MAE 109 - ELECTROMAGNETISM

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Part III Plane Wave In Vacuum

- I. Introduction
- II. Electromagnetism Maxwell Equation(Spatial/temporal equations)
- III. Plane Wave in vacuum
- IV. Wave propagation in matter
- V. Boundary conditions (Reflection/refraction)

Results | Electromagnetism - Maxwell Equation(Spatial/temporal equations)

(1) $\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\varepsilon_0}$

Local Gauss Law (electric flux density)

(2) $\vec{\nabla} \cdot \vec{B} = 0$

General Magnetism Law

(3) $\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$

Faraday Law (relationship between Electric and Magnetic Field)

(4) $\overrightarrow{V} \times \overrightarrow{B} = \mu_0 \left(\overrightarrow{j} + \epsilon_0 \frac{\partial \overrightarrow{E}}{\partial t} \right)$ Ampere law (current flow in a wire creating a magnetic field)

E & B Resolution

$$\Delta \vec{E} - \varepsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2} = \frac{1}{\varepsilon_0} \vec{\nabla} \rho + \mu_0 \frac{\partial \vec{J}}{\partial t}$$

$$\Delta \vec{B} - \varepsilon_0 \mu_0 \frac{\partial^2 \vec{B}}{\partial t^2} = -\mu_0 \vec{\nabla} \times \vec{J}$$



Electromagnetic **Wave Equation?**

Summary

- I. Introduction
- II. Electromagnetism Maxwell Equation(Spatial/temporal equations)
- III. Plane Wave in vacuum
 - 1. From D'Alembert Equation to Helmoltz equation
 - 2. Propagation of the EM plane wave in vacuum
 - 3. Harmonic Solution
 - 4. Velocity of the travelling Wave
 - 5. Polarization of the EM Wave in vacuum
- IV. Wave propagation in matter
- V. Boundary conditions (Reflection/refraction)

III.1. From D'Alembert Equation to Helmoltz equation

Wireless radio systems

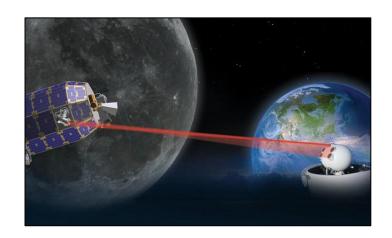


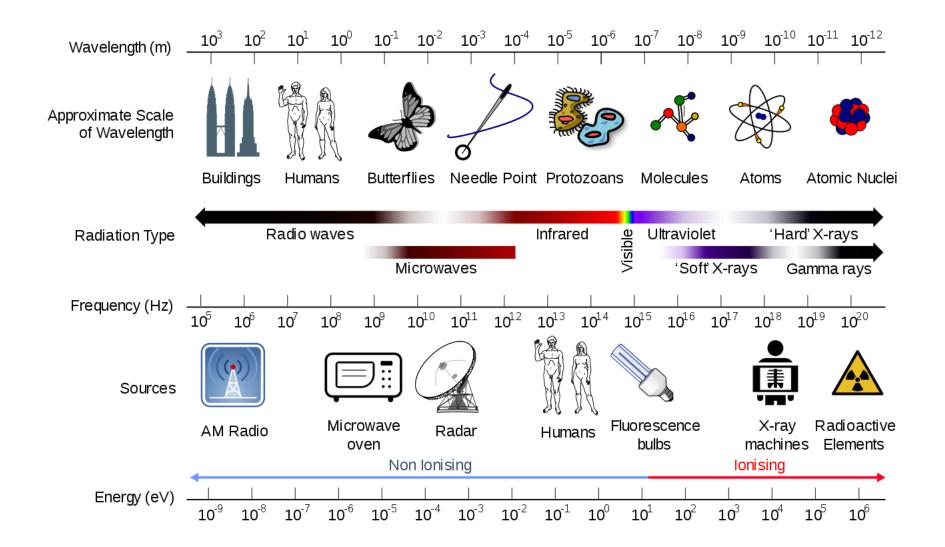


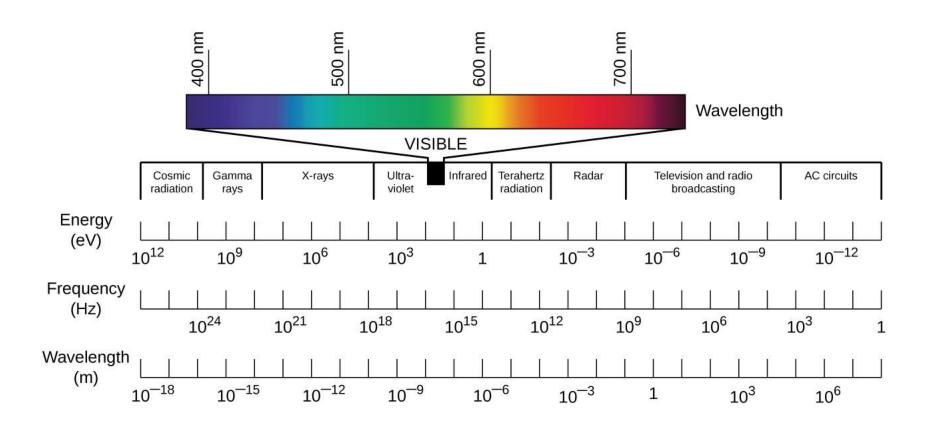
Electromagnetic propagation in the channel

Free space optics









III.1 From D'Alembert Equation to Helmoltz equation

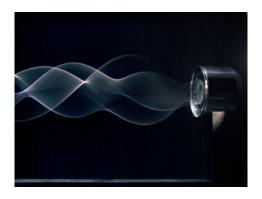
What is an Electromagnetic Field and Electromagnetic WAVE?

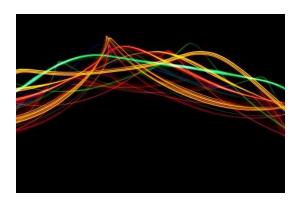
How to demonstrate the propagation?



Different kind of waves:

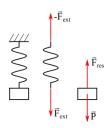
- Mechanical wave (fluid, acoustic)
- Electromagnetic Wave (Radio, Optics)
- Matter waves (De Broglie Wave / Wave nature of the matter).

