

Tangent Vectors
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Parallel Transport
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Covariant Derivatives
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Connections
oooooooooooo

What are
Connections?

Z. Virgilio

What are connections?

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Tangent Vectors
Parallel Transport
Covariant
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Vector Fields

The simplest idea of a vector field is a field of arrows. The idea is best explained in a picture:

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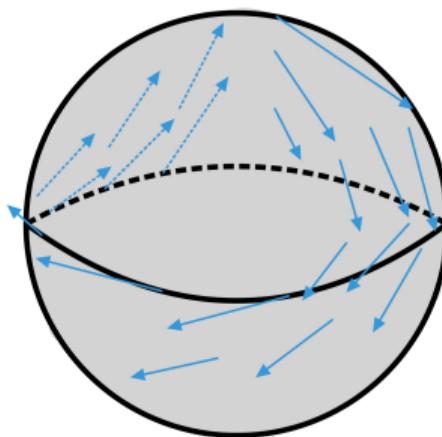
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Vector Fields

The simplest idea of a vector field is a field of arrows. The idea is best explained in a picture:



We have S^2 with a collection vectors drawn on it, all tangent to the surface at their base point.

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Vector Fields

More algebraically, given a field of arrows, we want to differentiate any function in the direction of the arrows.

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For the simple case of \mathbb{R}^n , given $f : \mathbb{R}^n \rightarrow \mathbb{R}$ and v a vector field on \mathbb{R}^n we can form this directional derivative, which we write

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Vector Fields

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For the simple case of \mathbb{R}^n , given $f : \mathbb{R}^n \rightarrow \mathbb{R}$ and v a vector field on \mathbb{R}^n we can form this directional derivative, which we write

$$vf$$

For x_1, \dots, x_n the coordinates of \mathbb{R}^n , and partial derivatives $\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n}$, we write

$$vf = v^1 \frac{\partial}{\partial x_1} f + \cdots + v^n \frac{\partial}{\partial x_n} f$$

and so we write $v = v^1 \frac{\partial}{\partial x_1} + \cdots + v^n \frac{\partial}{\partial x_n}$

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Vector Fields

In the formula

$$v = v^1 \frac{\partial}{\partial x_1} + \cdots + v^n \frac{\partial}{\partial x_n}$$

the partial derivatives aren't differentiating anything. The right hand side can be fed a function at any time, however.

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Vector Fields

In the formula

$$v = v^1 \frac{\partial}{\partial x_1} + \cdots + v^n \frac{\partial}{\partial x_n}$$

the partial derivatives aren't differentiating anything. The right hand side can be fed a function at any time, however.

But this leaves us with an identification between v , the vector field, and an operator that takes directional derivatives. Despite this seeming abuse of notation, this is a very convenient notion.

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Vector Fields

In general, for a manifold M , vector fields are defined as functions

$$v : C^\infty(M) \rightarrow C^\infty(M)$$

satisfying the following:

$$v(f + g) = v(f) + v(g)$$

$$v(\alpha f) = \alpha v(f)$$

$$v(fg) = v(f)g + fv(g)$$

This can be thought of as the basic rules all directional derivatives should follow.

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This can be thought of as the basic rules all directional derivatives should follow.

We shall denote the collection of vector fields on a manifold M by:

$$\text{Vect}(M)$$

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Tangent Vectors

Going back to our initial idea of a vector field assigning an arrow to each point of M , we call this kind of arrow a *tangent vector*.

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Tangent Vectors

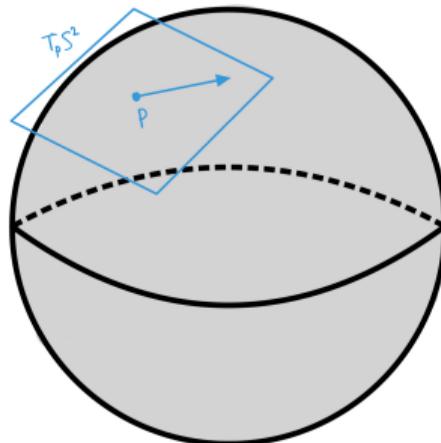
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Tangent Vectors

Going back to our initial idea of a vector field assigning an arrow to each point of M , we call this kind of arrow a *tangent vector*. For example:



The plane of all vectors tangent to S^2 at the point p is called $T_p S^2$, *the tangent plane at p* .

