

[3]

## 1. What is Olber's paradox and how was it resolved?

The Olber's paradox was that even though the universe was assumed to be infinitely large, eternally old and static; the night sky appeared dark. Ideally if the universe assumptions were true then the sky would always appear as hot as the surface of the sun at all times as eternally old implies that all the light has reached us.

Few way of resolving involved correcting the observations. Firstly, our line of sight does not tend to infinity, it may be blocked by a non luminous object. Also the assumption of eternally old was challenged as we can assume that some light has not yet reached us.

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2. A galaxy emits light of wavelength of  $1.21567 \times 10^{-7} \text{m}$ . When observed at earth, the wavelength is  $1.21445 \times 10^{-7} \text{m}$ . If the comoving distance of the galaxy is 2.5 million light years then find the magnitude and direction of the velocity of the galaxy with respect to the earth. 1 light year = 0.3067 parsec. [3]

Given, emitted  $\lambda_e = 1.21567 \times 10^{-7} \text{m}$   
 observed  $\lambda_o = 1.21445 \times 10^{-7} \text{m}$

comoving distance  $r = 2.5 \times 10^6 \text{ ly} = 0.767 \times 10^6 \text{ pc}$

The redshift is given by  $z = \frac{\lambda_{\text{obs}} - \lambda_{\text{em}}}{\lambda_{\text{em}}} = \frac{-0.00122 \times 10^{-7} \text{m}}{1.21567 \times 10^{-7} \text{m}} \approx -10^{-3}$

using Hubble's eq.  $v = H_0 r$

since  $z$  is -ve, the galaxy is moving closer to the earth.

$= 72 \text{ km s}^{-1} \text{ pc}^{-1} \times r$

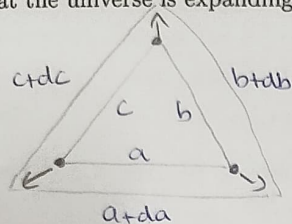
$= 72 \times 0.767 \times 10^6 \text{ km s}^{-1}$

$= 55.2 \times 10^6 \text{ km s}^{-1}$  towards the earth.

Ans.

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3. Three galaxies which are not in the same line form an acute triangle. Using the cosmological principle, and the fact that the universe is expanding, show that the expansion is according to the Hubble's law. [6]



cosmological principle  $\rightarrow$  universe is isotropic and homogeneous  
universe is expanding.

$\rightarrow$  triangle similarity is preserved due to cosmological principle.

$\hookrightarrow$  ratio of sides is constant.

$$\Rightarrow \frac{a+da}{a} = \frac{b+db}{b} = \frac{c+dc}{c} = 1 + \frac{da}{a} = k$$

~~all sides expand by same factor~~

~~expanding~~

$$\text{Let } v_1 = \frac{da}{dt} \quad v_2 = \frac{db}{dt} \quad v_3 = \frac{dc}{dt}$$

$$\text{we have } 1 + \frac{da}{a} \Rightarrow \frac{da}{dt a} = \frac{(k-1)}{dt}$$

$$\text{similarly } 1 + \frac{db}{b} \Rightarrow \frac{db}{dt b} = \frac{(k-1)}{dt}$$

$$1 + \frac{dc}{c} \Rightarrow \frac{dc}{dt c} = \frac{(k-1)}{dt}$$

} all are same.

$$\therefore \frac{\dot{a}}{a} = \frac{\dot{b}}{b} = \frac{\dot{c}}{c} = H_0 \text{ (Hubble's constant)}$$

$$\text{and } \frac{da}{dt} = a \left( \frac{k-1}{dt} \right) = H_0 a \quad \rightarrow v_1 = H_0 a$$

$$\text{similarly } v_2 = H_0 b \quad v_3 = H_0 c$$

} Hubble's law.

Ans

4. Suppose the universe was expanding at much faster rate in the past than it is now. If the present Hubble constant is  $H_0$ , explain if the universe older or younger than  $H_0^{-1}$ . [2]

Say  $H_0'$  is the Hubble const in the faster expanding universe, with  $a' > a$

$\hookrightarrow$  present universe

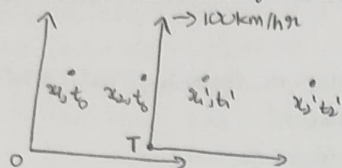
$$\therefore H_0' > H_0 \quad \text{as } H_0 = \frac{\dot{a}}{a}$$

$$\Rightarrow \frac{1}{H_0'} < \frac{1}{H_0}$$

$\Rightarrow t' < t$   $\therefore$  The new universe would be younger than our universe.