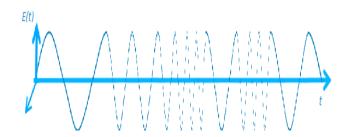




III.1 Analogy between Electromagnetic wave and mecanical waves Longitudinal & tranverse Wave

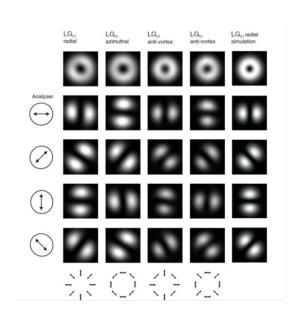
Longitudinal Wave as a Spring





Transverse Wave – Vibrating-Wire





Part III.1 Wave Equation

Harp string moving: D'Alembert Equation

Wave / stationary (SteadyState) phenomenon





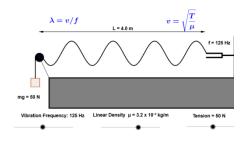
D'Alembert Equation (1730) for a vibrating string :

$$\Delta \vec{p} = \frac{1}{c^2} \frac{\partial^2 \vec{p}}{\partial t^2}$$

 $\overrightarrow{\boldsymbol{p}}$ is a vector corresponding to the pressure in acoustic

c is the speed of the sound in the air

Geogebra Activity on LMS: Vibrating String

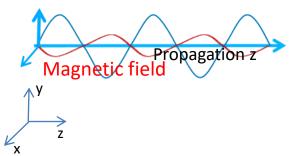


What about the Electromagnetism?

Part III.1 Plane wave in Vaccum: Vacuum Propagation

(No charge) (No current) Permittivity $arepsilon_0$ and Permeability μ_0

Electric field



$$\Delta \vec{E} - \varepsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2} = \frac{1}{\varepsilon_0} \vec{\nabla} \rho + \mu_0 \frac{\partial \vec{J}}{\partial t}$$

$$\Delta \vec{B} - \varepsilon_0 \mu_0 \frac{\partial^2 \vec{B}}{\partial t^2} = -\mu_0 \vec{\nabla} \times \vec{J}$$

General Wave equation

$$\Delta \vec{E} - \varepsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$
$$\Delta \vec{B} - \varepsilon_0 \mu_0 \frac{\partial^2 \vec{B}}{\partial t^2} = 0$$

$$\Delta \vec{B} - \varepsilon_0 \mu_0 \frac{\partial^2 \vec{B}}{\partial t^2} = 0$$

Travelling Wave

$$\Delta \vec{E} = \varepsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$
$$\Delta \vec{B} = \varepsilon_0 \mu_0 \frac{\partial^2 \vec{B}}{\partial t^2}$$

$$\Delta \vec{B} = \varepsilon_0 \mu_0 \frac{\partial^2 \vec{B}}{\partial t^2}$$



$$\Delta \vec{p} = \frac{1}{c^2} \frac{\partial^2 \vec{p}}{\partial t^2}$$

III.1 Plane wave in Vaccum: Wave Equation

$$\Delta \vec{E} = \varepsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\Delta \vec{B} = \varepsilon_0 \mu_0 \frac{\partial^2 \vec{B}}{\partial t^2}$$



D'Alembert Equation of E, B Field

$$\Delta \vec{A} = \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2}$$

$$c = \frac{1}{\sqrt{(\varepsilon_0 \mu_0)}}$$

$$\begin{split} \varepsilon_0 &= 8,85418782 \times 10^{-12} \, \mathrm{F \, m^{-1}} \, (or \, A^2. \, s^4. \, kg^{-1}. \, m^{-3}) \, , \\ \mu_0 &= 4\pi \cdot 10^{-7} \ H/m \, (or \, kg \cdot m \cdot A^{-2} \cdot s^{-2}) \end{split}$$

c is a velocity (m/s) Which Velocity?

The same for E and B Field.

E & B are Coupled and propagate with the same Velocity

III.1. Velocity of Electromagnetic Wave: speed of the light



Huygens Jupiter moon eclipse (1675) 220000km/s



Bradley Light Aberration (1729) 301000km/s



Fizeau Foucault
Rotating mirrors (1849)
298000km/s



Michelson Interferometer (1926) 299796km/s

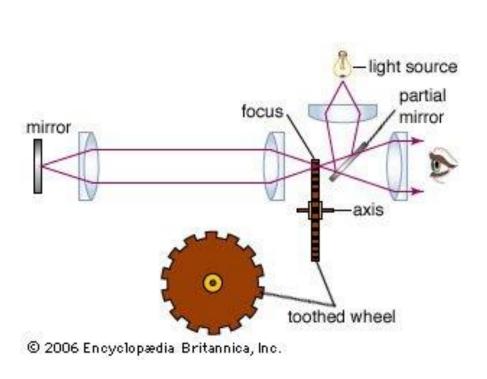
Maxwell 1864: Dynamical Theory of the Electromagnetic Field (Nature of Light and

Electromagnetics): c=31074000m. $s^{-1}=\frac{1}{\sqrt{\varepsilon_0\mu_0}}$

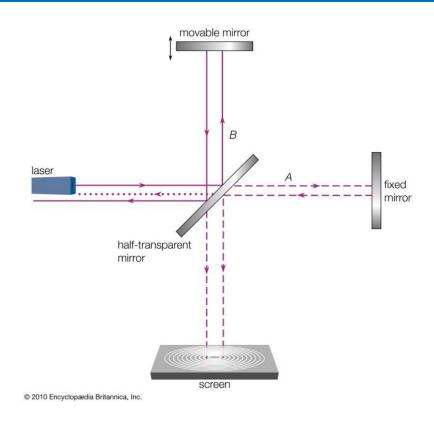
Today (General Conference on weights and measures – CGPM) **299792.458 km/s**

We assume $c \approx 3 \cdot 10^8 m/s$

III.1 speed of the light Experiments



Armand Fuzeau 1849 (5% difference) Léon Foucault 1850 (1% difference)



A.A Michelson Late 1870's (0,02 % difference)

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III. 1 From D'Alembert to Helmholtz Equation: Travelling Wave Equation

$$\Delta \vec{E} = \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\Delta \vec{B} = \frac{1}{c^2} \frac{\partial^2 \vec{B}}{\partial t^2}$$

