilities for stars to thin disk, thick disk, or halo populations based on their Galactic space velocities. We required all thin disk stars to have a TD/D probability > 80%, providing us with a final sample of 202,510 stars, shown in Figure 1.

With this full thin disk sample, we are able to characterize the vertical metallicity gradient as a function of Galactocentric radius as well as the radial metallicity gradient as a function of height above the plane. To investigate the chemical gradients and azimuthal variations in the plane of the disk, we further sub-sample the 202,510 stars in the following section.

## 2.1. Planar Thin Disk Sample

The density of stars in the stellar disk of the Milky 338 Way is thought to noticeably decline at about  $\sim 15.5$ 339 340 kpc away from the Galactic center (Momany et al. 2006; Reylé et al. 2009; Carraro et al. 2010). The inner radii of the Galactic disk are thought to be dominated by the bar, typically believed to extend out to  $\sim 3.5~\mathrm{kpc}$ 344 (Hammersley et al. 1994; Wegg et al. 2015; Lucey et al. 345 2023). Thus to minimize the chemical and dynamical 346 effects driven by the bar in the most central region, and 347 the drop in the density of stars in the furthest regions, we 348 adopted a Galactocentric radius cut of 3.5 < R < 15.5349 kpc to our final thin disk sample. We place a 0.3 kpc  $_{350}$  cut on  $Z_{max}$  to ensure that we are only selecting bona 351 fide thin disk stars with orbits confined to the Galactic 352 plane. After these final selection cuts, we have a planar 353 thin disk sample of 32,768 stars with the median uncertainties for  $T_{\text{eff}}$ , log g, and [Fe/H] being 33K, 0.07 dex, and 0.03 dex, respectively.

#### 3. METHODS

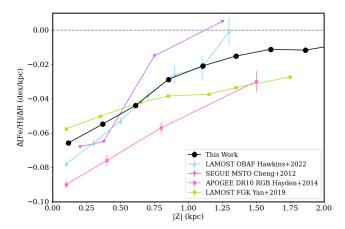


Figure 3. The radial metallicity gradient as a function of absolute vertical height Z above (and below) the plane. The gradient derived in this work is represented by the black dots, compared to a variety of other studies that use different tracers [Cheng et al. (2012) (purple triangles), Hawkins (2022) (blue triangles), Hayden et al. (2014) (pink squares), Yan et al. (2019) (green squares)]. Consistently,  $\Delta[Fe/H]/\Delta R$  starts off at its most negative in the plane and shallows out with greater distances from the disk. A vertical line is plotted at  $\Delta[Fe/H]/\Delta R$ =0 to illustrate where the radial gradient is no longer negative. The points derived in this work lie generally in the middle of the other studies conducted.

**Table 1.** Radial Metallicity Gradient Parameters for |Z| bins in the full thin disk sample

Z	$\Delta {\rm [Fe/H]}/\Delta {\rm R}$	$\sigma\Delta { m [Fe/H]}/\Delta { m R}$	$N_{\rm stars}$
(kpc)	(dex/kpc)	(dex/kpc)	
0.1	-0.0672	0.0004	40434
0.3	-0.0579	0.0005	42167
0.5	-0.0492	0.0006	35415
0.7	-0.0389	0.0007	25742
0.9	-0.0274	0.0009	14874
1.1	-0.0218	0.0011	9051
1.3	-0.0153	0.0013	5780
1.5	-0.0132	0.0015	3573
1.7	-0.0126	0.0016	2298
1.9	-0.0107	0.0018	1580

The behavior of the metallicity of stars throughout the Galactic disk can be best characterized by a negative linear gradient (e.g. Janes 1979; Rolleston et al. 2000). Using our planar thin disk sample, we derived the radial and vertical metallicity gradients stars as follows.

The gradients are initially represented as linear func-363 tions in which the gradients are represented by  $m_R$ 

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 $m_z$  in Equations 1 and 2. In this,  $[Fe/H]_R$  is the metallicity at a certain Galactocentric radius (R) and  $m_R = \Delta [Fe/H]/\Delta R$ . To account for the metallicity at 367 the Galaxy's center, we introduce the term  $b_R$ :

$$[Fe/H]_R = m_R R + b_R \tag{1}$$

Similarly, we model the vertical metallicity gradient are as:

$$[Fe/H]_Z = m_Z|Z| + b_Z \tag{2}$$

where  $m_z$  and  $b_z$  are constants.

To find the gradients, we fit linear models to the equa374 tions above with scipy.stats.linregress (Virtanen
375 et al. 2020) which calculates a linear least-squares re376 gression for two sets of measurements (in our case Galac377 tocentric radius/height and [Fe/H]). This returns the
378 slope of the line (metallicity gradient), the standard er379 ror of the gradient, and the y-intercept (metallicity at
380 the Galactic center/midplane). When computing the
381 vertical gradient, we use our full thin disk sample with382 out a radius or height cut. The resulting gradients are
383 in Section 4.1. In Tables 1 (and 2), we compute the
384 linear (vertical) metallicity gradients in radial (Galactic
385 height) bins.

We use our 1D models to create 2D metallicity residual maps to search for potential signature of azimuthal variations. We subtract the model abundances (in which the abundances are exactly equal to the linear gradient) from the observed [Fe/H] of each star in our sample. If the linear gradient is the only chemical feature of the disk, we would expect to see stochastic noise. If there is structure in this noise, then there are other processes driving the chemistry of the disk. We plot these residuals and investigate this azimuthal substructure in Section 4.2. When comparing the metallicity excess to azimuth, it is important to note that we define the line of sight from the Galactic center through the Sun to be an azimuth of  $\phi(\pi) = 1$ .

