

FILTERING FREQUENCIES IN A SHIFT-AND-INVERT LANCZOS ALGORITHM FOR THE DYNAMIC ANALYSIS OF STRUCTURES*

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Abstract. The shift-and-invert Lanczos algorithm is a commonly used solution procedure to compute the eigenpairs of large, sparse eigenvalue problems that arise when approximating the elastic dynamic response of large structures under the influence of seismic forces. Not all eigenvectors are equally important to the response when the structure is exposed to a mass-dependent external force of the form $g(t)Mb$, where M is the mass matrix of the system and b the rigid body vector. Structural engineers select eigenvectors x_j , $j = 1, \dots, \ell$, such that their cumulative mass participation, measured as $\sum_{j=1}^{\ell} (x_j^T Mb)^2 / (b^T Mb)$, is above a target threshold ξ . We show that when the starting vector for the unshifted Lanczos algorithm is the spatial distribution vector b , the Lanczos procedure can be used to provide an estimate of the cumulative mass participation. This allows us to identify intervals containing eigenvalues whose eigenvectors have a large contribution to the cumulative mass participation and filter out intervals containing eigenvalues whose eigenvectors have a negligible contribution. We use this information to devise a sequence of shifts $\sigma_1, \dots, \sigma_p$ for the shift-and-invert Lanczos algorithm as well as a stopping criterion for the iteration with shift σ_i so that the cumulative mass participation of the computed eigenvectors reaches the required level ξ . Numerical experiments on real engineering problems show that our approach computes up to 70% fewer eigenvectors and requires fewer shifts, on average, than the more general shifting strategy proposed by Ericsson and Ruhe [*Math. Comp.*, 35 (1980), pp. 1251–1268] together with its modification presented in Grimes, Lewis, and Simon [*SIAM J. Matrix Anal. Appl.*, 40 (1994), pp. 228–272].

Key words. shift-and-invert Lanczos algorithm, symmetric generalized eigenvalue problem, shifting strategy, structural analysis, orthogonal polynomials

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1. Introduction. A structural dynamics problem consists of finding the response of a structure, for instance, a building or a bridge, given some dynamic loading. Such problems may be written in the form of a system of second order differential equations,

$$(1.1) \quad M\ddot{u}(t) + D\dot{u}(t) + Ku(t) = f(t),$$

that results from the finite element discretization of the equation of motion, together with some initial conditions. The mass matrix $M \in \mathbb{R}^{n \times n}$ is usually symmetric positive semidefinite (denoted by $M \geq 0$), the stiffness matrix $K \in \mathbb{R}^{n \times n}$ is symmetric positive definite ($K > 0$), the damping matrix D is symmetric and often positive definite, $u(t)$ is the displacement, and $f(t)$ is the time-dependent external load on the given structure. Here, we concentrate on external forces of the form

$$(1.2) \quad f(t) = g(t)Mb,$$

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where $g(t)$ is a scalar function and $0 \neq b \in \mathbb{R}^n$ is the spatial distribution vector (also called rigid body vector or spatial vector of loading patterns). External forces of the form (1.2) are particular to earthquake loading, where $g(t)$ is the input earthquake acceleration.

Projection methods are usually employed to reduce the dimension n of the system (1.1). These methods consist of constructing a matrix $X_\ell \in \mathbb{R}^{n \times \ell}$ of full rank and transforming (1.1) into the reduced system

$$(1.3) \quad X_\ell^T M X_\ell \ddot{v}(t) + X_\ell^T D X_\ell \dot{v}(t) + X_\ell^T K X_\ell v(t) = g(t) X_\ell^T M b,$$

which is then solved for $v(t)$. An approximate solution to (1.1) is obtained as $u(t) \approx X_\ell v(t)$. In the mode superposition method, the columns x_1, \dots, x_ℓ of X_ℓ are eigenvectors corresponding to finite eigenvalues of the associated generalized eigenvalue problem (GEP)

$$(1.4) \quad (K - \lambda M)x = 0,$$

where we assume generalized proportional damping, so that the damping matrix D in (1.1) is diagonalizable by the matrix of eigenvectors of (1.4). If the columns of X_ℓ are M -orthonormal, i.e., $x_i^T M x_j = \delta_{ij}$ with δ_{ij} the Kronecker delta, then the reduced system (1.3) can be rewritten as ℓ decoupled second order differential equations,

$$(1.5) \quad \ddot{v}_j(t) + 2\zeta_j \omega_j \dot{v}_j(t) + \omega_j^2 v_j(t) = g(t) x_j^T M b, \quad j = 1, \dots, \ell,$$

where $\omega_j = \sqrt{\lambda_j}$, $\lambda_j > 0$ is an eigenvalue of (1.4) with corresponding eigenvector x_j and $x_i^T D x_j = 2\zeta_j \omega_j \delta_{ij}$ for some $\zeta_j \geq 0$ [20, Chap. 18].

Using the mode superposition method to solve the system of second order differential equations (1.1) in the time domain remains an expensive process both in terms of time and memory. Also, the information provided by the solution is usually excessive for the purpose of our engineering application. The method of response spectrum analysis (RSA) has been developed as a reliable but much cheaper method for the analysis of structures [3, Chap. 13, Part B]. As for the mode superposition method, the method of RSA requires a partial solution to the GEP (1.4). But once enough eigenpairs have been computed to adequately describe the behavior of the structure subjected to some earthquake ground motion, then instead of integrating the decoupled equations in (1.5), the process of RSA uses the computed eigenpairs to approximate the peak responses of the degrees of freedom of the structure using well-defined combination rules [3, sect. 13.7]. Here we concentrate on the computation of the eigenpairs.

Not all eigenpairs are equally important to a given system of differential equations, and it is clear from (1.5) that the response v_j depends on both the frequencies ω_j and the magnitude of $x_j^T M b$, which is called the *mass participation factor* of x_j . They satisfy

$$(1.6) \quad \sum_{j=1}^n \frac{(x_j^T M b)^2}{b^T M b} = 1$$

(see section 2.1). Within the context of seismic analysis and design, adhering to international standards [7, sect. 4.3.3.3.1] structural engineers aim to achieve at least 90% of mass participation in (1.6) using only a subset of the eigenvectors (see the justification in section 2.1). This leads to the following problem.

PROBLEM 1.1. For a given proportion $\xi \in (0, 1)$ and a spatial distribution vector b , find the smallest number of M -orthonormal eigenvectors x_{i_k} , $k = 1, \dots, \ell$, of (1.4) such that

$$(1.7) \quad \sum_{k=1}^{\ell} \frac{(x_{i_k}^T M b)^2}{b^T M b} \geq \xi,$$

where $\{i_1, \dots, i_{\ell}\} \subseteq \{1, \dots, n\}$.

Problem 1.1 is easy to solve if we can compute all the eigenvectors of the GEP (1.4), but this is not feasible for problems of large dimensions. Note that the solution may not be unique. It is usually the eigenvectors corresponding to small eigenvalues that contribute the most to the total mass participation (see, for example, [2] or section 2.1), so in previous work Problem 1.1 was relaxed to the following problem.

PROBLEM 1.2. For a given proportion $\xi \in (0, 1)$ and a spatial distribution vector b , find the M -orthonormal eigenvectors x_i , $i = 1, \dots, \ell$ associated with the ℓ smallest eigenvalues of (1.4), where ℓ is the smallest integer such that

$$(1.8) \quad \sum_{i=1}^{\ell} \frac{(x_i^T M b)^2}{b^T M b} \geq \xi.$$

Note that adding as an extra constraint an upper bound on the eigenvalues or frequencies may be necessary in practice—see section 3.4.

Recovering sufficient mass participation can be a challenge depending on the geometry of the structure and orientation of b . The vector b represents the “rigid body” deformation of the structure in cartesian x -, y -, and z -directions. For instance a given structure may have a geometry that makes it easy to solve Problem 1.2 with only a handful of smallest eigenvectors (i.e., ℓ is small) when b is along the x -direction but might require a large value of ℓ when b is along z . This is graphically illustrated in Figure 1, which shows the first and twentieth eigenvectors, x_1 and x_{20} , overlaid as displacements on the structure’s geometry. As can be seen from the diagram in the middle, x_1 is a sway mode wherein large parts of the structure deform in the Cartesian x -direction. By contrast, x_{20} (right of middle) is a bouncing mode and therefore the structure deforms significantly in the z -direction. The mass participations of these eigenvectors are tabulated in Table 1 for b in the x - and z -directions.

A natural approach to solve Problem 1.2 is to apply the shift-and-invert Lanczos algorithm (see [1], [6], [21] and references therein) to $K - \lambda M$ with or without shifts depending on how many eigenvectors will be needed to satisfy (1.8). Instead of using eigenvectors in (1.8), Wilson, Yuan, and Dickens [24] proposed to use the Ritz vectors resulting from a variant of the Arnoldi algorithm later known as the Wilson–Yuan–Dickens (WYD) algorithm [12], [23]. The use of the Lanczos vectors was later suggested by Nour-Omid and Clough [17]. But as noted in [2], there is no analysis to support the use of Lanczos vectors and Ritz vectors (unless they have converged to eigenvectors of $K - \lambda M$). In particular, we show in section 2 that under suitable conditions on $g(t)$ and the initial conditions for (1.1), the quantity $(1 - \xi)^2$ provides an upper bound on the relative error between the response $u(t)$ to the ODE (1.1) and its approximation by $X_{\ell} v(t)$ when the columns of X_{ℓ} are eigenvectors chosen such that (1.7) or (1.8) holds.