Propagation of the E and B Field at the same velocity

 \Rightarrow the speed of the light c

Light is an Electromagnetic Wave (UV, Visible or Infrared)



D'Alembert Equation resolution in Electromagnetism to

Helmholtz Equation (Harmonic Solution)

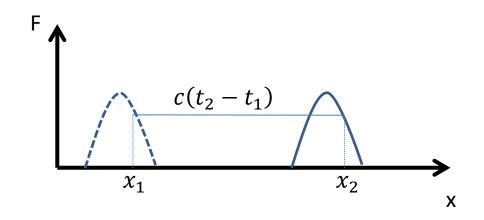
III. 2. Propagation of the EM plane wave in vacuum: One dimensions Wave Equation Resolution

1 Dimension wave Equation : $\frac{\partial^2 f}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 f}{\partial t^2}$

General solution

$$f(x,t) = F(ct - x) + G(ct + x)$$

F is a propagating signal (x direction) and G is a contrapropagating signal (-x direction)



Geogebra activity on LMS



Verify that
$$f(x,t) = F(ct-x) + G(ct+x)$$
 is a solution of Wave Equation if
$$\frac{\partial^2 f}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 f}{\partial t^2} = 0$$

We denote:

$$X = ct - x$$

and

$$Y = ct + x$$

Thus

$$x = 1/2(X - Y)$$
 and $ct = \frac{1}{2}(X + Y)$

Method:

Compute
$$\frac{\partial^2 f}{\partial x^2}$$
 and $\frac{1}{c^2} \frac{\partial^2 f}{\partial t^2}$

$$\frac{\partial f}{\partial x} = \frac{\partial f}{\partial X} \frac{\partial X}{\partial x} + \frac{\partial f}{\partial Y} \frac{\partial Y}{\partial x} = -\frac{\partial f}{\partial X} + \frac{\partial f}{\partial Y}$$

$$\frac{\partial^2 f}{\partial x^2} = \frac{\partial}{\partial x} \frac{\partial f}{\partial x} = \frac{\partial}{\partial x} \left(-\frac{\partial f}{\partial X} + \frac{\partial f}{\partial Y} \right)$$

$$= \frac{\partial}{\partial X} \left(-\frac{\partial f}{\partial X} \right) \frac{\partial X}{\partial x} + \frac{\partial}{\partial Y} \left(-\frac{\partial f}{\partial X} \right) \frac{\partial X}{\partial x} + \frac{\partial}{\partial X} \left(\frac{\partial f}{\partial Y} \right) \frac{\partial Y}{\partial x} + \frac{\partial}{\partial Y} \left(\frac{\partial f}{\partial Y} \right) \frac{\partial Y}{\partial x}$$



-1

-1



1

1

$$\frac{\partial^2 f}{\partial x^2} = \frac{\partial^2 f}{\partial X^2} + \frac{\partial^2 f}{\partial Y^2} + 2\frac{\partial^2 f}{\partial X \partial Y}$$

$$\frac{\partial f}{\partial ct} = \frac{\partial f}{\partial X} + \frac{\partial f}{\partial Y}$$

$$\frac{\partial^{2} f}{\partial (ct)^{2}} = \frac{\partial}{\partial (ct)} \frac{\partial f}{\partial (ct)} = \frac{\partial}{\partial (ct)} \left(\frac{\partial f}{\partial X} + \frac{\partial f}{\partial Y} \right)$$

$$= \frac{\partial}{\partial X} \left(\frac{\partial f}{\partial X} \right) \frac{\partial X}{\partial (ct)} + \frac{\partial}{\partial Y} \left(\frac{\partial f}{\partial X} \right) \frac{\partial X}{\partial (ct)} + \frac{\partial}{\partial X} \left(\frac{\partial f}{\partial Y} \right) \frac{\partial Y}{\partial (ct)}$$

$$+ \frac{\partial}{\partial Y} \left(\frac{\partial f}{\partial Y} \right) \frac{\partial Y}{\partial (ct)}$$

$$\frac{\partial^2 f}{\partial (ct)^2} = \frac{1}{c^2} \frac{\partial^2 f}{\partial t^2} = \frac{\partial^2 f}{\partial X^2} + \frac{\partial^2 f}{\partial Y^2} + 2 \frac{\partial^2 f}{\partial X \partial Y}$$





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Thus
$$\frac{\partial^2 f}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 f}{\partial t^2} = \mathbf{0}$$

III. 2. Propagation of the EM plane wave in vacuum: Electromagnetic Plane Wave

1 D solution of Wave equation in vacuum

Ttravelling wave in (z) Direction
$$\Rightarrow E(x, y, z, t) = E(x, y)(z - ct)$$

Thus $B(z, t) = B(x, y)(z - ct)$

Plane $z = z_0$ at $t = t_0$:

- identity of \vec{E}
- identity of \vec{B}

Plane Wave Definition

From Mawxell equation in Vacuum (no charges):

$$\vec{\nabla} \cdot \vec{E} = 0 \quad \& \vec{\nabla} \cdot \vec{B} = 0$$

E&B-Field No Component tangential to the propagation axis



$$E_z = 0$$

and

$$B_z = 0$$

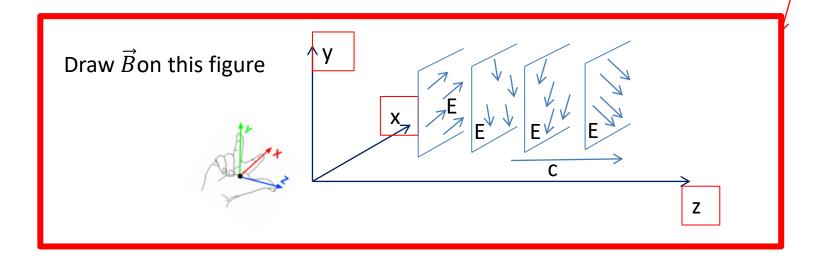
III. 2. Propagation of the EM plane wave in vacuum: EM Plane Wave

