

Lec 20

Particle in a 1D box

→ ground state ($n=1$)

$$\text{wave function } \psi_1(n) = \sqrt{\frac{2}{L}} \sin\left(\frac{\pi n}{L}\right)$$

Calculate $\langle n \rangle$, $\langle n^2 \rangle$, $\langle P \rangle$, $\langle P^2 \rangle$

$$\langle (\Delta n)^2 \rangle = \langle n^2 \rangle - \langle n \rangle^2$$

$$\langle (\Delta P)^2 \rangle = \langle P^2 \rangle - \langle P \rangle^2$$

$$\text{Note: } \langle A \rangle = \int_n \psi^*(n) \hat{A} \psi(n) dn$$

$$\int_n \psi^*(n) \psi(n) dn$$

$$\langle n \rangle = \int_0^L \frac{2}{L} \sin^2\left(\frac{\pi n}{L}\right) n dn$$

I₁

$$\int_0^L \frac{2}{L} \sin^2\left(\frac{\pi n}{L}\right) dn$$

I₂

$$I_2 = \frac{2}{L} \cdot \frac{1}{2} \int_0^L \left[1 - \cos \frac{2\pi n}{L} \right] dn$$

$$= \frac{1}{L} \int_0^L \left(1 - \cos \frac{2\pi n}{L} \right) dn$$

$$= \frac{1}{L} \left\{ [n]_0^L - \left[\frac{\sin \frac{2\pi n}{L}}{\frac{2\pi}{L}} \right]_0^L \right\}$$

$$= \frac{1}{L} \left\{ L - \frac{L}{2\pi} \sin 2\pi \right\} = 1$$

$$I_1 = \int_0^L \frac{2}{L} \sin^2 \left(\frac{\pi n}{L} \right) n \, dx$$

$$= \frac{2}{L} \left[\left[n \int \sin^2 \left(\frac{\pi n}{L} \right) dx \right]_0^L - \int_0^L \int \sin^2 \left(\frac{\pi n}{L} \right) dx \cdot du \right]$$

$$= \frac{2}{L} \left[\left[\frac{n}{2} \left(u - \frac{\sin \frac{2\pi n}{L}}{2\pi} \right) \right]_0^L - \int_0^L \left(\frac{n}{2} - \frac{1}{2} \frac{\sin \frac{2\pi n}{L}}{\frac{2\pi}{L}} \right) dx \right]$$

$$= \frac{2}{L} \left[\frac{L}{2} (L - 0) - \left[\frac{n^2}{4} + \frac{1}{2} \frac{\cos \frac{2\pi n}{L}}{\left(\frac{2\pi}{L} \right)^2} \right]_0^L \right]$$

$$= \frac{2}{L} \left[\frac{L^2}{2} - \frac{L^2}{4} - \frac{1}{2} \cdot \cancel{\frac{1}{4\pi^2}} + \frac{1}{2} \cancel{\frac{L^2}{4\pi^2}} \right]$$

$$= \frac{2}{L} \times \frac{L^2}{4} = \frac{L}{2}$$

$$\langle n^2 \rangle = \int_0^L n^2 \sin^2 \left(\frac{\pi n}{L} \right) dx = \frac{L^2}{6} \left(2 - \frac{3}{\pi^2} \right)$$

$$\langle P \rangle = \int_0^L \psi^*(x) \left(-i \hbar \frac{\partial}{\partial x} \psi(x) \right)$$

$$\frac{2}{L} \int_0^L \left(\frac{\pi n}{L} \right) \sqrt{\frac{2}{L}} \cos \left(\frac{\pi n}{L} \right) \cdot \frac{\pi}{L} dx$$

$$i \hbar \frac{2}{L} \cdot \frac{\pi}{L} \int_0^L \cos \left(\frac{\pi n}{L} \right) \sin \left(\frac{\pi n}{L} \right) dx$$

$$-i \hbar \frac{\pi}{L^2} \int_0^L \sin \left(\frac{2\pi n}{L} \right) dx = -i \frac{\pi \hbar}{L^2} \left[\frac{\cos \frac{2\pi n}{L}}{\frac{2\pi}{L}} \right]_0^L$$

$$= -i \frac{\pi \hbar}{L^2} \left[\frac{1}{\frac{2\pi}{L}} - \frac{1}{\frac{2\pi}{L}} \right] = 0$$

