

# Path Integral

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## 1 Quantum Mechanics case

Consider 1d propagator in non-relativistic quantum mechanics  $e^{-iHt}$  (we have set  $\hbar = 1$ ), we want to compute the matrix element

$$U(q', t'; q, t) = \langle q' | e^{-iH(t'-t)} | q \rangle \quad (1)$$

Suppose  $H$  has the form  $\frac{p^2}{2m} + V(q)$ , we firstly slice the time interval  $[t, t']$  into  $N$  slices (and let  $\Delta t = \frac{t'-t}{N}$ ), i.e.

$$U(q', t'; q, t) = \langle q' | \prod_{n=1}^N e^{-i(\frac{p_n^2}{2m} + V(q)) \frac{t'-t}{N}} | q \rangle \quad (2)$$

The insert the complete relation  $\int dq |q\rangle \langle q| = \int dp |p\rangle \langle p| = 1$

$$U(q', t'; q, t) = \langle q' | \prod_{n=1}^N \left( \int dp_n dq_n e^{-i\frac{p_n^2}{2m}\Delta t} |p_n\rangle \langle p_n|q_n\rangle e^{-iV(q_n)\Delta t} \langle q_n| + O(\Delta t^2) \right) | q \rangle \quad (3)$$

When  $N \rightarrow \infty$ , the  $O(\Delta t^2)$  term can be ignored (i.e. we are essentially consider the Lie algebra, which is commutative w.r.t. addition). Plugging in  $\langle p_n | q_n \rangle = (2\pi)^{-1/2} e^{ip_n q_n}$ , and notice that  $\langle q_n |$  pairs with  $|p_{n-1}\rangle$ ,  $\langle q_1 |$  pairs with  $|q\rangle$ , So

$$U(q', t'; q, t) = (2\pi)^{-N} \int dp_N dq_N e^{ip_N(q_N - q')} \prod_{n=1}^{N-1} \left[ \int dp_n dq_n e^{i(p_n q_n - p_n q_{n+1} - V(q_n)\Delta t - \frac{p_n^2}{2m}\Delta t)} \right] \delta(q_1 - q) \quad (4)$$

Integrate out  $p_n$  for each  $n \neq N$  gives

$$\int dp_n e^{i(p_n(q_n - q_{n+1}) - \frac{p_n^2}{2m}\Delta t)} = \int dp_n e^{-\frac{i\Delta t}{2m}(p_n - \frac{m}{\Delta t}(q_n - q_{n+1}))^2} e^{\frac{im}{2\Delta t}(q_n - q_{n+1})^2} = \left(\frac{2m\pi}{i\Delta t}\right)^{1/2} e^{\frac{im}{2\Delta t}(q_n - q_{n+1})^2} \quad (5)$$

Plugging in equation (4)

$$U(q', t'; q, t) = \left(\frac{m}{2\pi i\Delta t}\right)^{-(N-1)/2} \prod_{n=1}^N \left[ \int dq_n e^{i(\frac{m(q_n - q_{n+1})^2}{2\Delta t} - V(q_n)\Delta t)} \right] \delta(q' - q_N) \delta(q - q_1) \quad (6)$$

The delta functions are essentially fixing end points. Notice that  $\frac{m(q_n - q_{n+1})^2}{2\Delta t} \sim \frac{m\dot{q}_n^2}{2}\Delta t$ , so the exponential is essentially

$$\frac{m(q_n - q_{n+1})^2}{2\Delta t} - V(q_n)\Delta t = L(q_n, \dot{q}_n)\Delta t \quad (7)$$

Define the path integral measure as

$$\int D[q] = \lim_{N \rightarrow \infty} \prod_{n=1}^N \left( \frac{2\pi i(t' - t)}{Nm} \right)^{1/2} \int dq_n \quad (8)$$

Then the propagator is

$$U(q', t'; q, t) = \int_{q(t)=q}^{q(t')=q'} D[q] e^{i \int_t^{t'} L[q] dt} = \int_{q(t)=q}^{q(t')=q'} D[q] e^{i S[q]} \quad (9)$$