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Tangent Vectors

We formally define a tangent vector at a point $p \in M$ as a map

$$v_p : C^\infty(M) \rightarrow \mathbb{R}$$

by $v_p(f) = v(f)(p)$ (*recall $v(f)$ is again a $C^\infty(M)$ function*).

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by $v_p(f) = v(f)(p)$ (*recall $v(f)$ is again a $C^\infty(M)$ function*).

This definition of a tangent vector in fact looks like the arrows in the diagrams since:

$$(v_p + w_p)(f) = v_p(f) + w_p(f)$$

and

$$(\alpha v_p)(f) = \alpha v_p(f)$$

So they indeed form a vector space.

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Tangent Vectors

If we have a curve on our surface M , we can assign a tangent vector to each point of the curve.

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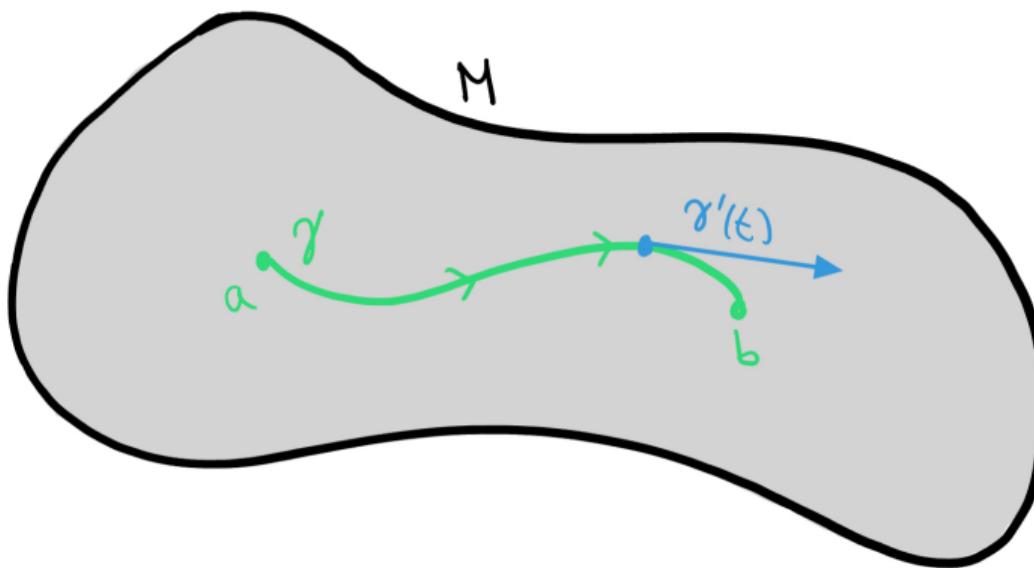
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Tangent Vectors

A curve γ is a function from \mathbb{R} (or some interval) to M that is *smooth*.

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A curve γ is a function from \mathbb{R} (or some interval) to M that is *smooth*. This means for any $f \in C^\infty(M)$, $f(\gamma(t))$ depends smoothly on t .

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The tangent vector $\gamma'(t)$ is defined for each tangent space $T_{\gamma(t)}M$, to send any function f to

$$\frac{d}{dt}f(\gamma(t))$$

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The tangent vector $\gamma'(t)$ is defined for each tangent space $T_{\gamma(t)}M$, to send any function f to

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In words, $\gamma'(t)$ differentiates functions in the direction that the curves γ is 'moving' at time t .

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Tangent Vectors

An important note for the rest of the talk:

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Tangent Vectors

An important note for the rest of the talk: each tangent space is a vector space, but there is no sensible way to add tangent vectors living in separate tangent spaces.