$$\langle P^2 \rangle = \frac{\hbar^2}{4L^2}$$

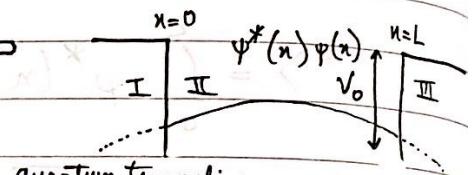
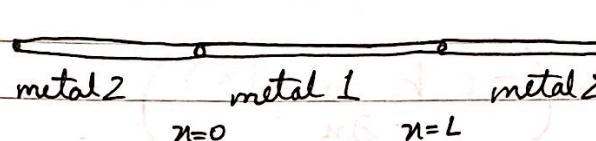
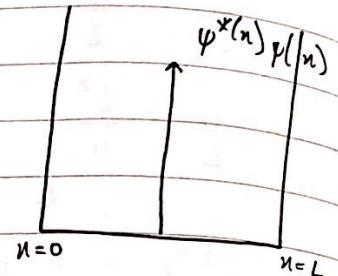
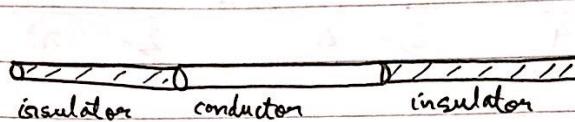
$\Delta n \equiv$ Uncertainty in n

$$= \sqrt{\langle n^2 \rangle - \langle n \rangle^2}$$

$\Delta p \equiv$ Uncertainty in p

$$= \sqrt{\langle p^2 \rangle - \langle p \rangle^2}$$

$\Delta n \cdot \Delta p \geq \frac{\hbar}{2} -$ Uncertainty Principle.



$$E < V_0, \quad H \psi(x) = E \psi(x)$$

$$E > V_0.$$

quantum tunneling
half sine wave
tail outside

Scanning tunneling microscopy (STM)
 $10^{-6} \text{ m} \rightsquigarrow 10^{-9} \text{ m}$

Region I : $V(n) = 0$

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial n^2} \psi_I(n) = E \psi_I(n)$$

$$\Rightarrow \psi_I(n) = C_1 e^{k_1 n} \quad k_1 = \sqrt{\frac{2m|E|}{\hbar^2}}$$

$$\psi_{III}(n) = C_3 e^{-k_1 n}$$

Boundary Conditions

$$\psi_I(n = -a) = \psi_2(n = -a)$$

[ψ should be single valued]

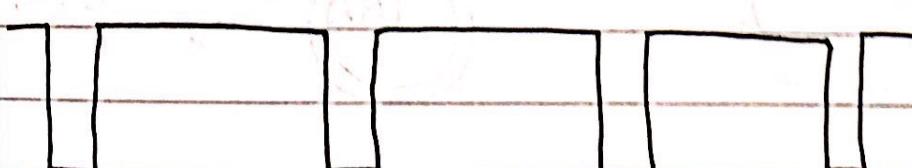
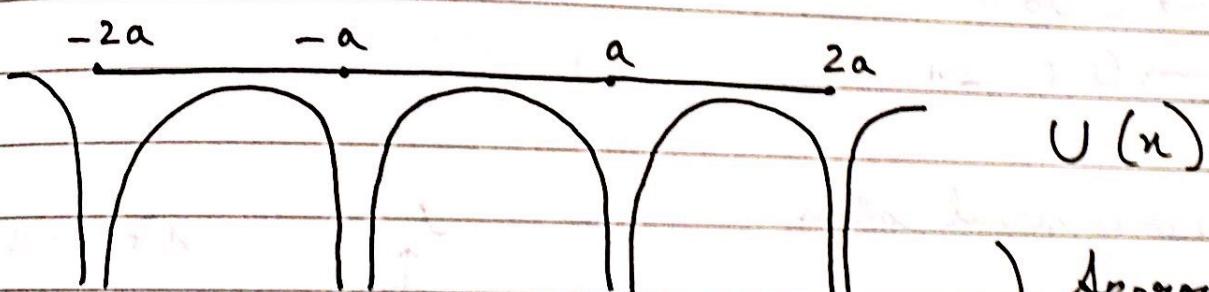
$$\frac{d}{dn} \psi_I(n = -a) = \frac{d}{dn} \psi_{II}(n = -a)$$

[ψ should be continuously differentiable]

Similarly :

$$\psi_{II}(n = a) = \psi_{III}(n = a)$$

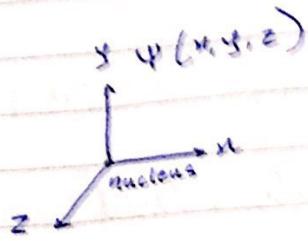
$$\frac{d}{dn} \psi_{II}(n = a) = \frac{d}{dn} \psi_{III}(n = a)$$



→ Approximation
as a series of
wells.