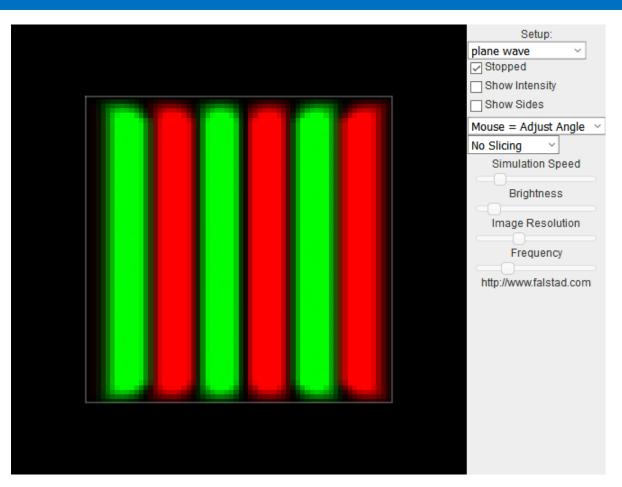
$$\overrightarrow{E(x,y,z,t)} = \overrightarrow{E_0(x,y,z)} (z-ct)$$

According to the figure propagation along z axis

$$\overrightarrow{E(x,y,z)}(z-ct) = \begin{pmatrix} E_x \\ E_y \\ E_z \end{pmatrix} = \begin{pmatrix} E_{0x}(z-ct) \\ E_{0y}(z-ct) \\ 0 \end{pmatrix}$$

$$\overrightarrow{B(x,y,z)}(z-ct) = \begin{pmatrix} B_x \\ B_y \\ B_z \end{pmatrix} = \begin{pmatrix} B_{0x}(z-ct) \\ B_{0y}(z-ct) \\ 0 \end{pmatrix}$$







to play the animation on wavebox of the plane wave according to the frequency

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Faraday Equation
$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$



See MemOperator
$$\left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z}\right) \cdot \overrightarrow{u_x} + \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x}\right) \cdot \overrightarrow{u_y} + \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y}\right) \cdot \overrightarrow{u_z}$$

$$\vec{\nabla} \times \vec{E(x, y, z)} = \begin{pmatrix} \frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} \\ \frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} \\ \frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \end{pmatrix} = \begin{pmatrix} -\frac{\partial E_y}{\partial z} \\ \frac{\partial E_x}{\partial z} \\ 0 \end{pmatrix} = \begin{pmatrix} -E_y \\ E_x \\ 0 \end{pmatrix} = \begin{pmatrix} -\frac{\partial B_x}{\partial t} \\ -\frac{\partial B_y}{\partial t} \\ 0 \end{pmatrix}$$

$$E_x = cB_y$$
 and $E_y = -cB_x$

$$ightharpoonup$$
 $|E| = c |B|$ & $\vec{E} \cdot \vec{B} = 0$ \Leftrightarrow $\vec{E} \perp \vec{B}$



FOR A PLANE WAVE IN VACUUM

$$\overrightarrow{E} \times \overrightarrow{B} = cB^2 \overrightarrow{u_z} = \frac{E^2}{c} \overrightarrow{u_z}$$

 \Rightarrow E & B are in the x,y plane

⇒ E & B are orthogonals

Propagation of a transversal wave along increasing x axis in vacuum:

$$\overrightarrow{E(x,y,z,t)} = \overrightarrow{E_0(x,y,z)} (x - ct) \qquad \overrightarrow{B(x,y,z,t)} = \overrightarrow{B_0(x,y,z)} (x - ct)$$

$$\overrightarrow{E(x,y,z)} = \begin{pmatrix} E_x \\ E_y \\ E_z \end{pmatrix} = \begin{pmatrix} \mathbf{0} \\ E_y \\ E_z \end{pmatrix} \quad \text{and} \quad \overrightarrow{B(x,y,z)} = \begin{pmatrix} B_x \\ B_y \\ B_z \end{pmatrix} = \begin{pmatrix} \mathbf{0} \\ B_y \\ B_z \end{pmatrix}$$

Relationship between E & B Maxwell Faraday

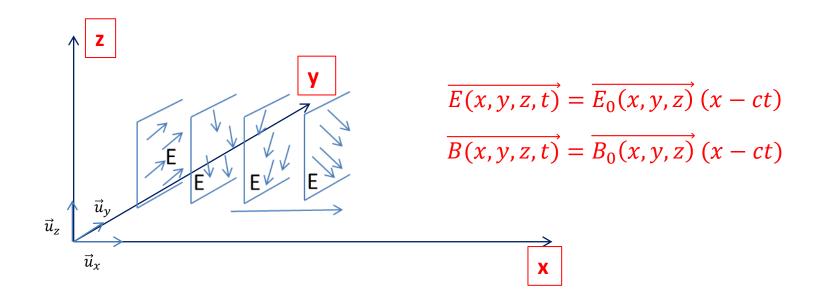
$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$-rac{\partial E_z}{\partial x} = -rac{\partial B_y}{\partial t}$$
 and $-rac{\partial E_y}{\partial x} = -rac{\partial B_z}{\partial t}$ $\Rightarrow E_z = -cB_y$ and $E_y = -cB_z$

$$\vec{E} \cdot \vec{B} = 0$$

 $\Rightarrow |\vec{E}| = c|\vec{B}|$ and $\vec{E} \times \vec{B} = cE^2 \overrightarrow{u_x}$

- Give the axis on the cartesian system
- Draw \vec{B} on this figure





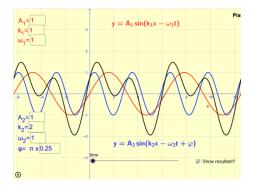
$$\vec{u}_x \times \vec{u}_y \text{=} \vec{u}_z$$

Assuming an harmonic solution allowing:

Generalization: Fourier domain (linearized), combination of waveform

$$f(x,t) = f_0 \cos \left[\omega \left(\pm \frac{x}{c} - t\right)\right] = f_0 \cos \left[\left(\pm kx - \omega t\right)\right]$$

- ω is the pulsation
- k is the Wave Number with $k = \frac{\omega}{c}$
- Geogebra activity on LMS: Superposition of two plane wave