Given an initial shift σ_1 , a common strategy to determine the sequence of shifts $\sigma_2, \sigma_3, \dots$ to be employed in a shift-and-invert Lanczos process for symmetric GEPs is to choose the new shift σ_i , $i > 1$, such that the largest converged eigenvalue λ_{\max} is

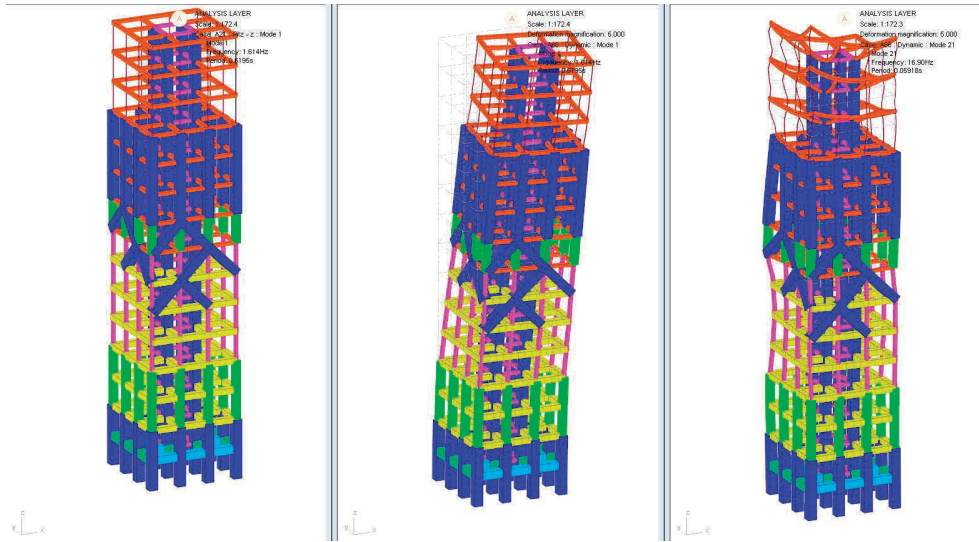


FIG. 1. The eigenvectors (modes of vibration) of a tall building structure. From left to right: the undeformed shape of the structure; its first eigenvector, x_1 , corresponding to the smallest eigenvalue, which is predominantly a “sway mode” (i.e., large mass participation factor in the Cartesian x -direction); and another eigenvector, x_{20} , which is a “bouncing mode” (large participation along cartesian z -direction). Deformations are exaggerated.

TABLE 1
Mass participation factors for M -orthonormal eigenvectors x_1 and x_{20} from Figure 1.

Eigenvector x_j	Mass participation factor $(x_j^T M b)^2 / (b^T M b)$	
	b along Cartesian x	b along Cartesian z
x_1	0.199	10^{-6}
x_{20}	10^{-8}	0.209

halfway between the old shift σ_{i-1} and the new shift σ_i , namely, $\sigma_i = 2\lambda_{\max} - \sigma_{i-1}$ [6]. This way, if the eigenvalues are roughly evenly distributed across the spectrum then the same number of eigenvalues is expected to converge to the left of the new shift as the number of eigenvalues that have converged to the right of the old shift. The Lanczos iteration with shift σ_i is then stopped when all eigenvalues between the shifts σ_{i-1} and σ_i have been found. One issue with this shifting strategy for Problem 1.2 is that the shift-and-invert Lanczos algorithm may return too many eigenvectors with small normalized mass participation. As a result, the number ℓ of returned eigenpairs may be unnecessarily large, whereas eigenvectors with large normalized mass participations corresponding to eigenvalues that are not in the lower end of the spectrum may not be detected.

Our main contribution is the presentation of a new shifting strategy for the shift-and-invert Lanczos algorithm together with a stopping criterion for the iteration with shift σ_i that are specifically designed to approximate the solution to Problem 1.1. For this, we use the theory of orthogonal polynomials to show that a few steps of the unshifted inverse Lanczos algorithm applied to $K^{-1}M$ with starting vector b provides, at no additional cost, information about the location of the eigenvalues whose corresponding eigenvectors have non-negligible mass participation and also helps to

identify intervals where eigenvalues have negligible mass participation. We use this information to devise a shifting strategy for the shift-and-invert Lanczos process so that condition (1.7) holds, albeit perhaps not for the smallest number ℓ of eigenvectors. This shifting strategy performs especially well in the cases where the eigenvectors with non-negligible mass participation correspond to larger eigenvalues while there are intervals of smaller eigenvalues whose corresponding eigenvectors have negligible contribution to the overall response of the structure. These are the cases that are the most problematic for available methods.

Numerical experiments performed on real structural engineering problems show an often large reduction in the number ℓ of eigenvectors computed to approximate the solution to Problem 1.1 using our new shifting strategy as opposed to the number of eigenvectors computed to approximate the solution to Problem 1.2 using the more general shifting strategy of Ericsson and Ruhe [6] together with its modification presented in Grimes, Lewis, and Simon [11]. The use of a new shift in the shift-and-invert Lanczos process has a cost since it leads to a new matrix factorization. For our set of test problems, our numerical experiments show that the number of shifts used with our new shifting strategy is, on average, smaller than with Ericsson and Ruhe's strategy.

Unlike previous attempts to solve Problems 1.1–1.2, we pay special attention to issues that can occur when the mass matrix M is singular. In this case, the Lanczos process proceeds with a quasi-inner product. In finite precision arithmetic, the computed Lanczos vectors and Ritz vectors have components in the null space of M and the magnitude of these unwanted components grows during the iterations. This can either delay or prevent convergence of the Ritz vectors [5]. These issues are overcome with the use of an appropriate starting vector and implicit filtering [16].

In the next section we give some preliminary material that includes a justification of Problem 1.1 and a discussion of issues that arise when using a quasi-inner product in a shift-and-invert Lanczos process, as well as possible remedies. We describe and justify our new shifting strategy in section 3 and illustrate its performance on a number of real structural engineering problems in section 4.

2. Preliminaries. Following [2], we show in section 2.1 that, under some assumptions on the initial conditions and the input function $g(t)$, choosing ξ close to 1 in (1.7) guarantees a small error between the exact response $u(t)$ and its approximation as a linear combination of the eigenvectors x_{i_k} , $k = 1, \dots, \ell$ satisfying (1.7).

We recall in section 2.2 the shift-and-invert Lanczos process for the GEP (1.4) and discuss issues related to the use of a quasi-inner product defined by the symmetric semidefinite matrix M .

2.1. Upper bound on the response error. The $n \times n$ mass matrix M in (1.1) is often singular in applications. As a result, the GEP (1.4) has the eigendecomposition

$$X^T(K - \lambda M)X = \Lambda - \lambda \begin{bmatrix} I_r & \\ & 0 \end{bmatrix},$$

where I_r is the $r \times r$ identity matrix with

$$r := \text{rank}(M) \leq n,$$

$\Lambda = \text{diag}(\lambda_1, \lambda_2, \dots, \lambda_n)$ has real positive diagonal entries displaying the r finite eigenvalues as λ_j , $j = 1, \dots, r$ (the remaining $n - r$ eigenvalues being at infinity), and X is a nonsingular matrix containing the corresponding eigenvectors x_1, \dots, x_n . Note

that the eigenvectors x_j , $j = r + 1, \dots, n$ associated with the eigenvalues at infinity belong to the null space of M , i.e.,

$$(2.1) \quad Mx_j = 0, \quad j = r + 1, \dots, n.$$

Let

$$(2.2) \quad u(t) = Xv(t) = \sum_{j=1}^n v_j(t)x_j,$$

where $v_j(t)$ is the j th entry of the vector $v(t)$. We assume generalized damping so that $X^T DX = 2\text{diag}(\zeta_1\omega_1, \dots, \zeta_n\omega_n)$ for some $\zeta_j \geq 0$ and $\omega_j = \sqrt{\lambda_j}$, $j = 1, \dots, n$, and rewrite (1.1) as

$$(2.3) \quad \begin{aligned} \ddot{v}_j(t) + 2\zeta_j\omega_j\dot{v}_j(t) + \omega_j^2 v_j(t) &= g(t)x_j^T Mb, & j = 1, \dots, r, \\ v_j(t) &= 0, & j = r + 1, \dots, n. \end{aligned}$$

The system of uncoupled equations (2.3) can then be solved by direct integration, yielding

$$(2.4) \quad \begin{aligned} v_j(t) &= (x_j^T Mb) \int_0^t \frac{e^{-\zeta_j\omega_j(t-s)}}{\tilde{\omega}_j} \sin(\tilde{\omega}_j(t-s)) g(s) ds, & j = 1, \dots, r, \\ v_j(t) &= 0, & j = r + 1, \dots, n, \end{aligned}$$

where we let $\tilde{\omega}_j = \omega_j(1 - \zeta_j^2)^{1/2}$ and assume for simplicity that $u(0) = \dot{u}(0) = 0$ so that $v(0) = \dot{v}(0) = 0$.

As mentioned in the introduction, in practical applications, n is large and it is infeasible to compute all n eigenpairs of (1.4) so the solution $u(t)$ in (2.2) is approximated instead by

$$(2.5) \quad \tilde{u}(t) = \sum_{j=1}^{\ell} v_j(t)x_j, \quad \ell \leq r.$$

We use the M -quasi vector norm,

$$\|y\|_M = \langle y, y \rangle_M^{1/2} = (y^T My)^{1/2},$$

to measure the relative error between the exact solution and its approximation,

$$(2.6) \quad \frac{\|u(t) - \tilde{u}(t)\|_M}{\|u(t)\|_M} = \frac{(\sum_{j=\ell+1}^r v_j^2(t))^{1/2}}{(\sum_{j=1}^r v_j^2(t))^{1/2}}.$$

We rewrite $v_j(t)$ in (2.4) as $v_j(t) = h_j(t)x_j^T Mb$. If the spatial distribution vector b has no components in the null space of M (this is usually the case), i.e., $b = \sum_{j=1}^r b_j x_j$, then

$$(2.7) \quad \sum_{j=1}^r (x_j^T Mb)^2 = \sum_{j=1}^r b_j^2 = b^T Mb,$$

which on using (2.1) is equivalent to (1.6). Now if we assume that $h_{\min} \leq |h_j(t)| \leq h_{\max}$ for $t > 0$, $j = 1, \dots, r$, and some positive scalars h_{\min}, h_{\max} , then it follows from (2.6) and (2.7) that

$$\begin{aligned} \frac{\|u(t) - \tilde{u}(t)\|_M}{\|u(t)\|_M} &\leq \frac{h_{\max}}{h_{\min}} \frac{(\sum_{j=\ell+1}^r (x_j^T M b)^2)^{1/2}}{(\sum_{j=1}^r (x_j^T M b)^2)^{1/2}} = \frac{h_{\max}}{h_{\min}} \left(1 - \frac{\sum_{j=1}^{\ell} (x_j^T M b)^2}{b^T M b}\right)^{1/2} \\ &\leq \frac{h_{\max}}{h_{\min}} (1 - \xi)^{1/2}, \end{aligned}$$

where ξ is as in (1.7). Thus, under the above assumptions on the initial conditions $u(0)$, $\dot{u}(0)$, and the functions h_j , choosing a ξ close to 1 guarantees a small relative error between the exact solution u and its approximation \tilde{u} . We refer to section 4 for an illustration of the above analysis.

