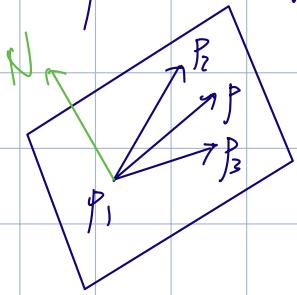


Lecture 3 plane

3 points to form a plane



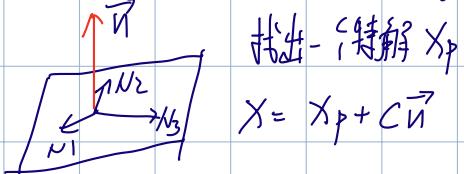
$$(1) \det(\vec{P_1P_2}, \vec{P_1P_3}, \vec{P_1P}) = 0 \quad \det(\vec{a}, \vec{b}, \vec{c}) = \begin{vmatrix} a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \\ c_1 & c_2 & c_3 \end{vmatrix}$$

$$(2) \vec{N} = \vec{P_1P_3} \times \vec{P_1P_2} \quad \text{解 } \vec{P_1P} \cdot \vec{N} = 0$$

\approx

$$Ax + By + Cz = D \quad (A, B, C) \text{ 为 normal vector } \vec{N}$$

若 $Ax + By + Cz = D$ 且 $\det(A) = 0$, 则 $\vec{N}_1, \vec{N}_2, \vec{N}_3$ are coplanar



Lecture 5 Equation of line

1. intersection of 2 planes

2. trajectory of a moving point parametric equation
ex. line through $Q_0 = (-1, 2, 2)$, $Q_1 = (1, 3, -1)$

$$\begin{cases} \overrightarrow{Q_0 Q(t)} = t \overrightarrow{Q_0 Q_1} = t \langle 2, 1, -3 \rangle \\ x(t) = -1 + 2t \\ y(t) = 2 + t \\ z(t) = 2 - 3t \end{cases} \quad Q(t) = Q_0 + t \overrightarrow{Q_0 Q_1}$$

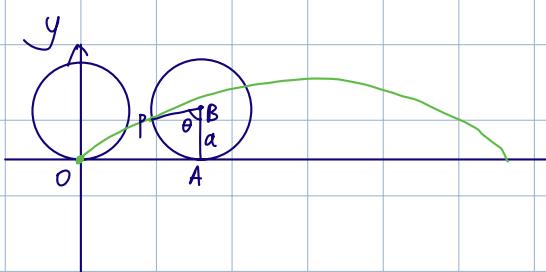


Application:

Intersection with plane? ex. $x + 2y + 4z = 7$ line through $(-1, 2, 2), (1, 3, -1)$

the plane divide the space into 2 part: $x + 2y + 4z > 7$ and $x + 2y + 4z < 7$ opposite side

$$x(t) + 2y(t) + 4z(t) = -t + 1 + 11 = 7 \quad t = \frac{1}{2} \quad \text{intersection point: } \left(0, \frac{5}{2}, \frac{1}{2}\right)$$



$$\begin{aligned} \overrightarrow{OP} &= \overrightarrow{OA} + \overrightarrow{AB} + \overrightarrow{BP} \\ \overrightarrow{OA} &= \langle a\theta, 0 \rangle \quad \overrightarrow{AB} = \langle 0, a \rangle \quad \overrightarrow{BP} = \langle -a\sin\theta, -a\cos\theta \rangle \\ &= x(\theta) \quad = y(\theta) \\ \overrightarrow{OP} &= \langle a\theta - a\sin\theta, a - a\cos\theta \rangle \end{aligned}$$

velocity ex $\vec{r}(t) = \langle t - \sin t, 1 - \cos t \rangle$

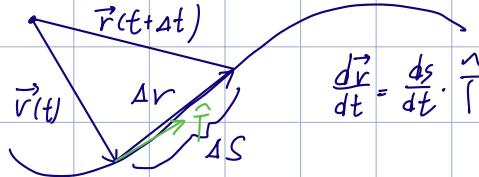
$$\vec{v} = \frac{d\vec{r}}{dt} = \left\langle \frac{dx}{dt}, \frac{dy}{dt}, \frac{dz}{dt} \right\rangle = \langle 1 - \cos t, \sin t \rangle$$

unit tangent vector $\hat{T} = \frac{\vec{v}}{|\vec{v}|}$

$$\vec{v} = \frac{d\vec{r}}{dt} = \frac{d\vec{r}}{ds} \cdot \frac{ds}{dt} = |\vec{v}| \quad , \quad \therefore \frac{d\vec{r}}{ds} = \frac{\vec{v}}{|\vec{v}|}$$

Velocity has direction, tangent to traj, \hat{T}
length: speed, $\frac{ds}{dt}$

$$\text{speed} = |\vec{v}|$$



Lecture 6 Partial derivative

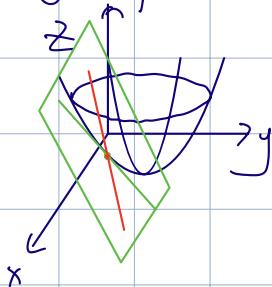
$$f = f(x, y) \quad \frac{\partial f}{\partial x} \Big|_{x_0, y_0} = \lim_{\Delta x \rightarrow 0} \frac{f(x_0 + \Delta x, y_0) - f(x_0, y_0)}{\Delta x}$$

$$\frac{\partial f}{\partial x} = \lim_{\Delta x \rightarrow 0} \frac{f(x + \Delta x, y) - f(x, y)}{\Delta x} = f_x \quad \frac{\partial f}{\partial y} = \lim_{\Delta y \rightarrow 0} \frac{f(x, y + \Delta y) - f(x, y)}{\Delta y} = f_y$$

Approximation:

If we change $x \rightarrow x + \Delta x$, $y \rightarrow y + \Delta y$, $z = f(x, y)$ then $\Delta z \approx f_x \Delta x + f_y \Delta y$

Tangent plane



$$L_1 = \begin{cases} z = z_0 + a(x - x_0) & \frac{\partial f}{\partial x}(x_0, y_0) = a \\ y = y_0 \end{cases}$$

$$L_2 = \begin{cases} z = z_0 + b(y - y_0) & \frac{\partial f}{\partial y}(x_0, y_0) = b \\ x = x_0 \end{cases}$$

$$\text{plane: } z = z_0 + a(x - x_0) + b(y - y_0)$$

Application. Optimization problems

At a [local] min or max: $\frac{\partial f}{\partial x}(x_0, y_0) = \frac{\partial f}{\partial y}(x_0, y_0) = 0$ tangent plane is horizontal

Def: (x_0, y_0) is a critical point if $f_x(x_0, y_0) = 0$ and $f_y(x_0, y_0) = 0$

$$\text{ex: } f(x, y) = x^2 - 2xy + 3y^2 + 2x - 2y$$

$$\begin{cases} f_x = 2x - 2y + 2 = 0 \\ f_y = -2x + 6y - 2 = 0 \end{cases} \Rightarrow \begin{cases} x = -1 \\ y = 0 \end{cases} \text{ critical point } (x, y) = (-1, 0)$$

Lecture 8 second derivative

critical point $\begin{cases} \text{local max} \\ \text{local min} \\ \text{saddle} \end{cases}$ global max/min: either critical point or boundary

ex. $w = ax^2 + bxy + cy^2$

$$w = \frac{1}{4a} [4a^2 (x + \frac{b}{2a}y)^2 + (4ac - b^2)y^2]$$

3 cases: (1) $4ac - b^2 < 0 \Rightarrow$ saddle point

(2) $4ac - b^2 = 0 \quad w = x^2$

(3) $4ac - b^2 > 0 \quad \begin{array}{l} \text{if } a > 0, \text{ minimum} \\ \text{if } a < 0, \text{ maximum} \end{array}$

