

Theoretical and practical investigation into the use of a bio-inspired “click” mechanism for the flight motor of a micro air vehicle

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Abstract

Recently, flapping wing micro air vehicles have received great attention with the drive to make smaller and smaller devices. This paper describes a theoretical investigation and subsequent practical implementation of a specific type of flight motor structure for this type of micro air vehicle that uses a “click” mechanism to improve mechanical efficiency. Diptera, which may use the mechanism, are the inspiration for this work. It builds on previous research into the “click” mechanism, which has been studied both from the biological and engineering points of view. It is difficult to capture the important fine details using a simple analytical model; hence, a multi-body dynamic software is used to model the device and to aid the design of a large-scale prototype. Force–deflection curves of the structure and the displacement response are obtained numerically and experimentally. The experimental and numerical results compare reasonably well, enabling the model to be used for further development and potential miniaturization of the flight motor structure. In a practical device, asymmetry occurs in the up- and down-stroke. The effects of this asymmetry are compared with previous results from analytical models. It is found that asymmetry offers a marginal improvement.

Keywords

Flapping wing, micro air vehicle, asymmetric, multi-body dynamic analysis, energy ratio

Date received: 10 June 2016; accepted: 27 January 2017

Introduction

Research into micro air vehicles (MAVs) has received great attention recently, because of the drive to make smaller and smaller devices. It has been an important area in biomimetic research for about two decades.^{1–3} Recently, some MAVs have been reduced to the scale of a fly.⁴ By comparing two different flight motor mechanisms, such as the flapping wing or the rotary wing, it has been found that the rotary wing has a higher power density (power to weight ratio).⁵ However, for a very small scale MAV with high maneuverability, the flapping wing system may still be attractive.^{5,6} A high degree of understanding of both aerodynamic theory and the dynamics of the flight motor mechanism are necessary to improve the performance of flapping wing MAV.

In the search for a suitable flight motor for a flapping wing MAV, some researchers have turned to

nature for inspiration. Many models for the flight mechanism of diptera have been proposed, one of which is the “click” mechanism, an overview of which is given in Overview of a simple model of the “click” mechanism section. The “click” mechanism was first

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proposed in Boettiger and Furshpan,⁷ Pringle⁸ and Thomson and Thompson⁹ derived a mathematical model for the “click” mechanism. It is called “click” mechanism because when the excitation force on the wing structure of diptera reaches a critical value, the wings suddenly “click” from an up to a down position at a high velocity. This interesting phenomenon is due to the bistable characteristic of the structure, and in engineering parlance is often called a “snap-through”. Alexander and Bennet-Clark¹⁰ analyzed the benefit of the “click” mechanism in the storage of elastic strain energy. Miyan and Ewing^{11,12} expressed doubt about the existence of the “click” mechanism using morphological analysis. However, a revised model by Miyan and Ewing¹³ shows a weak version of “click” mechanism in an inherently asymmetric structure. The new model attracted the interest of researchers and was reported by Bennet-Clark¹⁴ and Gronenberg.¹⁵ Ennos analyzed the distortion of the whole thorax and gave a novel linear model.¹⁶ The new model given in Miyan and Ewing^{11,12} was also criticized by Pfu who modified the “click” mechanism with a gear change structure.¹⁷ Recently, Pfu built a mechanical model with both “click” and gear mechanisms to test the plausibility of the gear change mechanism.¹⁸ The gear mechanism^{18,19} was put together with a clutch mechanism recently by Deora et al.²⁰

Based on the flight mechanism described in Thomson and Thompson,⁹ Brennan et al.²¹ and Tang and Brennan²² analyzed a simple mechanical model of the “click” mechanism with both linear and quadratic damping, to determine if it gives any mechanical advantage. Motivated by this work, Chin and Lau^{23,24} and Lau et al.²⁵ built a flapping-wing MAV, incorporating such a mechanism, and recently, Harne and Wang²⁶ used a compressed beam mechanism in a “click” mechanism for an MAV. Kok et al.²⁷ used a quasi-steady blade-element method to analyze the asymmetrical flapping on the aerodynamic efficiency of half-span wing of the MAV with “click” mechanism which was built by Chin and Lau,^{23,24} observing greater lift generation when the wing was excited below the natural frequency of the system.

The aim of this paper is to build on previous theoretical work, investigating the practical implementation of “click” mechanism, and to determine whether it offers any significant improvement in the mechanical efficiency in a practical device. A large-scale device is studied (because of resource limitations), but the practical device and the accompanying numerical model can be used in further studies of smaller scale devices. A practical issue that has not been addressed hitherto, and is of particular interest in this paper, is that a flight motor structure with a “click” mechanism may have unavoidable geometrical asymmetry between the

up- and down-stroke. One element of this research that sets it apart from previous work is that the flight motor structure is studied both theoretically and experimentally, resulting in a validated model that can be used for further development work and research.