Finally, we note that the factor $1/\omega_j$ in (2.4) suggests that the eigenvectors corresponding to smaller eigenvalues ω_j^2 (i.e., lower frequencies) are more likely to have a larger contribution to the response $u(t)$ in (2.2) than those corresponding to the higher frequencies.

2.2. Shift-and-invert Lanczos process with semidefinite inner product.

Applying the Lanczos algorithm with an M -quasi inner product to compute approximate eigenpairs of the definite pencil $K - \lambda M$ with $M \geq 0$ was first suggested by Scott [22]. A nontrivial null space of M may slow down the convergence of Ritz pairs and other choices for a norm may be considered [13], however the use of the M -quasi inner product restricts, at least in finite precision, the computation of eigenvectors to those that are in the range of M . These are exactly the eigenvectors corresponding to the finite eigenvalues. As is discussed later, the numerical issues arising from semidefiniteness of M can be detected and corrected at a reasonable cost. The M -quasi inner product is also used in quantifying the contribution of an eigenvector to the overall response of a structure, given a particular loading, and is therefore required to define the mass participation factor of an eigenvector.

For the Lanczos algorithm with an M -quasi inner product, given a shift σ near the eigenvalues of interest and a starting vector w , this Lanczos procedure constructs a matrix Q_k whose columns, called the Lanczos vectors, form an M -orthonormal basis for the k th order Krylov subspace

$$\mathcal{K}_k(A, w) = \text{span}\{w, Aw, \dots, A^{k-1}w\} = \text{span}\{q_1, q_2, \dots, q_k\},$$

where $A = (K - \sigma M)^{-1}M$. Given the first $k \geq 1$ M -orthonormal Lanczos vectors $q_1 = w/\|w\|_M, \dots, q_k$ and $q_0 = 0$, the next Lanczos vector q_{k+1} is obtained from the three-term recurrence

$$(2.8) \quad \beta_{k+1}q_{k+1} = (K - \sigma M)^{-1}Mq_k - \alpha_kq_k - \beta_kq_{k-1},$$

where $\alpha_k = q_k^T M(K - \sigma M)^{-1}Mq_k$, $\beta_1 = \|w\|_M = (w^T M w)^{1/2}$, and β_k for $k > 1$ is such that $\|q_k\|_M = (q_k^T M q_k)^{1/2} = 1$. The three-term recurrence (2.8) can be rewritten in matrix form as

$$(2.9) \quad (K - \sigma M)^{-1}MQ_k = Q_kT_k + \beta_{k+1}q_{k+1}e_k^T,$$

where $Q_k = [q_1, \dots, q_k] \in \mathbb{R}^{n \times k}$ and $T_k \in \mathbb{R}^{k \times k}$ is symmetric tridiagonal with

$$(T_k)_{jj} = \alpha_j, \quad j = 1, \dots, k, \quad (T_k)_{j,j+1} = \beta_{j+1}, \quad j = 1, \dots, k-1.$$

Since, by construction, q_{k+1} is M -orthogonal to the columns of Q_k and $Q_k^T M Q_k = I_k$, it follows from (2.9) that

$$(2.10) \quad Q_k^T M (K - \sigma M)^{-1} M Q_k = T_k.$$

Now, for an eigenpair $(\theta_j^{(k)}, s_j^{(k)})$ of T_k , the pair

$$(2.11) \quad (\lambda_j, x_j) := \left(1/\theta_j^{(k)} + \sigma, Q_k s_j^{(k)} \right)$$

is called a Ritz pair for $K - \lambda M$ and is considered as an approximate eigenpair of $K - \lambda M$ if the scaled residual (which approximates the backward error for (λ_j, x_j) [14, Thm. 2.1])

$$(2.12) \quad \eta(\lambda_j, x_j) = \frac{\|(K - \lambda_j M)x_j\|_2}{(\|K\|_1 + |\lambda_j| \|M\|_1) \|x_j\|_2}$$

is below a given tolerance.

An algorithm implementing the three-term recurrence (2.8) is provided in Algorithm 2.1. This is essentially the algorithm provided in [18]. Breakdown occurs when $\beta_{k+1} \leq 0$ in step 10. The number ℓ of eigenpairs returned will depend of the convergence criterion in step 3. We discuss the latter in section 3.3. Some more comments follow.

Algorithm 2.1. Shift-and-invert Lanczos algorithm.

This algorithm takes as input two $n \times n$ symmetric matrices $M \geq 0$ and $K > 0$, a shift $\sigma \geq 0$ such that $K - \sigma M$ is nonsingular, and a starting vector $0 \neq w \in \text{range}((K - \sigma M)^{-1} M)$. It returns ℓ eigenpairs (converged Ritz pairs) (λ_{i_j}, x_{i_j}) , $j = 1, \dots, \ell$.

- 1 $q_0 = 0$, $z = Mw$, $\beta_1 = \sqrt{w^T z}$
 - 2 Factor $K - \sigma M$.
 - 3 for $k = 1, 2, \dots$, until convergence
 - 4 $z = z/\beta_k$, $q_k = w/\beta_k$
 - 5 $v = (K - \sigma M)^{-1} z - \beta_k q_{k-1}$
 - 6 $\alpha_k = v^T z$
 - 7 $w = v - \alpha_k q_k$
 - 8 Reorthogonalize w against q_1, \dots, q_k with respect to $\langle \cdot, \cdot \rangle_M$ if necessary.
 - 9 $z = Mw$
 - 10 $\beta_{k+1} = \sqrt{w^T z}$
 - 11 Compute the k eigenpairs $(\theta_j^{(k)}, s_j^{(k)})$, $j = 1, \dots, k$ of $T_k = \text{tridiag}(\beta, \alpha, \beta)$ with $\alpha = [\alpha_1, \dots, \alpha_k]$ and if $k > 1$, $\beta = [\beta_2, \dots, \beta_k]$.
 - 12 Check for convergence of the Ritz pairs $(\lambda_j, x_j) = (\frac{1}{\theta_j^{(k)}} + \sigma, Q_k s_j^{(k)})$,
 $j = 1, \dots, k$, where $Q_k = [q_1, \dots, q_k]$.
 - 13 end
-

- (a) **Choice of starting vector and implicit filtering.** When M is singular, Nour-Omid et al. [18, sect. 2.2] recommend choosing a starting vector w in the range of $(K - \sigma M)^{-1} M$. Indeed, since $\mathbb{R}^n = \text{range}((K - \sigma M)^{-1} M) \oplus \text{null}((K - \sigma M)^{-1} M)$ and $\text{null}((K - \sigma M)^{-1} M) = \text{null}(M)$, we have that

$$\text{range}((K - \sigma M)^{-1} M) \cap \text{null}(M) = \{0\}.$$

For an eigenvector x_i of $K - \lambda M$ with finite eigenvalue λ_i , we have that

$$(2.13) \quad x_i \in \text{range}(K^{-1}M) = \text{range}((K - \sigma M)^{-1}M).$$

If the starting vector w for the shift-and-invert Lanczos procedure is in the range of $(K - \sigma M)^{-1}M$, then so is the first Lanczos vector q_1 and the subsequent Lanczos vectors q_2, \dots, q_k if operations are performed in exact arithmetic. As a result, the Ritz vectors lie in $\text{range}((K - \sigma M)^{-1}M)$. Now if $w \notin \text{range}((K - \sigma M)^{-1}M)$, then the Ritz vectors may have unwanted components in the null space of M , which in turn slows down their convergence. Note that in finite precision arithmetic, even if $w \in \text{range}((K - \sigma M)^{-1}M)$, rounding errors prevent the computed Lanczos vectors from staying in the range of $(K - \sigma M)^{-1}M$. The unwanted components in $\text{null}(M)$ are mainly introduced when solving the linear systems with $K - \sigma M$. Since this operation is performed at each iteration, the accumulation can be rapid. This is set off at the beginning when the starting vector w is constructed as the product of $(K - \sigma M)^{-1}M$ with some $y \in \mathbb{R}^n$, since this operation itself can introduce a null space component. The starting vector can be put in the right space explicitly by forming a particular projection matrix [5, sect. 2.3], or solving the linear system to higher precision, however both methods come at a substantial cost. Instead, as suggested by Meerbergen [16], we apply an implicit filter that alters the starting vector implicitly, producing Lanczos vectors lying in $\text{range}((K - \sigma M)^{-1}M)$. This approach is briefly discussed below.

For $q = q_R + q_N \in \mathbb{R}^n$ with $q_R \in \text{range}(M)$ and $q_N \in \text{null}(M)$, we have that $\|q\|_M = \|q_R\|_M$, i.e., the components in the null space of M are undetectable by the M -norm. So $\|q_N\|_2$ can be arbitrarily large even when $\|q\|_M = 1$. Hence, if at step k of the Lanczos process we have

$$(2.14) \quad \|q_k\|_2 > \text{tol} \|q_1\|_2,$$

for some tolerance $\text{tol} \gg 1$ (we choose $\text{tol} = 10^4$); then we apply implicit filtering. For this let

$$(2.15) \quad \underline{T}_k = \underline{V}_k R_k$$

be the QR factorization of $\underline{T}_k = [T_k, \beta_{k+1}e_k]^T \in \mathbb{R}^{k+1 \times k}$ with $R_k \in \mathbb{R}^{k \times k}$ upper triangular and $\underline{V}_k \in \mathbb{R}^{k+1 \times k}$ with orthonormal columns. Then $\tilde{Q}_k = Q_{k+1}\underline{V}_k$ has M -orthonormal columns, and if we let

$$\tilde{\underline{T}}_{k-1} = R_k \underline{V}_{k-1},$$

then it is not difficult to show that the matrix \tilde{T}_{k-1} obtained from $\tilde{\underline{T}}_{k-1}$ by deleting its last row is tridiagonal and satisfies the Lanczos recurrence

$$(2.16) \quad (K - \sigma M)^{-1}M\tilde{Q}_{k-1} = \tilde{Q}_{k-1}\tilde{T}_{k-1} + \tilde{\beta}_k\tilde{q}_ke_{k-1}^T,$$

where $\tilde{\beta}_k$ is such that $\|\tilde{q}_k\|_M = 1$. It is shown in [16, Thm. 3.1] that $\text{range}(\tilde{Q}_k) = \text{range}((K - \sigma M)^{-1}M\tilde{Q}_k)$. The implicit filter implicitly pre-multiplies the Lanczos vectors by $K^{-1}M$, thereby removing any components in $\text{null}(M)$. Since the dimension of the projection space is reduced by one,

we continue the Lanczos algorithm by forming the next Lanczos vector \tilde{q}_{k+1} from \tilde{q}_k as well as updating T_k in step 11 with \tilde{T}_{k-1} . Note that the implicit filtering described above is essentially one step of the unshifted QR algorithm applied to T_k . A Lanczos vector q_k with a large 2-norm means that q_k has large components in the null space of M . As a result, the smallest eigenvalue of T_k becomes close to zero and the unshifted QR step pushes that very small eigenvalue to the bottom of the tridiagonal matrix. Indeed, if T_k had a zero eigenvalue then after one step of unshifted QR, the last row of T_k would be zero. The construction of \tilde{T}_{k-1} corresponds to a deflation of the eigenpair of T_k with smallest eigenvalue, thereby removing the Ritz pair in (2.11) with a large (possibly infinite) eigenvalue.