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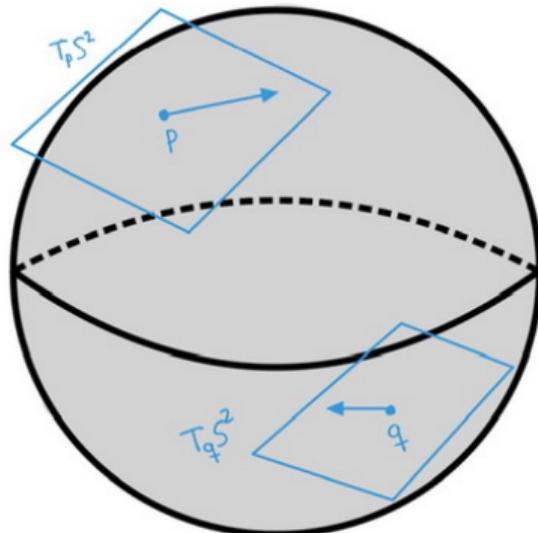
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Tangent Vectors

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Example

To each point (x, y) in \mathbb{R}^2 , we will define the basis to the tangent space to be given by:

$$e_1 = \begin{pmatrix} \cos(x) \\ \sin(x) \end{pmatrix}, \quad e_2 = \begin{pmatrix} -\sin(x) \\ \cos(x) \end{pmatrix}$$

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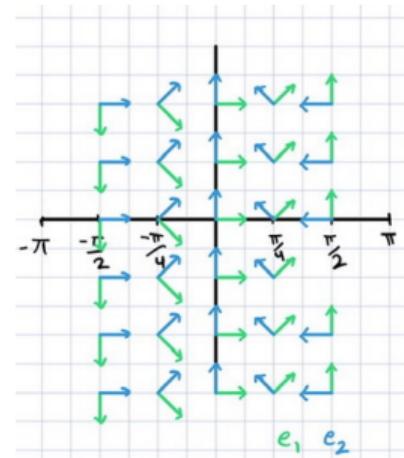
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$$e_1 = \begin{pmatrix} \cos(x) \\ \sin(x) \end{pmatrix}, \quad e_2 = \begin{pmatrix} -\sin(x) \\ \cos(x) \end{pmatrix}$$

So we have something along the lines of:



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Example

Now we consider a constant vector field given by:

$$v(x, y) = \begin{pmatrix} 2 \\ 1 \end{pmatrix}$$

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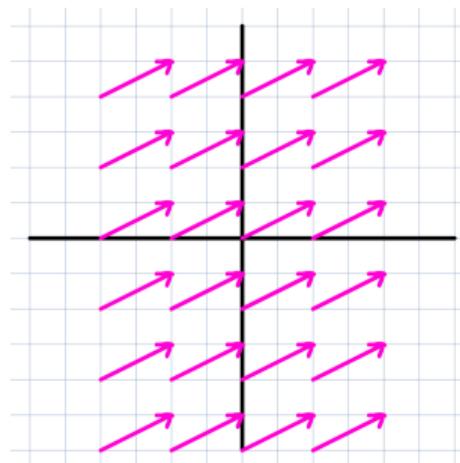
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Example

Now we consider a constant vector field given by:

$$v(x, y) = \begin{pmatrix} 2 \\ 1 \end{pmatrix}$$

Which looks like:



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Example

In our basis e_1, e_2 , even though our vector field is constant, the components change as x varies:

$$\begin{aligned}v(x, y) &= v^1 e_1 + v^2 e_2 \\&= (2 \cos x + \sin x) e_1 + (\cos x - 2 \sin x) e_2\end{aligned}$$

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So in this basis, we have:

$$v(x, y) = \begin{pmatrix} 2 \cos x + \sin x \\ \cos x - 2 \sin x \end{pmatrix}$$

which does not appear constant at all.