5. A train 1/2 Km long (as measured by an observer on the train) is traveling at a speed of 100 km/hr. Two lightning bolts strike the ends of the train simultaneously as determined by an observer on the ground. What is the time separation of the strikes as measured by an observer on the train? [5]



Ground frame - end of train  $\rightarrow (x_1)$  and  $(x_2)$  at time  $t_0$  (simultaneous)

Train frame - ends:  $(t_1', x_1')$   $(t_2', x_2')$

Also,  $x_2' - x_1' = 500\text{m}$  and  $v = 100 \frac{\text{km}}{\text{hr}} = \frac{100 \times 1000}{3600} = \frac{1000}{36} \text{ m/s}$

Speed of train

Assume relative vel along x axis only.

using Lorentz Transformation:  $x' = \begin{bmatrix} \gamma & -\gamma\beta \\ -\gamma\beta & \gamma \end{bmatrix} x$   $\beta = \frac{v}{c}$  and  $\gamma = \frac{1}{\sqrt{1-\beta^2}}$

$t_1' = \gamma t_0 - \gamma\beta x_1$

$x_1' = -\gamma\beta t_0 + \gamma x_1$

$t_2' = \gamma t_0 - \gamma\beta x_2$

$x_2' = -\gamma\beta t_0 + \gamma x_2$

now,  $x_2' - x_1' = \gamma(x_2 - x_1) = 500$

and  $t_2' - t_1' = -\gamma\beta(x_2 - x_1) = -\frac{\gamma\beta \times 500}{\gamma} = -500\beta$

$\beta = \frac{v}{c} = \frac{1000}{36 \times 3 \times 10^8} = \frac{1}{108 \times 10^5} \text{ s}$

$\therefore t_2' - t_1' = \frac{-500}{108 \times 10^5} = -4.65 \times 10^{-5} \text{ sec}$

$\therefore$  The front end ( $x_2'$ ) of the train strike is observed  $4.65 \times 10^{-5} \text{ s}$  before the back end ( $x_1'$ )

Ans.

6. Starting from the fluid equation

$$\frac{\partial p}{\partial t} + \frac{a}{a} (3p + 3P) = 0,$$

show that in a universe dominated by radiation only, the energy density scales as  $\rho \propto a^{-4}$ .

Dominated by radiation  $\Rightarrow P_M = P_A = 0$  ~~Pressure~~ PIP

Let the pressure energy density  $\rho = \omega$ .  $\rightarrow$  ~~PIP~~

then  $\frac{\partial p}{\partial t} + \frac{da}{da} (3p + 3P) = 0 \Rightarrow \frac{\partial p}{\partial t} + \frac{da}{da} (3\omega) = 0$

$\Rightarrow \frac{\partial p}{\partial t} + \frac{da}{da} (3\omega) = 0 \Rightarrow \frac{\partial p}{\partial t} = -\frac{da}{da} (3\omega)$

$\Rightarrow \ln(p) = -3(\omega H) \ln(a) + c = \ln(a^{-3(\omega H)})$

$\therefore p \propto a^{-3(\omega H)}$

$\Rightarrow p \propto a^{-3(\omega H)}$

for radiation dominated  $\Rightarrow P = \rho/3 \Rightarrow \omega = \frac{P}{\rho} = \frac{1}{3}$

$\Rightarrow p \propto a^{-3(1+\frac{1}{3})} = a^{-4} = a^{-4}$

$\therefore p \propto a^{-4}$

Ans.