- (b) **Testing for convergence of eigenpairs.** As explained earlier, we consider the Ritz pair

$$(\lambda_j, x_j) := \left(1/\theta_j^{(k)} + \sigma, Q_k s_j^{(k)} \right),$$

where $(\theta_j^{(k)}, s_j^{(k)})$ is an eigenpair of T_k , to have converged if its backward error $\eta(\lambda_j, x_j)$ in (2.12) is below a given tolerance.

We suggest to only compute $\eta(\lambda_j, x_j)$ when the Ritz pair (λ_j, x_j) is likely to have converged. This can be checked as follows. On using (2.9), we have that

$$\begin{aligned} (K - \lambda_j M)x_j &= (K - \sigma M)x_j - \frac{1}{\theta_j^{(k)}} Mx_j \\ &= \frac{1}{\theta_j^{(k)}} (K - \sigma M)(\theta_j^{(k)} x_j - (K - \sigma M)^{-1} Mx_j) \\ &= \frac{1}{\theta_j^{(k)}} (K - \sigma M)(Q_k T_k - (K - \sigma M)^{-1} M Q_k) s_j^{(k)} \\ &= -\frac{1}{\theta_j^{(k)}} (K - \sigma M) \beta_{k+1} q_{k+1} (e_k^T s_j^{(k)}) \end{aligned}$$

so that

$$\|(K - \lambda_j M)x_j\|_2 \leq \frac{\|K\|_2 + |\sigma| \|M\|_2}{|\theta_j^{(k)}|} |e_k^T s_j^{(k)}| \|\beta_{k+1}\| \|q_{k+1}\|_2.$$

If $s_j^{(k)}$ is normalized such that $\|s_j^{(k)}\|_2 = 1$, then

$$\|x_j\|_M^2 = x_j^T M x_j = s_j^{(k)T} Q_k^T M Q_k s_j^{(k)} = s_j^{(k)T} s_j^{(k)} = 1$$

and

$$1 = \|x_j\|_M^2 = |x_j^T M x_j| \leq \|x_j\|_2 \|M x_j\|_2 \leq \|x_j\|_2^2 \|M\|_2 \leq \sqrt{n} \|x_j\|_2^2 \|M\|_1,$$

where we used Cauchy–Schwarz for the first inequality. Without loss of generality, we can assume that $\|M\|_1 = \|K\|_1 = 1$ (if not replace $K - \lambda M$ by $K/\|K\|_1 - \tilde{\lambda} M/\|M\|_1$ with $\tilde{\lambda} = \lambda \|M\|_1/\|K\|_1$). Hence, $\|x_j\|_2 \geq n^{-1/2}$ and $\|K\|_2 + |\sigma| \|M\|_2 \leq n^{1/2} (\|K\|_1 + |\sigma| \|M\|_1)$. Now if $|\lambda_j| \approx |\sigma|$, then

$$(2.17) \quad \eta(\lambda_j, x_j) \lesssim n^{1/4} |e_k^T s_j^{(k)}| \|\beta_{k+1}\| \|q_{k+1}\|_2 / |\theta_j^{(k)}|.$$

All the quantities in the above approximate upper bound are readily available during the Lanczos steps including $\|q_{k+1}\|_2$ (see point (a) and (2.14)). Hence, we suggest to only compute $\eta(\lambda_j, x_j)$ when the upper bound in (2.17) is below, say, 10 times the tolerance on the backward error.

3. Shifting strategies for an approximate solution to Problems 1.1–1.2.

A considerable number of eigenvectors may be required to solve Problem 1.2 for problems where there are large normalized mass participation factors

$$(3.1) \quad m(x_j) := \frac{(x_j^T M b)^2}{b^T M b}$$

corresponding to eigenvalues λ_j that lie further away from the lower end of the spectrum. In this case, the Lanczos algorithm (Algorithm 2.1 with $\sigma = 0$ as we are aiming for the small eigenvalues first) becomes increasingly slow and memory intensive. So if condition (1.8) is not satisfied by the converged eigenvectors after a given number of, say, k_{\max} steps of Algorithm 2.1 with zero shift, we then restart Algorithm 2.1 with a sequence of shifts. As mentioned in the introduction, Ericsson and Ruhe [6] propose to use the following sequence:

$$(3.2) \quad \sigma_1 = \lambda_{\max}^{(0)} + \frac{\lambda_{\max}^{(0)}}{2}, \quad \sigma_i = \sigma_{i-1} + 2\delta_{i-1}, \quad i \geq 2,$$

where

$$\delta_{i-1} = \lambda_{\max}^{(i-1)} - \sigma_{i-1},$$

and $\lambda_{\max}^{(i)}$ is the largest converged eigenvalue when using shift σ_i . Following the suggestion of Grimes, Lewis, and Simon [11], we choose δ_{i-1} to be the maximum of its previous value and the distance between the old shift σ_{i-1} and the largest converged eigenvalue $\lambda_{\max}^{(i-1)}$, that is,

$$(3.3) \quad \delta_{i-1} = \max(\delta_{i-2}, \lambda_{\max}^{(i-1)} - \sigma_{i-1}),$$

with $\delta_0 = \lambda_{\max}^{(0)}/2$. Algorithm 2.1 with shift σ_i is stopped when all eigenvalues in the interval $[\sigma_{i-1}, \sigma_i]$ have converged, or once condition (1.8) is satisfied, in which case we terminate the computation. The shift-and-invert Lanczos algorithm with this shifting strategy, which we refer to as SIL, approximates the solution to Problem 1.2.

An issue with Problem 1.2 is that its solution may include many eigenvectors with very small or negligible normalized mass participation factors. In what follows, we show that Algorithm 2.1 with $\sigma = 0$ and starting vector w equal to the spatial distribution vector b provides at almost no additional cost information about where the eigenvalues associated with eigenvectors of non-negligible normalized mass participation lie, while at the same time identifies the parts of the spectrum that do not contribute much to the total mass participation. We use this information to devise a new shifting strategy.

3.1. Estimating the cumulative mass participation. Let q_1, q_2, \dots be the Lanczos vectors generated by the three-term recurrence (2.8) with $\sigma = 0$ and $0 \neq q_1 \in \text{range}(K^{-1}M)$. Each Lanczos vector q_i can be written as

$$(3.4) \quad q_i = p_i(K^{-1}M)q_1,$$

where p_i is called the i th Lanczos polynomial. It is well known from the theory of orthogonal polynomials [8, Chap. 2], [10, Chap. 4] that these polynomials are orthogonal with respect to the inner product defined in terms of the Riemann–Stieltjes integral

$$(3.5) \quad \langle p_i, p_j \rangle_\phi := \int_a^b p_i(\mu) p_j(\mu) d\phi(\mu),$$

where $a \leq \mu_{\min}$, and $b \geq \mu_{\max}$, with μ_{\min} and μ_{\max} the smallest and the largest eigenvalues of $K^{-1}M$, respectively (i.e., $\mu_{\min} = 1/\lambda_{\max}$ and $\mu_{\max} = 1/\lambda_{\min}$, where λ_{\max} and λ_{\min} are the largest and smallest eigenvalues of $K - \lambda M$). The distribution function $\phi(\mu)$ is a step function with jumps at the eigenvalues μ_i of $K^{-1}M$ and is given by

$$(3.6) \quad \phi(\mu) = \sum_{i=1}^n \phi_i^2 h(\mu - \mu_i),$$

where

$$h(t) = \begin{cases} 1, & t \geq 0, \\ 0, & \text{otherwise} \end{cases}$$

is the Heaviside function and the ϕ_i 's are the coefficients of the first Lanczos vector q_1 when expressed in the M -orthonormal basis x_1, \dots, x_r for $\text{range}(K^{-1}M)$, namely,

$$(3.7) \quad q_1 = \sum_{i=1}^r \phi_i x_i,$$

(and $\phi_i = 0$ for $i = r+1, \dots, n$).

It turns out that the step function $\phi(\mu)$ coincides with the cumulative mass participation sum when the starting vector for the Lanczos algorithm with zero shift is the spatial distribution vector b in (1.2) with $b \in \text{range}(K^{-1}M)$, as we now show.

PROPOSITION 3.1. *Let $0 \neq b$ be the spatial distribution vector in (1.2) and assume that $b \in \text{range}(K^{-1}M)$. If $w = b$ in Algorithm 2.1 with $\sigma = 0$, then for the step function $\phi(\lambda)$ in (3.6) we have that*

$$\phi_i^2 = m(x_i), \quad i = 1, \dots, n,$$

where $m(x_i)$ is the normalized mass participation of the eigenvector x_i in (3.1).