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Example

We have:

$$v^1 = 2 \cos x + \sin x \quad v^2 = \cos x - 2 \sin x$$

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Example

We have:

$$v^1 = 2 \cos x + \sin x \quad v^2 = \cos x - 2 \sin x$$

So we can compute how the components change as x changes by taking the partial derivatives:

$$v_x^1 = -2 \sin x + \cos x$$

$$v_x^2 = -\sin x - 2 \cos x$$

Of course, $v_y^1 = v_y^2 = 0$.

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Of course, $v_y^1 = v_y^2 = 0$.

So how do we tell if the vector field is truly changing, or if this is just an artefact of our basis?

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Example

In \mathbb{R}^2 (and \mathbb{R}^n in general) there is in fact a notion of how to add vectors between different tangent spaces:

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Example

In \mathbb{R}^2 (and \mathbb{R}^n in general) there is in fact a notion of how to add vectors between different tangent spaces: *slide the vector along a straight line (a geodesic) between two points, keeping the angle constant between the vector and the line. Add the vectors once they are in the same tangent space*

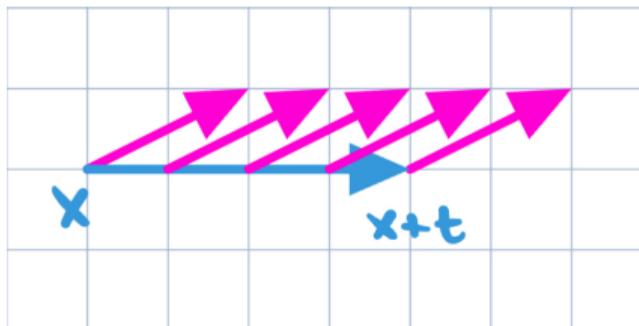
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Example

Due to this we can make sense of the notion a 'true' derivative:

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Example

Due to this we can make sense of the notion a 'true' derivative: We differentiate along the path $(x + t, y)$:

$$\frac{dv}{dt} = \lim_{\epsilon \rightarrow 0} \frac{v(x + \epsilon, y) - v(x, y)}{\epsilon}$$

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Example

Due to this we can make sense of the notion a 'true' derivative: We differentiate along the path $(x + t, y)$:

$$\frac{dv}{dt} = \lim_{\epsilon \rightarrow 0} \frac{v(x + \epsilon, y) - v(x, y)}{\epsilon}$$

But this is just the x -partial derivative of v :

$$\begin{aligned}\frac{dv}{dt} &= \frac{d}{dt}(v^1 e_1 + v^2 e_2) \\ &= v_x^1 e_1 + v^1 (e_1)_x + v_x^2 e_2 + v^2 (e_2)_x \\ &= (v_x^1 - v^2) e_1 + (v^1 + v_x^2) e_2 = 0\end{aligned}$$

Example

$$\frac{dv}{dt} = (v_x^1 - v^2) e_1 + (v^1 + v_x^2) e_2 = 0$$

Notice that in each component e_i , there are two terms:

- 1 the partial derivative of the component, v_x^i
- 2 a correction term, which turns out to be v^2 for e_1 and v^1 for e_2

Example

$$\frac{dv}{dt} = (v_x^1 - v^2)e_1 + (v^1 + v_x^2)e_2 = 0$$

Notice that in each component e_i , there are two terms:

- 1 the partial derivative of the component, v_x^i
- 2 a correction term, which turns out to be v^2 for e_1 and v^1 for e_2

The notion of the 'true' derivative sees that the vector field is constant, despite the changing components. We wish to arrive at this idea for manifolds more complicated than \mathbb{R}^n .

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Parallel Transport

To make sense of this notion of true derivative on a more general manifold, we need to be able to transport a vector from one tangent space to another.

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Parallel Transport

To make sense of this notion of true derivative on a more general manifold, we need to be able to transport a vector from one tangent space to another. This is the idea of *parallel transport*.

The way we parallel transport a tangent vector $w \in T_p M$ to $T_q M$ for some nearby point q is to simply place w on q . There is no reason for this to be a tangent vector to q , but if we consider M to be living in some larger ambient space, we can 'rotate' w until it lies in $T_q M$.