Harmonic solution / Maxwell equation: complex notation

$$f(x,t) = f_0 e^{\left[i(\omega t - kx)\right]} + f_0 e^{\left[i(\omega t + kx)\right]} = f_0 e^{\left[-2i\pi\left(\frac{t}{T} \pm \frac{x}{\lambda}\right)\right]}$$

Propagative travelling wave (x>0)

Contra-propagative travelling wave (x<0)



According to the sign convention fixed by the observer

$$\lambda = \frac{2\pi}{k} = cT$$
 spatial period or wavelength $T = \frac{2\pi}{\omega}$ is the period ω

$$\vec{E} = \vec{E_0} \cos(kz - \omega t)$$
 $\vec{B} = \vec{B_0} \cos(kz - \omega t)$

$$\overrightarrow{E_0}$$
 and $\overrightarrow{B_0}$ Constant & Real such as $\overrightarrow{E_0} = \begin{pmatrix} E_{0x} \\ E_{0y} \\ E_{0z} \end{pmatrix}$ and $\overrightarrow{B_0} = \begin{pmatrix} B_{0x} \\ B_{0y} \\ B_{0z} \end{pmatrix}$

By considering 1 component of the E-Field (1 Dimension) = E-Field Exist only in one direction (x):

$$\vec{E} = \overrightarrow{E_0} \cos(kz - \omega t) = \begin{pmatrix} E_{0x} \cos(kz - \omega t) \\ 0 \\ 0 \end{pmatrix}$$

Computation of the wave equation: $\Delta \vec{E} - \epsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2} = \vec{0}$

Wave equation
$$\Delta \vec{E} - \epsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2} = \vec{0}$$

$$\Delta \vec{E} = \left(\frac{\partial^{2} E_{x}}{\partial x^{2}} + \frac{\partial^{2} E_{x}}{\partial y^{2}} + \frac{\partial^{2} E_{x}}{\partial z^{2}}\right) \overrightarrow{u_{x}} + \left(\frac{\partial^{2} E_{y}}{\partial x^{2}} + \frac{\partial^{2} E_{y}}{\partial y^{2}} + \frac{\partial^{2} E_{y}}{\partial z^{2}}\right) \overrightarrow{u_{y}} + \left(\frac{\partial^{2} E_{z}}{\partial x^{2}} + \frac{\partial^{2} E_{z}}{\partial y^{2}} + \frac{\partial^{2} E_{z}}{\partial z^{2}}\right) \overrightarrow{u_{z}}$$

$$\Delta \vec{E} = \frac{\partial^{2} E_{x}}{\partial z^{2}} \overrightarrow{u_{x}} = E_{0x} \frac{\partial}{\partial z} \left(\frac{\partial \cos(kz - \omega t)}{\partial z}\right) \overrightarrow{u_{x}} = -kE_{0x} \frac{\partial \sin((kz - \omega t))}{\partial z} \overrightarrow{u_{x}} = -k^{2} E_{0x} \cos(kz - \omega t) \overrightarrow{u_{x}}$$

$$\frac{\partial^{2}\vec{E}}{\partial t^{2}} = \frac{\partial}{\partial t} \left(\frac{\partial \vec{E}}{\partial t} \right) = \frac{\partial}{\partial t} \left(\frac{\partial E_{0x} cos(kz - \omega t)}{\partial t} \right) = \omega E_{0x} \frac{\partial sin(kz - \omega t)}{\partial t} \overrightarrow{u_{x}} = \omega^{2} E_{0x} cos(kz - \omega t) \overrightarrow{u_{x}}$$

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$$\vec{\nabla} \times \vec{E} = \begin{pmatrix} \frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} \\ \frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} \\ \frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \end{pmatrix} = \begin{pmatrix} 0 \\ \frac{\partial E_x}{\partial z} \\ 0 \end{pmatrix}$$

$$|\vec{\nabla} \times \vec{E}| = \begin{pmatrix} \frac{\partial \mathbf{E}_{\mathbf{z}}}{\partial \mathbf{y}} - \frac{\partial \mathbf{E}_{\mathbf{y}}}{\partial \mathbf{z}} \\ \frac{\partial \mathbf{E}_{\mathbf{x}}}{\partial \mathbf{z}} - \frac{\partial \mathbf{E}_{\mathbf{z}}}{\partial \mathbf{x}} \\ \frac{\partial \mathbf{E}_{\mathbf{y}}}{\partial \mathbf{x}} - \frac{\partial \mathbf{E}_{\mathbf{x}}}{\partial \mathbf{y}} \end{pmatrix} = \begin{pmatrix} \mathbf{0} \\ \frac{\partial E_{\mathbf{z}}}{\partial z} \\ 0 \end{pmatrix} \qquad \begin{vmatrix} \frac{\partial \vec{B}}{\partial t} = -\begin{pmatrix} B_{0x} \frac{\partial \cos(kz - \omega t)}{\partial t} \\ B_{0y} \frac{\partial \cos(kz - \omega t)}{\partial t} \\ B_{0z} \frac{\partial \cos(\omega t - kz)}{\partial t} \end{pmatrix} = \omega \sin(kz - \omega t) \begin{pmatrix} B_{0x} \\ B_{0y} \\ B_{0z} \end{pmatrix}$$

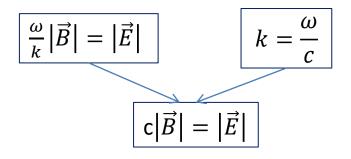
$$B_{0x}=0$$
 and $B_{0x}=0$ $\overrightarrow{B}=egin{pmatrix} 0 \ B_{0y}cos(kz-\omega t) \ 0 \end{pmatrix}$ $\overrightarrow{E}\cdot\overrightarrow{B}=0$

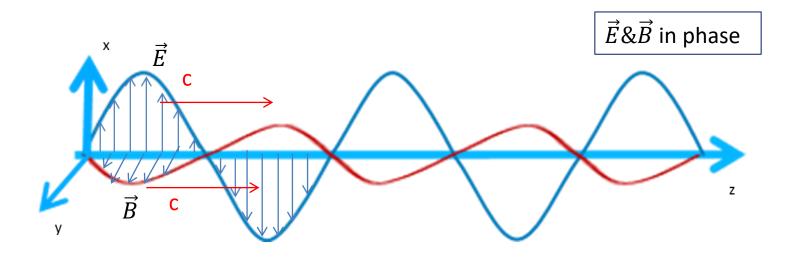
$$E_{0x} \cdot k \cdot \sin(kz - \omega t) = B_{0y} \cdot \omega \cdot \sin(kz - \omega t)$$