Proof. The first Lanczos vector is given by $q_1 = b/\|b\|_M$ and, for $i = 1, \dots, r$,

$$m(x_i) = \frac{(x_i^T M b)^2}{\|b\|_M^2} = \frac{(x_i^T M q_1 \|b\|_M)^2}{\|b\|_M^2} = (x_i^T M q_1)^2 = \left(x_i^T M \sum_{j=1}^n \phi_j x_j \right)^2 = \phi_i^2.$$

For $i = r+1, \dots, n$, $x_i \in \text{null}(M)$ so that $Mx_i = 0$ and hence $m(x_i) = 0 = \phi_i^2$. \square

If the nonzero eigenvalues μ_j , $j = 1, \dots, r$, of $K^{-1}M$ are ordered by decreasing values then it follows from (3.6) and Proposition 3.1 that

$$\phi(\mu_j) = \sum_{i=1}^j m(x_i), \quad j \leq r.$$

Let $(\theta_i^{(k)}, s_i^{(k)})$, $i = 1, \dots, k$, be the eigenpairs of the tridiagonal matrix T_k resulting from k steps of the unshifted Lanczos algorithm (Algorithm 2.1 with $\sigma = 0$) and ordered such that

$$\theta_1^{(k)} \geq \theta_2^{(k)} \geq \dots \geq \theta_k^{(k)}.$$

The Lanczos polynomials p_i in (3.4) are not only orthogonal with respect to the inner product (3.5), they are also orthogonal with respect to the inner product

$$(3.8) \quad \langle p_i, p_j \rangle_{\tau_k} = \int_a^b p_i(\mu) p_j(\mu) d\tau_k(\mu), \quad 1 \leq i, j \leq k$$

induced by the step function

$$(3.9) \quad \tau_k(\mu) = \sum_{i=1}^k \tau_{k,i}^2 h\left(\mu - \theta_i^{(k)}\right),$$

where $\tau_{k,i}$ is the first entry of the eigenvector $s_i^{(k)}$, namely,

$$(3.10) \quad \tau_{k,i} = e_1^T s_i^{(k)}$$

(see [8, Chap. 2], [10]). As a result, $\tau_k(\mu)$ and the distribution function $\phi(\mu)$ in (3.6) have the same modified moments up to degree $2k - 1$, namely,

$$(3.11) \quad \langle 1, p_i \rangle_\phi = \langle 1, p_i \rangle_{\tau_k}, \quad i = 1, \dots, 2k - 1.$$

In turn, by the following theorem due to Karlin and Shapley [15, Thm. 22.2], it follows that $\tau_k(\mu)$ serves as a good approximation to $\phi(\mu)$.

THEOREM 3.2. *If $\phi(\mu)$ in (3.6) and $\tau_k(\mu)$ in (3.9) have the same modified moments up to degree $2k - 1$, then, if not identically zero, the difference function $\phi(\mu) - \tau_k(\mu)$ has $2k - 1$ sign changes.*

Thus if τ_k and ϕ do not coincide, the vertical and horizontal steps of τ_k will intersect ϕ exactly $2k - 1$ times. This theory has been used, for instance, in estimating eigenvalue distribution [8], and in constructing polynomial preconditioners [9]. In our case, we use the step function τ_k , obtained after k steps of the Lanczos algorithm, as an approximation to the cumulative mass participation sum, i.e.,

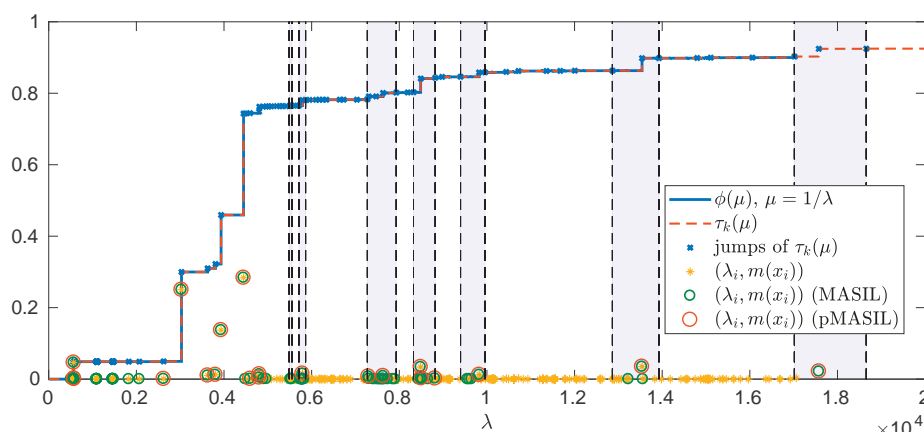
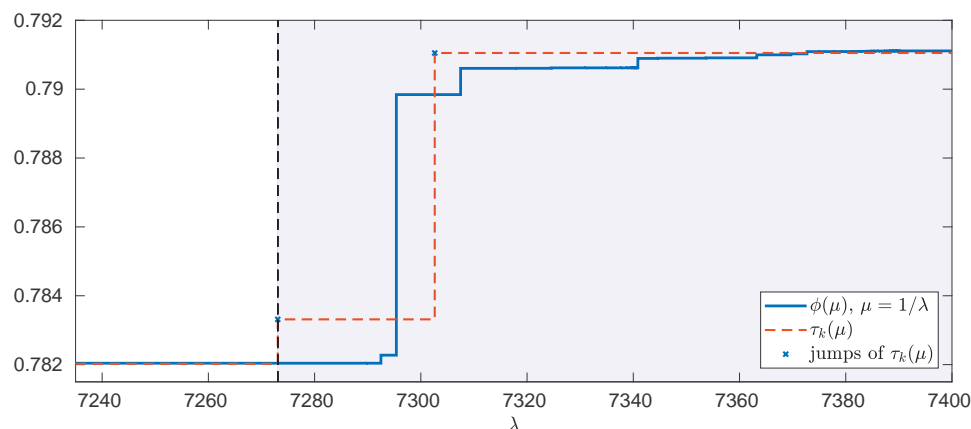
$$\tau_k(\mu_\ell) \approx \phi(\mu_\ell) = \sum_{i=1}^{\ell} m(x_i), \quad 1 \leq \ell \leq n.$$

This is illustrated in Figure 2(a)–(b) for a real structural engineering problem called **chilled** problem. On plot (a), the step functions τ_k and ϕ appear to coincide at least to the eye but plot (b), which is a closeup of plot (a) around $\lambda = 7.3 \times 10^3$ shows that the step functions indeed intersect. We will return to Figure 2 at the end of section 3.2.

3.2. A new shifting strategy. The step function $\tau_k(\mu)$ in (3.9) is readily available after k steps of Algorithm 2.1 with $\sigma = 0$ and starting vector $w = b$. Intuitively, the eigenvalues corresponding to the eigenvectors with largest normalized mass participation should lie under tall and narrow steps of $\phi(\lambda)$, whereas short and wide steps would indicate intervals of eigenvalues corresponding to eigenvectors of negligible mass participation (see Figure 2 for illustration). To make this formal, denote by $\Psi(c, d)$ the total mass participation of eigenvectors whose corresponding eigenvalues lie in the interval $[c, d]$, that is,

$$(3.12) \quad \Psi(c, d) = \sum_{j \in \mathcal{J}} m(x_j), \quad \mathcal{J} = \{j \in \{1, \dots, n\} : \lambda_j \in [c, d]\},$$

where $\lambda_j = 1/\mu_j$ is the j th largest eigenvalue of $K - \lambda M$. It follows from Theorem 3.2 that $\phi(\mu)$ will either cross or coincide with the vertical steps of $\tau_k(\mu)$. Therefore, the change in $\phi(\mu)$ over an interval between any two jumps of $\tau_k(\mu)$ can be bounded from below and above in terms of the heights of the vertical steps of $\tau_k(\mu)$. In particular, we have that for $1 < i < j < k$,

(a) chilled problem in the x -direction, $n = 93445$, $k_{\max} = 200$.

(b) Closeup of the above stair graph.

FIG. 2. Step functions $\phi(\mu)$ and $\tau_k(\mu)$ for the **chilled** problem in the x -direction. The small yellow stars in (a) correspond to the smallest eigenvalues λ_i and their normalized mass participation factors $m(x_i)$ such that $\sum_i m(x_i) \geq 0.9$. The green and red circles correspond to the pairs $(\lambda_i, m(x_i))$ obtained by mass accumulating SIL (MASIL) and MASIL with purging step (pMASIL), respectively, and such that $\sum_i m(x_i) \geq 0.9$. The shading shows the intervals selected by the new shifting strategy.

$$(3.13) \quad \sum_{\ell=i+1}^{j-1} \tau_{k,\ell}^2 \leq \Psi \left(\frac{1}{\theta_i^{(k)}}, \frac{1}{\theta_j^{(k)}} \right) \leq \sum_{\ell=i-1}^{j+1} \tau_{k,\ell}^2.$$

We can now make use of the bounds in (3.13) to identify the union of intervals of smallest total length over which we are guaranteed to satisfy

$$(3.14) \quad \sum_{j=1}^{\ell} m(x_{i_j}) \geq \xi$$

for some $\xi \in (0, 1)$. Since we require that the total mass participation over that union of intervals is at least ξ , we must look at the lower bounds of (3.13). Let us denote by γ_{ij} the following ratios

$$\gamma_{ij} := \frac{\sum_{\ell=i+1}^{j-1} \tau_{k,\ell}^2}{\frac{1}{\theta_j^{(k)}} - \frac{1}{\theta_i^{(k)}}},$$

so that, by (3.13),

$$\gamma_{ij} \leq \frac{\Psi\left(\frac{1}{\theta_i^{(k)}}, \frac{1}{\theta_j^{(k)}}\right)}{\frac{1}{\theta_j^{(k)}} - \frac{1}{\theta_i^{(k)}}}.$$

Assuming that the eigenvalues are distinct and roughly evenly distributed, large γ_{ij} indicate high relative mass participation in the interval $[1/\theta_i^{(k)}, 1/\theta_j^{(k)}]$, and small γ_{ij} indicate low relative mass participation. Since the lower bound for $\Psi(1/\theta_i^{(k)}, 1/\theta_{i+1}^{(k)})$ is 0, we look at intervals of the form $[1/\theta_{i-1}^{(k)}, 1/\theta_{i+1}^{(k)}]$. Thus to simplify the following discussion, define γ_i to be

$$\gamma_i := \gamma_{i-1,i+1} = \frac{\tau_{k,i}^2}{\frac{1}{\theta_{i+1}^{(k)}} - \frac{1}{\theta_{i-1}^{(k)}}}.$$