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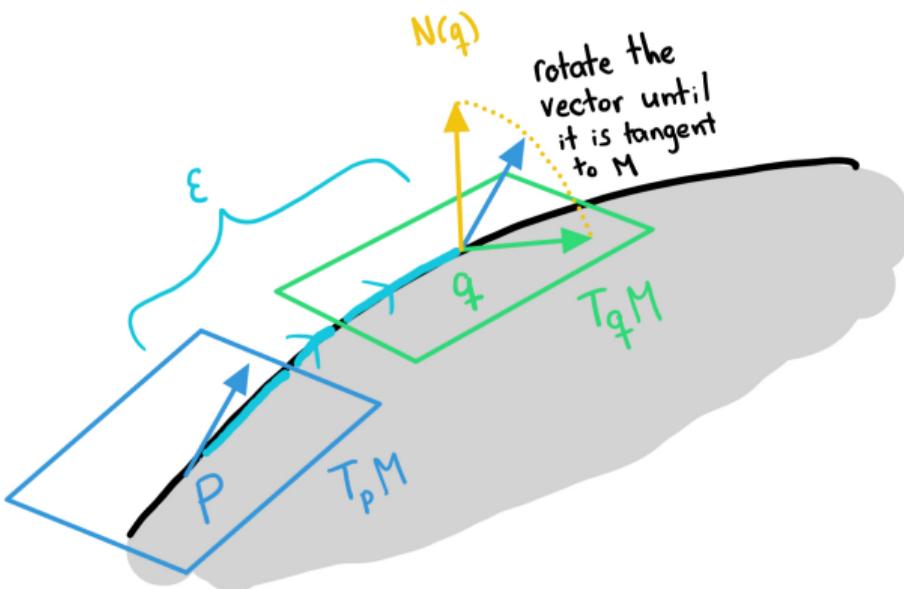
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A picture is worth a thousand words:



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Parallel Transport

In general, parallel transport is not unique and depends on the path connecting two points:

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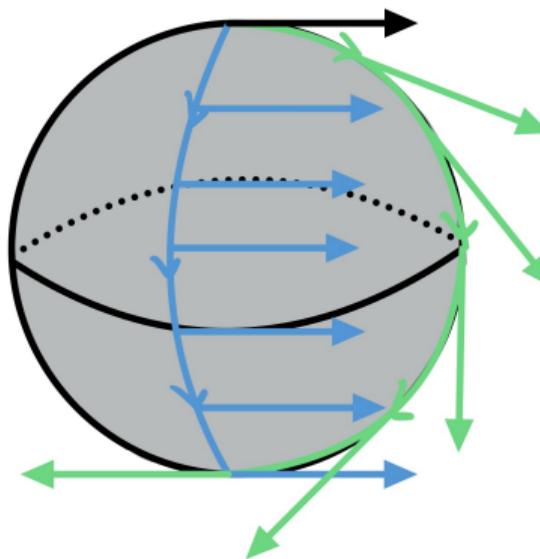
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Parallel Transport

This is true even for a path which is a closed loop:

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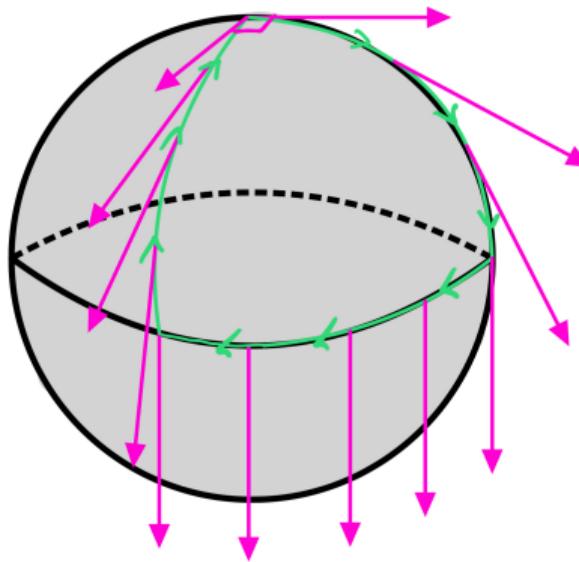
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Parallel Transport

Some important properties of parallel transport:

- 1 If two vectors in the same tangent space are parallel transported along the same curve, then the angle between them remains constant
- 2 From an initial point and an arbitrary initial tangent vector v (velocity), parallel transport of v along v yields a unique geodesic arising from these conditions.