E & B are in Phase

 $\mathsf{E} \perp \& \; \mathsf{B} \bot \mathsf{To} \; \mathsf{the} \; \mathsf{propagation} \; \mathsf{Axis} \; \mathsf{z}$

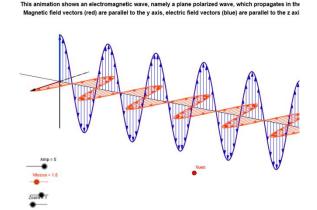
 $\vec{E} \times \vec{B}$ in the direction of the propagation Axis z





Geogebra activity





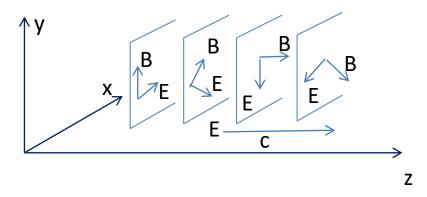
• Wave Impedance in vacuum definition : η_0

$$\begin{split} \mathbf{c} \big| \vec{B} \big| &= \big| \vec{E} \big| \quad \Rightarrow \quad \mathbf{c} \mu_0 \big| \vec{H} \big| = \big| \vec{E} \big| \qquad \Rightarrow \quad \eta_0 = \frac{|\vec{E}|}{|\vec{H}|} = \sqrt{\frac{\mu_0}{\epsilon_0}} \\ \varepsilon_0 &= 8,85418782 \times 10^{-12} \; \mathrm{F \; m^{-1}} \; (or \; A^2. \, s^4. \, kg^{-1}. \, m^{-3}) \; , \, \mu_0 = 4\pi \cdot 10^{-7} \quad H/m \; (or \; kg \cdot m \cdot A^{-2} \cdot s^{-2}) \end{split}$$

$$\eta_0 \approx 377 \sqrt{\frac{kg \cdot m \cdot A^{-2} \cdot s^{-2}}{A^2 \cdot s^4 \cdot kg^{-1} \cdot m^{-3}}} \approx 377 \ m^2 \cdot kg \cdot s^{-3} \cdot A^{-2} \approx 377 \Omega$$

Evolution of the EM field Between 2 Wave Planes (Wavefront) : the behavior of $\vec{E} \& \vec{B}$ Field in a transverse plan (x,y) along the propagation axis.

$$E_x \neq 0, E_y \neq 0, B_x, \neq 0$$
 $B_y \neq 0$, $E_z, = 0$ $B_z = 0$



Monochromatic Plane Wave Propagation

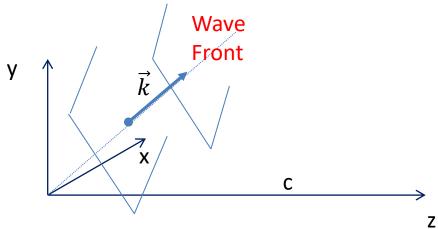
Plane Wave: The Phase of the Wave is the same at all points, on each plane \perp to the propagation wave.

1. considering of $\vec{E} \& \vec{B}$ Field in a transverse plan (x,y) along the propagation axis.

$$E_x \neq 0$$
, $E_y \neq 0$, $B_x \neq 0$, $B_y \neq 0$, $E_z = 0$, $B_z = 0$.

$$\vec{E} = \overrightarrow{E_0}\cos(kz - \omega t) \& \vec{B} = \vec{B}\cos(kz - \omega t)$$

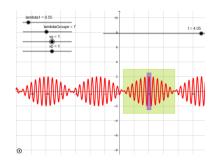
1. In a real system, the plane waves are constructed with arbitrary direction of propagation.



What is the velocity of the travelling wave?

- ☐ Phase velocity: velocity of the crest of the travelling
- ☐ Group velocity: velocity of the envelope
- > Due to superposition of travelling waves with different wavelengths
 - Geogebra activity on LMS
 - Phase and group velocity





Group velocity





Velocity of the travelling wave : Phase Velocity

Phase velocity: velocity of the crest of the travelling wave

$$\vec{E} = \overrightarrow{E_0} \cos(kz - \omega t)$$

$$\omega t - kz$$
 is constant such as $\frac{d(kz - \omega t)}{dt} = 0$

$$k\frac{dz}{dt} - \omega \frac{dt}{dt} = \frac{\omega}{c} \frac{dz}{dt} - \omega = 0 \qquad \Rightarrow \frac{dz}{dt} = \boldsymbol{v_{\varphi}} = c$$

In vacuum, $v_{\varphi} = c$

For Two travelling waves with different wavelenghts

$$\overrightarrow{E_1} = \overrightarrow{E_0} \cos(k_1 z - \omega_1 t)$$
 with $v_{\varphi} = \frac{\omega_1}{k_1}$

$$\overrightarrow{E_2} = \overrightarrow{E_0} \cos(k_2 z - \omega_2 t)$$
 with $v_{\varphi} = \frac{\omega_2}{k_2}$

III.4. Velocity of the travelling wave

Group velocity : velocity of the envelope due to the superposition of travelling waves with different wavelenghts

$$\overrightarrow{E_T} = \overrightarrow{E_1} + \overrightarrow{E_2} = \overrightarrow{E_0} \cos(k_1 z - \omega_1 t) + \overrightarrow{E_0} \cos(k_2 z - \omega_2 t)$$

$$\cos(A) + \cos(B) = 2\cos\left(\frac{A-B}{2}\right)\cos\left(\frac{A+B}{2}\right)$$

$$\overrightarrow{E_T} = \overrightarrow{E_1} + \overrightarrow{E_2} = 2\overrightarrow{E_0}\cos\left(\frac{k_1 - k_2}{2}z - \frac{\omega_1 - \omega_2}{2}t\right)\cos\left(\frac{k_1 + k_2}{2}z - \frac{\omega_1 + \omega_2}{2}t\right)$$