The linear gradient is a useful approximation for [Fe/H], however this does not necessarily hold true for all elements (see Appendix B). To attempt to account for any non-linearity in our sample of elements, we employed a 'running median' method. We began by calulating the median [X/Y] abundances in 0.2 kpc radial bins within the bounds of  $\sim 6-15$  kpc. In these bins, we then took the median [X/Y] value and following the linear method, subtracted Data-Median to quantify how the individual stars may deviate azimuthally from the median [X/Y] abundance. However, when this approach was taken, the radial bins did not sample azimuthal angles isotropically due to APOGEE only probing one part of the disk. Thus, for the context of this work we only

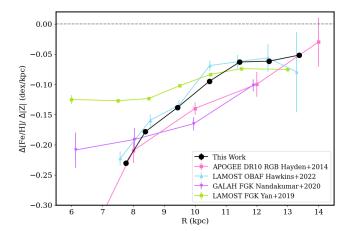


Figure 4. The vertical metallicity gradient as a function of Galactocentric radius. The gradient derived in this work is depicted by black dots, whereas the other colored points represent the gradient as determined by different tracers in other studies [Hayden et al. (2014) (pink squares), Hawkins (2022) (blue traingles), Nandakumar et al. (2020) (purple triangles), Yan et al. (2019) (green squares)]. The vertical metallicity gradient is at its most negative closest to the Galactic center and shallows out  $(\Delta [Fe/H]/\Delta Z$  approaches 0) as distance increases.

**Table 2.** Vertical Metallicity Gradient Parameters for R bins in the full thin disk sample

R	$\Delta { m [Fe/H]}/\Delta { m Z}$	$\sigma\Delta { m [Fe/H]}/\Delta { m Z}$	$N_{\rm stars}$
(kpc)	(dex/kpc)	(dex/kpc)	
7.5	-0.2307	0.0026	29634
8.5	-0.1785	0.0022	49208
9.5	-0.1383	0.0025	33837
10.5	-0.0951	0.0026	22099
11.5	-0.0632	0.0028	14729
12.5	-0.0617	0.003	8025
13.4	-0.0518	0.0034	3090

414 select elements that behave monotonically with respect 415 to Galactocentric radius and move forward with the lin-416 ear gradient method in Section 4.3.

## 4. RESULTS AND DISCUSSION

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In this section we present our radial and vertical metallicity gradients (Section 4.1), as well as azimuthal [Fe/H] variations throughout the thin disk (Section 4.2). We uncover azimuthal variations in elements beyond Fe in Section 4.3. To investigate which mechanisms may be responsible for these variations, we bin our sample

by age to explore the age dependence on azimuthal variations (Section 4.4) and close out the section by examining the link between metallicity excess and several dynamical properties (Section 4.5).

#### 4.1. Metallicity Gradients

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Our first step in this analysis was to measure the radial 429  $(\Delta [Fe/H]/\Delta R)$  and vertical  $(\Delta [Fe/H]/\Delta Z)$  metallicity 430 gradients in our planar thin disk sample that have been 432 previously observed with different tracer populations. 433 Employing the methods outlined in Section 3, we ob-434 tained a radial metallicity gradient with our planar thin disk sample of  $\sim -0.066 \pm 0.0004 \, \mathrm{dex/kpc}$  (Figure 2) as well as a vertical metallicity gradient of  $\sim -0.164 \pm 0.001$ 437 dex/kpc. Overall our gradients match within reason to 438 other recently calculated gradients. Hawkins (2022) de-439 rived a radial metallicity gradient of  $\sim -0.078 \pm 0.001$ 440 dex/kpc in the Galactic disk and a vertical metallicity gradient of  $\sim -0.15 \pm 0.01$  dex/kpc with an OBAF-type 442 stellar sample from the LAMOST survey. Using the six-443 teenth data release of APOGEE, Eilers et al. (2022) determined a radial gradient of  $\sim -0.057 \pm 0.001$  dex/kpc. 445 Imig et al. (2023) calculated the running median ra-446 dial metallicity gradients with APOGEE DR17 ASPCAP parameters and found  $\Delta [Fe/H]/\Delta R = -0.064 \pm 0.001$ 448 dex/kpc in their low- $\alpha$  sample at  $|Z| \leq 0.25$ , which is in good agreement with our gradient of  $\sim -0.066$  dex/kpc. 450 We note one of the main distinctions in our sample com-451 pared to other works in the literature is the cut we im- $_{452}$  posed on  $Z_{max}$  as opposed to present day height above 453 the plane.

The radial metallicity gradient is illustrated in Figure 2. The metal enhancement of the inner Galaxy as compared to the outer Galaxy could point to the 'inside-out' formation theory (Larson 1976; Kobayashi & Nakasato 2011). Starting at the inner Galaxy (i.e., R  $\sim$ 3.5 kpc), the metallicity decreases linearly — for every kiloparsec moving towards the outer disk, the global metallicity decreases by  $\sim 0.07$  dex.

Figure 3 explores how  $\Delta [Fe/H]/\Delta R$  varies with height above (and below) the Galactic plane. We separate the full thin disk sample into 0.2 kpc steps of —Z— and plot these results along with a select few other studies that use different tracer populations: LAMOST OBAF-type stars (Hawkins 2022), SEGUE main sequence turn-off stars (Cheng et al. 2012), APOGEE red giant stars (Hayden et al. 2014), and LAMOST FGK-type stars (Yan et al. 2019). Regardless of tracer population, all studies how a clear trend that as the distance from the Galactic plane increases, the radial metallicity gradient becomes more shallow, approaching zero. Physically, this means that the radial metallicity gradient is a promi-

nent feature in the disk, but with increasing height from the Galactic plane the radial gradient is no longer the dominant observed relation and the vertical metallicity gradient starts to take over. We tabulate these results in Table 1. This work's  $\Delta[\text{Fe/H}]/\Delta R$  vs. —Z— trend lies generally in the center of the other studies' trends. When compared with Hayden et al. (2014) which used the same survey as this work, merely an earlier data relasse (DR10 compared to our DR17), our work seems to be shifted up by  $\sim 0.01$  dex. This is not surprising due to the results of Jönsson et al. (2018) showing that APOGEE data releases can have systematic differences of less than 0.05 dex.

Recent studies have shown that there is a vertical metallicity gradient that varies as a function of Galactocentric radius. Önal Taş et al. (2016) found that the radial gradient is flat within 0.5 - 1 kpc of the plane and then becomes positive greater than 1 kpc away from the plane, suggesting that there is a vertical metallicity gradient in the Galaxy. To derive the vertical gradient, we applied the same methodology used for the radial metallicity gradient, obtaining  $\Delta [{\rm Fe}/{\rm H}]/\Delta {\rm Z} \sim -0.164 \pm 0.001$  dex/kpc.