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Covariant Derivatives

Consider a geodesic G on some surface M . Let $G'(t) = v(t)$, the velocity of the curve. Suppose there is a vector field on G , $w(t)$. Parallel transport a tangent vector $w(p) \in T_p M$ along G .

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Since G is a geodesic, it can be 'unrolled' into a flat strip.

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Since G is a geodesic, it can be 'unrolled' into a flat strip. We can then measure the how much $w(t)$ is changing as it moves along G compared to the parallel transport of $w(p)$. This is what is meant by differentiating a vector field $w(t)$ in the direction $v(t)$.

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Covariant Derivatives

Consider a geodesic G on some surface M . Let $G'(t) = v(t)$, the velocity of the curve. Suppose there is a vector field on G , $w(t)$. Parallel transport a tangent vector $w(p) \in T_p M$ along G .

Since G is a geodesic, it can be 'unrolled' into a flat strip. We can then measure the how much $w(t)$ is changing as it moves along G compared to the parallel transport of $w(p)$. This is what is meant by differentiating a vector field $w(t)$ in the direction $v(t)$. It is denoted:

$$D_v w$$

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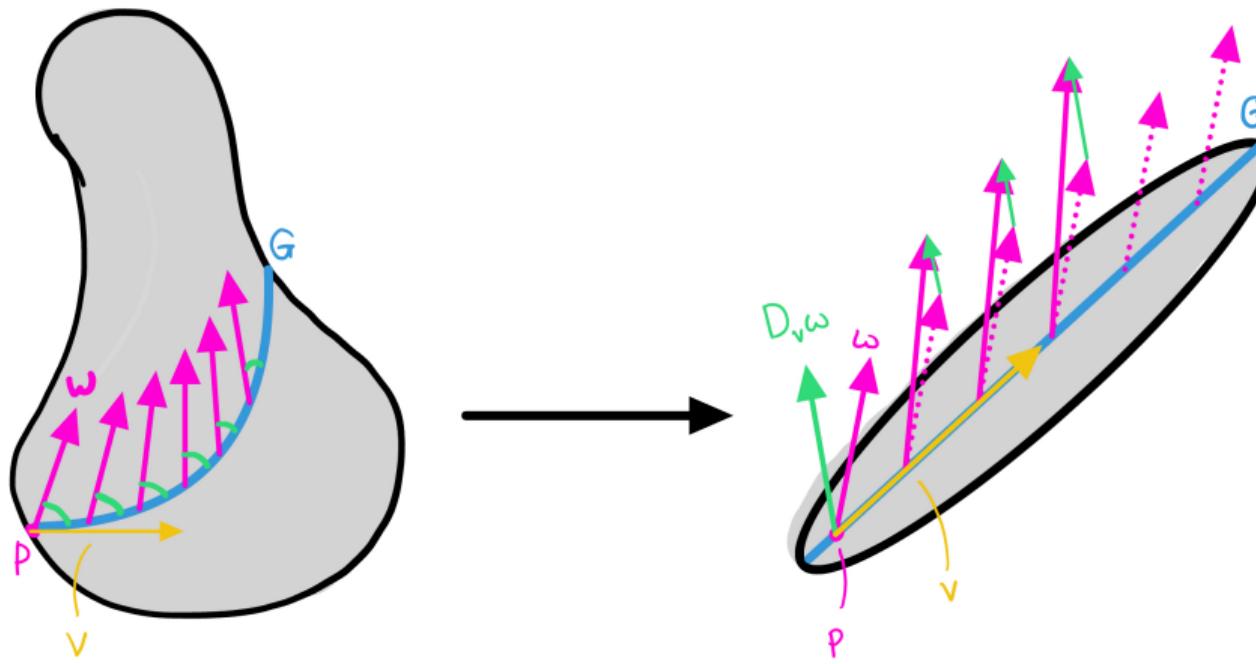
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Once again, pictures are superior:



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$D_v w$ is quantifying the rate at which $w(t)$ is changing as it moves in the direction of v .