Lec 21

Atoms



$\psi^*(x, y, z) \psi(x, y, z)$: Probability density of finding electron at (x, y, z)

Potential energy $V(x, y, z) \neq 0$.

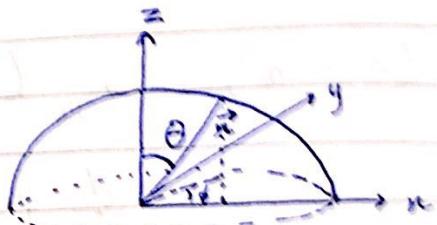
- electron-nucleus interaction is not zero.

$$\psi(x, y, z) \sim \psi(r, \theta, \phi)$$

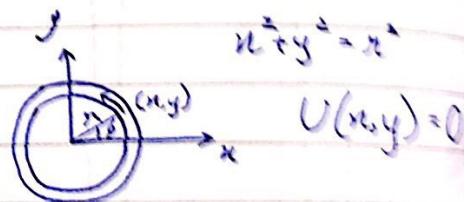
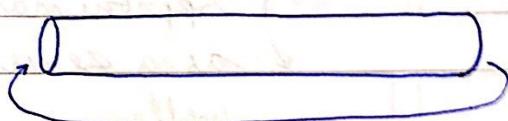
$$r \rightarrow 0 \text{ to } \infty$$

$$\theta \rightarrow 0 \text{ to } \pi$$

$$\phi \rightarrow 0 \text{ to } 2\pi$$



- One-Dimensional atom



$$\text{Schrodinger Equation: } -\frac{\hbar^2}{2m} \left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right] \psi(x, y) = E \psi(x, y)$$

Cartesian to Polar : $x = r \cos \phi, y = r \sin \phi$

r is a constant

$$\phi \in [0, 2\pi)$$

$$\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right] = \left[\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2} \right]$$

To Derive:

$$\frac{\partial f(r,y)}{\partial r} = \frac{\partial f}{\partial r} \cdot \frac{\partial r}{\partial r} + \frac{\partial f}{\partial y} \cdot \frac{\partial y}{\partial r}$$

- for this problem: r is fixed.

$$\frac{\partial^2}{\partial r^2} + \frac{\partial^2}{\partial y^2} = \frac{1}{r^2} \cdot \frac{\partial^2}{\partial \phi^2}$$

Schrodinger equation in polar coordinates:

$$-\frac{\hbar^2}{2m} \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2} \psi(r, \phi) = E \psi(r, \phi)$$

|||

$$-\frac{\hbar^2}{2m} \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2} \bar{\psi}(\phi) = E \bar{\psi}(\phi)$$

$$\psi(r, \phi) = R(r) \bar{\psi}(\phi)$$

$I = m r^2 \Rightarrow$ Moment of Inertia

$$-\frac{\hbar^2}{2I} \cdot \frac{\partial^2}{\partial \phi^2} \bar{\psi}(\phi) = E \bar{\psi}(\phi)$$

$$\frac{\partial^2}{\partial \phi^2} \bar{\psi}(\phi) = -\frac{2IE}{\hbar^2} \bar{\psi}(\phi)$$

→ general solution $\bar{\psi}(\phi) = A e^{im_1 \phi}$

$$\bar{\psi}(\phi + 2\pi) = \bar{\psi}(\phi)$$

$$A e^{im_1(\phi + 2\pi)} = A e^{im_1 \phi}$$

$$e^{im_1 2\pi} = 1 ; m_1 = 0, \pm 1, \pm 2, \pm 3, \dots$$

Substitute: $\Phi(\phi) = Ae^{im\phi}$ in the Schrödinger equation & determine E .
 $\Rightarrow E_{m_L} = \frac{m_L^2 \hbar^2}{2I}$ $m_L = 0, \pm 1, \pm 2, \dots$