To quantify how the vertical gradient changes as a 499 function of Galactocentric radius, we follow a similar 500 methodology to the evaluation of the radial gradient 501 varying with Galactic height, this time splitting our full  $_{502}$  thin disk sample of stars into radial bins of 1 kpc (Table 503 2). In Figure 4, we show our derived vertical metal-504 licity gradient using the APOGEE red giants, denoted 505 by black circles. Other works that we compared to in-506 clude vertical gradients obtained from APOGEE red gi-507 ant stars (Hayden et al. 2014), LAMOST OBAF-type 508 stars (Hawkins 2022), GALAH FGK-type stars (Nan-509 dakumar et al. 2020), and LAMOST FGK-type stars 510 (Yan et al. 2019). We find that the vertical metallicity 511 gradient is heavily correlated with Galactocentric radius <sub>512</sub> in that  $\Delta [Fe/H]/\Delta Z$  approaches zero with increasing 513 distance from the Galactic center.

## 4.2. Azimuthal Variations in $\Delta [Fe/H]/\Delta R$

In this section, we aim to characterize the angledependent deviations from the linear radial metallicity gradient in our planar thin disk sample and investigate any correlations with these deviations and the
spiral arms. The discernible radial metallicity gradient
has been recognized and characterized for decades (e.g.,
Mayor 1976; Andrievsky et al. 2002; Magrini et al. 2009;
Boeche et al. 2013, 2014; Cunha et al. 2016). With increased sample size and precision provided by advances
in instrumentation throughout the years, we are now

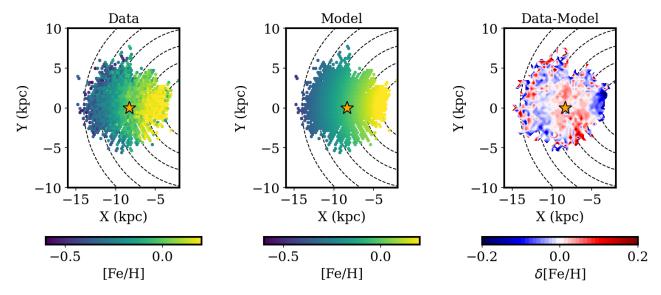


Figure 5. Left: The full *APOGEE* planar thin disk sample colored by the metallicity. Middle: Each position of the data points colored by a model gradient of  $\Delta [\text{Fe/H}]/\Delta R = -0.066 + 0.546 \text{ dex/kpc}$ . Right: the residuals of the observed [Fe/H] abundances and the linear model abundances.

<sup>525</sup> able to peel back the gradient to see if there is any under-<sup>526</sup> lying structure beneath this strong signature.

To look beyond the metallicity gradient, we follow the 527 528 steps outlined in Section 3 and compute the model linear  $\Delta [Fe/H]/\Delta R$  gradient throughout the disk. We subtract 529 530 this model off of the data to search for any structure in the residuals. When applied, we find variations in the 531 [Fe/H] abundances that correlate with azimuthal angle. 532 We plot the top-down view of the Milky Way with 533 534 our planar thin disk data in X and Y colored by [Fe/H] 535 in Panel 1 of Figure 5. We illustrate the model lin-536 ear metallicity gradient in the second panel of Figure and in the final panel we show the residuals that are 538 found after subtracting the model from the data (labelled  $\delta$ [Fe/H]).

There is a blue-red-blue-red pattern in  $\delta$ [Fe/H] shown 540 the final panel of Figure 5. The red bins repre-542 sent areas of the Milky Way that are more metal-rich 543 than the model predicts and the blue bins are where 544 the model over-estimates the stellar metallicities. This 545 oscillating pattern has been observed previously with 546 different stellar surveys, such as Gaia DR3 stars in Fig-547 ure 2 of Poggio et al. (2022), with similar results to this 548 work. Hawkins (2022) use LAMOST OBAF stars as well as Gaia DR3 stars, uncovering azimuthal variations in 550 both of these populations, showing similar patterns in which the solar-neighborhood is more metal-rich (red) and the inner/outer galaxy are more metal-poor (blue). To probe different formation pathways for this oscil-554 lating pattern, we first compare these results with the

555 location of the spiral arms. Khoperskov et al. (2018)

from the center of the Galaxy may travel along the spifrom the center of the Galaxy may travel along the spifrom the center of the Galaxy may travel along the spifrom the center of the Galaxy may travel along the spifrom the center of the Galaxy may travel along the spifrom the center of the Galaxy may travel along the spifrom appear more metal rich. Khoperskov et al. (2018) show that azimuthal [Fe/H] variations in spiral galaxies would from arise naturally if there is an initial negative radial metallicity gradient due to the migration of stars. To detect any correlation with the spiral arms, we plot our  $\delta$ [Fe/H] from Figure 6.

The black solid lines in the bottom panel of Figure 6 567 represent the spiral arms as determined by Reid et al. 568 (2019) using the parallaxes and proper motions associ-569 ated with high-mass star-forming regions using a Very 570 Long Baseline Array, European VLBI Newtork, and 571 the Japanese VLBI Exploration of Radio Astrometry 572 project. They locate multiple arm segments with pitch <sub>573</sub> angles ranging from 7° to 20° with the widths of the 574 arms increasing with distance from the Galactic center. 575 The black contours in the top panel of Figure 6 repre-576 sent the spiral arm locations determined by Poggio et al. 577 (2021) using the over-density of upper main sequence 578 stars determined by Gaia DR3. These results were con-579 sistent with some of the Reid et al. (2019) arms, such as 580 the Sagittarius-Carina spiral arm, while the geometry of <sub>581</sub> arms with Galactic longitudes from 180° to 270° were 582 significantly different from other spiral arm models. The 583 metal-rich and metal-poor portions of the Milky Way in 584 our work are not fully encompassed by either determi-585 nation of the spiral arm locations, thus we look towards

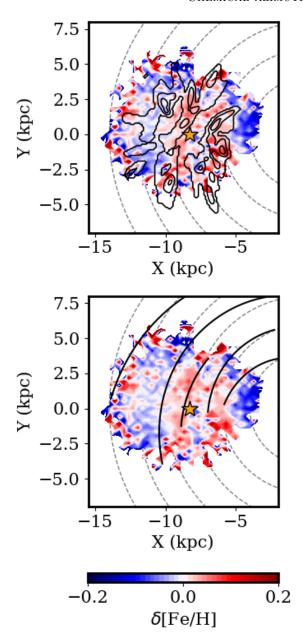


Figure 6. The azimuthal metallicity variations ( $\delta$ [Fe/H] in the last panel of Figure 5) as compared to different determinations of the spiral arms. The black contours in the top panel represent the spiral arms as derived by Poggio et al. (2021) using main sequence stars in *Gaia* and the solid black lines in the bottom panel are the spiral arms as determined by Reid et al. (2019) using high-mass star-forming regions.