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$D_v w$ is quantifying the rate at which $w(t)$ is changing as it moves in the direction of v .

If we are not moving along a geodesic, the picture remains the same, as long as we consider only very short segments.

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One more fact that is obvious from this picture is that $D_v w = 0$ if and only if w is parallel transported along v .

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We can now extend D to be defined on the entirety of a manifold M , so now v and w are vector fields and $D_v w$ is map that assigns to each $x \in M$ the tangent vector corresponding to $D_{v(x)} w(x)$.

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Connections

We can now extend D to be defined on the entirety of a manifold M , so now v and w are vector fields and $D_v w$ is map that assigns to each $x \in M$ the tangent vector corresponding to $D_{v(x)} w(x)$.

In this way, we may consider D as a map

$$D : \text{Vect}(M) \times \text{Vect}(M) \rightarrow \text{Vect}(M)$$

However, the inputs are not symmetric.

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D satisfies the following properties for $v_1, v_2, w_1, w_2 \in \text{Vect}(M)$ and $f \in C^\infty(M)$ and α a scalar:

$$D_{v_1}(\alpha w_1) = \alpha D_{v_1} w_1$$

$$D_{v_1}(w_1 + w_2) = D_{v_1} w_1 + D_{v_1} w_2$$

$$D_{v_1}(fw_1) = v(f)w_1 + fD_{v_1} w_1 \quad (*)$$

$$D_{v_1+v_2} w_1 = D_{v_1} w_1 + D_{v_2} w_1$$

$$D_{fv_1} w_1 = fD_{v_1} w_1 \quad (\dagger)$$

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$$D_{v_1+v_2} w_1 = D_{v_1} w_1 + D_{v_2} w_1$$

$$D_{fv_1} w_1 = fD_{v_1} w_1 \quad (\dagger)$$

In particular, for the second input D behaves like a derivative ($(*)$ this is known as the *Leibnitz rule*), while D is linear with respect to functions in the first input (\dagger)

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So D is collection of all the information about differentiation vector fields and hence parallel transport.

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So D is collection of all the information about differentiation vector fields and hence parallel transport.

What about a more general setting? So far we have been exclusively looking at the tangent bundle. This connection is in fact canonical and called the *Levi-Civita connection*.

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In more generality, we don't necessarily have the same notion of angles to measure when vectors are parallel or not, or be able to rotate vectors. We have more general *vector bundles*.