In general: look at second derivative

$$\frac{\partial^2 f}{\partial x^2} = f_{xx} \quad f_{xy} = \frac{\partial^2 f}{\partial x \partial y} = \frac{\partial^2 f}{\partial y \partial x} = f_{yx} \quad \frac{\partial^2 f}{\partial y^2} = f_{yy}$$

second derivative test

At a critical point (x_0, y_0) of f

$$\text{let } A = f_{xx}(x_0, y_0), B = f_{xy}(x_0, y_0), C = f_{yy}(x_0, y_0)$$

If $AC - B^2 > 0$ and $A > 0$, $f(x_0, y_0)$ is a local minimum

If $AC - B^2 > 0$ and $A < 0$, local maximum

If $AC - B^2 < 0$ saddle

If $AC - B^2 = 0$ can't conclude

proof: quadratic approximation

$$\begin{aligned} \Delta f &= f_x(x - x_0) + f_y(y - y_0) + \frac{1}{2} f_{xx}(x - x_0)^2 + f_{xy}(x - x_0)(y - y_0) + \frac{1}{2} f_{yy}(y - y_0)^2 \\ &= \frac{1}{2} f_{xx}(x_0, y_0)(x - x_0)^2 + f_{xy}(x_0, y_0)(x - x_0)(y - y_0) + f_{yy}(x_0, y_0)(y - y_0)^2 \end{aligned}$$

In degenerate case, depends on higher derivative

Lecture 9 Least Square

$$y = \hat{a}x + \hat{b}$$

deviation for each point: $y_i - (\hat{a}x_i + \hat{b})$

$$\text{Minimise } P(a, b) = \sum_{i=1}^n [y_i - (ax_i + b)]^2 = \sum_{i=1}^n [y_i^2 + (ax_i + b)^2 - 2y_i(ax_i + b)]$$

$$\frac{\partial P}{\partial a} = \sum_{i=1}^n [2(ax_i + b) \cdot x_i - 2x_i y_i] = 0 \Leftrightarrow a \sum x_i^2 + b \sum x_i = \sum x_i y_i$$

$$\frac{\partial P}{\partial b} = \sum_{i=1}^n [2(ax_i + b) - 2y_i] = 0 \Leftrightarrow a \sum x_i + b n = \sum y_i$$

solve for (a, b)

$$\left\{ \begin{array}{l} \hat{a} = \frac{n \sum x_i y_i - \sum x_i \sum y_i}{n \sum x_i^2 - (\sum x_i)^2} \\ \hat{b} = \frac{\sum y_i - a \sum x_i}{n} \end{array} \right.$$

Lecture 11 Differential

Total differential

$$f(x,y,z) \quad df = f_x dx + f_y dy + f_z dz = \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} dy + \frac{\partial f}{\partial z} dz$$

Important: df is NOT Δf

can do ① encode how change in x, y, z affect f

② placeholder for small variation $\Delta x, \Delta y, \Delta z$

to get approximation formula $\Delta f \approx f_x \Delta x + f_y \Delta y + f_z \Delta z$

③ divide by something like Δt to get a rate of change when $x(t), y(t), z(t)$

$$\frac{df}{dt} = f_x \frac{dx}{dt} + f_y \frac{dy}{dt} + f_z \frac{dz}{dt}$$

$$\text{ex } W = x^2 y + z, \quad x = t, \quad y = e^t, \quad z = \sin t$$

$$\frac{dw}{dt} = 2xy \cdot \frac{dx}{dt} + x^2 \cdot \frac{dy}{dt} + 1 \cdot \frac{dz}{dt}$$

$$= 2te^t + t^2 e^t + \cos t = (t^2 + 2t)e^t + \cos t$$

Prove chain rule with more variables

$$W = f(x, y) \quad x = x(u, v) \quad y = y(u, v)$$

$$W = f(x(u, v), y(u, v)), \quad \text{what about } \frac{\partial w}{\partial x}, \frac{\partial w}{\partial y}$$

$$dw = f_x dx + f_y dy = f_x(x_u du + x_v dv) + f_y(y_u du + y_v dv)$$

$$= (\underbrace{f_x x_u + f_y y_u}_{\text{has to be } \frac{\partial w}{\partial u}}) du + (\underbrace{f_x x_v + f_y y_v}_{\text{has to be } \frac{\partial w}{\partial v}}) dv$$

$$\frac{\partial f}{\partial u} = \frac{\partial f}{\partial x} \cdot \frac{\partial x}{\partial u} + \frac{\partial f}{\partial y} \cdot \frac{\partial y}{\partial u}$$

$$\frac{\partial f}{\partial v} = \frac{\partial f}{\partial x} \cdot \frac{\partial x}{\partial v} + \frac{\partial f}{\partial y} \cdot \frac{\partial y}{\partial v}$$

Lecture 12 Gradient

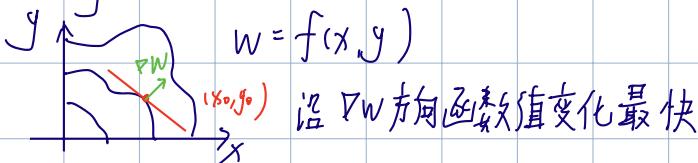
$$w = f(x, y, z)$$

$\nabla w = \langle w_x, w_y, w_z \rangle$ gradient of w at some point (x, y, z)

$$\frac{d\vec{r}}{dt} = \left\langle \frac{dx}{dt}, \frac{dy}{dt}, \frac{dz}{dt} \right\rangle$$

$$\frac{dw}{dt} = \nabla w \cdot \frac{d\vec{r}}{dt}$$

Theory: $\nabla w \perp$ level surface ($w = \text{constant}$)



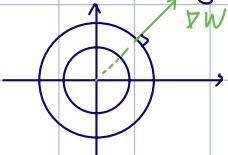
$$\text{ex1. } w = a_1 x + a_2 y + a_3 z$$

$$\nabla w = (a_1, a_2, a_3)$$

level surface: $a_1 x + a_2 y + a_3 z = C$ plane with normal vector ∇w

ex2. $w = x^2 + y^2$ level curve $w = C$ is a circle $x^2 + y^2 = C$

$$\nabla w = (2x, 2y)$$



∇w always points at higher value of a function

Proof: Take a curve $\vec{r} = \vec{r}(t)$ that stays on the level $w = C$

$$w(x, y, z) = C \quad \text{Take a motion on the surface } w(x, y, z) = C$$

$v = \frac{d\vec{r}}{dt}$ is tangent to the plane $w(x, y, z) = C$

By chain rule: $\frac{dw}{dt} = \nabla w \cdot \frac{d\vec{r}}{dt} = \nabla w \cdot v = 0$ because $w(t) = C = \text{constant}$

so $\nabla w \perp v$, This is true for any motion on $w = C$

v can be any vector tangent to the surface

so ∇w is perpendicular to the tangent plane to the level

Applications find the tangent plane to surface $x^2 + y^2 - z^2 = 4$ at $(2, 1, 1)$?