Overview of a simple model of the “click” mechanism

The flight motor described in Thomson and Thompson⁹ and the related simplified mechanical model described in Brennan et al.²¹ and Tang and Brennan²² are shown in Figure 1. Figure 1(a) to (c) shows the three successive positions of the wing articulation during the course of a beat. Figure 1(d) to (f) shows the corresponding positions of the simplified model, in which the mass m at point C is the equivalent mass of the moving part of the whole mechanism given in Figure 1(a) to (c); BE and DF are the mass-less flexible links which represent the pleural apophysis (pa) and notum (n), respectively; BAC and CD are the mass-less rigid links which represent axillary sclerite 2 (ax2) and parascutum (p) respectively, and WA is the simple mass-less model of the wing. The harmonic excitation force is $f(t) = F \cos \omega t$, where F is the amplitude and ω is the angular frequency. It is applied at point C, while the input for the insect model is applied to the scuteller lever (sl).

The displacement of the mass from the position in which BCD is horizontal is given by y , and the lateral displacement of the cantilevers BE and DF is denoted by x . The simplified mechanism is assumed to be symmetric as shown by Brennan et al.²¹ and Tang and Brennan,²² which means that $BC = CD = l$ and the distance between vertical undeformed link BE and point C (DF and C) is b . The equation of motion for the simplified model is then given by

$$m\ddot{y} + f(\dot{y}) + 2k\left(-1 + \frac{b}{\sqrt{l^2 - y^2}}\right)y = F \cos \omega t \quad (1)$$

where k is the bending stiffness of each of the two mass-less flexible links BE and DF, $f(\dot{y})$ is the damping force induced by aerodynamic and mechanical damping and the overdots denote the differentiation with respect to time. The derivation of the non-linear restoring force is given in Brennan et al.²¹ and is presented in Appendix 1 for ease of reference.

Recent research on a flapping flight mechanism that utilizes the “click” mechanism has shown that, in certain circumstances, it has some advantages compared to that which uses a linear mechanism.^{21–27} This is primarily for small devices where the natural frequency of the system is too high for resonance to be utilized as the

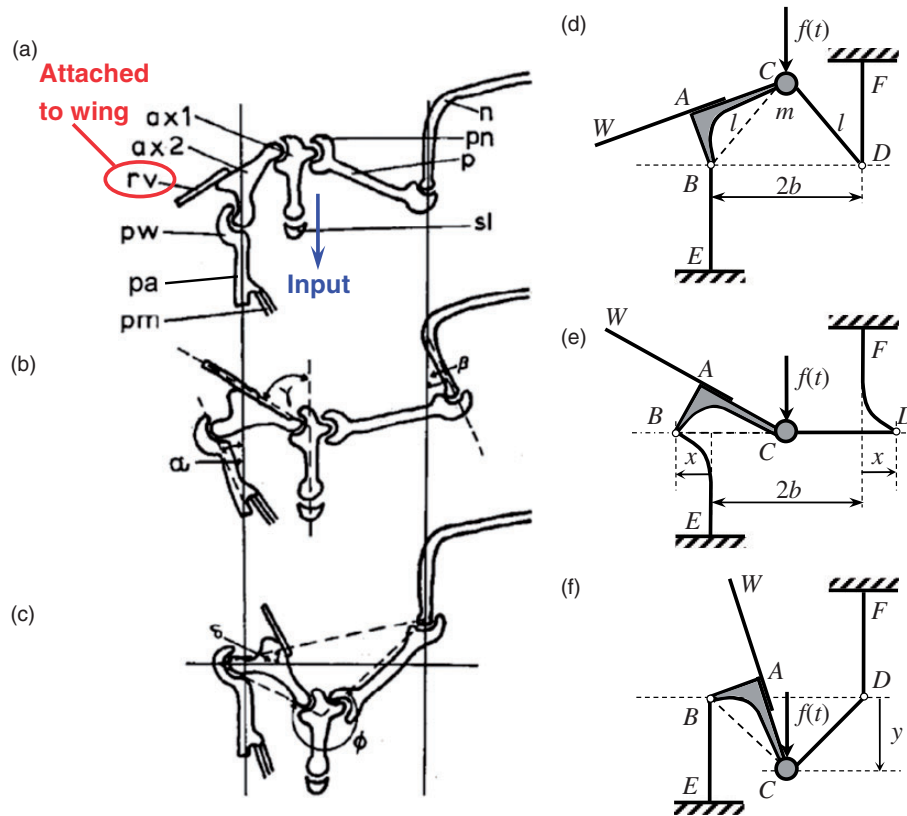


Figure 1. Diagram showing the flight motor (a–c) and physical model (d–f) of an insect and three successive positions of the wing articulation during the course of a beat, from the down position (a or d) to the up position (c or f). In (a–c), $ax1$, $ax2$ – axillary sclerites 1 and 2; n – notum; p – parascutum; pm – pleurosternal muscle; pn – anterior notal process; pw – wing process; rv – base of radial vein; sl – section through scutellar lever ((a–c) After Thomson and Thompson⁹). pa – pleural apophysis (not shown in (Pringle⁸)). In (d–f), m : equivalent mass; BE , DF : mass-less flexible link; ABC , CD : mass-less rigid link; WA : wing.