586 other processes that may have induced this azimuthal 587 substructure.

If the spiral arm locations do not obviously coincide with the azimuthal substructure, this begs the question, what is causing these variations? Carr et al. (2022) quantified the effect that a Sagittarius (Sgr)-like dwarf galaxy would have on the kinematics and chemsitry of 593 the Milky Way-like disk upon first interaction. They did 594 this by painting the particles in their N-body simulation 595 of the disk with an negative radial [Fe/H] gradient. The 596 passage of Sgr through the plane will cause radial rear-597 rangement and disrupt stellar orbits. If the heating of <sup>598</sup> orbits is due to a non-axisymmetric feature, such as Sgr. 599 migration and mixing signatures will manifest as an ap-600 proximate quadrupole in chemical azimuthal variations 601 across the disk. In their present day snapshot of a sim-602 ulated Milky Way being perturbed by a Sgr-like dwarf 603 galaxy, Carr et al. (2022) found that the maximum de-<sub>604</sub> parture from the  $\Delta [Fe/H]/\Delta R$  gradient occurs in the 505 solar annulus on the side of disk closest to the current 606 position of Sgr, while the minimum is found adjacent 607 to Sgr in the outermost annuli. These azimuthal varia-608 tions in metallicity will be dependent on the model of Sgr 609 chosen. It is interesting to note that in our sample the 610 solar annulus appears to be more metal-rich, although 611 we lack a complete view of the Galaxy to classify this 612 definitively.

While external influences like Sgr can cause azimuthal 614 metallicity variations, studies have also shown that these 615 features can arise through secular evolution with struc-616 tures like the Galactic bar. Filion et al. (2023) sim-617 ulated the effects of radial rearrangement in a barred 618 galaxy and find substantial changes in the radii of stars 619 that, when paired with the negative radial metallicity 620 gradient, will induce azimuthal substructure similar to 621 what we find with the APOGEE stars (see also Di Mat-622 teo et al. 2013). The radial rearrangement of stars due 623 to the bar can be split into 3 zones to characterize the 624 dynamics of the stars. The effects driven by the bar 625 would be most prominent in the inner Galaxy where or-626 bits evolve inward due to the angular momentum loss of 627 the stellar orbits. Their intermediate zone is composed 628 of stars moving both inwards and outwards producing 629 no mean radial evolution. In the outer zone of their 630 simulations, orbits generally evolve outwards (likely due 631 to the net effect of the bar moving angular momentum 632 outwards) and the trends are less aligned with the bar 633 angle. Any migration of stars across the disk will pro-634 duce azimuthal metallicity variations. In line with their 635 findings, we see that the inner-disk stars of our sample 636 have a lower metallicity than predicted by the gradient 637 alone.

To discern the responsible mechanisms for this substructure, it is crucial to confirm that this pattern is not unique to iron but extends to other elements as well. While the azimuthal [Fe/H] variations have been quantified before using *Gaia* (Hawkins 2022; Poggio et al. 2022), *APOGEE* gives us access to a suite of other elements. To confirm that this substructure is present in

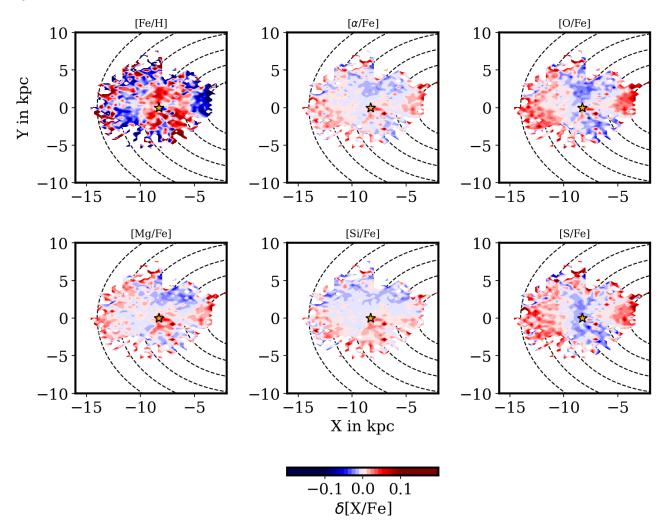


Figure 7. The Data-Model panel for the elements we deem to behave monotonically with respect to Galactocentric radius. The contours are colored by the shared [X/Fe] colorbar to illustrate the varying intensities of the deviations from the linear radial gradient. The second panel represents the average  $[\alpha/Fe]$  abundance, with the individual  $\alpha$ -elements in the following panels. The elements with the most saturated contours, such as [Fe/H], showcase the most exaggerated deviation from the radial gradients. The  $\alpha$ -elements appear to be loosely anti-correlated with [Fe/H], following predictions from the difference in timescales between events that mainly produce  $\alpha$ -elements (Type II SN) and events that mainly produce [Fe/H] (Type Ia SN).

645 elements aside from iron, we evaluate other azimuthal 646 chemical variations in the following section.

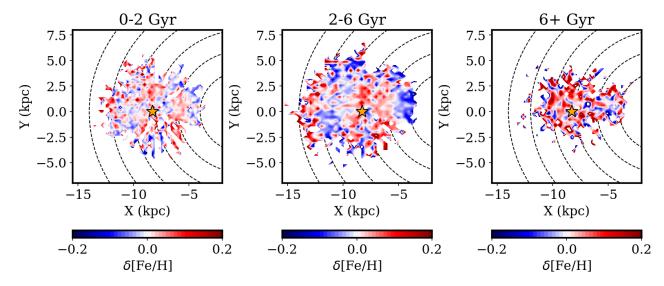
## 4.3. Azimuthal Variations in Other Elements

Here, we explore azimuthal variations in [Fe/H] and  $\alpha$ elements (O, Mg, Si, S, Ca) in our planar thin disk sample. We choose to specifically highlight the  $\alpha$ -elements
because the abundances of these elements in stars appear to behave monotonically with respect to Galactocentric radius, similar to [Fe/H]. For this methodology,
we used our planar thin disk sample and plotted the median abundances of these elements in 0.2 kpc radial bins
in Figure 11. As not all of these elements are monotonic with respect to Galactocentric radius, we focus on
the  $\alpha$ -elements. We fit linear gradients to these elements
and subtracted off the linear model (explained further in

Section 3) to look for structure in the residuals (similar to [Fe/H]), shown in Figure 7.