Level set $w = 4$ when $w = x^2 + y^2 - z^2$

gradient $\nabla w = (2x, 2y, -2z) = (4, 2, -2)$, which is the normal vector to the plane $4(x-2) + 2(y-1) - 2(z-1) = 0 \Rightarrow 4x + 2y - 2z = 8$

Another way: $dw = 2x dx + 2y dy - 2z dz = 4 dx + 2 dy - 2 dz$ at $(2, 1, 1)$

$$dw \approx 4dx + 2dy - 2dz$$

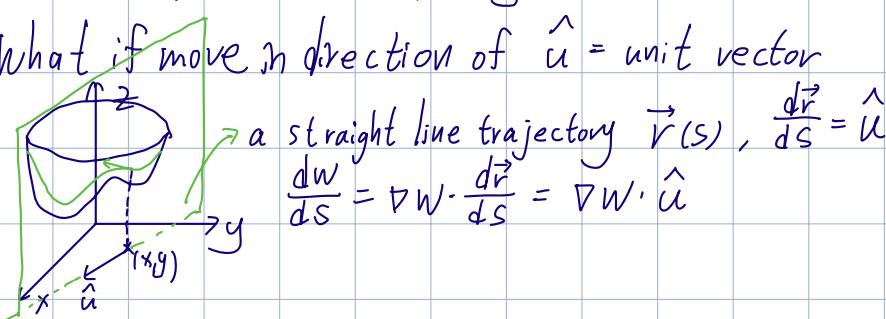
when $dw = 0$; its tangent plane: $4dx + 2dy - 2dz = 0$

$$4(x-2) + 2(y-1) - 2(z-1) = 0$$

Directional derivatives

$w = w(x, y) \rightsquigarrow$ know $\frac{\partial w}{\partial x}, \frac{\partial w}{\partial y}$

What if move in direction of \hat{u} = unit vector



Def: If $\hat{u} = \langle a, b \rangle$ (unit vector)

$$\begin{cases} x(s) = x_0 + a s \\ y(s) = y_0 + b s \end{cases}$$

$$\frac{dw}{ds}|_{\hat{u}} = \nabla w \cdot \hat{u} \quad \text{components of } \nabla w \text{ in direction of } \hat{u}$$

Directional derivative

Geometric view: slice the graph by the plane of the unit vector, slope = $\frac{dw}{ds}|_{\hat{u}}$

$$\frac{dw}{ds} \Big|_{\hat{u}} = \nabla w \cdot \hat{u} = |\nabla w| \cdot |\hat{u}| \cdot \cos \theta$$

$= |\nabla w| \cdot \cos \theta$

A diagram shows a horizontal vector labeled ∇w pointing to the right. A second vector, labeled \hat{u} , originates from the same point and makes an angle θ with the ∇w vector. The length of \hat{u} is indicated by a vertical tick mark.

\Rightarrow be maximal when $\cos \theta = 1 \Rightarrow \hat{u} = \operatorname{dir}(\nabla w)$

Conclusion: $\nabla w = \operatorname{dir}$ of fastest increase of w

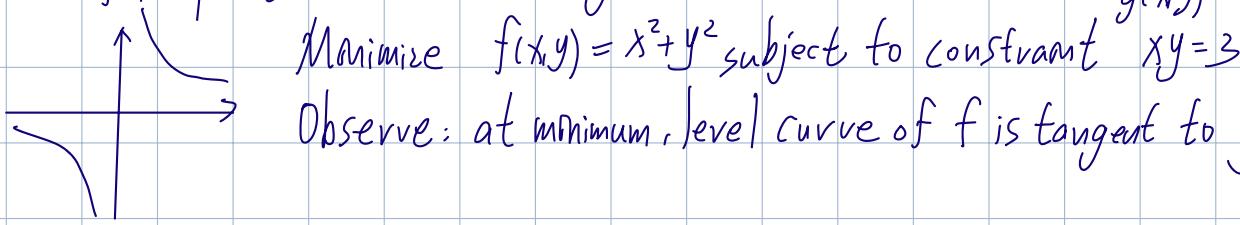
$$|\nabla w| = \frac{dw}{ds} \Big|_{\hat{u}} = \operatorname{dir}(\nabla w)$$

minima when $\cos \theta = -1 \quad \hat{u} = \operatorname{dir}(\nabla w)$

when $\frac{dw}{ds} \Big|_{\hat{u}} = 0 \quad \cos \theta = 0 \quad \theta = 90^\circ \quad \hat{u} \perp \nabla w \Leftrightarrow \hat{u} \text{ tangent to level}$

Lecture 13 Lagrange multipliers

min/max a function $f(x,y,z)$ where x,y,z are not independent, $g(x,y,z) = c$
ex: find point closest to origin on $xy=3$



When $f(x,y)$'s level curve is tangent to $g(x,y)$'s level curve, normal vectors parallel

$\nabla f \parallel \nabla g$ so $\nabla f = \lambda \nabla g$

$\Rightarrow \begin{cases} f_x = \lambda g_x \\ f_y = \lambda g_y \\ xy = 3 \end{cases}$

from $\nabla f = \lambda \nabla g$
from the constraint

Why is this method valid?

At constraint min/max, many directions along $g=c$, the rate of change of f must be 0

For any \hat{u} tangent to $g=c$, we must have $\frac{df}{ds}|_{\hat{u}} = 0$
 $\nabla f \cdot \hat{u} = 0$, so any such \hat{u} is $\perp \nabla f$, so $\nabla f \perp$ level set of g
know $\nabla g \perp$ level set of g , so $\nabla f \parallel \nabla g$

WARNING: method don't tell whether it's maximum or minimum (or a saddle)
CAN'T use second derivative test

Lecture 14 Non-independent variables

$f(x, y, z)$ where $g(x, y, z) = C$

then $z = z(x, y)$ $\frac{\partial z}{\partial x}$? $\frac{\partial z}{\partial y}$?

ex $x^2 + yz + z^3 = 8$ at $(2, 3, 1)$

Take differential: $2x dx + z dy + (y + 3z^2) dz = 0$

at $(2, 3, 1)$ $4dx + dy + 6dz = 0$

If we view $z = z(x, y)$ $dz = -\frac{1}{6}(4dx + dy)$

$$\frac{\partial z}{\partial x} = -\frac{z}{3}/1 = -\frac{z}{3} \quad \frac{\partial z}{\partial y} = -\frac{1}{6}/1 = -\frac{1}{6}$$

(y constant $dy = 0$) (x constant $dx = 0$)

In general:

$g(x, y, z) = C$ then $dg = g_x dx + g_y dy + g_z dz = 0$

Solve for dz : $dz = -\frac{g_x}{g_z} dx - \frac{g_y}{g_z} dy$

so for $\frac{\partial z}{\partial x}$, set $dy = 0$ $dz = -\frac{g_x}{g_z} dx \quad \frac{\partial z}{\partial x} = \frac{-g_x}{g_z}$

ex. $f(x, y) = x + y \quad \frac{\partial f}{\partial x} = 1$

$x = u$ $y = u + v$, then $f = 2u + v \quad \frac{\partial f}{\partial u} = 2$

$x = u$, but $\frac{\partial f}{\partial x} \neq \frac{\partial f}{\partial u}$, because $\frac{\partial f}{\partial x}$ means keeping y constant, $\frac{\partial f}{\partial u}$ means keeping v constant

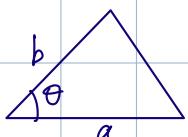
Need clearer Notations

$(\frac{\partial f}{\partial x})_y$ = keep y constant

$(\frac{\partial f}{\partial u})_v$ = keep v constant

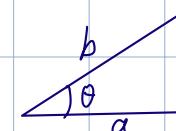
$$\text{so } \underbrace{(\frac{\partial f}{\partial x})_y}_1 \neq \underbrace{(\frac{\partial f}{\partial x})_v}_2 = \underbrace{(\frac{\partial f}{\partial u})_v}_2$$

ex. area of a triangle



$$A = \frac{1}{2}ab \sin \theta$$

Assume it's a right triangle



$$a = b \cos \theta$$

constraint

Rate of A respect to θ

(1) treat a, b, θ as independent $\frac{\partial A}{\partial \theta} = (\frac{\partial A}{\partial \theta})_{a,b} = \frac{1}{2}ab \cos \theta$ loss the constraint

(2) keep a constant $b = b(a, \theta) (= \frac{a}{\cos \theta})$ to keep right angle $(\frac{\partial A}{\partial \theta})_a$

(3) keep b constant, $a = a(b, \theta) \quad (\frac{\partial A}{\partial \theta})_b$

Compute $(\frac{\partial A}{\partial \theta})_a$

Method 0: solve for b then substitute

$$= \frac{\partial}{\partial \theta} \left(\frac{1}{2} a \cdot \frac{a}{\cos \theta} \cdot \sin \theta \right) = \frac{1}{2} a^2 \sec^2 \theta$$