III.4. Velocity of the travelling wave

envelope

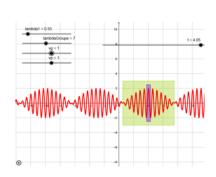
carrier

$$\overrightarrow{E_T} = \overrightarrow{E_1} + \overrightarrow{E_2} = 2\overrightarrow{E_0}\cos\left(\frac{k_1 - k_2}{2}z - \frac{\omega_1 - \omega_2}{2}t\right)\cos\left(\frac{k_1 + k_2}{2}z - \frac{\omega_1 + \omega_2}{2}t\right)$$

 $\overrightarrow{E_T} = \overrightarrow{E_1} + \overrightarrow{E_2} = 2\overrightarrow{E_0}\cos\left(\frac{k_1 - k_2}{2}z - \frac{\omega_1 - \omega_2}{2}t\right)\cos\left(\frac{k_1 + k_2}{2}z - \frac{\omega_1 + \omega_2}{2}t\right)$ Assume $k_1 \approx k_2 \approx k$ and $\omega_1 \approx \omega_2 \approx \omega$ \Longrightarrow 2 velocities : $v_{\varphi} = \frac{\omega}{k}$ and $v_g = \frac{\omega_1 - \omega_2}{k_1 - k_2} = \frac{d\omega}{dk}$

The carrier is the higher frequency signal $f_1 + f_2$

The envelop is the lower frequency signal: $f_1 - f_2$



III.4. Velocity of the travelling wave

What is the velocity of the travelling wave?

$$\overrightarrow{E_T} = \overrightarrow{E_1} + \overrightarrow{E_2} = 2\overrightarrow{E_0}\cos\left(\frac{k_1 - k_2}{2}z - \frac{\omega_1 - \omega_2}{2}t\right)\cos\left(\frac{k_1 + k_2}{2}z - \frac{\omega_1 + \omega_2}{2}t\right)$$

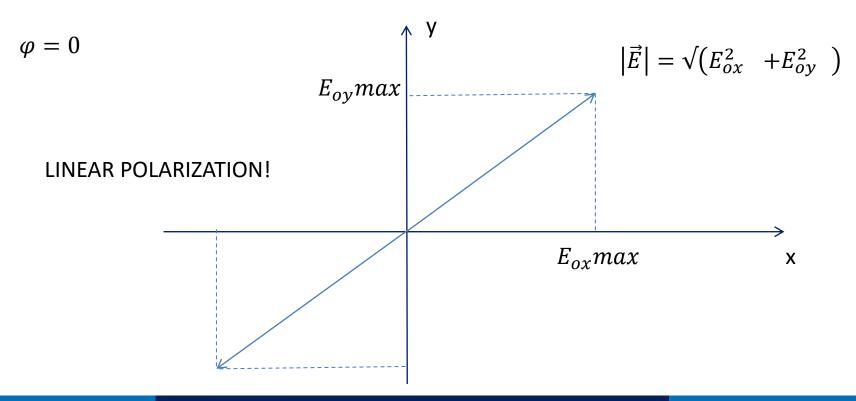
Assume $k_1 \approx k_2 \approx k$ and $\omega_1 \approx \omega_2 \approx \omega \implies$ 2 velocities : $v_{\varphi} = \frac{\omega}{k}$ and $v_g = \frac{\omega_1 - \omega_2}{k_1 - k_2} = \frac{d\omega}{dk}$

In the current case, v_{arphi} is the velocity of the carrier and v_{g} is the velocity of the modulating signal.

 $oldsymbol{v}_g$ is the velocity of the group (or packet) and corresponds to the energy flow.

III.5 Polarization of the EM Wave: Linear Polarization

$$\vec{E} = \begin{pmatrix} E_{0x} \cos(\omega t - kz) \\ E_{0y} \cos(\omega t - kz + \varphi) \\ 0 \end{pmatrix}$$



III.5 Polarization of the EM Wave: Elliptical Polarization

$$\varphi = \frac{\pi}{2} \qquad \qquad \vec{E} = E_{ox} \; \cos(\omega t - kz) \overrightarrow{e_x} + E_{oy} \; \cos\left(\omega t - kz + \frac{\pi}{2}\right) \overrightarrow{e_y}$$

$$\vec{E} = E_{ox} \; \cos(\omega t - kz) \overrightarrow{e_x} + E_{oy} \; \sin(\omega t - kz) \overrightarrow{e_y}$$
 ELLIPTICAL POLARIZATION!
$$E_{oy} max$$

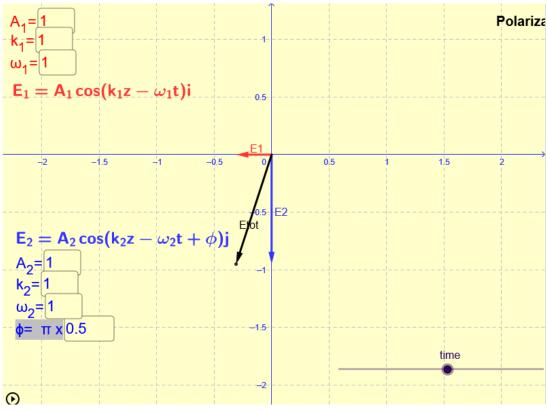
$$E_{ox} max \qquad \times$$
 CIRCULAR POLARIZATION if $E_{ox} = E_{oy}$!

III.5 Polarization of the EM Wave: Activity

 Geogebra activity on LMS: Change the value of φ and observe the animation according to

the time

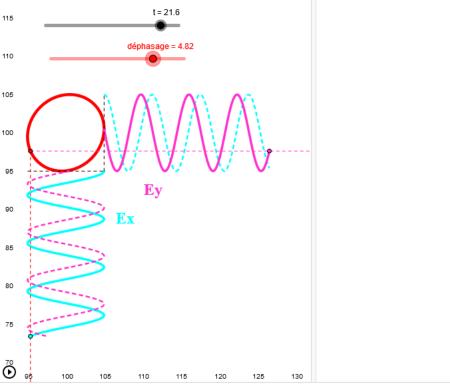




III.5 Polarization of the EM Wave: Activity

Geogebra activity, change the phase between the component X and Y of the E field to draw linear, elliptical and circular wave



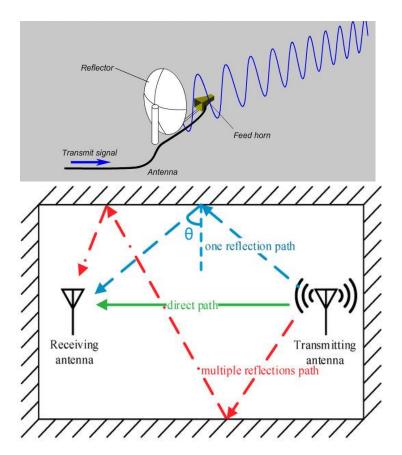


III.5 Polarization of the EM Wave: Application in Microwave

Wave's polarisation

media).