wing flapping frequency. In such a case, it is necessary to flap the wings at a frequency well below the natural frequency, and in this case, the “click” mechanism has advantages. This was demonstrated in Brennan et al.²¹ and Tang and Brennan²² by showing that, for a highly damped system, the peak kinetic energy of the wing root normalized by power input to the mechanism is larger than that of the linear system, as shown in Figure 2(a)^{21,22} for a system in which the excitation frequency is 1/10 of the natural frequency (the definition of the ratio of peak kinetic energy to work input is given in Appendix 2). This is because the peak velocity of the wings is increased by the “click” mechanism compared to the linear system for a given input frequency, providing a larger lift force. To illustrate the way in which the “click” mechanism changes the velocity of the wings, the normalized velocity is plotted in Figure 2(b) as a function of non-dimensional time for $\hat{F} = 1.0$ (see Appendix 2 and Brennan et al.²¹ and Tang and Brennan²²).

Model of the flight motor structure with a “click” mechanism

An example of a flight motor structure with a “click” mechanism is depicted in Figure 3; a photograph is shown in Figure 3(a) and a schematic in Figure 3(b). It should be noted that this system was not intended to fly. It was a first attempt to study the practicalities of designing and testing a structure, which incorporates the “click” mechanism, and therefore it is much larger than a practical flapping wing MAV would be. The frame of the flight motor structure (pleural apophysis) is made from 0.5 mm thick carbon fiber-reinforced polymer (CFRP). The hinge for the flight motor wing, marked hinge I, consists of two 55 μm thick layers of polyimide film tape, which are bonded together, and hinge II consists of a single layer. The wings are made from polyethylene terephthalate polyester (PET) film and are 12.5 μm thick. The two pieces of CFRP, which are used for the tergum, are connected

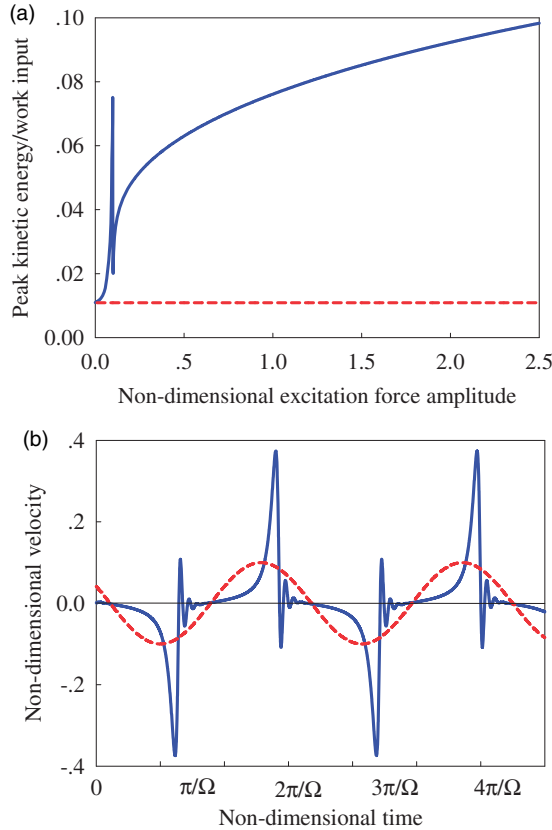


Figure 2. Comparison between the linear system (dashed line) and the system with the “click” mechanism (solid line). (a) Ratio of peak kinetic energy to work input as a function of excitation force amplitude, (b) non-dimensional velocity when $\hat{F} = 1.0$.

by one layer of polyimide film tape. A polyimide cylinder, which acts as the tergum hinge, is fixed on the top of the film. The driving mechanism is a DC motor, which can develop a maximum torque of 0.92 mNm, and connects to the tergum hinge through a crank and connecting rod mechanism. The dimensions and the masses of the flight motor structure are given in Figure 3(c) and Table 1, respectively, and the approximated simple mechanical model is given in Figure 3(d).

Although it is possible to derive a detailed analytical model of the flight motor structure, it would be highly complicated. A more appropriate approach is to develop a multi-body dynamic model using commercial software to predict and study the behavior of the system, and here ADAMS (MSC Software Corp, USA) is used. The pleural apophysis, pleura process, tergum, tergum hinge, and wing frames are modelled as rigid bodies with a density of 1758 kg/m^3 . The wing film is also modelled as a rigid body with a density of 1400 kg/m^3 .