2 systematic methods

Method 1: differentials

keep a fixed: $da=0$

constraint $a = b \cos \theta$ $da = \cos \theta db - b \sin \theta d\theta = 0$

$db = b \tan \theta d\theta$ (To find when θ change, how b change)

function $A = \frac{1}{2} ab \sin \theta$

$$\begin{aligned} dA &= \frac{1}{2} b \sin \theta da + \frac{1}{2} a \sin \theta db + \frac{1}{2} ab \cos \theta d\theta \\ &= \frac{1}{2} a \sin \theta \cdot b \tan \theta d\theta + \frac{1}{2} ab \cos \theta d\theta = (\frac{1}{2} ab \sin \theta \tan \theta + \frac{1}{2} ab \cos \theta) d\theta = \frac{1}{2} ab \sec \theta d\theta \end{aligned}$$

$$\text{so } (\frac{\partial A}{\partial \theta})_a = \frac{1}{2} ab \sec \theta$$

Summary: (1) write dA in terms of $da, db, d\theta$

(2) set $da = 0$

⊗ (3) differentiate constraint \Rightarrow solve for db in terms of $d\theta$

) plug in

Method 2: Chain Rule $(\frac{\partial}{\partial \theta})_a$ in formula for A

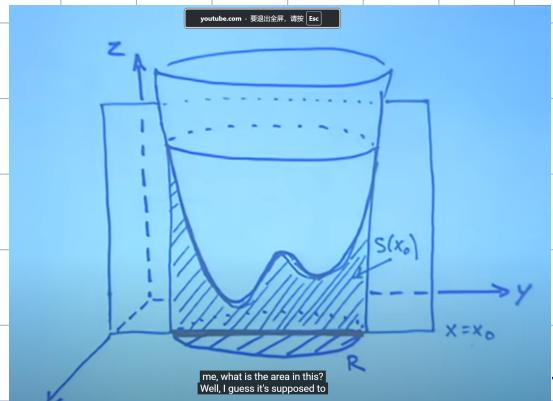
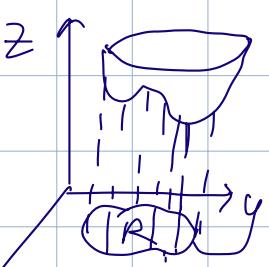
$$(\frac{\partial A}{\partial \theta})_a = A_\theta (\underbrace{\frac{\partial \theta}{\partial \theta}}_1)_a + A_a (\underbrace{\frac{\partial a}{\partial \theta}}_0)_a + A_b (\underbrace{\frac{\partial b}{\partial \theta}}_{\text{use constraint}})_a$$

Lecture 15 Double integral

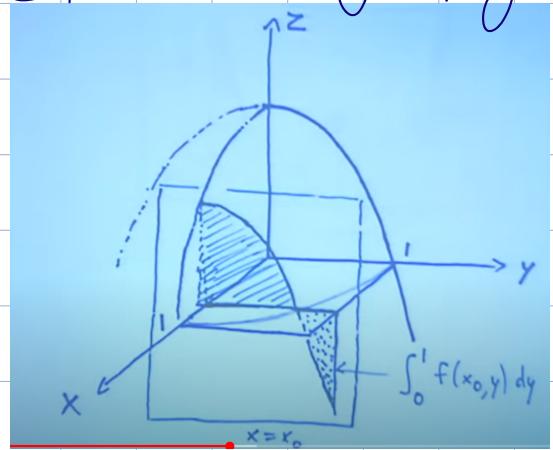
one-variable integral: area of graph

Double integral: volume below the graph

over a region R in xy plane $\iint_R f(x,y) dA$



ex1. $z = 1 - x^2 - y^2$ region $0 \leq x \leq 1$, $0 \leq y \leq 1$



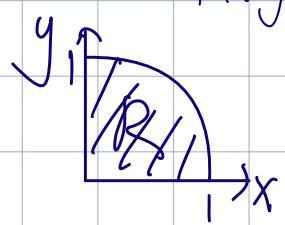
move a plane from one axis, take slices
add all the $A \cdot dx$

$$V = \int_{x_{\min}}^{x_{\max}} S(x) dx$$

$$\text{For given } x, S(x) = \int_{y_{\min}(x)}^{y_{\max}(x)} f(x,y) dy$$

$$\iint_R f(x,y) dA = \int_{x_{\min}}^{x_{\max}} \left[\int_{y_{\min}(x)}^{y_{\max}(x)} f(x,y) dy \right] dx$$

ex2: integral in such R



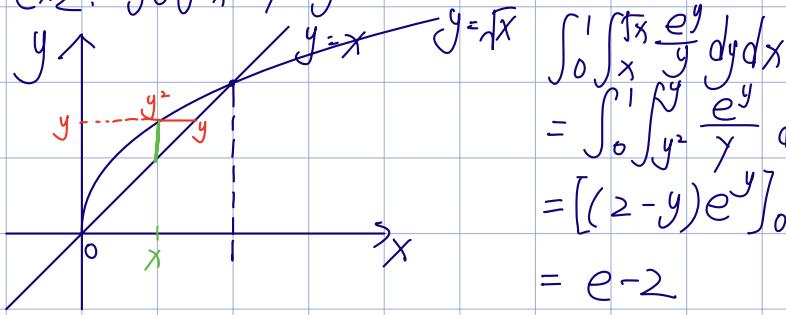
$$\begin{aligned} \iint_R (1-x^2-y^2) dy dx &= \int_0^1 [(1-x^2)\sqrt{1-x^2} - \frac{1}{3}(1-x^2)\sqrt{1-x^2}] dx \\ &= \int_0^1 \frac{2}{3}(1-x^2)^{\frac{3}{2}} dx = \frac{2}{3} \int_0^{\pi/2} \cos^4 \theta d\theta \\ &= \frac{2}{3} \cdot \frac{3}{16} \pi = \frac{1}{8} \pi \end{aligned}$$

Easier in Polar coordinates

Exchanging order

$$\text{ex1. } \int_0^1 \int_0^2 dx dy = \int_0^2 \int_0^1 dy dx$$

$$\text{ex2. } \int_0^1 \int_x^2 \frac{e^y}{y} dy dx$$

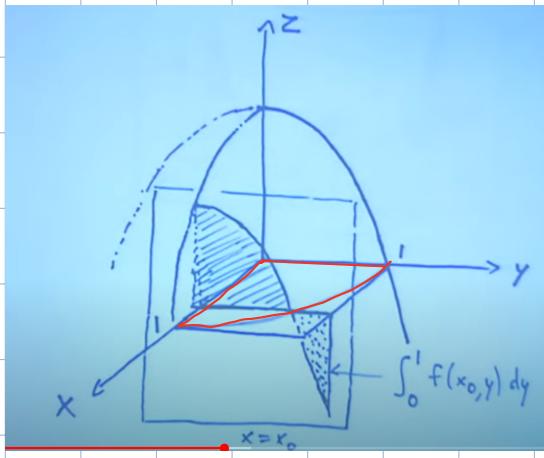


$$\begin{aligned} & \int_0^1 \int_x^2 \frac{e^y}{y} dy dx \\ &= \int_0^1 \int_{y^2}^y \frac{e^y}{y} dy dx = \int_0^1 (1-y) e^y dy \\ &= [(2-y)e^y]_0^1 \\ &= e-2 \end{aligned}$$

$$\bar{x} = \frac{1}{\text{Mass}} \iint x f(\theta, r) dy dx$$

$$\bar{y} = \frac{1}{\text{Mass}} \iint y f(\theta, r) dy dx$$

Lecture 17 Polar Coordinates



$x = r \cos \theta, y = r \sin \theta$

$\Delta A \approx \Delta r \cdot r \Delta \theta$

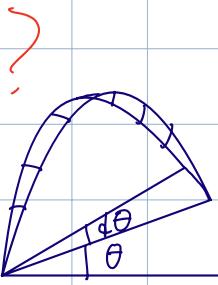
$dA = r dr d\theta$

$$\frac{1}{2} \Delta \theta ((r + \Delta r)^2 - r^2)$$

$$= \Delta \theta r dr$$

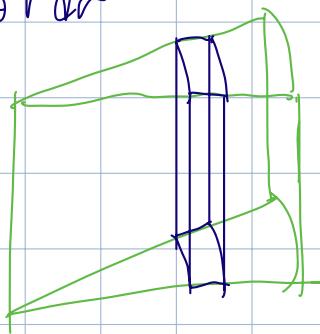
$$\begin{aligned} & \int_0^{\pi/2} \int_0^1 f(\theta, r) r dr d\theta \\ &= \int_0^{\pi/2} \int_0^1 (1-r^2) r dr d\theta \\ &= \int_0^{\pi/2} \left(\frac{1}{2}r - \frac{1}{4}r^4 \right) \Big|_0^1 d\theta \\ &= \int_0^{\pi/2} \frac{1}{4} d\theta = \frac{\pi}{8} \end{aligned}$$