In Figure 7, we show azimuthal variations with ranging intensities for all elements in our analysis, with variations detectable on at least the  $\sim 0.05$  dex-level. The second panel in this figure represents the average  $[\alpha/\text{Fe}]$ abundance, while the individual  $[\alpha/\text{Fe}]$  abundances are found in the following panels. We find that the  $\alpha$ elements are loosely anticorrelated with [Fe/H], which aligns with expected Galactic chemical evolution (e.g., Tinsley 1980); in the disk where there is a  $\delta[\text{Fe/H}]$  ex-

The anti-correlation between [Fe/H] and  $[\alpha/\text{Fe}]$  can be explained by the astrophysical processes (and timescales) that are largely responsible for the produci-



**Figure 8.** This figure is similar to Panel 3 of Figure 5 increasing age groups of 0-2, 3-6, and 6+ Gyr from left to right with 8730, 18493 and 5545 stars respectively. The color bars are constant throughout all panels, thus the youngest age group has the least amount of contrast and represents the smallest deviations from the linear gradient. Conversely, the oldest age group has the most saturated colors due to the larger variations from the modelled linear gradient.

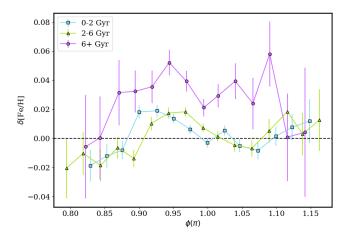


Figure 9. This figure shows  $\delta[\text{Fe/H}]$  as a function of azimuthal angle  $(\phi)$ . For reference,  $\phi(\pi)=1$  is the line of sight from the Sun towards the Galactic center. The youngest age group, 0-2 Gyr (blue squares), displays the smallest deviations from zero in  $\delta[\text{Fe/H}]$ . The intermediate age group, 2-6 Gyr (green triangles), is the largest sample and has noticable deviations from zero. The oldest age group, 6+ Gyr (purple circles), varies the most dramatically in  $\delta[\text{Fe/H}]$  azimuthally. This depiction matches the theorized azimuthal variations, discussed further in Section 4.4.

ton of these elements. At early times, Type II SNe were able to effectively disperse  $\alpha$ -elements (e.g., O, Mg, Si, S, and Ca) as well as small amounts of Fe that were synthesized during the lifetimes of the first generations of stars. As the Galaxy continues to evolve, Type Ia SNe will then dominate, causing the  $[\alpha/\text{Fe}]$  abundances to lower due to more Fe-peak (Fe, Cr, Mn, Co, Ni) elements be-

 $^{682}$  ing made available to the next generation of stars (e.g.,  $^{683}$  Tinsley 1980; Greggio & Renzini 1983; Woosley & Hoff- $^{684}$  man 1992; Matteucci & Recchi 2001; Gonzalez et al.  $^{685}$  2011). Following this sequence of events, wherever there  $^{686}$  are [Fe/H]-excess regions, there will be [ $\alpha$ /Fe]-defecit re- $^{687}$  gions and vice versa.

The presence of azimuthal variations in other elements confirms the existence of some process that is causing the observed chemical substructure throughout the Galactic disk. When looking for the responsible process, Filion et al. (2023) noted that focusing on younger populations ( $\sim 1-2$  Gyr) in the inner disk would help discern the role of secular evolution. Other simulations have used exclusively older stars to identify the cause of azimuthal chemical variations (Di Matteo et al. 2013). Consequently, we delve into the possible age-dependencies observed in azimuthal metallicity variations in the following section.

## 4.4. Azimuthal Metallicity Variations by Age

In this section, we aim to quantify the relationship between stellar age has and the magnitude of the azimuthal
substructure found in Section 4.2. We are particularly
interested in whether the amplitude of the azimuthal
[Fe/H] variations are stronger in older or younger populations. We additionally aim to explore how the [Fe/H]
gradient and azimuthal variations change as a function
of age. For reference, it is important to note the difference between look-back time and present-day age. Lookback time is primarily used in simulations, where it is
the time elapsed from the 'final'/present time back to a

previous point in time, used to analyze historical states. In observations, we are limited to the present-day age to the object, or the total age of the object from when it formed to the present time. Thus, we aim to use the present-day ages of these stars and compare any trends in the ages with simulated works.

Using present-day age, the Debattista et al. (2024) 719 simulations found that removing stars younger than 2 720 Gyr still produces azimuthal variations. This implies 721 that azimuthal variations are not solely primordial in 722 origin because older populations exhibit these variations 723 as well. Bellardini et al. (2021) demonstrated that az-<sub>724</sub> imuthal scatter increases with increasing look-back time. 725 This would indicate that stars that are born at earlier 726 times will be born with stronger azimuthal variations 727 due to the inhomogeneous interstellar medium (ISM) 728 from which they form. While these two hypotheses ap-729 pear to be contradictory, we cannot probe look-back 730 times observationally so we cannot compare directly to 731 Bellardini et al. (2021), but we can compare our results 732 to Debattista et al. (2024). Consequently, in the rest of 733 this section we make and test predictions about the ob-734 served relationship between azimuthal metallicity variations and stellar present-day age.

If the azimuthal variations are exclusively natal in origin (i.e., they are a result of variations in the gas-phase
abundances at the time a population was born), then
we would predict that younger populations would exhibit strong metallicity variations due to these populations forming more recently. However, if azimuthal variations are exclusively the result of dynamical processes,
we would predict that older populations would show
stronger azimuthal variations because these stars would
have had more time to interact with non-axisymmetric
features in the Galaxy that would cause the rearrangement of stars, leading to the observed azimuthal metallicity variations. To distinguish between these predictions, we applied the present-day ages from APOGEEastroNN to separate our sample by stellar age.