- When use LP or/and CP plane wave in RF?
 - ➤ LP (Linearly polarized)
 Emitting and receiving antenna systems fixed.
 - Receiver mobile, in multipath and fading environments and in communication with space vehicles above the earth's ionosphere (see later in dispersive



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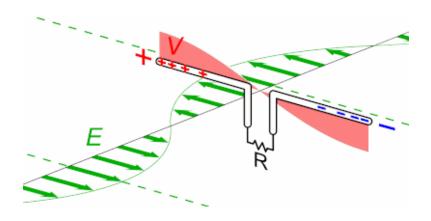
> Several polarizations: Polarization reconfigurability to improve the quality of the wireless link, optimize the link reliability

III.5 Polarization of the EM Wave: Application in Microwave

Wave's polarisation

How is created the electromagnetic wave?

By antenna (to be studied the next semester)

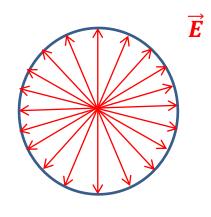


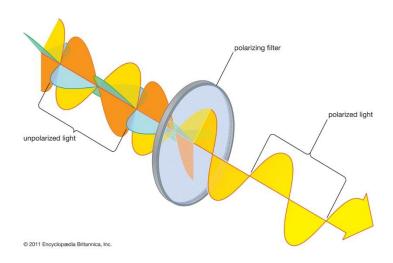
Linearly polarized antenna

https://fr.wikipedia.org/wiki/Antenne dipolaire

III.5 Polarization of the EM Wave: Application in optics

A light source is not polarized. To control the direction of the E field, polarizer (film or crystal) and polarisation controller (fiber optics or crystal).





https://www.britannica.com/science/light/Unpolarized-light

The polarization of the light is very sensitive whatever the source and the application. A control of polarization is required to maintain the properties of the optical beam but has an impact on the total field energy propagation especially in photonics (fiber optics communication, quantum optics, optical sensing).

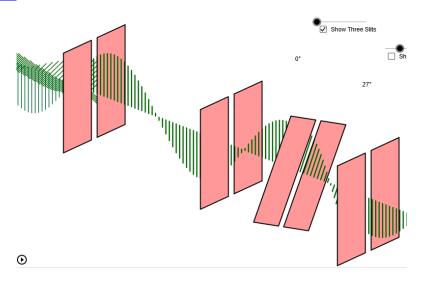
III.5 Polarization of the EM Wave: Activity

What is the polarization of the light if the second polarizer is horizontal?

Check you answer with geogebra activit



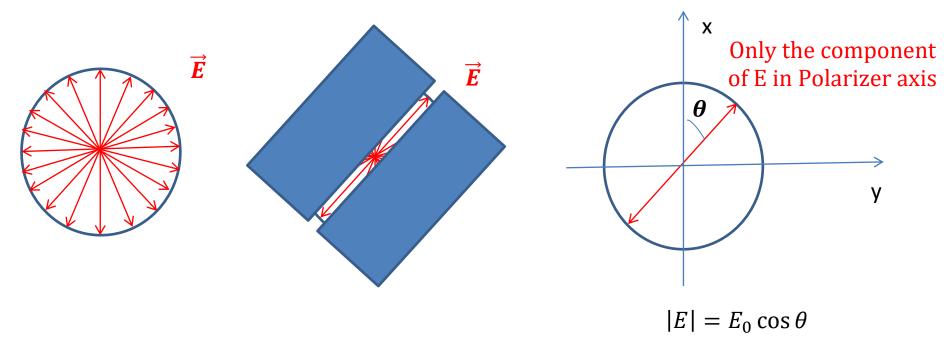
Fiber optic polarization controller (twist the fiber)





III.5 Polarization of the EM Wave: Application in optics

The polarized optical field follows the Malus Law:



The intensity of the field is : $I = |E|^2$

From the Malus Law: $I \sim E_0^2 cos^2 \theta$

III. Conclusion Plane Wave in Vacuum

From Maxwell Equation in vacuum (without charge and current), Wave Equation

$$\Delta \vec{E} = \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\Delta \vec{B} = \frac{1}{c^2} \frac{\partial^2 \vec{B}}{\partial t^2}$$

IN VACUUM

The wave number is $k = \frac{\omega}{c}$

From Maxwell Equation : $\vec{E} \perp \vec{k}$, $\vec{B} \perp \vec{k}$

$$\vec{E} \perp \vec{B}$$

$$\mathsf{c}|\vec{B}| = |\vec{E}|$$

Wave impedance in vacuum

$$\eta_0 = \frac{|\vec{E}|}{|\vec{H}|} = \sqrt{\frac{\mu_0}{\epsilon_0}} = 377\Omega$$

Harmonic wave

Wavelength
$$\lambda = \frac{2\pi}{k} = cT$$

Period
$$T = \frac{2\pi}{\omega}$$

Frequency
$$f = \frac{\omega}{2\pi}$$

Velocity

Light speed : $c \approx 3 \cdot 10^8 \text{m/s}$

Phase velocity
$$v_{\varphi} = \frac{\omega}{k}$$

Group velocity $v_g = \frac{d\omega}{dk}$

Polarization

Linear polarization: $\varphi=0~or~\pi$

Circular polarization:

$$\varphi = \pm \frac{\pi}{2}$$
 and $E_{ox} = E_{oy}$

Elliptical polarization: each other

III. Conclusion plane Wave in Vacuum

Polarization

Malus Law, Intensity $I \sim E_0^2 cos^2 \theta$



How can we propagate energy with an electromagnetic Field?

Part IV