???



$$\iint f(\theta, r) dr \cdot r d\theta \cdot \frac{1}{2}$$

面積 A
体積 dV



$$\iint f(\theta, r) dr \cdot \frac{1}{2} r d\theta$$

S
dV

Application

(1) convert single integral to double integral



$$\text{Area} = \iint_R 1 \cdot dA$$

(2) Average value of $f = \bar{f} = \frac{1}{\text{Area}(R)} \iint_R f dA$

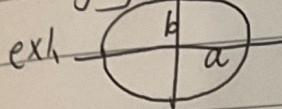
weighted average of f with density δ : $\bar{f} = \frac{1}{\text{Mass}(R)} \iint_R f \delta dA$

(3) Center of mass of a (planar) object (with density δ)



$$(\bar{x}, \bar{y}) \quad \bar{x} = \frac{1}{\text{Mass}} \iint_M x \delta dA, \bar{y} = \frac{1}{\text{Mass}} \iint_M y \delta dA$$

changing variables in $\iint \frac{x}{a} + \frac{y}{b} \leq 1$



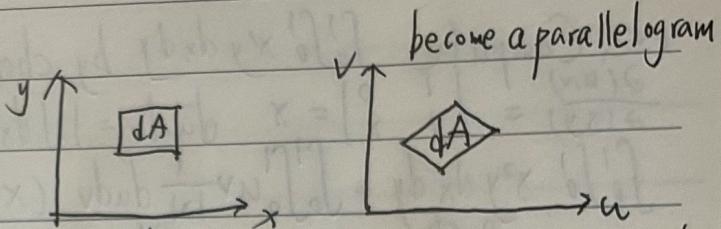
$$\iint_{\frac{x}{a} + \frac{y}{b} \leq 1} dx dy = ab \iint_{u+v \leq 1} du dv = \pi ab$$

set $\frac{x}{a} = u$ $\frac{y}{b} = v$

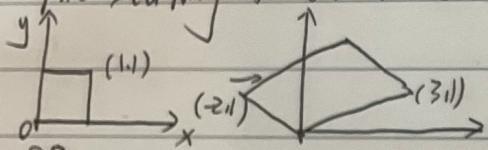
In general: finding scaling factor between $dx dy \leftrightarrow du dv$

$$\begin{cases} u = 3x - 2y \\ v = x + y \end{cases}$$

$$dA = dx dy \text{ and } dA' = du dv$$



The scaling factor here doesn't depend on choice of rectangle



$$A' = \begin{vmatrix} 3 & 1 \\ -2 & 1 \end{vmatrix} = 5$$

$$\text{so } dA' = 5dA \quad du dv = 5 dx dy$$

$$\iint f(x,y) dx dy = \iint g(u,v) \frac{1}{5} du dv, \text{ bound also needs change}$$

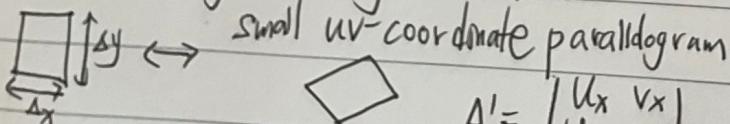
↑ Only for linear transformation

General case: $u = u(x, y)$ $v = v(x, y)$

$$\Delta u \approx du \approx u_x \cdot \Delta x + u_y \cdot \Delta y$$

$$\Delta v \approx dv \approx v_x \cdot \Delta x + v_y \cdot \Delta y$$

Small rectangle in x-y coordinates



$$A' = \begin{vmatrix} u_x & v_x \\ u_y & v_y \end{vmatrix}$$

can be seen as linear transformation, which changing scalar depends on $u_x, u_y, v_x, v_y (x_0, y_0)$

$$dA' = \begin{vmatrix} u_x & v_x \\ u_y & v_y \end{vmatrix} dA \quad \text{or} \quad dA' = \begin{vmatrix} u_x & u_y \\ v_x & v_y \end{vmatrix} dA$$

$$\text{Jacobian} : J = \frac{\partial(u, v)}{\partial(x, y)} = \begin{vmatrix} u_x & u_y \\ v_x & v_y \end{vmatrix}$$

$$du dv = |J| dx dy = \left| \begin{vmatrix} u_x & u_y \\ v_x & v_y \end{vmatrix} \right| dx dy$$

ex1 Polar coordinate

$$x = r \cos \theta \quad y = r \sin \theta$$

$$\frac{\partial(x, y)}{\partial(r, \theta)} = \begin{vmatrix} x_r & x_\theta \\ y_r & y_\theta \end{vmatrix} = \begin{vmatrix} \cos \theta & -r \sin \theta \\ \sin \theta & r \cos \theta \end{vmatrix} = r^2$$

$$dx dy = |r| dr d\theta = r dr d\theta$$

ex2. Compute $\int_0^1 \int_0^1 x^2 y dx dy$ by changing to $u=x$ $v=xy$

$$\frac{\partial(u, v)}{\partial(x, y)} = \begin{vmatrix} 1 & 0 \\ y & x \end{vmatrix} = x \quad du dv = |x| dx dy = x dx dy$$

$$\int_0^1 \int_0^1 x^2 y dx dy = \int_0^1 \int_0^1 uv \frac{1}{x} du dv \quad (\times)$$

$$= \int_0^1 \int_0^1 v du dv = \int_0^1 v dv = \frac{1}{2} \quad (\times)$$

For bounds

$\iint_{\text{region}} v du dv$ or $dv du$?



$\iint v du dv$

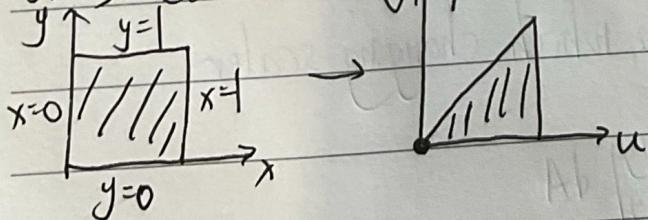
u change, v constant
 $xy = C$

when v is constant, find scope of u

$$y=1 \quad v=x, \quad u=v \quad \int_v^1 v du$$

$$= \int_0^1 \int_v^1 v du dv$$

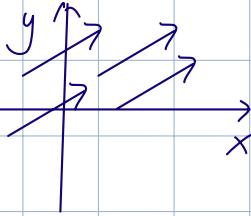
or switch to u-v picture



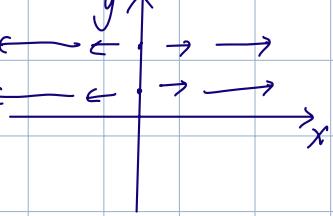
Lecture 19 Vector fields

$\vec{F} = M\vec{i} + N\vec{j}$, M, N are functions of x, y

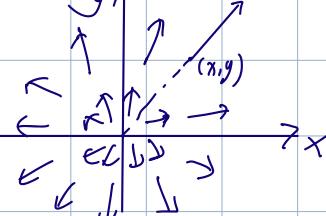
ex: $\vec{F} = 2\vec{i} + \vec{j}$



$\vec{F} = x\vec{i}$



$\vec{F} = xy\vec{i} + y\vec{j}$

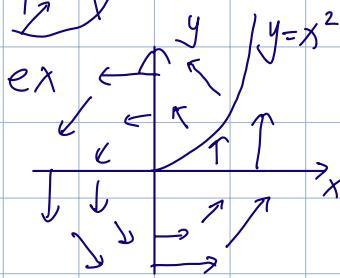


line integral

$$W = \vec{F} \cdot d\vec{r} \quad \text{Along a trajectory } C \quad \vec{F} = M\vec{i} + N\vec{j}$$



$$W = \int_C \vec{F} \cdot d\vec{r} = \int_{t_1}^{t_2} \vec{F} \cdot \frac{d\vec{r}}{dt} dt = \int_C M dx + N dy$$