[5]

$$\frac{\partial p}{\partial t} + \frac{a}{a} (3p + 3P) = 0$$

$$\frac{\partial p}{\partial t} + \frac{da}{da} (3\omega + 3) = 0$$

$$\frac{\partial p}{\partial t} + \frac{da}{da} (3)(\omega H) = 0$$

$$\left( \omega \frac{\partial p}{\partial t} \right) = -3(\omega H) \frac{da}{da}$$

$$\omega \ln(p) = \ln(a) \left( -\frac{3\omega H}{\omega} \right)$$

$$p \propto a^{-3(\omega H)}$$

$$\omega \ln\left(\frac{p}{\rho}\right)$$

$$\frac{p}{\rho} \propto a^{-4}$$

7. Consider a light source at a redshift  $z = 3$  in an Einstein-de Sitter universe, where  $\Omega_M = 1.0$  and  $\Omega_\Lambda = 0$ ,  $\Omega_R = 0$ ,  $\Omega_K = 0$ , i.e., the universe is dominated by non-relativistic matter only. Calculate the distance the light has travelled to reach us. You can use  $H_0 = 70 \text{ km/s/Mpc}$ . [7]

Einstein de sitter universe,  $\Omega_M = 1, \Omega_\Lambda = 0, \Omega_R = 0, \Omega_K = 0$

$$2 = 3 = 1 + \frac{a(t_0)}{a(t_e)}$$

$$\Rightarrow \frac{a(t_0)}{a(t_e)} = 2$$

$$\frac{\rho_0}{\rho_e} = \left( \frac{a(t_e)}{a(t_0)} \right)^3 = \frac{1}{8}$$

$$\Rightarrow \rho_0 = \frac{1}{8} \rho_e$$

0  $\rightarrow$  present  $e \rightarrow$  emitted.

$$\text{at the time of emission } \rho_M^0 = 1 = \frac{\rho_M^0}{\rho_M^0} \Rightarrow \rho_e = \rho_M^0 = \rho_e$$

currently

$$\rho_M = \frac{\rho_0}{\rho_e} = \frac{1}{8} \frac{\rho_e}{\rho_e} = \frac{1}{8} \quad H^2 = H_0^2 \Omega_M$$

$$\Rightarrow H^2 = H_0^2 \left( \rho_M^0 \left( \frac{a_0}{a} \right)^3 \right) = \frac{H_0^2}{8}$$

$$H^2 = H_0^2 \left( \rho_M^0 \left( \frac{a_0}{a} \right)^3 \right) = H_0^2 \left( 1 \cdot \left( \frac{1}{2} \right)^3 \right)$$

$$= \frac{H_0^2}{8}$$

$$\Rightarrow H = \frac{H_0}{2\sqrt{2}}$$

$$\frac{a_0}{a} = \frac{a(t_e)}{a(t_0)} = \frac{1}{2}$$

$$t = \int dt = \int \frac{dx}{H_0 x \sqrt{\Omega_M x^{-3}}} \quad \left( x = \frac{1}{1+z} \right)$$

$\Rightarrow$  just  $\Omega_R = \Omega_K = \Omega_\Lambda = 0, \Omega_M = 1$

$$= \int_0^1 \frac{dx}{H_0 x^{1/2}} = \frac{1}{H_0} \int_0^1 x^{-1/2} dx = \frac{1}{H_0} \left[ \frac{x^{1/2}}{1/2} \right]_0^1$$

$$= \frac{1}{H_0} \left[ \frac{1}{1/2} - 0 \right] = \frac{2}{3H_0}$$

$$t = \frac{2}{3H_0} = \frac{\text{dist}}{\text{speed} \rightarrow c = 1}$$

$$\Rightarrow \text{dist} = \frac{2}{3H_0} = \frac{2}{3 \times 70 \text{ km/s/Mpc}}$$

$$d = \frac{2}{3 \times 3.25 \times 10^{-14} \text{ m}}$$

$$= \frac{20}{3 \times 3.25} \times 10^{13} \text{ m}$$

$$= 2.05 \times 10^{13} \text{ m}$$

Ans.

$$= \frac{2}{3} \times 3.25 \times 10^{14} \text{ m}$$

$$= 2.16 \times 10^{14} \text{ m}$$

$$\frac{1000}{3.25 \times 9.46 \times 10^{15}} = 30.8$$

$$\frac{100}{30.8 \times 10^{14}} = 3.25$$



8. The definition of the horizon distance is

$$d_h = a(t_0) \int_0^{r_h} \frac{dr}{\sqrt{1 - Kr^2}},$$

where  $(r_h, 0, 0)$  is any point in the particle horizon, and  $t_0$  is the present age of the universe. Show that in a radiation dominated universe  $d_h = 2t_0$  (in the unit  $c = 1$ ). How do you explain the fact that  $d_h$  is greater than the length of the path that light travels in time  $t_0$ .