We separated the *APOGEE* giants into 3 distinct age groups to explore the effects that age has on departures from the metallicity gradient. The age bins we selected are 0-2 Gyr, 2-6 Gyr, and 6+ Gyr, with 8730, 18493 and 5545 stars, respectively. The strongest radial metallicity gradient appears in the middle age group with  $\Delta [\text{Fe/H}]/\Delta R \sim -0.0774 \pm 0.0005 \, \text{dex/kpc}$ , outranking the youngest ( $\Delta [\text{Fe/H}]/\Delta R \sim -0.0674 \pm 0.0006 \, \text{dex/kpc}$ ) and oldest ( $\Delta [\text{Fe/H}]/\Delta R \sim -0.0484 \pm 0.0012 \, \text{dex/kpc}$ ) age groups.

The correlations we find with radial metallicity gradient and stellar age are consistent with the findings of Wang et al. (2019) who used the LAMOST survey to

<sup>764</sup> illustrate that stars of 4-6 Gyr in age exhibit a steeper respectively gradient than either younger or older stars. Similar work respectively has been done with a CoRoT and APOGEE sample by respectively Anders et al. (2017) where their middle-aged population respectively of stars had the steepest radial metallicity gradient, followed by the youngest then oldest population. This age-gradient trend is consistent with a systematic offset of respectively respectively. The consistent with a systematic offset of respectively consistent with a systematic offset of respectively.

After following the same methodology as Section 4.2 and subtracting the data from the linear gradient model, Figure 8 shows the  $\delta$ [Fe/H] residuals for each of the 3 age groups. Visually, the signatures seem to be the strongest in the panel containing the oldest stars, while the deviations lessen in intensity with decreasing age.

In Figure 9 we plot azimtuhal angle on the X axis and metallicity excess on the Y axis for each of the different age bins. We show that the oldest age group has the most extreme variations in terms of  $\delta[\text{Fe/H}]$ , with the youngest stars showing minimal azimuthal variations. Interestingly, there seems to be a phase offset in the curves among the different age bins. This phase offset is also seen in Figure 12 of Carr et al. (2022). For all age groups, the variations are minimized at  $\phi(\pi)=1$ , which most likely due to the sample selection and the number of stars located along the line of sight from the Sun towards the center of the Galaxy. The future releases of SDSS V will provide better radial and azimuthal coverage to alleviate this problem.

We find that the strongest azimuthal [Fe/H] variations 794 are apparent in solar-age stars. Additionally, the metal-795 licity excess is largest in the oldest stars (6+ Gyr) and 796 similar in the youngest (0-2 Gyr) and intermediate aged 797 (2-6 Gyr) stars, with a phase offset that differentiates 798 between them. These results are in line with the simu-799 lated results of Debattista et al. (2024). The authors 800 found that azimuthal metallicity variations were still 801 present when excluding younger populations, whereas 802 we find the strongest azimuthal metallicity variations 803 in our older populations. Although azimuthal vartia-804 tions have been quantified in younger populations (e.g., 805 the OBA sample in Hawkins 2022), the presence of the 806 strongest azimuthal variations in our older populations 807 suggest an important contribution from dynamical pro-808 cesses to the creation of azimuthal metallicity variations 809 throughout the Galactic disk. Thus, in the following 810 section, we delve into possible dynamical origins.

#### 4.5. Linking Chemistry to Dynamics

The driving mechanisms behind the chemical azimuthal substructure we see in the Galactic disk can be predominantly categorized into either radial migra-

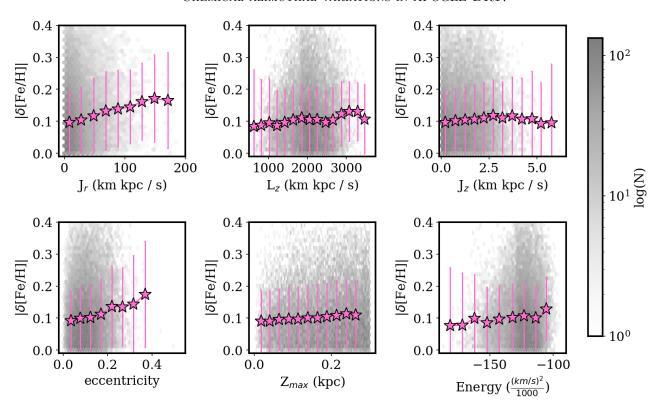


Figure 10. The correlation between absolute metallicity excess ( $|\delta[\text{Fe/H}]|$ ) on the y axes and different dynamical properties on the x axes. The hexagonal bins are colored by density with the pink stars representing the running medians in each panels with error bars showing the standard deviation. In the  $J_r$  and eccentricity panels, we see a clear increasing trend hinting that stars with high eccentricities and high radial actions contribute to the variations with the largest magnitudes.

815 tion/churning or kinematic heating/blurring. Radial 816 migration (Sellwood & Binney 2002) (also sometimes 817 known as cold torquing), hereafter referred to as churn-818 ing, is a dynamical process which altering stellar or-819 bits through the interactions with non-axisymmetric fea-820 tures in the disk. More specifically, churning results in change in the stellar orbits by changing the angular 822 momenta of the orbits without adding any excess en-823 ergy (e.g., stars on circular orbits experiencing churning will not see a change in eccentricity, only radii) (Carr 825 et al. 2022). Blurring (also known as kinematic heating) 826 heats an orbit without an increase in angular momenta of stellar orbits (i.e., there is an increase in eccentricity). Churning and blurring could be due to interactions with the spiral arms of the Galaxy (Jenkins & Binney 1990), 830 interactions with the bar (Schönrich & Binney 2009; Filion et al. 2023), and interactions with a satellite galaxy 832 like Sgr (Carr et al. 2022).

In this section, we aim to provide the initial basis of linking the azimuthal variations with Galactic dynamical cal properties to help elucidate the mechanisms responsible for these observed chemical signatures. While the goal of this paper is not to define the exact mechanisms responsible for azimuthal chemical variations, this ini-

tial linking to dynamics will motivate future studies. The dynamical properties of these stars could point to whether churning or blurring has a greater impact on the magnitude of the azimuthal variations. Thus we plot the determined radial action  $(J_r)$ , angular momentum  $(L_z)$ , vertical action  $(J_z)$ , orbital eccentricity (e), maximum height above the Galactic plane  $(Z_{max})$ , and total energy to plot against absolute metallicity excess  $(|\delta[Fe/H]|)$ .