$\vec{F} = -y\vec{i} + x\vec{j}$

$C: x = t$

$y = t^2 \quad (0 \leq t \leq 1)$

$$\begin{aligned} \int_C \vec{F} \cdot d\vec{r} &= \int_0^1 (-y, x) \cdot (1, 2t) dt \\ &= \int_0^1 (-t^2, t) \cdot (1, 2t) dt \\ &= \int_0^1 t^2 dt = \frac{1}{3} \end{aligned}$$

Special case, when \vec{F} is gradient of a function $f(x, y)$

Fundamental theorem of Calculus for line integral

$$\int_C \nabla f \cdot d\vec{r} = f(P_1) - f(P_0)$$



proof: $\int_C \nabla f \cdot d\vec{r} = \int_C f_x dx + f_y dy = \int_C df = f(P_1) - f(P_0)$

proof: $\int_C \nabla f \cdot d\vec{r} = \int_C f_x dx + f_y dy \quad C: x = x(t), y = y(t)$

$$= \int_C \left(f_x \frac{dx}{dt} + f_y \frac{dy}{dt} \right) dt = \int_{t_0}^{t_1} \frac{df}{dt} dt = f(x(t_1), y(t_1)) \Big|_{t_0}^{t_1} = f(P_1) - f(P_0)$$

Only works if \vec{F} is a gradient field

(1) Path-independence $\int_{C_1} \nabla f \cdot d\vec{r} = \int_{C_2} \nabla f \cdot d\vec{r}$ if C_1 and C_2 have same start/end points

\Leftrightarrow (2) \vec{F} is a gradient field

\Leftrightarrow (3) $M dx + N dy$ is an exact differential (can be expressed in df)

Lecture 2 | Exact differential

if $\vec{F} = \nabla f$, $M = f_x$ $N = f_y$, then $f_{xy} = f_{yx} \Rightarrow M_y = N_x$

Then: if $\vec{F} = \langle M, N \rangle$ defined, differentiable everywhere

and $M_y = N_x$, then \vec{F} is a gradient field

ex. $\vec{F} = -y\hat{i} + x\hat{j}$

$\frac{\partial M}{\partial y} = -1$ $\frac{\partial N}{\partial x} = 1$ are not the same, not a gradient field

ex. $\vec{F} = (4x^2 + axy)\hat{i} + (3y^2 + 4x^2)\hat{j}$

$\frac{\partial M}{\partial y} = ax = \frac{\partial N}{\partial x} = 8x$ $a=8$

Then, how do we find the potential (only if $M_y = N_x$)

(1) Computing line integral

$$\int_C \vec{F} d\vec{r} = f(x_1, y_1) - f(0, 0) \text{ choose a constant}$$
$$\int_C \vec{F} d\vec{r} = (4x^2 + 8xy) dx + (3y^2 + 4x^2) dy$$
$$C_1: y=0 \quad dy=0 \quad \int_{C_1} \vec{F} d\vec{r} = \int_0^{x_1} 4x^2 dx = \frac{4}{3} x_1^3$$
$$C_2: y: 0 \rightarrow y_1 \quad x=x_1 \quad dx=0 \quad \int_{C_2} \vec{F} d\vec{r} = \int_0^{y_1} (3y^2 + 4x_1^2) dy = y_1^3 + 4x_1^2 y_1$$
$$f(x_1, y_1) = \frac{4}{3} x_1^3 + y_1^3 + 4x_1^2 y_1 + C \quad f(x, y) = \frac{4}{3} x^3 + y^3 + 4x^2 y + C$$

(2) Antiderivatives

Want to solve $\begin{cases} f_x = 4x^2 + 8xy & \textcircled{1} \\ f_y = 3y^2 + 4x^2 & \textcircled{2} \end{cases}$

$\textcircled{1}: f = \frac{4}{3} x^3 + 4x^2 y + g(y)$

$\Rightarrow f_y = 4x^2 + g'(y) \quad \text{match with } \textcircled{2} \quad g'(y) = 3y^2 \quad g(y) = y^3 + C$

$f = \frac{4}{3} x^3 + 4x^2 y + y^3 + C$

\vec{F} is a conservative field, then $\oint_C \vec{F} d\vec{r} = 0$

Definition $\text{curl}(\vec{F}) = Nx - My$ and $Nx = My \Rightarrow \text{curl}(\vec{F}) = 0$

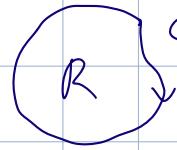
For a velocity field, curl measures rotation component of motion

ex. $\vec{F} = \langle -y, x \rangle$ $\text{curl}(\vec{F}) = \frac{\partial}{\partial x}(x) - \frac{\partial}{\partial y}(-y) = 2$

curl measures (2π) angular velocity of rotation component of velocity field

in a force field, curl measure torque $\frac{\text{torque}}{\text{moment of inertia}} = \frac{d}{dt}(\text{angular velocity})$

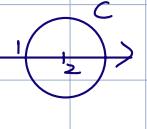
Lecture 22 Green Theorem

 If C is a closed curve enclosing a region R , counter-clockwise

\vec{F} vector field defined and differentiable in R

$$\text{Then } \oint_C \vec{F} \cdot d\vec{r} = \iint_R \operatorname{curl}(\vec{F}) dA$$

$$\oint_C M dx + N dy = \iint_R (N_x - M_y) dA$$

ex. 

$$\begin{aligned} & \oint_C ye^{-x} dx + (\frac{1}{2}x^2 - e^{-x}) dy \\ &= \iint_R (x + e^{-x} - e^{-x}) dA \\ &= 2\pi \end{aligned}$$

Proof: Special case: $\operatorname{curl}(\vec{F}) = 0$

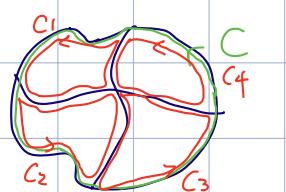
$$\oint_C \vec{F} \cdot d\vec{r} = \iint_R \operatorname{curl}(\vec{F}) dA = 0 \text{ if } \operatorname{curl}(\vec{F}) = 0$$

Proof: $\oint_C M dx + N dy = \iint_R (N_x - M_y) dA$

Observe: first prove $\oint_C M dx = \iint_R -M_y dA$ (special case when $N=0$)

Similarly: $\oint_C N dy = \iint_R N_x dA$ ($M=0$)

Observe 2: decompose R into simple regions

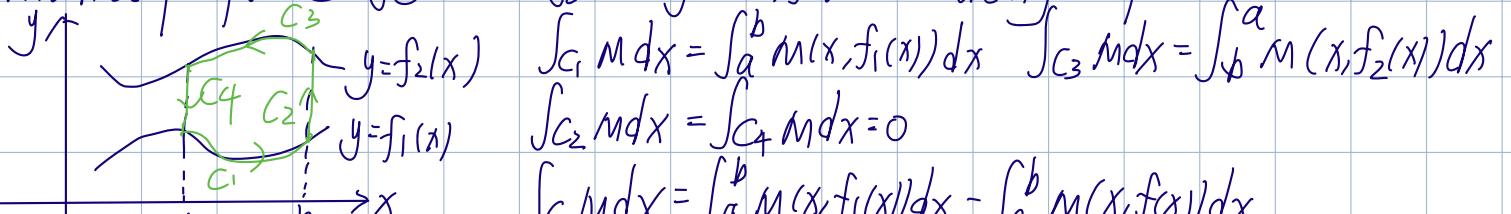


$$\begin{aligned} \oint_C M dx + N dy &= \sum_{i=1}^n \oint_{C_i} M dx + N dy \\ \iint_R (N_x - M_y) dA &= \sum_{i=1}^n \iint_{R_i} (N_x - M_y) dA \end{aligned}$$