Suppose that after k steps of the Lanczos algorithm the first ℓ Ritz pairs (λ_ν, x_ν) , $\nu = 1, \dots, \ell$ of $K - \lambda M$ have converged. Let us denote by ξ_ℓ the sum of their mass participation factors, that is,

$$\xi_\ell = \sum_{\nu=1}^{\ell} m(x_\nu).$$

To construct a union of intervals of smallest total length over which we are guaranteed to attain the remaining mass participation $\xi - \xi_\ell$, we pick the s largest γ_i , say, $\gamma_{i_1}, \dots, \gamma_{i_s}$, $i_\nu \in \{\ell + 1, \dots, k - 1\}$ such that

$$(3.15) \quad \sum_{\nu=1}^s \tau_{k,i_\nu}^2 \geq \xi - \xi_\ell.$$

The wanted union of intervals is then $\bigcup_{\nu=1}^s [1/\theta_{i_\nu-1}^{(k)}, 1/\theta_{i_\nu+1}^{(k)}]$. By merging the neighboring and the overlapping intervals we can construct a set of $s' \leq s$ disjoint intervals

$$(3.16) \quad [1/\theta_{i_\nu}^{(k)}, 1/\theta_{j_\nu}^{(k)}], \quad \nu = 1, \dots, s'.$$

We then choose the shifts σ_ν to be the midpoints of those intervals, namely,

$$\sigma_\nu = \frac{1/\theta_{j_\nu}^{(k)} + 1/\theta_{i_\nu}^{(k)}}{2}, \quad \nu = 1, \dots, s'.$$

We end the shift-and-invert Lanczos iteration with shift σ_ν whenever the sum of the mass participation of the converged Ritz vectors with eigenvalues in $[1/\theta_{i_\nu}^{(k)}, 1/\theta_{j_\nu}^{(k)}]$ attains the minimum in (3.13), i.e., $\sum_{\ell=i_\nu+1}^{j_\nu-1} \tau_{k,\ell}^2$.

Remark 3.3. The inequality (3.15) may not hold if even $s = k - \ell + 1$ although in practice, it is usually satisfied after a small number of steps k . If (3.15) does not hold then we can increase k , or reduce ξ if possible.

As an illustration, let us look at Figure 2(a). The jumps of $\tau_k(\mu)$ correspond to the points $(1/\theta_i^{(k)}, \tau_k(\theta_i^{(k)}))$, $i = 1, \dots, k$ with $k = k_{\max}$. There are 34 more jumps outside of plot (a) corresponding to those $\theta_i^{(k)}$ such that $1/\theta_i^{(k)} > 2 \times 10^4$. The plot shows as small yellow stars the normalized mass participation factors $m(x_i)$ of the eigenvectors x_i corresponding to the ℓ smallest eigenvalues λ_i , solving Problem 1.2. Although not visible on plot (a), the eigenvalues of the **chilled** problem are clustered: for example, the smallest eigenvalue is around 507, there are 47 eigenvalues in the interval $[535, 538]$ and 22 in $[580.1, 580.7]$. The shaded regions correspond to the $s' = 7$ disjoint intervals in (3.16). They define seven shifts for the **chilled** problem. The nonshaded areas correspond to those intervals in which no more unconverged eigenvalues associated with eigenvectors of non-negligible mass participation remain. These intervals are ignored by our approach after the initial Lanczos iterations with the zero shift.

3.3. An algorithm for the approximate solution to Problem 1.1. Given two $n \times n$ matrices $M \geq 0$ and $K > 0$, a spatial distribution vector $b \in \text{range}(K^{-1}M)$, a proportion $\xi \in (0, 1)$, and a maximum number of iterations k_{\max} our algorithm for the approximate solution to Problem 1.1 goes through the following steps.

step 1 Call Algorithm 2.1 with $\sigma = 0$, $w = b$, and the implicit filtering turned off.

Stop the Lanczos iterations at step k when either

- (a) the converged Ritz vectors x_j , $j = 1, \dots, \ell_0$ are such that $\sum_{j=1}^{\ell_0} m(x_j) \geq \xi$, or
- (b) k has reached a number of k_{\max} iterations (or larger if $\sum_{j=1}^k \tau_{k,j}^2 \geq \xi$ is not satisfied).

If (a) holds then return the converged Ritz vectors x_j , $j = 1, \dots, \ell_0$ as an approximate solution to Problem 1.1. End the algorithm.

If (b) holds then

- if $\|q_k\|_2 \leq \text{tol}\|q_1\|_2$, then save the converged Ritz vectors x_j , $j = 1, \dots, \ell_0$, let $\xi_{\ell_0} = \sum_{j=1}^{\ell_0} m(x_j)$, and proceed to **step 2** with the converged and unconverged Ritz pairs $(1/\theta_i^{(k)}, x_i)$, $i = 1, \dots, k$.
- if $\|q_k\|_2 > \text{tol}\|q_1\|_2$, then keep the computed Ritz pairs $(1/\theta_i^{(k)}, \tilde{x}_i)$, $i = 1, \dots, k$ for **step 2**. Apply implicit filtering as described in section 2.2(a) and continue the Lanczos iterations (Algorithm 2.1). Save the converged Ritz vectors x_j , $j = 1, \dots, \ell_0$, let $\xi_{\ell_0} = \sum_{j=1}^{\ell_0} m(x_j)$, and proceed to **step 2**.

step 2 Construct a sequence of disjoint intervals $[1/\theta_{i_\nu}^{(k)}, 1/\theta_{j_\nu}^{(k)}]$, $\nu = 1, \dots, s'$, as discussed in section 3.2, using the k converged and unconverged Ritz pairs from **step 1**. Compute the sequence of shifts $\sigma_\nu = (1/\theta_{j_\nu}^{(k)} + 1/\theta_{i_\nu}^{(k)})/2$, $\nu = 1, \dots, s'$.

step 3 For $\nu = 1, \dots, s'$, call Algorithm 2.1 with $\sigma = \sigma_\nu$ and $w = (K - \sigma_\nu M)^{-1} M \tilde{w}$, where \tilde{w} is the sum of the unconverged Ritz vectors from the previous step. Apply implicit filtering when $\|q_k\|_2 > \text{tol}\|q_1\|_2$. Stop the Lanczos iterations when either

- (i) the converged Ritz vectors x_j , $j = p+1, \dots, p+\ell_\nu$ with $p = \ell_0 + \sum_{i=1}^{\nu-1} \ell_i$, are such that

$$\xi_{\ell_\nu} := \sum_{j=p+1}^{p+\ell_\nu} m(x_j) \geq \sum_{\ell=i_\nu+1}^{j_\nu-1} \tau_{k,\ell}^2,$$

or

(ii) $\bar{\xi} := \xi_{\ell_0} + \sum_{j=1}^{\nu} \xi_{\ell_j} \geq \xi$.
 Stop the for loop when (ii) holds. Return the converged Ritz vectors x_j ,
 $j = 1, \dots, p$ with $p = \ell_0 + \sum_{i=1}^{\nu} \ell_i$ as an approximate solution to Problem 1.1.
 End the algorithm.

Implicit filtering cannot be used in the first call to Algorithm 2.1 in **step 1** since this would alter the starting vector. Our selection of shifts and their computation in **step 2** rely on the theory described in sections 3.1–3.2, which requires the first Lanczos vector to be $b/\|b\|_M$. We do not think it is possible to recover the approximation to the cumulative mass participation sum after implicit filtering. In **step 1(b)** the condition $\sum_{j=1}^k \tau_{k,j}^2 \geq \xi$ is usually satisfied after a small number of steps k (see Remark 3.3). The choice of the maximum number of iterations k_{\max} is important. If k_{\max} is too small, then the approximation $\tau_k(\lambda)$ of $\phi(\lambda)$ is too rough and leads to large intervals and shifts that are not close enough to eigenvalues with eigenvectors that have large mass participation. On the other hand, a too large k_{\max} can lead to unnecessary computations, in particular of converged eigenvectors with negligible mass participation but can also result in too many shifts being identified by our shifting strategy. This last point is illustrated in Figure 3 for a real engineering problem. The plot shows that the number of shifts (or intervals in (3.16)) determined by our shifting strategy increases as the number k of unshifted Lanczos steps increases. In practice, we found that taking $k_{\max} = 200$ is a reasonable choice. For tol we choose $\text{tol} = 10^4$, and we consider that the eigenpair (λ_j, x_j) has converged if $\eta(\lambda_j, x_j) \leq nu$, where u is the machine precision and $\eta(\lambda_j, x_j)$ is defined in (2.12).

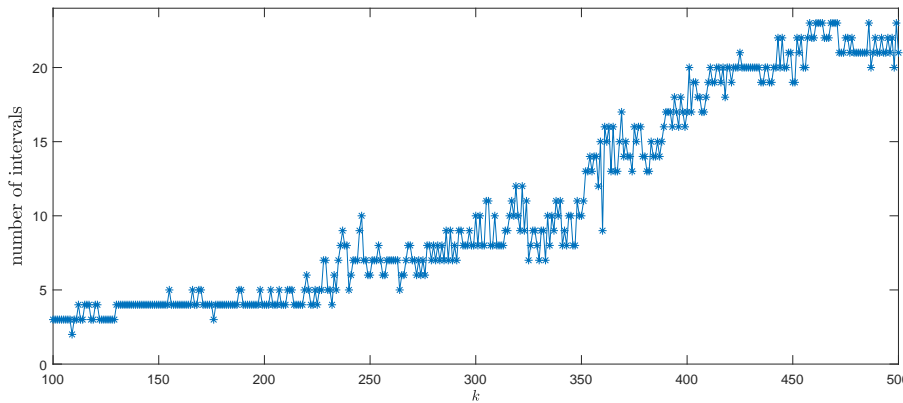


FIG. 3. Number of intervals s' in (3.16) as a function of the number k of unshifted Lanczos steps for the `local_modes` problem in the z -direction.