The two pairs of wing hinges are modelled as flexible elements with a Young's modulus of 2.5 GPa and a density of 1400 kg/m^3 . Wing hinge I, which consists of

two bonded layers of polyimide film tape, is modelled as two parallel polyimide films, each $55 \mu\text{m}$ thick. Wing hinge II, which behaves as a rigid revolute joint, is modelled by setting the Young's modulus of the material to be one eighth of the actual Young's modulus. To simulate the elastic characteristics of the flexible joints, the wing hinge elements are analyzed using ANSYS (ANSYS, Inc., USA). The first two natural frequencies and mode shapes of the wing hinge elements are saved in an MNF file and imported into ADAMS. To simulate the aerodynamic force, a damping element is created in ADAMS and located vertically between the tergum hinge and the pleural apophysis. Using the multi-body model, the restoring force, the dynamic response, and the energy consumption can be determined. The natural frequency of the whole model, estimated using ADAMS, was found to be about 12.5 Hz. This is the primary frequency of the whole structure when the moving parts of the structure oscillate around one of the stable equilibrium points as shown in Figure 1(d) or (f).

Force–deflection characteristic

The force–deflection characteristic of the flight motor structure depicted in Figure 3(a), as seen by the electric motor, was measured using an Instron Micro Tester 5848 (Instron, Norwood, MA, USA) with a 5 N load cell and 20 Hz sample rate. The test-rig is shown in Figure 4. As shown in the inset, the pleural apophysis was fixed, and the tergum hinge was connected by a metal rod to the moving part of the Instron Micro Tester, which was excited at a rate of 18 mm/min. During the test, the motor was replaced by a solid plastic block to maintain the operational shape of the flight motor structure frame. The wings were also removed. As well as measuring the force applied to the tergum and the resulting displacement, a video of the structural deformation was taken using an IXUS 210 Canon camera as shown in Figure 4. Because the test process was quasi-static, the camera was easily synchronized with the load step of the Instron Micro Tester, and the deformation of structure, force load and displacement of the structure could be captured at the same time.

The measured applied force on the tergum and the displacement of the tergum are plotted together to give the force–deflection curve shown in Figure 5 as a solid blue line. Also shown in this figure are some pictures (video-stills) of the flight motor structure, at positions labelled (A) to (G). Superimposed on the measured force–deflection curve is the predicted force–deflection characteristic of the model of the structure shown in Figure 3(b), and the corresponding deformation of the structure, both calculated using the multi-body

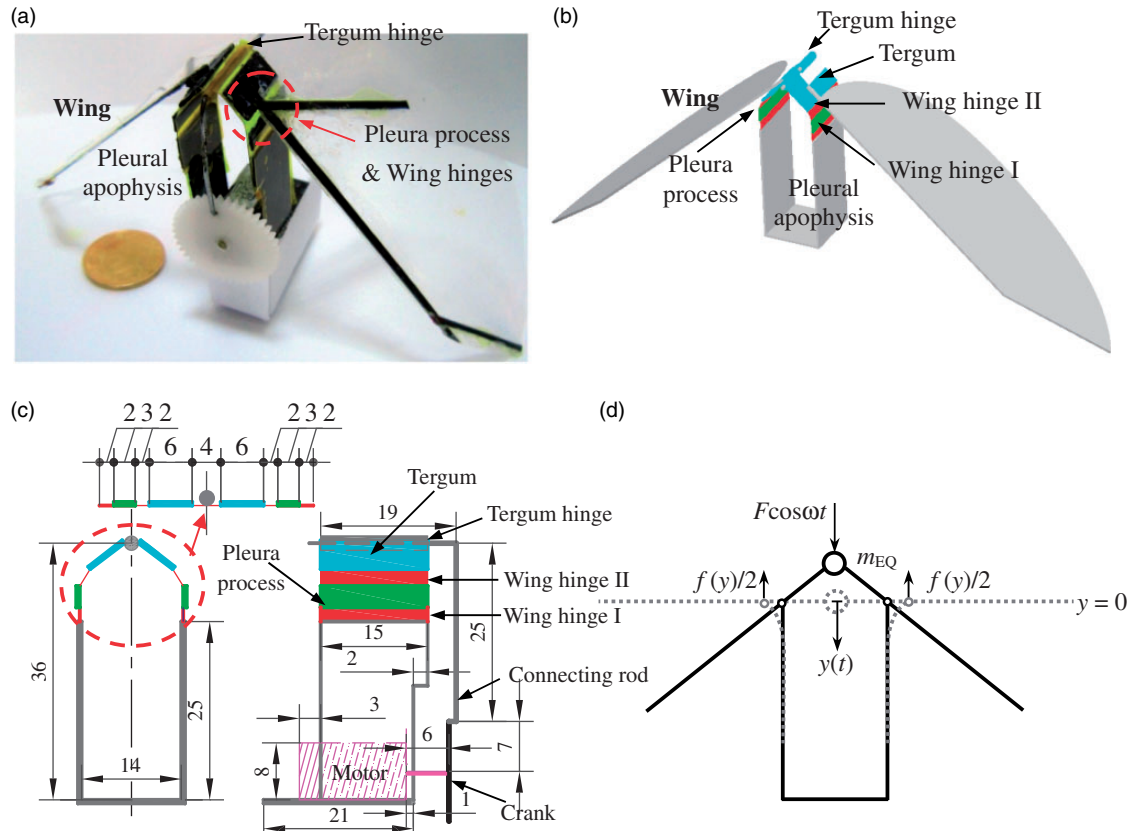


Figure 3. Flight motor structure with click mechanism (a) photograph with 50 Chinese cents, (b) multi-body dynamic model, (c) dimensions of the structure (units in mm), (d) simple lumped parameter model (dashed line, unstable equilibrium position), and $f(y)$ is the nonlinear restoring force given by the two vertical uniform cantilever beams.