Figure 10 is a six panel plot showing the aforementioned parameters versus  $|\delta[\text{Fe/H}]|$ . The bins are colored by density with the pink stars representing the running medians in each panels with error bars showing the standard deviation. There is a slight positive trend with  $|\delta[\text{Fe/H}]|$  and  $J_r$ , suggesting that stars with a larger metallicity excess tend to have higher radial actions. This correlation with radial action follows the predictions of the Debattista et al. (2024) simulation<sup>2</sup>. There are no obvious trends with the angular momen-

 $<sup>^2</sup>$  Note that those authors suggest that  ${\rm J}_r$  should be computed as a time-averaged quantity. We do not investigate the impact of time-averaging, and leave further investigation of this to future work.

tum, vertical action, maximum height, and total energy of stars when comparing to  $|\delta[\text{Fe/H}]|$ .

With the first and third panels of Figure 10, we can make inferences about the relative importance of churning and blurring assuming that there is an initial negative metallicity gradient and no natal azimuthal variations. If churning is the primary mechanism causing azimuthal variations, we would predict that the largest  $|\delta[\text{Fe/H}]|$  would be present at smaller eccentricities. This is because churning would largely cause stars to migrate from nearly circular orbits to nearly circular orbits of a different radius. If blurring is the primary mechanism causing azimuthal variations, we would expect the largest  $|\delta[\text{Fe/H}]|$  values at larger eccentricities (and radial actions) because blurring is generally characterized by an increase in eccentricity.

In the first and third panels of Figure 10, we see larger  $|\delta[{\rm Fe/H}]|$  at higher eccentricities and radial actions. This implies that blurring (heating of orbits) contributes a non-negligible amount to the mechanisms that are causing these observed azimuthal metallicity variations. In reality, multiple Galactic mechanisms are likely responsible for the azimuthal metallicity variations and more work is required to pin down mechanism responsible for the variations observed in this work and others (e.g., Hawkins 2022; Poggio et al. 2022; Imig et al. 2023).

While the goal of this paper is not to define the exact mechanisms responsible, this initial linking to dynamism ics will motivate future work on mock observations of simulations to directly compare different dynamical processes. We show that working in orbit-space, rather than present-day positions, can provide valuable insight into the dynamical properties that are driving the chemistry of the Galactic disk.

## 5. SUMMARY

Stellar spectroscopists have been using chemical car-893 894 tography to map the chemistry of the Milky Way for the 895 past decade (Hayden et al. 2015). With the chemistry 896 of stars, we are able to discern what observable patterns 897 are present throughout the Galactic disk and identify 898 the key processes that contribute to Galactic evolution. 899 Some of the most prominent patterns throughout the 900 disk are the negative vertical and radial metallicity gra-901 dients (Mayor 1976; Andrievsky et al. 2002; Magrini 902 et al. 2009; Luck & Lambert 2011; Bergemann et al. 903 2014; Xiang et al. 2015; Yan et al. 2019; Hawkins 2022). These gradients support theories such as 'inside-out' for-905 mation, which would cause the stars in the Galactic cen-906 ter to form first and ignite rapid star formation, with 907 the outer Galaxy forming at later times. This will lead 908 to the inner Galaxy having more metal-rich stars and

the outer Galaxy to be populated by metal-poor stars. Hiding under the radial  $\Delta [{\rm Fe/H}]/\Delta {\rm R}$  gradient are azimuthal variations on the scale of  $\sim 0.1$  dex (Poggio et al. 2022; Hawkins 2022). To constrain the properties of these azimuthal variations, we aim to confirm the [Fe/H] azimuthal variations in *APOGEE* DR17, characterize how the variations interplay with stellar age, identify if azimuthal variations exist in elements other than [Fe/H], and attempt to link the azimuthal variations to dynamical properties.

To accomplish these tasks, we model the linear radial metallicity gradient in the planar thin disk and subtract this off from the actual stellar metallicities to look for azimuthal substructure in [Fe/H] and  $\alpha$ -elements. We then separate our sample by age in to 3 subsets: 0-2 Gyr, 2-6 Gyr, and 6+ Gyr. We then quantified how metallicity excess ( $\delta$ [Fe/H]) changes with respect to radial action ( $J_r$ ), angular momentum ( $L_z$ ), vertical action ( $J_z$ ), orbital eccentricity (e), maximum height above the Galactic plane ( $Z_{max}$ ), and total energy. Based on these methods, we present the following results:

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- 1. A radial metallicity gradient ( $\Delta [Fe/H]/\Delta R$ ) of  $\sim -0.066 \pm 0.0004$  dex/kpc is found throughout the kinematic thin disk ( 3.5 < R < 15.5 kpc and  $Z_{max} < 0.3$  kpc) of the Milky Way. We additionally find a vertical metallicity gradient ( $\Delta [Fe/H]/\Delta Z$ ) of  $\sim -0.164 \pm 0.001$ . Both gradients are in fairly good agreement with previous studies that use different tracer populations (Section 4.1).
- 2. Azimuthal variations are found throughout the disk in [Fe/H]. These deviations were quantified by modelling the linear gradient and subtracting the model from the data to illustrate the [Fe/H]-poor to [Fe/H]-rich oscillating azimuthal pattern in the disk. These results are consistent with other azimuthal studies that have used different datasets or tracer populations (Eilers et al. 2022; Hawkins 2022). While previous studies have found a correlation with azimuthal metallicity variations and the location of the spiral arms, we find some deviations to that in this work. We argue this is worth further investigation over a wider range of azimuthal angles (Hackshaw et al. in prep) (Section 4.2).
- 3. Azimuthal substructure is seen in elements beyond iron in  $[\alpha/\text{Fe}]$ . While [Fe/H] exhibits the strongest variations from the gradient, the  $\alpha$ -elements (O, Mg, Si, S) still show azimuthal variations on the

 $\sim 0.1$  dex-level. These  $\alpha$ -abundances manifest as anti-correlations with [Fe/H], following the predictions by the production mechanisms between  $\alpha$  and Fe-peak elements (Section 4.3).

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- 4. Azimuthal substructure varies by stellar age. Older populations (6+ Gyr) exhibit more extreme deviations from the radial metallicity gradient than seen in younger (0-2 Gyr) and intermediate (2-6) populations, which suggests that the azimuthal metallicity variations largely stem from dynamical processes rather than natal processes. This age dependence can be coupled with mock observed simulations to identify why the older populations are disproportionately affected by the cause of these azimuthal variations (Section 4.4).
- 5. There appear to be discernible trends with absolute metallicity excess ( $|\delta[\text{Fe/H}]|$ ) and different dynamical properties. There is a positive trend between  $|\delta[\text{Fe/H}]|$  and  $J_r$ , as well as  $|\delta[\text{Fe/H}]|$  and eccentricity. This implies that stars with high eccentricities and high radial actions contribute to the azimuthal variations with the largest magnitudes, hinting that blurring is an important dynamical process in the production of azimuthal [Fe/H] variations. Further investigation of these links can help definitively characterize the cause of azimuthal metallicity variations (Section 4.5).