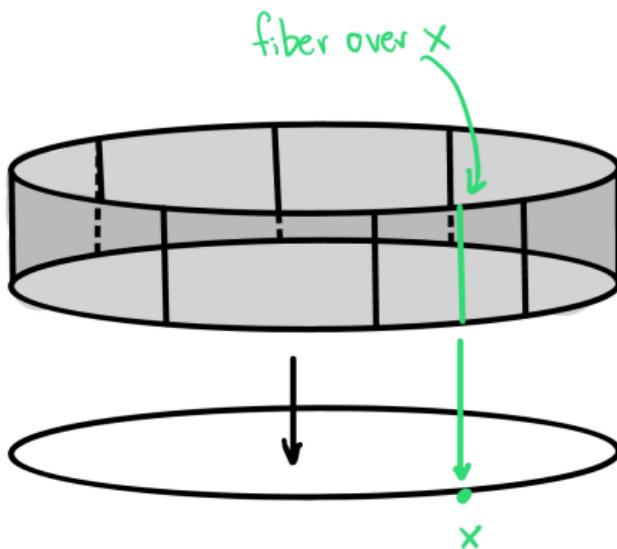
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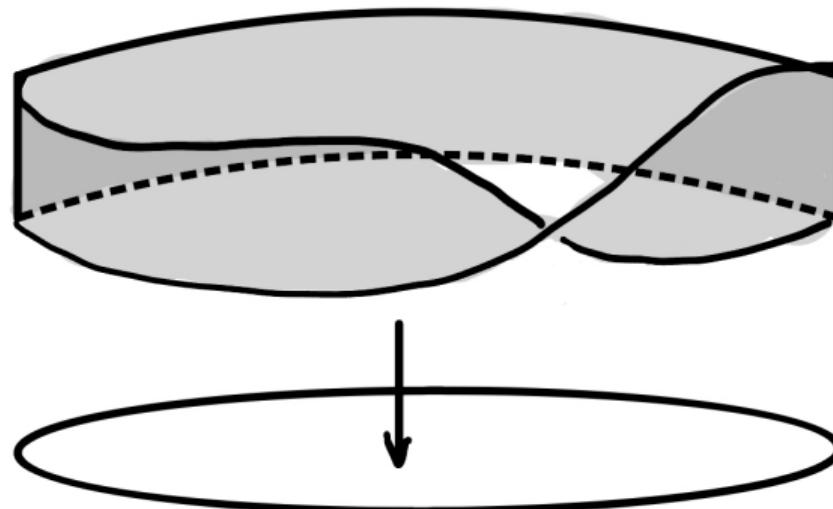
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A non-trivial example of a vector bundle is the mobius bundle:



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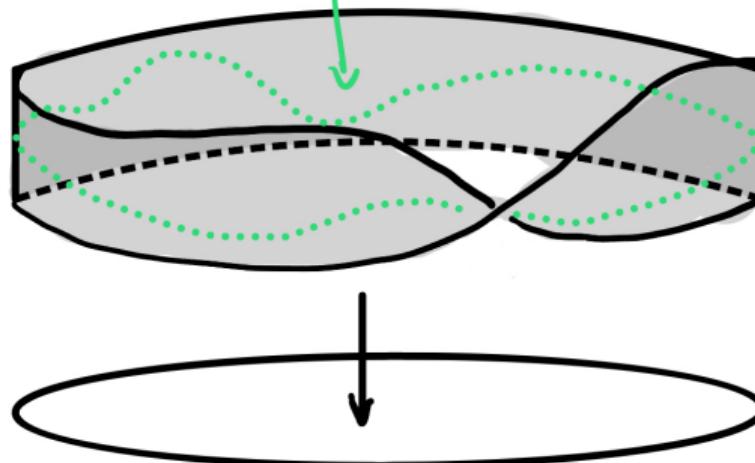
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A section of vector bundle:

Section of the
vector bundle



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Relating the original example

In the prior case we had:

$$D : \text{Vect}(M) \times \text{Vect}(M) \rightarrow \text{Vect}(M)$$

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Relating the original example

In the prior case we had:

$$D : \text{Vect}(M) \times \text{Vect}(M) \rightarrow \text{Vect}(M)$$

Now for a vector bundle E over M , it becomes:

$$D : \text{Vect}(M) \times \Gamma(E) \rightarrow \Gamma(E)$$

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Relating the original example

We replace vector fields in the second input with sections of our bundle, denote $\Gamma(E)$. So for $v \in \text{Vect}(M)$ and $s \in \Gamma(E)$ we can calculate the covariant derivative:

$$\begin{aligned} D_v s &= D_{v^j \partial_j} s \\ &= v^j D_{\partial_j} s \\ &= v^j (D_{\partial_j} (s^i e_i)) \\ &= v^j ((\partial_j s^i) e_i + A_{ji}^k s^i e_k) \\ &= v^j (\partial_j s^i + A_{jk}^i s^k) e_i \end{aligned}$$

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Relating the original example

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which bears a striking resemblance to the 'true' derivative from earlier:

$$\frac{dv}{dt} = (v_x^1 - v^2) e_1 + (v^1 + v_x^2) e_2$$

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Thank You!

The end!

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