Table 1. Masses of the structure (see Figure 3).

Components	Mass (g)
Wings	0.45
Pleura process and wing hinges	0.35
Tergum hinge	0.07
Pleural apophysis	0.97
Crank and connecting rod	1.03
Motor	4.00
Total	6.87

dynamic model. To obtain these quantities, the tergum hinge joint of the model was excited with quasi-static displacement excitation.

As shown in Figure 5, it can be seen that the force–deflection characteristic obtained experimentally and that obtained from the multi-body model, agree reasonably well. The main difference between the two is the hysteretic behavior that is clearly evident in the

experimental structure due to the structural damping in the hinges, which is not seen in the simulations, as this was not incorporated in the multi-body model. The asymmetric behavior of the flight motor structure is also clearly evident, with the structure being much stiffer on the up-stroke compared with the down-stroke.

Displacement response

The setup for dynamic displacement measurement is shown in Figure 3(a). The flight motor structure was driven by the electric motor with a rotation speed of 162 r/min (2.7 Hz). As in the previous experiment, the camera was used to capture the deformation of flight motor structure. To extract the displacement response of the tergum from the measured images, the computer program Tracker was used.²⁸ The displacement is shown in Figure 6 together with the response from the multi-body dynamic model. In the model, the damping coefficient was set to 1.5 Ns/m and the excitation force amplitude to 0.2 N, which are used to simulate the damping effect and the excitation load on the

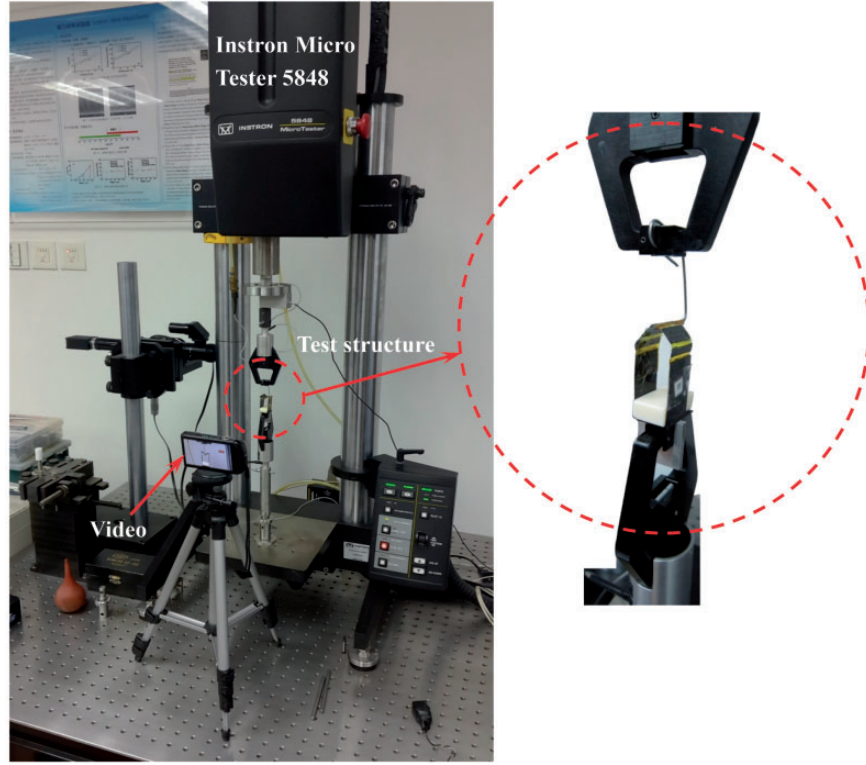


Figure 4. Measurement of the force–deflection characteristic of the structure.

moving parts of the prototype structure respectively. It can be seen that the numerical results and the measured response compare reasonably well. This means that the model can be used to investigate some of the features of the MAV structure with “click” mechanism, and this is done in the following sections.