3.4. Cutoff frequency. An issue with the algorithm described in section 3.3 may arise if the intervals returned contain eigenvalues that are too large, either because they do not represent the frequencies of interest to engineers, or because they cannot be considered reliable due to limited validity of the underlying finite element mesh. This can be prevented by confining the search for eigenvalues to the interval $(0, \lambda^*]$ for some cutoff frequency λ^* . The theory developed in this section provides insight into how a cutoff frequency can affect the solutions to Problems 1.1 and 1.2. In particular, let j be the smallest integer for which the following holds for a given ξ

$$(3.17) \quad \sum_{i=1}^j \tau_{k,i}^2 \geq \xi,$$

where $\tau_{k,i}$ is as in (3.10). Due to the bounds given in (3.13), we then have that

$$(3.18) \quad \Psi\left(0, \frac{1}{\theta_{j-1}^{(k)}}\right) \leq \xi \leq \Psi\left(0, \frac{1}{\theta_{j+1}^{(k)}}\right),$$

that is, the eigenvalue λ_ℓ in (1.8), at which the threshold ξ is exceeded by computing eigenvectors corresponding to the smallest eigenvalues, must lie in the interval $[1/\theta_{j-1}^{(k)}, 1/\theta_{j+1}^{(k)}]$. Thus for $\lambda^* < 1/\theta_{j-1}^{(k)}$, the required mass participation cannot be reached even if all eigenvalues are computed in the interval $(0, \lambda^*]$, and only when $\lambda^* \geq 1/\theta_{j+1}^{(k)}$, we are certain that the condition (1.8) can be satisfied by computing some or all eigenvalues that lie in the interval $(0, \lambda^*]$.

Note that when $\lambda^* \geq 1/\theta_{j+1}^{(k)}$, we are also certain that intervals as given in (3.16) can be constructed from the Ritz values $1/\theta_i^{(k)}$ which lie entirely in the interval $(0, \lambda^*]$. For λ^* large enough, namely, $\lambda^* \geq 1/\theta_{j_s'}^{(k)}$, where $1/\theta_{j_s'}^{(k)}$ is as in (3.16), the intervals returned by the algorithm described in section 3.3 would not differ from those returned by the same algorithm if no cutoff frequency is introduced. The intervals returned will differ only if $\lambda^* \in [1/\theta_{j+1}^{(k)}, 1/\theta_{j_s'}^{(k)})$. As λ^* approaches $1/\theta_{j+1}^{(k)}$, one would expect the intervals to become larger and fewer of them to be returned. For $\lambda^* = 1/\theta_{j+1}^{(k)}$, no intervals containing eigenvalues whose eigenvectors have negligible mass participation would be identified and SIL algorithm with the shifting strategy (3.2) is likely to be more efficient.

4. Numerical experiments. For our numerical experiments we used matrices M and K , and spatial distribution vectors b provided by Arup Group Limited that were constructed by the finite element software package Oasys GSA [19] from models of real structural engineering problems. These are listed in Table 2. Our numerical experiments are performed with MATLAB. As mentioned in the introduction, b is always associated with a direction, either x , y , or z . So for a given problem, we will have different solutions depending on the chosen direction since they correspond to different spatial distribution vectors.

We compare the following approaches to solve our problem:

- SIL: shift-and-invert Lanczos algorithm with the Ericsson and Ruhe shifting strategy (3.2)–(3.3),
- MASIL: mass accumulating shift-and-invert Lanczos algorithm with the new shifting strategy described in section (3.3).

When no shifts are used, SIL and MASIL are identical, except for their starting vectors: SIL uses $w = K^{-1}Mb$ as suggested in [24] and MASIL uses $w = b$.

In practice, the cumulative sum of the mass participation factors of the eigenvectors returned by SIL and MASIL is always slightly larger than the wanted proportion ξ . So we have the possibility to remove some of the eigenvectors with small mass participation from the list returned by SIL or MASIL, as long as the cumulative sum of the mass participation factors of the remaining eigenvectors is exactly ξ or just above. It can be seen from (2.4) that the contribution to the solution of an eigenvector x_j depends not only on the magnitude of its mass participation factor, but also on

TABLE 2

List of test problems together with their size n and the number k_{\max} of unshifted Lanczos steps for SIL and the new method MASIL. The last four columns list the number of eigenvectors needed to satisfy the mass participation condition (3.14) with $\xi = 0.9$ for SIL and MASIL, and their purged versions pSIL and pMASIL. The number of shifts used by SIL and MASIL is provided inside brackets.

Problem (direction)	n	k_{\max}	SIL (shifts)	MASIL (shifts)	pSIL	pMASIL
local_modes (z)	51,348	200	984 (12)	365 (3)	589	308
ccnb (x)	57,152	200	148 (1)	77 (3)	64	55
chilled (x)	93,445	200	455 (1)	101 (7)	69	21
chilled (y)		100	117 (1)	32 (3)	6	9
chilled (y)		200	49 (0)	49 (0)	6	6
TT (y)	131,835	200	494 (7)	407 (8)	204	151
TT (z)		200	1173 (21)	992 (6)	540	430

the magnitude of the corresponding eigenvalue. We therefore choose to purge those eigenvectors x_j , for which the quantities

$$\frac{|x_j^T M b|}{\omega_j}$$

are the smallest, thereby targeting to purge the eigenvectors with small mass participation and large eigenvalues. We refer to this small modification of SIL and MASIL as pSIL and pMASIL, respectively, where the “p” stands for extra purging step. An illustration of the purging step can be seen in Figure 2(a), where the pairs $(\lambda_i, m(x_i))$ from MASIL are shown as small green circles and those kept by pMASIL are shown as large red circles.

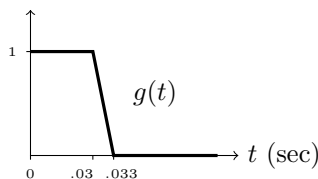
In Table 2, we compare the number of computed eigenvectors required to satisfy the mass participation condition (3.14) with $\xi = 0.9$. Some directions for the problems are excluded from the table. This is either because shifts were not employed (e.g., **ccnb** in the y - or the z -direction, where the 90% mass participation was reached in fewer than k_{\max} steps and no shift), or too many eigenvectors were required to achieve 90% mass participation, exceeding time and memory constraints (e.g., for the **chilled** problem in z -direction and $\xi = 0.75 \ll 0.9$, SIL returns 6,959 eigenvectors whereas MASIL returns 4,201 eigenvectors).

The new shifting strategy allows us to exclude intervals containing eigenvalues whose eigenvectors have a negligible mass participation and shift in the middle of intervals containing eigenvalues with eigenvectors of large mass participation. As a result, the number of eigenvectors returned by MASIL can be much smaller than that returned by SIL. This is, for example, the case for the **local_modes** problem in the z -direction and the **chilled** problem in the x -direction. The number of shifts used by SIL or MASIL depends on the problem and the shifting strategy employed. Notice the large number of shifts employed by SIL for the **local_modes** problem in the z -direction and for the **TT** problem in the z -direction. For MASIL, the number of shifts needed is known beforehand so if this number is too high, then there is always the possibility to increase the value of k_{\max} or to see if there is a $k < k_{\max}$ that leads to larger search intervals but fewer of them. We reported results for two different values of k_{\max} for the **chilled** problem in the y -direction. For $k_{\max} = 100$, SIL uses only one shift but returns 117 eigenvectors whereas MASIL uses 3 shifts but returns only 32 eigenvectors. No shifts are needed if we increase k_{\max} to 200.

Now comparing the number of eigenvectors returned by pSIL to that of SIL shows that a large proportion of the eigenvectors computed by SIL have a negligible mass participation and can be purged away. The last column in Table 2 shows that MASIL still returns many eigenvectors with negligible mass participation that can be removed while still maintaining the condition (3.14), but the reduction is not as drastic as for SIL. Note that the number of eigenvectors returned by pMASIL is lower or equal than those returned by pSIL.

Let us now look at the quality of the approximation $\tilde{u}(t) = \sum_{j=1}^p v_j(t)x_j$ to the response $u(t)$ of (1.1) when the p eigenvectors x_j are computed by SIL, MASIL, pSIL, or pMASIL. For the time dependent external load $f(t) = g(t)Mb$ we use

$$g(t) = \begin{cases} 1, & 0 \leq t \leq 0.03, \\ 11 - \frac{10^3}{3}t, & 0.03 \leq t \leq 0.033, \\ 0, & 0.033 \leq t \leq 1 \end{cases}$$



for the loading function g as suggested in [4, Fig. 2]. We employ 2% damping which consists of setting $\zeta_j = 0.02$ in (2.4) for all j . For the relative error between the exact solution $u(t)$ and its approximation $\tilde{u}(t)$, the analysis in section 2.1 tells us to expect

$$(4.1) \quad \frac{\|u(t) - \tilde{u}(t)\|_M}{\|u(t)\|_M} \lesssim \sqrt{1 - \xi} \approx 0.3.$$

We do not have access to $u(t)$. Therefore as a reference we use the solution obtained from SIL with $\xi = 0.99$. The relative errors for the approximate responses $\tilde{u}(t)$ returned by SIL, pSIL, MASIL, and pMASIL with $\xi = 0.9$ are shown in Figure 4 for the `local_modes` problem in the z -direction in (a) and for the `ccnb` problem in the x -direction in (b). The relative errors agree with (4.1). We note that for the `local_modes` problem, the pMASIL solution with 308 eigenvectors is almost as good as the pSIL solution which requires 589 eigenvectors. For the `ccnb` problem MASIL algorithm is comparable to or outperforms SIL algorithm with significantly fewer eigenvectors, and the pMASIL solution is almost as good as the pSIL solution for this problem, with fewer eigenvectors.

We plot in Figure 5(a)-(b) the i th entry of the reference solution $u(t)$ and its approximations obtained by SIL, pSIL, MASIL, and pMASIL for the `ccnb` and `TT` problems. We chose the index i for which the reference solution has the largest amplitude. For the `TT` problem in the y -direction in plot (a), all solutions agree with the reference solution even the pMASIL solution which uses significantly fewer eigenvectors. Although (4.1) holds for all the approximate solutions, the SIL and pSIL solutions do not agree as well with the reference solution for the `ccnb` problem in the x -direction in plot (b), as the MASIL and pMASIL solutions, thereby suggesting a better selection of eigenvectors representing the solution.