(go twice along boundary, cancelled out)

Cut R into pieces that for $a \leq x \leq b$, $f_1(x) \leq y \leq f_2(x)$

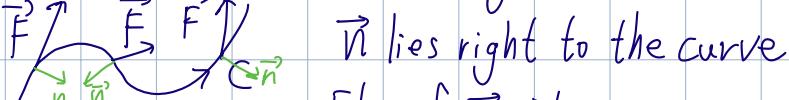
Main step: prove $\oint_C M dx = \iint_R -M_y dx$ if R vertically simple



$$\iint_R -M_y dA = - \int_a^b \int_{f_1(x)}^{f_2(x)} \frac{\partial M}{\partial y} dy dx = - \int_a^b [M(x, f_2(x)) - M(x, f_1(x))] dx = \int_a^b [M(x, f_1(x)) - M(x, f_2(x))] dx$$

Lecture 23 Flux

Flux: another line integral: $\int_C \vec{F} \cdot \hat{n} \cdot d\vec{s}$



Flux: $\int_C \vec{F} \cdot \hat{n} \cdot d\vec{s}$ measure normal component of \vec{F}

Interpretation: For \vec{F} a velocity field, Flux measure how much fluid passes through C

Calculation: $d\vec{r} = \hat{T} d\vec{s} = \langle dx, dy \rangle$

\hat{n} is \hat{T} rotate 90° clockwise: $\hat{n} d\vec{s} = \langle dy, -dx \rangle$

$$\text{if } \vec{F} = \langle P, Q \rangle \quad \int_C \vec{F} \cdot \hat{n} \cdot d\vec{s} = \int_C \langle P, Q \rangle \cdot \langle dy, -dx \rangle = \int_C -Q dx + P dy$$

Green Theorem for flux

If C encloses a region R counterclockwise and $\vec{F} = \langle P, Q \rangle$ defined in R
then $\int_C \vec{F} \cdot \hat{n} \cdot d\vec{s} = \iint_R \operatorname{div} \vec{F} dA$ (divergence of \vec{F})

$$\operatorname{div} \langle P, Q \rangle = P_x + Q_y$$

$$\int_C -Q dx + P dy = \iint_R (P_x + Q_y) dA \quad \text{the same as Green Theorem}$$

ex. $\vec{F} = x\hat{i} + y\hat{j}$ C = circle of $r=a$

$$\operatorname{div} \vec{F} = \frac{\partial}{\partial y}(y) + \frac{\partial}{\partial x}(x) = 2$$

$$\int_C \vec{F} \cdot \hat{n} \cdot d\vec{s} = \iint_R 2 dA = 2\pi a^2$$

Interpretation

(1) how much the flow is "expanding"

(2) the source rate = amount of fluid added to the system per unit time and area

line integral: $\int_C \vec{F} \cdot d\vec{r} = \int_C M dx + N dy$

line integral Flux: $\int_C \vec{F} \cdot \hat{n} \cdot d\vec{s} = \int_C -Q dx + P dy$

$\operatorname{curl} \vec{F} = N_x - M_y \quad \int_C \vec{F} \cdot d\vec{r} = \iint_R (N_x - M_y) dA$

$\operatorname{div} \vec{F} = P_x + Q_y \quad \int_C \vec{F} \cdot \hat{n} \cdot d\vec{s} = \iint_R (P_x + Q_y) dA$

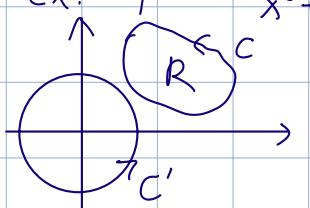
$\langle dx, dy \rangle \cdot \hat{T} \cdot d\vec{s}$

$\langle dy, -dx \rangle \cdot \hat{n} \cdot d\vec{s}$

More validity of Green's Theorem

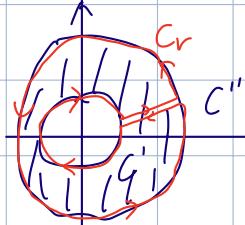
Only work if \vec{F} (and derivative) defined everywhere

ex. $\vec{F} = \frac{-y\hat{i} + x\hat{j}}{x^2 + y^2}$ For not $x=y=0$, $\text{curl } \vec{F} = 0$



$$\oint_C \vec{F} d\vec{r} = \iint_R \text{curl } \vec{F} dA = 0$$

$\oint_{C'} \vec{F} d\vec{r} = ?$ can't use Green theorem directly

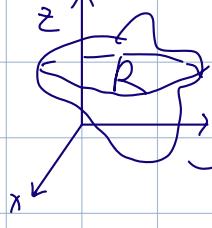


$$\oint_{C''} \vec{F} d\vec{r} - \oint_{C'} \vec{F} d\vec{r} = \iint_R \text{curl } \vec{F} dA$$

take line integral for C_r to prove

If $\text{curl } \vec{F} = 0$ and domain where \vec{F} defined is simply connected, \vec{F} is conservative and is a gradient

Lecture 25 Triple Integrals



$$\iiint_R f(x, y, z) dz dy dx$$

ex. region between $z = x^2 + y^2$ $z = 4 - x^2 - y^2$

$$\begin{aligned} \text{volume} &= \iiint_R dV = \int \int \int_{x^2+y^2}^{4-x^2-y^2} dz dy dx \\ &= \int_{-12}^{12} \int_{\sqrt{12-x^2}}^{\sqrt{4-x^2-y^2}} \int_{x^2+y^2}^{4-x^2-y^2} dz dy dx \\ &= \int_{-12}^{12} \int_{\sqrt{12-x^2}}^{\sqrt{4-x^2-y^2}} (4-x^2-y^2) dy dx \end{aligned}$$

find whenever $z_{\text{bottom}} < z_{\text{top}}$ $x^2 + y^2 \leq 4 - x^2 - y^2$, $x^2 + y^2 \leq 2$

project

switch to polar coordinate

Better: Polar coordinates

$$\int_0^{2\pi} \int_0^{\sqrt{2}} \int_{r^2}^{4-r^2} dz \cdot r dr d\theta \quad \text{cylindrical coordinate}$$

$$\begin{aligned} &\begin{array}{l} z \uparrow \\ r \\ (r, \theta, z) \quad x = r \cos \theta \\ \theta \\ r \\ z \uparrow \\ y \\ y = r \sin \theta \end{array} \\ &dx dy dz = dV = r dr d\theta dz \end{aligned}$$

Average value of $f(x, y, z)$ on x, y, z

$$\bar{f} = \frac{1}{\text{Vol}(R)} \iiint_R f dV \quad \text{or with weighted } \bar{f} = \frac{1}{\text{Mass}(R)} \iiint_R f \delta dV$$

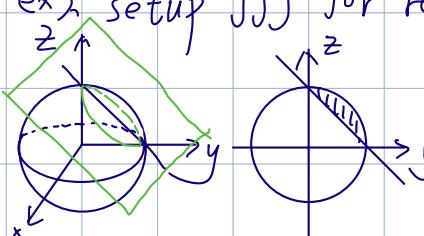
$$\text{Center of Mass } (\bar{x}, \bar{y}, \bar{z}), \quad \bar{x} = \frac{1}{\text{Mass}} \iiint_R x \delta dV \quad \bar{y} = \frac{1}{\text{Mass}} \iiint_R y \delta dV \quad \bar{z} = \frac{1}{\text{Mass}} \iiint_R z \delta dV$$