Effects of asymmetry

To investigate the effects of the asymmetry of the device, the velocity of the tergum was first calculated using the multi-body dynamic model and is shown in Figure 7(a). Further analysis is possible by developing a single-degree-of-freedom (SDOF) approximate model of the system, which is given by

$$m_{EQ}\ddot{y} + c\dot{y} + f(y) = F\cos\omega t \quad (2)$$

in which the restoring force is approximated by a sixth order polynomial given by $f(y) = \sum_{n=0}^6 k_n y^n$, where the stiffness coefficients are determined by curve fitting to the force–deflection characteristic from the multi-body model. They were found to be $k_0 = -0.0033$ N, $k_1 = -27.7$ N/m, $k_2 = -4.1 \times 10^3$ N/m², $k_3 = 5.9 \times 10^5$ N/m³, $k_4 = -3.1 \times 10^8$ N/m⁴, $k_5 = 5.1 \times 10^9$ N/m⁵, $k_6 = 3.7 \times 10^{12}$ N/m⁶. The damping coefficient c is set to be the same as in the multi-body model and has a

value of 1.5 Ns/m, and the mass has a value of 10 g, which is estimated using the equivalence of the total kinetic energies of the multi-body model given by²⁹

$$\frac{1}{2}m_{EQ}\dot{y}^2 = \sum_{i=1}^N \frac{1}{2}m_i\dot{y}_i^2 + \sum_{j=1}^M \frac{1}{2}J_j\dot{\theta}_j^2 \quad (3)$$

where $m_i\dot{y}_i^2/2$ is the kinetic energy of the i th ($i=1, \dots, N$) moving body of the multi-body flight motor structure model, and $J_j\dot{\theta}_j^2/2$ is the kinetic energy of the j th ($j=1, \dots, M$) rotating body of the multi-body flight motor structure model. The root mean square error between the two sides of equation (3) is used to minimize the error in the estimated equivalent mass.

The velocity response of the simplified SDOF model is shown in Figure 7(b). By comparing Figure 7(a) and (b), it can be seen that the simplified SDOF model captures the general behavior of the multi-body model. Note, however, that the multi-body model was necessary to capture the fine details of the system, from which the lumped mass and stiffness characteristics were determined. By retaining only the odd order terms in the approximate expression for the restoring force, a symmetric system can be obtained. The velocity of this system is shown in Figure 7(c). By comparing Figure 7(b) and (c), it can be seen that the asymmetry

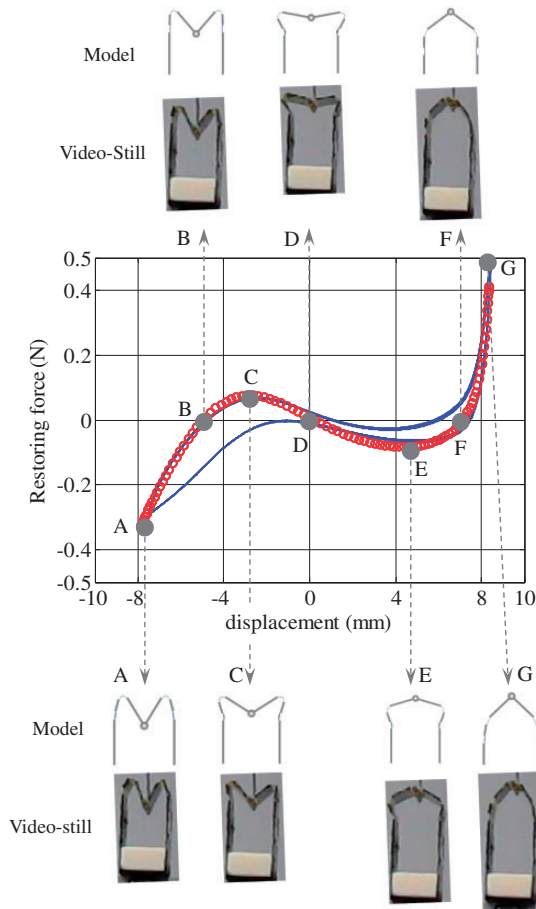


Figure 5. Measured and calculated restoring force. Circle \circ , multi-body model; solid line —, experimental results. (a) and (g), tergum down and up points; (b) and (f), stable equilibrium points; (c) and (e), local maximum force points; (d), unstable equilibrium point.

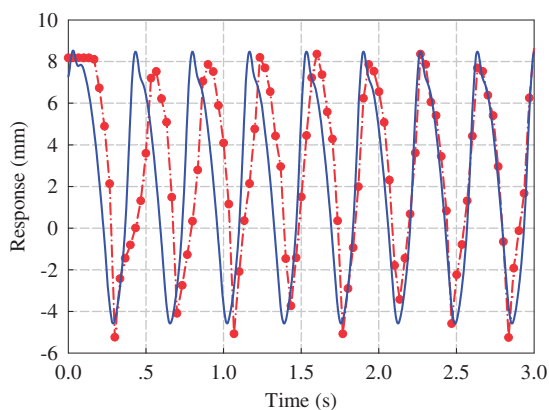


Figure 6. Experimental and simulated displacement response of the tergum. Simulations from the multi-body model (solid line) and experimental results using Tracker (dashed-dotted line with thick dot).

