

Nonlinear Legendre Spectral Finite Elements for Wind Turbine Blade Dynamics



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Motivation

- Beam model currently used in FAST
 - · Euler-Bernoulli beam model with shortening effect
 - Two degree-of-freedoms
 - Assumed-mode method
- Beam models used in other wind turbine tools
 - Multibody-formulation
 - Linear beam models
 - Constraints introduced between linear beams to describe large deflections and rotations
 - Finite element method

Objective

- Objective: create efficient high-fidelity beam models for wind turbine blade analysis that can
 - Capture geometrical nonlinearity systematically
 - Capture anisotropic and heterogeneous behavior of composite materials rigorously
 - Modeling moving beams (translation and rotation)
 - Achieve the speed of computational design without significant loss of accuracy comparing to the ultimate accuracy obtained by 3D nonlinear FEA
 - Compatible with the FAST modularization framework





NWTC static and dynamic blade tests showing typical large, elastic deflections.

- Implementation
 - Geometrically Exact Beam Theory (GEBT)

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 - Successfully applied to simulation of fluid dynamics, geophysics, elastodynamics
 - Limited usage in structural dynamics
 - FAST Modularization Framework (Jonkman, 2013)
 - State-space formulation for tight-coupling scheme
 - Time integrator for first-order PDEs
- Result: BeamDyn, which can be used as a structural module of FAST

Governing Equation

$$\frac{\dot{\underline{h}} - \underline{F'} = \underline{f}}{\dot{\underline{g}} + \dot{\overline{u}}\underline{h} - \underline{M'} - (\widetilde{x}'_0 + \widetilde{u}')\underline{F} = \underline{m}}$$

Constitutive Equation

$$\left\{ \frac{\underline{h}}{\underline{g}} \right\} = \underline{\mathcal{M}} \left\{ \underline{\underline{u}} \right\} \\
 \left\{ \underline{\underline{F}} \right\} = \underline{\underline{\mathcal{C}}} \left\{ \underline{\underline{\epsilon}} \right\}$$

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- Elastic couplings are captured

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Strain Measures

$$\left\{ \underline{\underline{\epsilon}} \right\} = \left\{ \underline{\underline{x}}_0' + \underline{\underline{u}}' - (\underline{\underline{R}} \, \underline{\underline{R}}_0) \overline{\iota}_1 \right\}$$

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- Geometrically exact: deformed beam geometry is represented exactly
- Small strains

State-space Formulation

Governing Equation

$$\underline{\underline{\mathfrak{M}}}\,\underline{a}+f(\underline{q},\underline{v},t)=0$$

$$\underline{q}^T = \left[\underline{u}^T \ \underline{p}^T\right]$$
> State-Space Form

$$\underline{\boldsymbol{v}}^T = \begin{bmatrix} \underline{\dot{\boldsymbol{u}}}^T & \underline{\boldsymbol{\omega}}^T \end{bmatrix} \qquad \qquad \underline{\boldsymbol{a}}^T = \begin{bmatrix} \underline{\ddot{\boldsymbol{u}}}^T & \underline{\dot{\boldsymbol{\omega}}}^T \end{bmatrix}$$

$$\underline{\underline{A}}\,\dot{\hat{\underline{x}}}(t)=\mathfrak{f}(\dot{\hat{\underline{x}}}(t),t)$$

$$\underline{\underline{x}}(t) \equiv \left\{ \underline{\underline{q}}(t) \right\} \qquad \underline{\underline{\underline{q}}}(\underline{\hat{\underline{x}}}(t)) = \begin{bmatrix} \underline{\underline{\underline{D}}} & \underline{\underline{\underline{M}}} \end{bmatrix}$$

$$\underline{\underline{\underline{D}}}(\underline{\hat{x}}(t)) = \int_0^1 \underline{\underline{\underline{M}}}^T \begin{bmatrix} \underline{\underline{I}} & \underline{\underline{\underline{0}}} \\ \underline{\underline{\underline{0}}} & \underline{\underline{\underline{H}}} \end{bmatrix} \underline{\underline{\underline{M}}} dx_1 \qquad \qquad \mathfrak{f}(\underline{\hat{x}}(t), t) = \left\{ \underbrace{\int_0^1 \underline{\underline{\underline{M}}}^T \underline{\underline{V}} dx_1}_{\underline{\underline{F}}(\underline{\hat{x}}(t), t)} \right\}$$

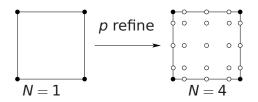
Second-order Adams-Moulton (AM2)

$$\underline{\underline{A}}_{k+1}(\underline{\hat{x}}_{k+1} - \underline{\hat{x}}_k - \frac{\Delta t}{2}\underline{\hat{x}}_k) = \frac{\Delta t}{2}\mathfrak{f}(\underline{\hat{x}}_{k+1}, t_{k+1})$$

Legendre Spectral Finite Elements

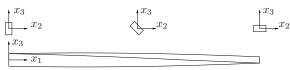
LSFE methods combine the geometric flexibility of the FE method with the accuracy of global spectral methods.