Finally, we investigate the sensitivity of the cumulative mass participation factors with respect to the quality of the eigenvector approximations and the choice of starting vectors for SIL. For a fair comparison, we only allow one shift $\sigma = 0$. The results are summarized in Table 3 for the `local_modes` problem in the z -direction and the `TT` problem in the y -direction. Note that we had to lower the threshold to $\xi = 0.75$ to ensure that the required eigenvectors were computed without the need to restart the algorithm with nonzero shifts. Recall that MASIL and pMASIL would return the results in the first row of the table. We expect the eigenvectors computed by SIL with

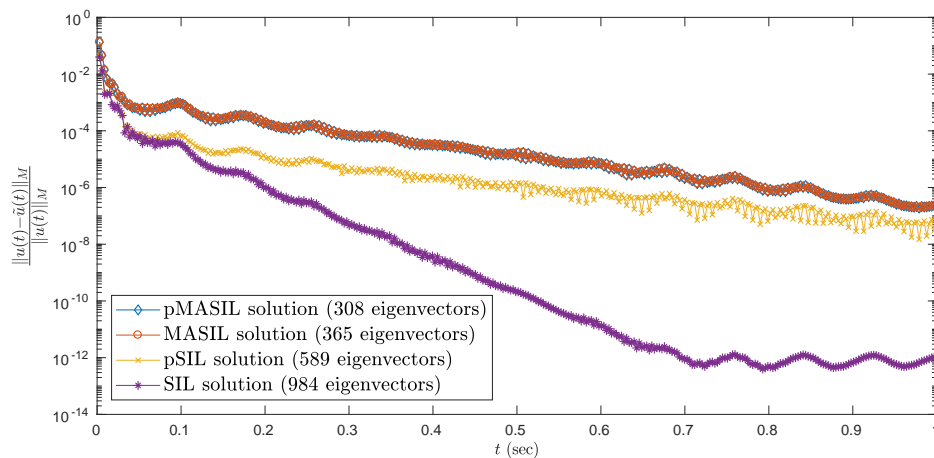
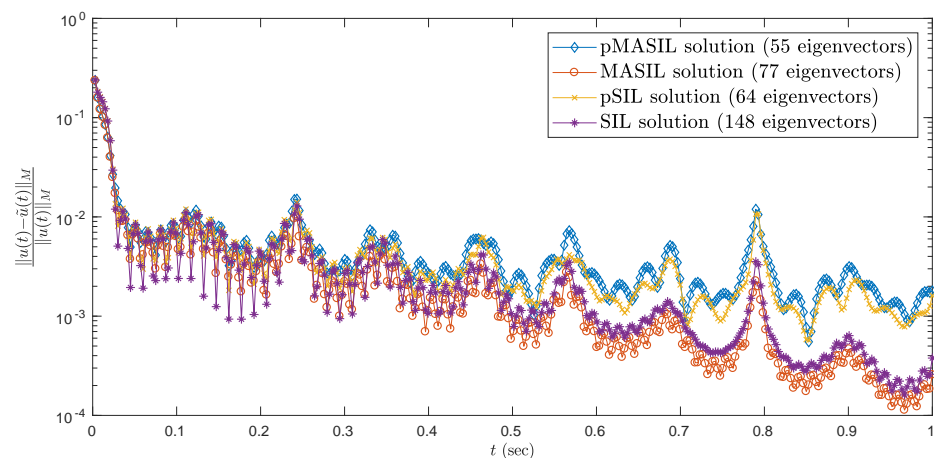
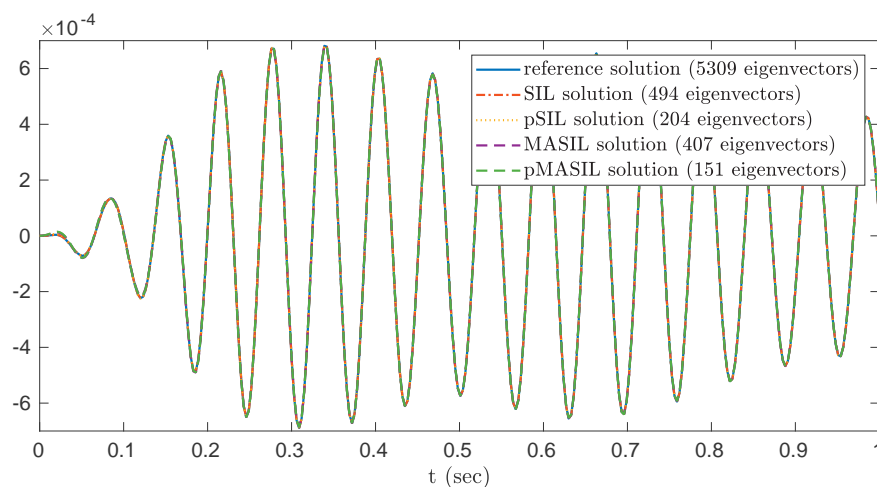
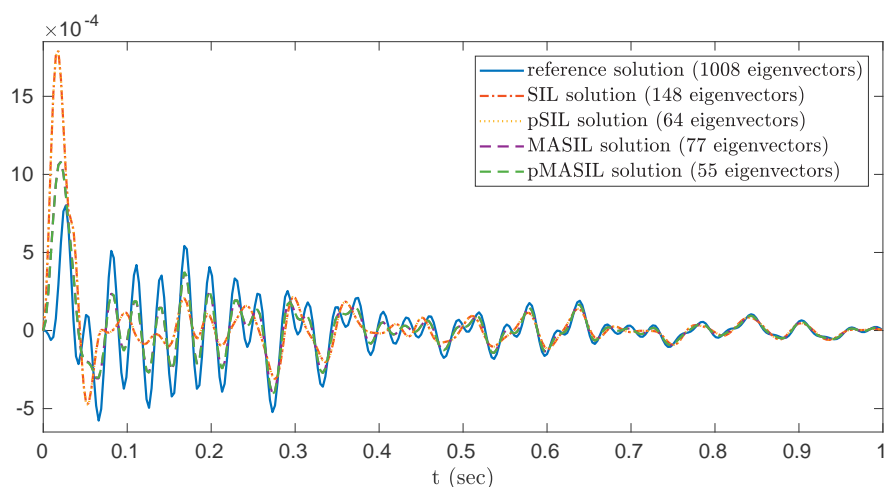
(a) `local_modes` problem in the z -direction, $k_{\max} = 200$.(b) `ccnb` problem in the x -direction, $k_{\max} = 200$.

FIG. 4. Relative error between the reference solution $u(t)$ and its approximation $\tilde{u}(t)$ obtained from SIL, MASIL, pSIL, and pMASIL with a proportion of $\xi = 90\%$ for the total mass participation. (SIL with $\xi = 0.99$ was used to construct the reference solution requiring 3,643 eigenvectors for the `local_modes` problem and 1,008 eigenvectors for the `ccnb` problem.)

tolerance nu on the backward error of converged eigenpairs to be better approximations to the exact eigenvectors than those computed using the larger tolerance $n\sqrt{u}$. The starting vector affects the convergence of SIL and hence that of the cumulative mass participation factors. This explains the differences in the number eigenvectors required to satisfy the mass participation condition (1.8). Note that the differences are not as pronounced when the tolerance is set to nu . In particular, after purging the eigenvectors with negligible mass participation, the number of eigenvectors returned by pSIL with tolerance nu is more or less the same for all seven starting vectors, an exception being for the `local_modes` problem and the starting vector generated as `rng(2);K\M*randn(n,1)`. For the latter, an eigenvector of relatively large mass

(a) TT problem in the y -direction, $i = 21,943$.(b) cnb problem in the x -direction, $i = 14,802$.FIG. 5. Plot of the i th entry of the response vector $u(t)$ (reference solution) and its approximations obtained by SIL, pSIL, MASIL, and pMASIL.

participation did not converge. As a result, more eigenvectors were required by pSIL to make up for the missing mass participation. The computed eigenvectors with tolerance $n\sqrt{u}$ are bad approximations to the exact eigenvectors, in particular when using $b, K^{-1}Mb, (K^{-1}M)^2b$ as starting vectors. This results in mass participation factors that are larger than they should be, leading to a cumulative mass participation that reaches the threshold ξ too quickly. This shows that it is important to use eigenvectors that have converged with a small backward error to compute the cumulative mass participation before checking the mass participation condition (1.8).

5. Conclusion. We have shown that if the Lanczos process is applied to $K^{-1}M$ with starting vector equal to the spatial distribution vector b , then the Lanczos polynomials are orthogonal with respect to the inner product induced by a step function

TABLE 3

Number of eigenvectors required to satisfy the mass participation condition (1.8) with $\xi = 0.75$ when computed by SIL and pSIL using different starting vectors and convergence tolerances on the backward error for the computed eigenpairs for the `local_modes` problem in the z -direction and the TT problem in the y -direction. Here u denotes the machine precision.

Starting vector	local_modes ($n = 51,348$)				TT ($n = 131,835$)			
	tolerance nu		tolerance $n\sqrt{u}$		tolerance nu		tolerance $n\sqrt{u}$	
	SIL	pSIL	SIL	pSIL	SIL	pSIL	SIL	pSIL
<code>b</code>	94	65	92	65	164	35	3	2
<code>K\M * b</code>	94	65	4	4	163	35	2	1
<code>(K\M)² * b</code>	94	65	8	8	163	35	3	1
<code>K\M * ones(n, 1)</code>	98	64	91	54	168	35	23	14
<code>rng(0); K\M * randn(n, 1)</code>	96	65	97	62	171	35	142	53
<code>rng(1); K\M * randn(n, 1)</code>	97	64	97	56	171	35	164	98
<code>rng(2); K\M * randn(n, 1)</code>	97	79	93	66	172	35	158	74

ϕ that coincides with the cumulative mass participation sum, that is, the quantity we are interested in to solve Problems 1.1–1.2. The Lanczos polynomials are also orthogonal with respect to an inner product induced by a step function τ_k that, unlike ϕ , is readily available at step k of the Lanczos process. The step function τ_k offers an approximation to ϕ and hence to the cumulative mass participation sum. The eigenvalues of $K - \lambda M$ lie on the positive real line, and we use τ_k to identify intervals containing eigenvalues whose corresponding eigenvectors have non-negligible mass participation as well as intervals containing eigenvalues whose eigenvectors have negligible contribution to the cumulative mass participation. We use this information to construct a sequence of shifts $\sigma_1, \sigma_2, \dots, \sigma_p$ for the SIL algorithm as well as a stopping criterion for the SIL steps with shift σ_i , $i = 1, \dots, p$ so that (1.7) holds. The numerical experiments we performed on real engineering problems show that our approach computes up to 70% fewer eigenvectors and requires fewer shifts, on average, compared with the shifting strategy proposed by Ericsson and Ruhe and its slight modification by Grimes, Lewis, and Simon.

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