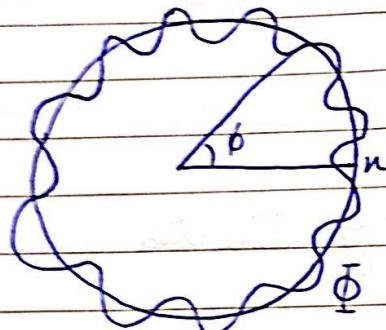
$m_L = \pm 2$ 2nd excited
 $m_L = \pm 1$ 1st excited
 $m_L = 0$ Ground

$$E_{m_L=0}$$

find A : normalize: $\Phi(\phi) = A e^{im_L \phi}$

$$\int_0^{2\pi} \Phi^*(\phi) \Phi(\phi) d\phi = 1 \rightarrow A = \frac{1}{\sqrt{2\pi}}$$

$$\therefore \Phi(\phi) = \frac{1}{\sqrt{2\pi}} e^{im_L \phi}; m_L = 0, \pm 1, \pm 2, \dots$$



Angular Momentum

$$\vec{L} = \vec{r} \times \vec{p}$$

$$i\hat{L}_x + j\hat{L}_y + k\hat{L}_z = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ n & y & z \\ P_x & P_y & P_z \end{vmatrix}$$

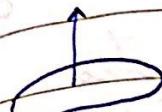
$$\hat{L}_z = nP_y - yP_n, \quad \hat{L}_z = \hat{i} \left(-i\hbar \frac{\partial}{\partial y} \right) - \hat{j} \left(-i\hbar \frac{\partial}{\partial n} \right)$$

Polar: $\hat{L}_z = \frac{\hbar}{i} \frac{\partial}{\partial \phi}$

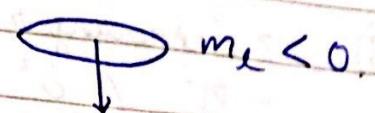
$\hat{L}_z \phi(\phi) = m_e \hbar \phi(\phi)$

Angular Momentum $= m_e \hbar$
 ↳ quantised

$$m_e = 0, \pm 1, \pm 2, \dots$$



$$m_l > 0$$



Angular Momentum
 quantum number
 $l = 0, 1, 2, \dots$

Particle on a sphere



$n \rightarrow$ fixed

$\theta \rightarrow 0$ to π

$\phi \rightarrow 0$ to 2π

$$-\frac{\hbar^2}{2m} \left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right] \psi(n, y, z)$$

$$(\hat{H} + n) \psi(n, y, z) = E \psi(n, y, z)$$

$\tan \theta = r$

$\cos \theta \sin \theta = y$

$\cos \theta = z$

$$\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \equiv -\nabla^2$$

$\hat{H} = -\nabla^2 + V$

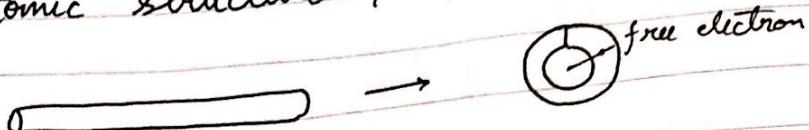
Angular momentum \vec{L} commutes with \hat{H}

$$\hat{L}_x \psi = i \hbar \psi$$

$$\hat{L}_y \psi = j \hbar \psi$$

Lec 22 Atomic Structure / Spectra.

Recap:



1) Polar : $(x, y) \rightarrow (r, \phi)$

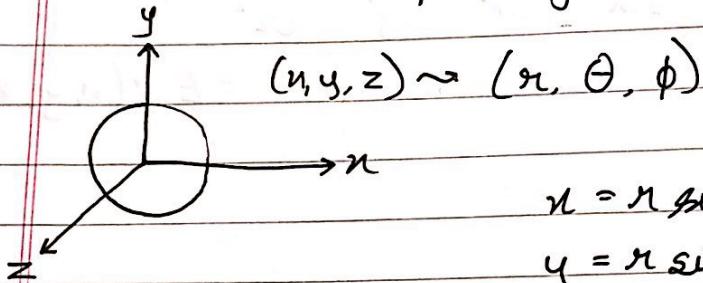
2) $\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \rightarrow \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2}$

3) \rightarrow Schrodinger equation in polar coordinates :

$$-\frac{\pi^2}{2m} \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2} \right) \Phi(\phi) = E \Phi(\phi)$$

4) Apply the boundary condition : $\Phi(\phi + 2\pi) = \Phi(\phi)$
 $\Rightarrow \Phi(\phi) = \frac{1}{\sqrt{2\pi}} e^{im\phi}$, m being the magnetic quantum number.

- Particle on a sphere of radius r .



$$x = r \sin \theta \cos \phi$$

$$y = r \sin \theta \sin \phi$$

$$z = r \cos \theta$$

Step 1: express : $\nabla^2 \equiv \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$

in terms of r, θ, ϕ .