Azimuthal chemical variations are an apparent feature of the thin disk of our Galaxy. However, the cause of this substructure is unknown, with hints lying among the age dependence of these variations, azimuthal variations in elements beyond iron, and the dynamics of these stars. To answer this open question, we recommend deeper exploration of azimuthal variations with data like the impending release of SDSS-V, providing better radial and azimuthal coverage of the Galaxy. Including the orbit-space perspective rather than only the positions is a key step to understanding what drives these processes. Connecting the observable signatures in the Galactic disk with with mock simulated observations will bring us one step closer conclusively identifying what drives Galactic evolution.

1001 Software: astropy (Astropy Collaboration et al. 1002 2022, 2013), astronn (Leung & Bovy 2019a), gala 1003 (Price-Whelan et al. 2020), PyAstronomy (Czesla et al.

2019), scipy (Virtanen et al. 2020), matplotlib.pyplot (Hunter 2007), numpy (van der Walt et al. 2011)

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SDSS is managed by the Astrophysical Research Con-1036 sortium for the Participating Institutions of the SDSS Collaboration, including the Carnegie Institution for 1038 Science, Chilean National Time Allocation Committee (CNTAC) ratified researchers, the Gotham Participa-1040 tion Group, Harvard University, Heidelberg University, 1041 The Johns Hopkins University, L'Ecole polytechnique 1042 fédérale de Lausanne (EPFL), Leibniz-Institut für As-1043 trophysik Potsdam (AIP), Max-Planck-Institut für As-1044 tronomie (MPIA Heidelberg), Max-Planck-Institut für 1045 Extraterrestrische Physik (MPE), Nanjing University, 1046 National Astronomical Observatories of China (NAOC), 1047 New Mexico State University, The Ohio State Univer-1048 sity, Pennsylvania State University, Smithsonian Astro-1049 physical Observatory, Space Telescope Science Institute 1050 (STScI), the Stellar Astrophysics Participation Group, 1051 Universidad Nacional Autónoma de México, University 1052 of Arizona, University of Colorado Boulder, University 1053 of Illinois at Urbana-Champaign, University of Toronto, 1054 University of Utah, University of Virginia, Yale Univer-1055 sity, and Yunnan University.

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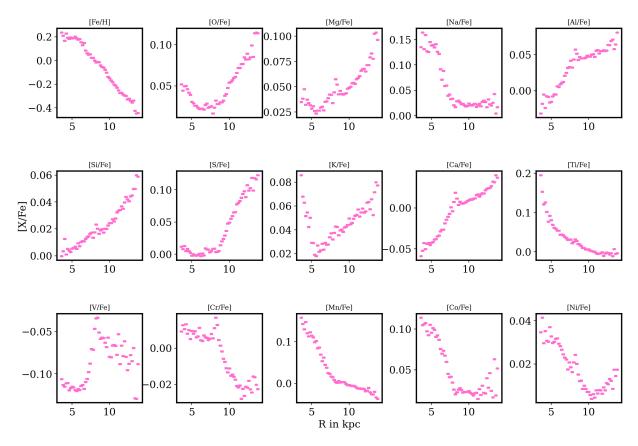


Figure 11. The running median [X/Fe] abundances in 0.2 kpc radial bins for 15 elements in our sample. The high-fidelity elements ([Fe/H], [O/Fe], and [Mg/Fe]) are featured on the top row. A number of these elements would not be best represented by a linear gradient, so we leave a majority of them out of our analysis. The elements included when searching for azimuthal variations in the linear gradient are the approximately linear elements: Fe, O, Mg, Si, S. To look at the  $\alpha$ -element behavior more holistically, we averaged the  $\alpha$ -element abundances in Figure 7.

1056 APPENDIX

## A. THIN DISK SAMPLE DETERMINATION

The thick-to-thin disk probability ratios (TD/D) were found for each star using the assumption that the space velocities U, V, and W have Gaussian distributions (Bensby et al. 2014),

$$f = k \cdot \exp\left(-\frac{(U_{\rm LSR} - U_{\rm asym})^2}{2\sigma_{\rm U}^2} - \frac{(V_{\rm LSR} - V_{\rm asym})^2}{2\sigma_{\rm V}^2} - \frac{W_{\rm LSR}^2}{2\sigma_{\rm W}^2}\right)$$
(A1)

where the normalization factor k is given by

$$k = \frac{1}{(2\pi)^{3/2} \sigma_{\mathcal{U}} \sigma_{\mathcal{V}} \sigma_{\mathcal{W}}} \tag{A2}$$

The characteristic velocity dispersions are represented by  $\sigma_{\rm U}, \sigma_{\rm V}$  and  $\sigma_{\rm W}$  and  $U_{\rm asym}$  are the asymmetric drifts. Finally the thick (TD) to thin disk (D) ratio is found with

$$TD/D = \frac{X_{TD}}{X_{D}} \cdot \frac{f_{TD}}{f_{D}} \tag{A3}$$

We characterize thin disk stars as any star with a thin disk probability  $\geq 80\%$ .

#### B. ELEMENTAL GRADIENTS

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This appendix contains the [X/Fe] median abundance trends in 0.2 kpc bins in Figure 11. Note the y-axis of the panels vary, but the elemental families show general trends as expected, such as the Fe-peak (Cr, Mn, Fe, Co, and Ni) exhibiting negative gradients. The  $\alpha$ -elements (O, Mg, Si, S, Ca) all have positive trends. These gradients match well with shape and scale to other studies that have looked at abundance trends throughout the Galactic disk (for example, Figure 7 in Eilers et al. 2022).

Due to the fact that azimuths are not sampled isotropically in the 0.2 kpc radial bins, the methodology of calculating  $\Delta$  [Fe/H] from the abundance medians is inaccurate. Thus, we only choose elements with roughly linear and monotonic shapes to complete our analysis (Fe, O, Mg, Si, S).

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