[6]

as particle horizon is based on the light  $ds^2 = 0$  property,

$$\int_0^{r_h} \frac{dr}{\sqrt{1 - Kr^2}} = \int_0^{t_0} \frac{dt}{a(t)} \quad \hookrightarrow \quad ds^2 = -dt^2 + a(t)^2 \left( \frac{1}{1 - Kr^2} \right) dr^2 + a(t)^2 (d\theta^2 + d\phi^2)$$

$$\Rightarrow \frac{a(t) dr}{\sqrt{1 - Kr^2}} = dt \quad \text{only radial distance}$$

$\therefore$  horizon distance

$$d_h = a(t_0) \int_0^{t_0} \frac{dt}{a(t)}$$

For radiation dominated universe,

$$a(t) = \left( 2 \sqrt{\frac{8\pi G \rho_0 c^4}{3}} t - t^2 \right)^{1/2} \quad \text{for } k=1$$

$$t_0 \quad \text{for } k=0$$

$$t^2 \quad \text{for } k=-1$$

$$\therefore a(t) \propto t^{1/2}$$

$$\text{Let } a(t) = t^{1/2}$$

$$\text{then } \int_0^{t_0} t^{-1/2} dt = \left[ \frac{t^{1/2}}{1/2} \right]_0^{t_0} = 2t_0^{1/2}$$

$$\text{also } a(t_0) = t_0^{1/2}$$

$$\therefore d_h = a(t_0) \int_0^{t_0} \frac{dt}{a(t)} = a(t_0) \int_0^{t_0} t^{-1/2} dt = t_0^{1/2} \cdot 2t_0^{1/2} = 2t_0 \quad \text{--- Ans.}$$

For light we would assume that  $d_h = t_0$  and in this case  $d_h = 2t_0$ , the reason being that for horizon distance the measurements are simultaneous.

--- Ans.

9. The number density of photons in thermal equilibrium at temperature  $T$  and between frequency  $\nu$  and  $\nu + d\nu$  is

$$n_T(\nu) d\nu = \frac{8\pi\nu^2 d\nu}{\exp(h\nu/k_B T) - 1}.$$

Show that the equilibrium temperature  $T$  scales as inversely with the scale factor.

[5]

Given,  $n_T(\nu) d\nu = \frac{8\pi\nu^2 d\nu}{\exp(h\nu/k_B T) - 1}$

There is no change in number density ~~per~~ in a comoving volume.

$$\therefore n_T(\nu, t) a(t)^3 d\nu = n_{TL}(\nu_L, t_L) a(t_L)^3 d\nu_L$$

$\hookrightarrow$  Equilibrium temperature.

we know,  $\frac{\nu_L}{\nu} = \frac{a(t)}{a(t_L)} = \frac{d\nu_L}{d\nu}$

$$\rightarrow \frac{8\pi\nu^2 d\nu \cdot a(t)^3}{\exp\left(\frac{h\nu}{k_B T}\right) - 1} = \frac{8\pi\nu_L^2 d\nu_L \cdot a(t_L)^3}{\exp\left(\frac{h\nu_L}{k_B T_L}\right) - 1}$$

$$\text{now, } \frac{\nu^2 d\nu \cdot a(t)^3}{\nu_L^2 d\nu_L a(t_L)^3} = \frac{\nu_L^2}{\nu_L^2} \cdot \frac{d\nu}{d\nu_L} \cdot \frac{a(t)^3}{a(t_L)^3} = \frac{a(t_L)^2}{a(t)} \cdot \frac{a(t_L)}{a(t)} \cdot \frac{a(t)^3}{a(t_L)} = 1$$

$$\therefore \exp\left(\frac{h\nu}{k_B T}\right) - 1 = \exp\left(\frac{h\nu_L}{k_B T_L}\right) - 1$$

$$\Rightarrow \frac{\nu}{T} = \frac{\nu_L}{T_L} \Rightarrow \frac{T}{T_L} = \frac{\nu}{\nu_L} = \frac{a(t_L)}{a(t)}$$

$$\therefore T \propto \frac{1}{a}$$

— Ans.