- Solution improved through increased basis polynomial order (p-refinement)
- ► LSFEs employ Lagrangian interpolant shape functions with nodes at Gauss-Lobatto-Legendre (GLL) points
- Exponential convergence rates for sufficiently smooth solutions



Example 1: Initially Twisted Beam

Sketch of Initially Twisted Beam



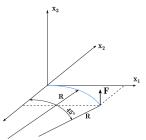
Result

Table: Comparison of tip displacements of an initially twisted beam

	u_1 (m)	<i>u</i> ₂ (m)	u_3 (m)
BeamDyn	-1.132727	-1.715123	-3.578671
ANSYS	-1.134192	-1.714467	-3.584232
Percent Error	0.129%	0.038%	0.155%

Example 1: Initially Curved Beam

Sketch of Initially Curved Beam



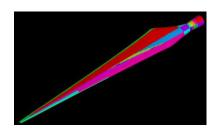
Result

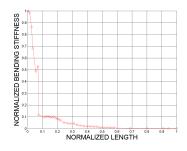
Table: Comparison of tip displacements of an initially curved beam

	u_1 (inches)	u ₂ (inches)	u_3 (inches)
BeamDyn (one LSFE)	-23.7	13.5	53.4
Bathe-Bolourchi ?	-23.5	13.4	53.4

Example 2: CX-100

▶ Sketch of CX-100



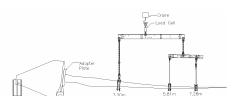


Sectional Properties at 2.2 m

$$C = 10^{3} \times \begin{bmatrix} 193,000 & -75.4 & 12.2 & -75.2 & -1970 & -3500 \\ -75.4 & 19,500 & 4,760 & 62.6 & 67.3 & 11.3 \\ 12.2 & 4,760 & 7,210 & -450 & 17.0 & 2.68 \\ -75.2 & 62.6 & -450 & 518 & 1.66 & -1.11 \\ -1,970 & 67.3 & 17.0 & 1.66 & 2,280 & -879 \\ -3,500 & 11.6 & 2.68 & -1.11 & -875 & 4,240 \end{bmatrix}$$

Example 2: CX-100 (Continued)

Static Test Configuration



Deflection

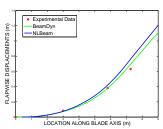


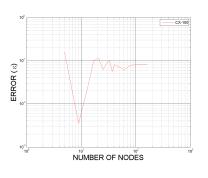
Table: Experimental and BeamDyn simulation results for the CX-100 static test

Saddle

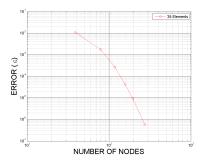
	u_3 at saddle #1 (m)	u_3 at saddle #2 (m)	u_3 at saddle #3 (m)
Experimental	0.083530	0.381996	0.632460
BeamDyn	0.072056	0.381074	0.698850
Percent Error	13.74%	0.24%	10.5%

Example 2: Convergence Study

Validation



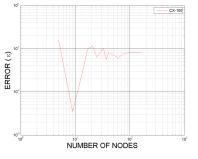
Verification

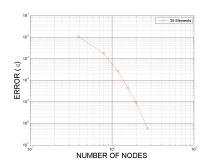


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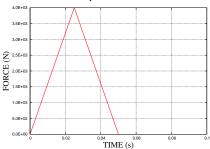




- Sharp gradients in sectional properties
- Erratic data

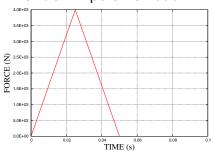
Example 3: Damping Effect

Cantilever Beam Under Impulsive Load



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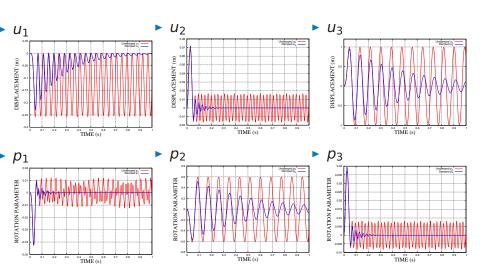
Viscous Damping

$$\underline{\underline{f}}_{d} = \underline{\underline{\mu}} \, \underline{\underline{\mathcal{C}}} \, \left\{ \begin{matrix} \dot{\boldsymbol{\epsilon}} \\ \dot{\boldsymbol{\kappa}} \end{matrix} \right\}$$

