

# NOTES ON ETLING VELOCITY PROFILE FOR EKMAN/PRANDTL LAYER

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## 1. BASIC VELOCITY PROFILES

### 1.1. Velocity profile in the Prandtl layer.

$$(1) \quad u_p(z) = \frac{u_\star}{\kappa} \log \frac{z^+}{z_0^+}$$

In coordinates of the geostrophic wind (for a veering angle of  $\alpha_0 = \alpha(z=0)$ )

$$(2) \quad u_p(z) = \frac{u_\star}{\kappa} \log \frac{z^+}{z_0^+} \begin{pmatrix} \cos \alpha_0 \\ \sin \alpha_0 \end{pmatrix}$$

### 1.2. Velocity profile in the Ekman layer.

$$(3) \quad \begin{pmatrix} u_{ek} \\ v_{ek} \end{pmatrix} (z) = G\sqrt{2} \sin \alpha_0 \begin{pmatrix} \frac{1}{\sqrt{2} \sin \alpha_0} - e^{-\tilde{z}} \cos(\tilde{z} + \frac{\pi}{4} - \alpha_0) \\ e^{-\tilde{z}} \sin(\tilde{z} + \frac{\pi}{4} - \alpha_0) \end{pmatrix}$$

with the scaled Ekman height

$$\tilde{z} = \frac{z - z_P}{D},$$

where  $z_P$  is the Prandtl layer thickness (here assumed as 10% of the BL thicknes) and  $D$  the Ekman layer thickness. For  $D$ , there exists a matching based on the constant diffusivity within the Prandtl layer

$$(4) \quad D = \sqrt{2 \frac{K_m}{f}} = \sqrt{2 \frac{\kappa u_\star z_p}{f}} = \sqrt{2 \kappa \delta z_p}$$

with  $\delta = u_\star/f$ . Using  $z_p = 0.1\delta$ , we obtain

$$D = \sqrt{0.2 \kappa \delta} \approx 0.3\delta$$

## 2. MATCHING POINT

The profiles shall be matched at the upper end of the Prandtl layer ( $z = z_P$ ). Given the geostrophic wind, Prandtl layer height  $z_p$  and height offset (in the case of a rough surface, the roughness height; for aerodynamically smooth surfaces a scaled version of the profile offset parameter), the profiles can be matched to yields the remaining parameters  $u_\star$  and  $\alpha_0$ .

$$(5a) \quad \frac{u_\star}{\kappa} \log \frac{z_P}{z_0} \cos \alpha_0 = G\sqrt{2} \sin \alpha_0 \left[ \frac{1}{\sqrt{2} \sin \alpha_0} - e^{-\tilde{z}} \cos \phi \right]$$

$$(5b) \quad \frac{u_\star}{\kappa} \log \frac{z_P}{z_0} \sin \alpha_0 = G\sqrt{2} \sin \alpha_0 e^{-\tilde{z}} \sin \phi$$

where – utilizing the condition  $z = z_P$ , it is  $\phi = \pi/4 - \alpha_0$ . We can then eliminate  $\sin \alpha_0$  from the equation for the span-wise component to obtain

$$(6a) \quad \cos \alpha_0 = \frac{G}{u_\star} \sqrt{2} \lambda_z \sin \alpha_0 \left[ \frac{1}{\sqrt{2} \sin \alpha_0} - \cos \phi \right]$$

$$(6b) \quad \frac{u_\star}{G} = \sqrt{2} \lambda_z \sin \phi$$

where

$$\lambda_z = \frac{\kappa}{\log \frac{z_P}{z_0}}.$$

Next, we rewrite Eq. (6a) using Eq. (6b) as

$$(7a) \quad \cos \alpha_0 = \sqrt{2} \lambda_z \frac{G}{u_\star} \left( \frac{1}{\sqrt{2} \sin \phi} - \frac{\sin \alpha_0 \cos \phi}{\sin \phi} \right)$$

$$(7b) \quad \sin \phi \cos \alpha_0 + \cos \phi \sin \alpha_0 = \frac{1}{\sqrt{2}}$$

The LHS of the latter equation can be rewritten as

$$\sin(\phi + \alpha_0) = \sin(\pi/4)$$

such that the equation holds for any  $\alpha_0$ . Thus,  $\alpha_0$  becomes a parameter of the solution, and we can simply write

$$(7c) \quad \frac{u_\star}{G} = \frac{\sqrt{2} \kappa}{\log(z_P/z_0)} \sin\left(\frac{\pi}{4} - \alpha_0\right)$$

**2.1. Relation to the Reynolds number.** In a smooth flow, the scaling parameter of the logarithmic height is

$$\frac{1}{\kappa} \log z^+ + A = \frac{1}{\kappa} [\log z^+ - (-\kappa A)] = \frac{1}{\kappa} [\log z^+ - \log(e^{-\kappa A})] = \log(z^+/z_0^+)$$

where  $z_0^+ = e^{-\kappa A} \approx 1/8$ . This implies that  $z_P/z_0 = z_P^+/z_0^+ = e^{\kappa A} z_P^+ = e^{\kappa A} \gamma_P Re_\tau$ . Hence, the ratio  $z_P/z_0$  implicitly defines the Reynolds number  $Re_\tau$  as

$$(8) \quad Re_\tau = \frac{z_P}{z_0} \frac{1}{e^{\kappa A} \gamma_P} \sim Re_\tau.$$

**2.2. Linearization for  $\alpha_0$  in the trigonometric terms.** Based on a first guess for  $u_\star$  and  $\alpha_0$ , we can then estimate  $u_\star/G$  from Eq. (6b) and  $\alpha_0$  from Eq. (6a) and iterate this procedure to arrive at a solution for  $\alpha_0$  and  $u_\star$ .

A first guess may be obtained from a linearized version of the equations in the trigonometric terms. It is

$$(9a) \quad \sin \phi = \sin(\pi/4 - \alpha_0) \approx \sin(\pi/4) - \cos(\pi/4) \alpha_0 = \frac{1}{\sqrt{2}}(1 - \alpha_0)$$

$$(9b) \quad \cos \phi = \cos(\pi/4 - \alpha_0) \approx \cos(\pi/4) + \sin(\pi/4) \alpha_0 = \frac{1}{\sqrt{2}}(1 + \alpha_0)$$

$$(9c) \quad \sin \alpha_0 \approx \sin(0) + \alpha_0 = \alpha_0.$$

We thus obtain

$$(10) \quad -1 = \frac{G \lambda_z}{u_\star} \left[ \frac{1}{\alpha_0} - (1 - \alpha_0) \right]$$

$$(11) \quad \Rightarrow 0 = \alpha_0 + \frac{G \lambda_z}{u_\star} [1 - (\alpha_0 - \alpha_0^2)]$$