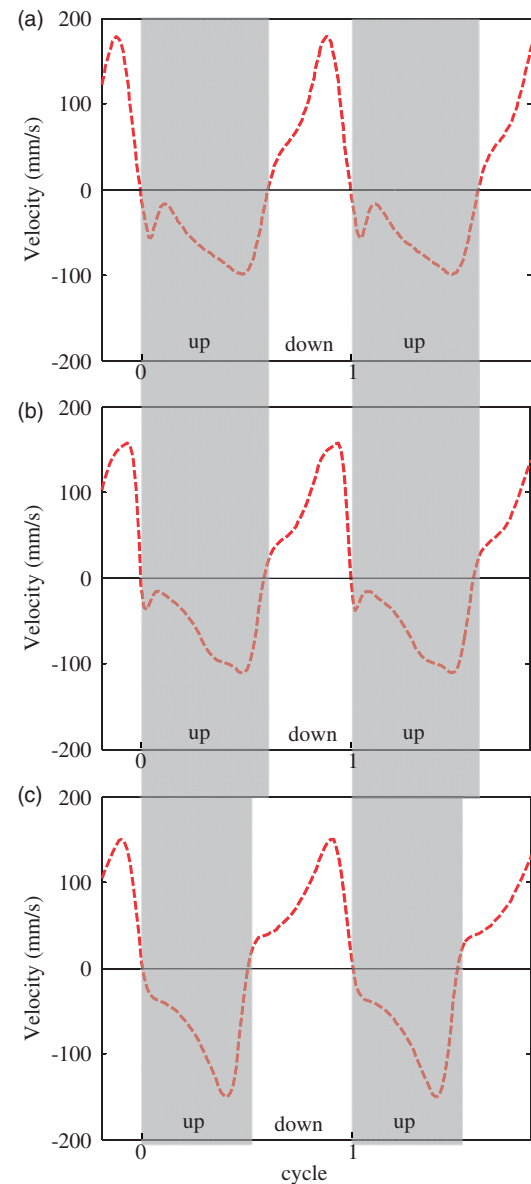


Figure 7. Time history response for velocity of (a) the tergum of multi-body model, (b) simple model given in equation (2), and (c) simple symmetric lumped model.

has the effect of inducing a larger peak velocity during the down-stroke. This may be considered advantageous as it may generate a larger lift force.

Energy considerations

As the simulations and the experimental results for both the stiffness characteristics and the dynamic response of the flight motor structure compare reasonably well, the multi-body dynamic model and simplified lumped parameter model can be used to further investigate some additional characteristics of the flight motor structure. Of particular interest is the ratio the

energy ratio of peak kinetic energy of the wings and the work input per cycle which is given by³⁰

$$\begin{aligned}\hat{E} &= \frac{\text{Peak kinetic energy}}{\text{Work input over one cycle}} \\ &= \frac{m(\dot{y}_{peak})^2/2}{\int_0^{2\pi/\omega} F \cos \omega t \dot{y} dt}\end{aligned}\quad (4)$$

This can be used as a metric to characterize the efficacy of the structure. For the simplified lumped parameter model shown in Figure 1(d), the peak kinetic energy $m(\dot{y}_{peak})^2/2$ can be obtained by taking the product of the lumped mass and the square of the maximum velocity of the mass. For the multi-body model, the peak kinetic energy can be obtained by choosing the maximum value of the sum of the total kinetic energy of all the moving components from the time histories. The work input over one cycle can be obtained by integrating the product of the excitation harmonic force and tergum velocity. This is calculated for a harmonic excitation force amplitude ranging from 0.01 to 0.3 N, by determining steady-state time histories from numerical integration of the equations of motion. Equation (4) is plotted in Figure 8, for four cases; the multi-body model, the approximate SDOF model, the symmetric approximate SDOF model and the linear system, which has a natural frequency, equivalent mass, and damping coefficient similar to those of the multi-body model.

It can be seen that all the systems which have a “click” mechanism perform better than the linear system. It can also be seen that the symmetric system and the asymmetric system behave in a qualitatively

similar way, but the asymmetry, which is hard to avoid in a practical flight motor design and observed in the wing motion of insects, offers a marginal improvement.

Conclusions

A bio-inspired flight motor structure with a “click” mechanism, which was realized using polyimide film and CFRP, has been investigated numerically and experimentally in this paper. Asymmetry, which occurs in the up- and down-stroke in the practical flight motor structure studied, has also been considered. The numerical analysis was carried out using both a multi-body dynamic model and a simple lumped parameter model. The experimental and numerical results for the restoring force and displacement response compared well, meaning that the multi-body dynamic and lump parameter models can capture the dynamic characteristics of the asymmetric flight motor structure reasonably well. The energy ratio of the kinetic energy to work input over one period of excitation was subsequently obtained from the numerical models. It was found that the symmetric and asymmetric structures have a similar performance when the energy ratio of the kinetic energy to work input over one period of excitation is chosen as the criterion, with the asymmetric structure being marginally better. Although a large-scale device has been studied in this paper (because of resource limitations), the advantages of the “click” mechanism have been demonstrated practically and the numerical model developed can be used in further studies of smaller scale devices.