— 5 —



10. Assume that the present universe is dominated by the cosmological constant  $\Lambda = 8\pi G\rho_\Lambda > 0$  only, i.e.,  $\rho_M = 0, \rho_R = 0, K = 0$ . At present the temperature of the CMB is  $T_0 = 2.7\text{K}$ . Calculate the time (from the present epoch) when the CMB temperature will be  $2\text{K}$ . You are given  $G = 6.67 \times 10^{-11} \text{m}^3 \text{kg}^{-1} \text{s}^{-2}$ ,  $\rho_\Lambda = \Omega_\Lambda \rho_{\text{crit}}$ ,  $\Omega_\Lambda = 0.685$ , and  $\rho_{\text{crit}} = 8.5 \times 10^{-27} \text{kg/m}^3$ . [5]

Given,  $\Lambda = 8\pi G\rho_\Lambda > 0$   $\rho_M = \rho_R = K = 0$   $T_0 = 2.7\text{K}$  time when CMB Temp will be  $2\text{K}$

using Friedmann Equation,  $\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3} \rho - \frac{K}{a^2}$

$$\rho = \rho_\Lambda, K = 0$$

$$\Rightarrow \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3} \rho_\Lambda = \frac{\Lambda}{3}$$

$$\Rightarrow \frac{\dot{a}}{a} = \sqrt{\frac{\Lambda}{3}} = \frac{da}{adt}$$

2nd Friedmann eq  $\Rightarrow \frac{\ddot{a}}{a} = -\frac{4\pi G}{3} (\rho_\Lambda - 3\rho_\Lambda)$  as  $\rho = \rho_\Lambda$

$$= \frac{8\pi G}{3} \rho_\Lambda = \frac{\Lambda}{3}$$

conformal time  $\rightarrow d\eta = \frac{dt}{a} \Rightarrow \eta = \int \frac{dt}{a(t)}$

$$\dot{a} = \frac{da}{dt} = \frac{da}{d\eta a} \Rightarrow \frac{\dot{a}}{a} = \frac{da}{d\eta a^2} = \sqrt{\frac{\Lambda}{3}}$$

$$\Rightarrow \frac{da}{a^2} = \sqrt{\frac{\Lambda}{3}} d\eta \quad \text{integrate both sides}$$

$$\Rightarrow -\frac{1}{a} + \frac{1}{a_i} = \sqrt{\frac{\Lambda}{3}} d\eta \quad a \propto T^{-1}$$

$$\Rightarrow -\frac{1}{a_f} + \frac{1}{a_i} = \sqrt{\frac{\Lambda}{3}} d\eta$$

$$\therefore -T_f + T_i = \sqrt{\frac{\Lambda}{3}} d\eta = 2.7 - 2.0 \text{ K} = 0.7 \text{ K}$$

$$\Rightarrow d\eta = 0.7 \times \sqrt{\frac{3}{\Lambda}}$$

$$\begin{aligned} \Lambda &= 8\pi G\rho_\Lambda \\ &= 1.68 \times 10^{-11} \times 0.685 \times 8.5 \times 10^{-27} \\ &= 796 \times 10^{-38} \end{aligned}$$

$$3/\Lambda = 3.77 \times 10^{35}$$

$$\sqrt{\frac{3}{\Lambda}} = 6.14 \times 10^{17}$$

$$\Rightarrow \text{time} = 0.7 \times \sqrt{\frac{3}{\Lambda}} = 4.3 \times 10^{17} \text{ sec}$$

$$= \frac{4.13 \times 10^{17}}{3.15 \times 10^7} \text{ years} = 1.3 \times 10^{10} \text{ years} \rightarrow 10 \text{ billion years}$$

Ans.



11. The number density of particles of type  $i$  at equilibrium temperature  $T$  is given by

$$n_i = g_i \left( \frac{m_i k_B T}{2\pi \hbar^2} \right)^{3/2} \exp \left( - \frac{m_i c^2}{k_B T} \right). \quad (1)$$