► RMS Error

$$arepsilon_{RMS} = \sqrt{rac{\sum_{k=0}^{n_{max}} [u_3^k - u_b(t^k)]^2}{\sum_{k=0}^{n_{max}} [u_b(t^k)]^2}}$$

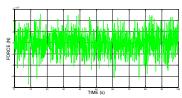
Example 3: Root forces and moments



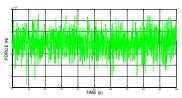
▶ NREL 5-MW Blade; Cantilevered at root

- NREL 5-MW Blade; Cantilevered at root
- White noise force applied at the free tip along flap direction

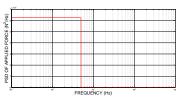
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- Time History of Applied Force



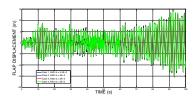
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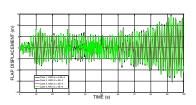
PSD of Applied Force



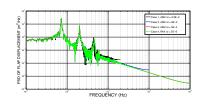
Flapwise Response



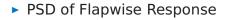
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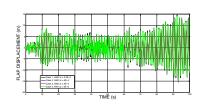


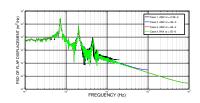
PSD of Flapwise Response



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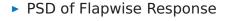


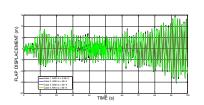


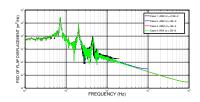


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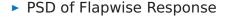


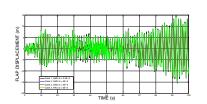


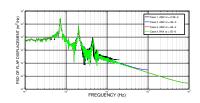


- For an implicit AM2 time step beyond 0.005 s, the solution is nearly identical to the fully resolved explicit RK4 solution
- For an AM2 time step of 0.025 s, the solution remains stable and tracks the other solutions, but error grows at higher frequencies

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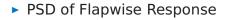


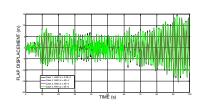


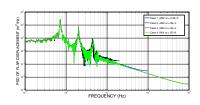


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- For an AM2 time step of 0.025 s, the solution remains stable and tracks the other solutions, but error grows at higher frequencies
- The spikes at 0.7 Hz and 2 Hz correspond to the first and second blade flapwise natural frequencies
- The spike above 5 Hz-above the frequency range of excitation-is brought about by nonlinear effects

Convergence Rate

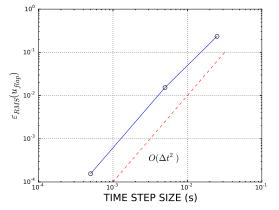


Figure: Normalized RMS error of flapwise displacement histories as a function of time step size for AM2 time integrator. The dashed line shows ideal second-order convergence.

Solver Statistics

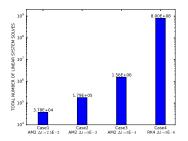


Figure: Total number of linear system solves.

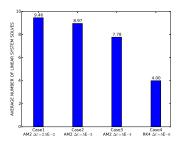


Figure: Average number of linear system solves per step.

Summary

Conclusion

- Based on geometrically exact beam theory, BeamDyn is capable of dealing with geometric nonlinear beam problems with arbitrary magnitude of displacements and rotations for both static and dynamic analyses
- Along with a preprocessor like PreComp or VABS, BeamDyn takes full elastic coupling effects into account
- The governing equations are reformulated into state-space form, thus, making it amendable into FAST for tight-coupling analysis
- The space is discretized by spectral finite elements, which
 is a p-version finite element, so that exponential
 convergence rate can be expected for smooth solutions
- Different time integrators have been implemented in BeamDyn; users will have options based on their needs
- BeamDyn is implemented following the programming requirements (data structures and interfaces) of the FAST modularization framework

Summary (Continued)

- ► Future Work
 - Coupling BeamDyn to FAST
 - Full-Turbine validation
 - Enhancement of numerical performance

Questions?

Acknowledgments

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