Moments of inertia

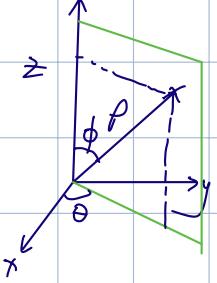
$$I = \iiint_R (\text{distance to the axis})^2 \delta dV \quad \text{ex. } I_z = \iiint_R r^2 \delta dV$$

ex3. setup \iiint for region $z > 1-y$ inside unit ball

$$\begin{aligned} &\begin{array}{l} z \uparrow \\ \text{unit ball} \\ z \uparrow \\ y \\ y \end{array} \\ &= \int_0^1 \int_{\sqrt{1-y^2}}^{\sqrt{1-x^2-y^2}} \int_{-y}^{1-y} dz dx dy \\ &(1-y) < \sqrt{1-x^2-y^2} \\ &(1-y)^2 < 1-x^2-y^2 \end{aligned}$$



Lecture 26 Spherical coordinate

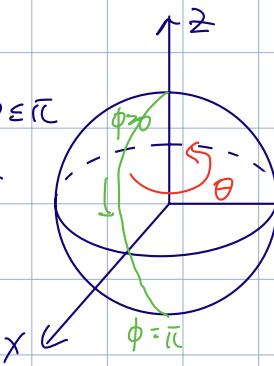


ϕ : go down from z -axis $0 \leq \phi \leq \pi$

θ : counterclockwise from x -axis

$$z = \rho \cos \phi \quad \rho = \sqrt{r^2 + z^2} = \sqrt{x^2 + y^2 + z^2}$$

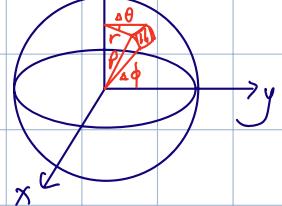
$$r = \rho \sin \phi \quad x = \rho \sin \phi \cos \theta \\ y = \rho \sin \phi \sin \theta$$



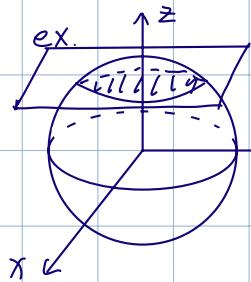
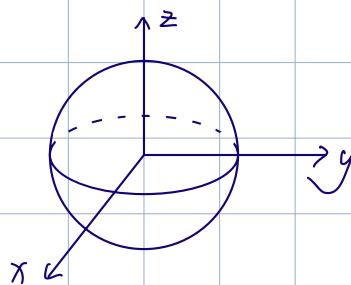
fix $\rho = a$

$\phi \leftrightarrow$ latitude
 $\theta \leftrightarrow$ longitude

Triple integral in spherical coordinate



$$dV = \rho d\phi \, r d\theta \, dr \\ = \rho \rho \sin \phi \, d\rho \, d\phi \, d\theta \\ = \rho^2 \sin \phi \, d\rho \, d\phi \, d\theta$$

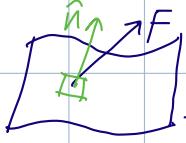


$$= \int_0^{2\pi} \int_0^{\pi/4} \int_{\frac{r}{2} \sec \theta}^1 (1 - \frac{\sqrt{2}}{2} \sec \theta) \rho^2 \sin \phi \, d\rho \, d\phi \, d\theta \\ = \frac{2\pi}{3} - \frac{5\pi}{6\sqrt{2}}$$

Lecture 27 Flux in space

flux through a surface; surface integral

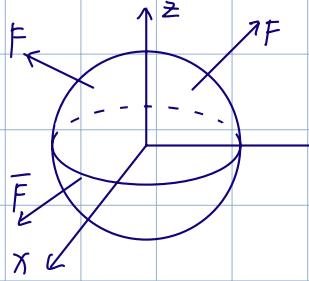
\vec{F} vector field, S a surface



2 choices for \hat{n} (orientation)

$$\text{Flux} = \iint_S \vec{F} \cdot \hat{n} dS \quad \text{Notation } d\vec{S} = \hat{n} dS$$

ex1. Flux of $\vec{F} = \langle x, y, z \rangle$ through sphere of radius a centered at O



$$\hat{n} = \frac{1}{a} \langle x, y, z \rangle$$

$$\vec{F} \cdot \hat{n} = |\vec{F}| = a$$

$$\iint_S \vec{F} \cdot \hat{n} dS = \iint_S a dS = 4\pi a^3$$

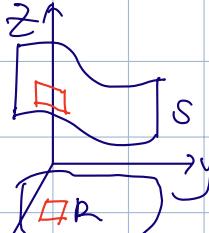
ex2. Same sphere $\vec{F} = z \hat{k}$

$$\vec{F} \cdot \hat{n} = \langle 0, 0, z \rangle \cdot \langle x, y, z \rangle \cdot \frac{1}{a} = \frac{z^2}{a}$$

$$\iint_S \vec{F} \cdot \hat{n} dS = \iint_S \frac{z^2}{a} dS$$

$$dS = \rho^2 \sin\phi d\phi d\theta = a^2 \sin\phi d\phi d\theta \quad z = a \cos\phi$$

$$\iint_S \frac{z^2}{a} dS = \int_0^{2\pi} \int_0^\pi \frac{a^2 \cos^2\phi}{a} \cdot a^2 \sin\phi d\phi d\theta = \frac{4}{3}\pi a^3$$



graph $z = f(x, y)$

$$\hat{n} dS = \pm \langle -f_x, -f_y, 1 \rangle \underbrace{dx dy}_{\text{Not } \hat{n}} \underbrace{dS}_{\text{Not } ds}$$

simplifies magically!

To set up bounds on $\iint \cdots dx dy$, look at shadow on $x-y$ plane

$$\pm \vec{u} \times \vec{v} = \Delta S \cdot \hat{n}$$

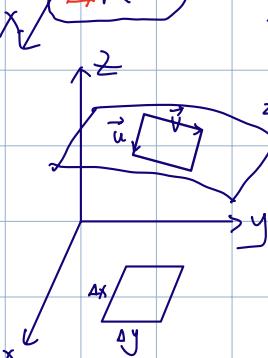
$$\approx f(x, y) + \Delta x \cdot f_x$$

\vec{u} : from $(x, y, f(x, y))$ to $(x + \Delta x, y, f(x + \Delta x, y))$

$$\begin{aligned} \vec{u} &= \langle \Delta x, 0, f_x \Delta x \rangle & \vec{v} &= \langle 0, \Delta y, f_y \Delta y \rangle \\ &= \langle 1, 0, f_x \rangle \Delta x & &= \langle 0, 1, f_y \rangle \Delta y \end{aligned}$$

$$\hat{n} \Delta S = \vec{u} \times \vec{v} = \begin{vmatrix} i & j & k \\ 1 & 0 & f_x \\ 0 & 1 & f_y \end{vmatrix} \Delta x \Delta y = \langle -f_x, -f_y, 1 \rangle \Delta x \Delta y$$

$$\hat{n} dS = \pm \langle -f_x, -f_y, 1 \rangle dx dy$$



ex, $\vec{F} = z\hat{k}$, $z = x^2 + y^2$ above the unit circle

$$= \iint_{\text{unit circle}} \langle 0, 0, z \rangle \cdot \langle -2x, -2y, 1 \rangle dy dx \quad \text{choose the unit vector upward}$$

$$= \int_{-1}^1 \int_{-\sqrt{1-x^2}}^{\sqrt{1-x^2}} x^2 + y^2 dy dx \quad \text{get rid of } z$$

$$= \int_0^{2\pi} \int_0^1 r^2 \cdot r dr d\theta = \int_0^{2\pi} \frac{1}{4} d\theta = \frac{\pi}{2}$$