- (a) Assuming the neutrons ( $i = n$ ) and the protons ( $i = p$ ) to be in thermal equilibrium, calculate the temperature and the age of the universe (since Big Bang) when  $n_n/n_p$  was 0.50. You may use  $g_p = g_n = 2$ . Neutron and proton mass difference  $(m_n - m_p)c^2 = 1.29 \text{ MeV} = 1.29 \times 10^6 \text{ eV}$ , and the ratio  $(m_n/m_p)^{3/2} = 1.002$ . The Boltzmann constant  $k_B = 8.6 \times 10^{-5} \text{ eV/K}$ . Reduced Planck constant  $\hbar = 6.582 \times 10^{-16} \text{ eV.s}$ . You may need the numerical value  $\ln(0.50/1.002) = 0.695145$ . [4]
- (b) After the neutrinos decouple and freeze out, the number densities of protons and neutrons evolve independently of each other. Since, the neutrons decay to protons ( $n \rightarrow p + e^- + \bar{\nu}$ ), the number density of neutrons deplete according to  $n_n(t) = n_n^{\text{eq}} e^{-t/\tau_n}$ , where  $n_n^{\text{eq}}$  is the equilibrium number densities at the moment of freezeout, and  $\tau_n = 880 \text{ s}$  is the neutron lifetime. Assuming  $n_n^{\text{eq}}/n_p = 1/5$  at the time of freezeout, calculate the value of  $X_n = n_n/(n_n + n_p)$  at time  $t = 250 \text{ s}$ . Hint: You may need the numerical value  $e^{-250/880} = 0.752698$ . [6]
- (c) Assume that all the neutrons end in helium during the BBN. Calculate the ratio of number density of helium  $n_{\text{He}}$  to hydrogen  $n_H = n_p$  at  $t = 250 \text{ s}$  and show that it is close to 0.07. Hint:  $e^{-250/880} = 0.752698$ . [3]

a) when neutrons and protons are in thermal equilibrium

$$\frac{n_n}{n_p} = \frac{g_n (m_n k_B T / 2\pi \hbar^2)^{3/2} \exp(-m_n c^2 / k_B T)}{g_p (m_p k_B T / 2\pi \hbar^2)^{3/2} \exp(-m_p c^2 / k_B T)} = \left( \frac{m_n}{m_p} \right)^{3/2} \exp \left( - \frac{c^2 (m_n - m_p)}{k_B T} \right)$$

$$= 1.002 \times \exp \left( - \frac{1.29 \times 10^6 \text{ eV}}{k_B T} \right) = \frac{1}{2} = 0.5$$

$$\Rightarrow \frac{-1.29 \times 10^6 \text{ eV}}{k_B T} = \ln \left( \frac{0.5}{1.002} \right) = -0.695145$$

$$\Rightarrow k_B T = \frac{0.695145}{1.002 \times 10^6 \text{ eV}} \times 1.29 \times 10^6 \text{ eV} = 1.86 \times 10^6 \text{ eV}$$

$$\Rightarrow T = \frac{1.86 \times 10^6 \text{ eV}}{k_B = 8.6 \times 10^{-5} \text{ eV/K}} = 0.216 \times 10^{11} \text{ K}$$

$$T = 2.1 \times 10^{10} \text{ K} \quad \text{Ans.}$$

$$\text{time } t = 0.994 \text{ sec} \left( \frac{10^{10} \text{ K}}{T} \right)^2$$

$$= 0.994 \left( \frac{1}{2.1} \right)^2 = 0.225 \text{ sec} \quad \text{Ans.}$$

b)  $n \rightarrow p + e^- + \bar{\nu}$  using Saha equation,  $\frac{n_n}{n_p} = \frac{n_p}{n_n} \Rightarrow \frac{n_n^{\text{eq}} e^{-t/\tau_n}}{n_n^{\text{eq}}} = 5$  at freezeout.

$$\text{at freezeout } \frac{n_n^{\text{eq}}}{n_p} = \frac{1}{5}$$

$\Rightarrow$  num density of protons evolve as  $\Rightarrow$

$$n_p(t) = n_p + (n_n^{\text{eq}} - n_n(t))$$

$$\therefore n_p(t) = 5 n_n^{\text{eq}} + n_n^{\text{eq}} - n_n(t) = 6 n_n^{\text{eq}} - n_n(t)$$

$$= n_n^{\text{eq}} (6 - e^{-t/\tau_n})$$

$$\text{at } 250 \text{ s, } x_n = \frac{n_n^{\text{eq}} e^{-250/880}}{n_n^{\text{eq}} e^{-250/880} + n_p^{\text{eq}} (6 - e^{-250/880})} = \frac{0.752698}{6} = 0.125 \quad \text{Ans.}$$