Acknowledgements

The authors thank Dr. Yao-Wei Chin, Dr. Rong Yang and Dr. Yihui Feng who helped during the experiments.

Declaration of conflicting interests

The author(s) declared no potential conflicts of interest with respect to the research, authorship, and/or publication of this article.

Funding

The author(s) disclosed receipt of the following financial support for the research, authorship, and/or publication of this article: The first author wishes to acknowledge the financial support from the National Natural Science Foundation of China (Grant Nos. 11202048 and 11672058), the Fundamental Research Funds for the Central Universities of China (Grant No. DUT15QY30).

References

1. Vincent JFV, Bogatyreva OA, Bogatyrev NR, et al. Biomimetics: its practice and theory. *J R Soc Interface* 2006; 3: 471–482.

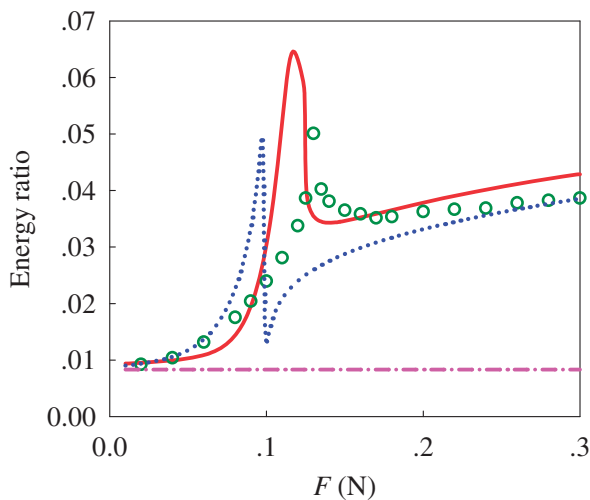


Figure 8. Energy ratio comparison at $\omega = 1.25$ Hz. Circle \circ , multi-body model; solid line —, simple model given in equation (2); dotted line \cdots , simple symmetric lumped model; dashed-dotted line $- \cdot -$, linear system.

point (Figure 1(e)) and bottom stable equilibrium point (Figure 1(f)). The mass-less flexible link BC (effective length of ABC) and CD have an angle of θ with respect to the horizontal line. The free-body diagram of Figure 9(a) is given in Figure 9(b), which shows the internal forces within the simplified physical model. The equations in this Appendix are derived from the left part of the free-body diagram, assuming the positive directions of the displacement and rotation angle shown in Figure 9(b). Due to symmetry, the right part can be analyzed in a similar way. The force relationship between f_x and f_y can be written as $f_y = -f_x \tan \theta$, and because of Hooke's law $f_x = kx$, where k is the bending stiffness of the mass-less flexible link DF.

The geometric relationships between the rigid and flexible links can be written as $x = l \cos \theta - b$ and $y = l \sin \theta$, which results in $x = -b + (l^2 - y^2)^{1/2}$ and $\tan \theta = y/(x + b)$.

Combining the geometric and force relationships gives

$$f_y = k \left[-1 + \frac{b}{\sqrt{l^2 - y^2}} \right] y \quad (5)$$

which shows that the force-deflection characteristic is nonlinear and is a function of the geometry of the device.

Appendix 2. Non-dimensional equation of motion for the simple “click” mechanism model

In this Appendix, the non-dimensional model of the simple model with a click mechanism is described.

Assuming that the damping force is given by $f(\dot{y}) = c\dot{y}$, where c is a linear viscous damping coefficient, equation (1) can be written in non-dimensional form as²²

$$u'' + \gamma u' - \alpha u + \alpha \beta \frac{u}{\sqrt{1 - u^2}} = \hat{F} \cos \Omega \tau \quad (6)$$

where $u = y/l$, $\gamma = c/m\omega_0$, $\omega_0^2 = 2k(1/\beta^2 - 1)/m$, $\beta = b/l$, $\hat{F} = F/ml\omega_0^2$, $\tau = \omega_0 t$, $\Omega = \omega/\omega_0$, $\alpha = \beta^2/(1 - \beta^2)$, and $(\bullet)'$ denotes differentiation with respect to non-dimensional time τ . Note that ω_0 is the natural frequency for small oscillations when the system is at either of its static equilibrium positions as shown in Figure 1(d) or (f).

The non-dimensional equation for the linear system is given by²²

$$u'' + \gamma u' + u = \hat{F} \cos \Omega \tau \quad (7)$$

When equations (6) and (7) are solved numerically with $\beta = 0.9$ and $\gamma = 1.46$, as given in Tang and Brennan,²² for different amplitudes of non-dimensional force \hat{F} , the ratio of non-dimensional peak kinetic energy to the work input over one cycle is given by

$$\hat{E} = \frac{(u'_{peak})^2/2}{\int_0^{2\pi/\Omega} \hat{F} \cos \Omega \tau u' d\tau} \quad (8)$$

the results of which are plotted in Figure 2(a).