Step 2 : rewrite the Schrodinger equation in terms of r, θ, ϕ .

$$\nabla^2 = \cancel{\frac{\partial^2}{\partial r^2}} + \frac{2}{r} \cancel{\frac{\partial}{\partial r}} + \frac{1}{r^2} \Lambda^2$$

$$\Lambda^2 \equiv \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right)$$

Schrödinger equation

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$$\nabla^2 \psi(r, \theta, \phi) = -\left(\frac{2IE}{\hbar^2}\right) \psi(r, \theta, \phi)$$

$$I = m\pi^2$$

Solution: Spherical harmonics $Y_{lm}(\theta, \phi)$
↳ 2 quantum numbers: l, m

$$Y_{lm}(\theta, \phi) = \sqrt{\frac{(2l+1)}{4\pi}} \frac{(l-m)!}{(l+m)!} P_{lm}(\cos \theta) e^{im\phi}$$

$$\psi(r, \theta, \phi) = Y_{lm}(\theta, \phi)$$

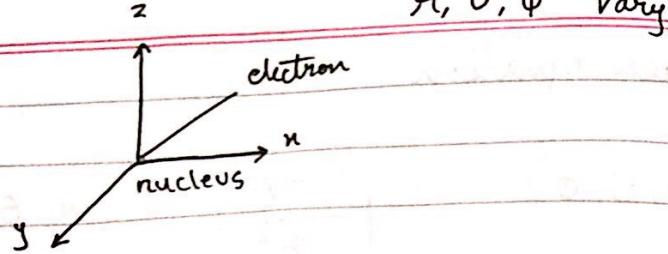
- here $P_{lm}(\cos \theta)$ is called the associated Legendre Polynomial
 $m \rightarrow$ magnetic quantum number $[0, \pm 1, \pm 2, \dots]$
 $l \rightarrow$ angular momentum quantum number $[0, 1, 2, \dots]$

Spherical Harmonics

$$\begin{array}{c|c|c|c} l & m & Y_{lm}(\theta, \phi) \\ \hline 0 & 0 & \frac{1}{2} \sqrt{\frac{1}{\pi}} \\ \hline 1 & 0 & \frac{1}{2} \sqrt{\frac{3}{\pi}} \cos \theta \\ \hline 1 & \pm 1 & \mp \sqrt{\frac{3}{2\pi}} \sin \theta e^{\pm i\phi} \end{array}$$

$l=0 \rightarrow$ s orbitals [spherical]

H-atom



r, θ, ϕ vary

$$U(r) = -\frac{1}{4\pi\epsilon_0} \frac{e^2}{r}$$

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \quad [\text{Derive}]$$

$$\rightarrow -\frac{\hbar^2}{2m} \left(\frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) \right) \psi(r, \theta, \phi)$$

$$\begin{aligned} & \# -\frac{\hbar^2}{2m} \left(\frac{1}{\sin\theta} \frac{\partial}{\partial\theta} \left(\sin\theta \frac{\partial}{\partial\theta} \right) + \frac{1}{\sin^2\theta} \frac{\partial^2}{\partial\phi^2} \right) \\ & - \frac{e^2}{4\pi\epsilon_0 r} \end{aligned} \quad \left. \begin{array}{l} \text{Kinetic energy} \\ \text{Potential energy} \end{array} \right\}$$

- Method of separation of variables:

$$\psi(r, \theta, \phi) = R(r) \cdot Y(\theta, \phi)$$

radial

angular

$$\Rightarrow -\frac{\hbar^2}{2m} Y(\theta, \phi) \left(\frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} R(r) \right) \right)$$

radial + angular = 0.

Using the particle on a sphere model.

$Y_{lm}(\theta, \phi)$: solution to angular part

$R_{nl}(r)$: solution to radial part

3 quantum numbers : $n = 1, 2, 3, \dots$

$l = 0, 1, 2, \dots, n-1$

$m = 0, \pm 1, \pm 2, \dots, \pm l$.

$$\Psi_{nlm} = R_{nl}(r) Y_{lm}(\theta, \phi)$$

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• Principal Quantum Number

n l m ψ

$$1 \quad 0 \quad 0 \quad \psi_{100} = \frac{1}{\sqrt{\pi}} \left(\frac{1}{a_0} \right)^{3/2} e^{-r/a_0}$$

$$r = \frac{a_0}{a_0} \rightarrow a_0 = \text{Bohr Radius.}$$

1s orbital