Cohort Project 2021, Team 4 Simulating quantum advantage with trapped ions

1 Introduction

In 2019 it was announced that the superconducting qubit based quantum computer Sycamore from Google outperformed state-of-the-art classical computers in sampling a pseudo-random quantum circuit probability distribution [1]. In this work we perform the same calculation using trapped ion gates simulated on a classical computer. We explore how the probability distribution depends on the depth (number of layers or gates) and width (number of qubits) of the circuit and how sensitive the resulting speckled pattern is to perturbations (errors in the circuit). Next we study the convergence of a perfect quantum random-circuit to a Porter-Thomas distribution and deviations from these distribution as the 2-qubit gate errors increase.

2 Task 1: Probability distribution of a quantum randomcircuit

We first simulate the probability distribution if the gates were to be error-free. Even if that would be the case we don't expect the probability distribution to be flat, i.e. every bit-string x to appear with the same frequency if the measurement is repeated a large number of times S, where S is the number of samples. The reason is that due to quantum interference some bit-strings occur much more often than others, forming a speckled pattern. The computation of such quantum distribution grows exponentially with the number of qubits on a classical computer.

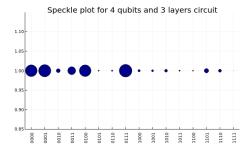
We apply 1 and 2-qubit random gates as in ref. [1] to entangle the qubits,

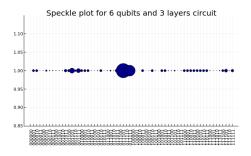
$$R(\theta, \varphi) = \begin{pmatrix} \cos\frac{\theta}{2} & -ie^{i\varphi}\sin\frac{\theta}{2} \\ -ie\sin\frac{\theta}{2} & \cos\frac{\theta}{2} \end{pmatrix}$$
 (1)

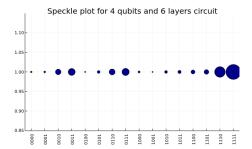
$$M(\theta) = \begin{pmatrix} \cos\Theta & 0 & 0 & -i\sin\Theta \\ 0 & \cos\Theta & -i\sin\Theta & 0 \\ 0 & -i\sin\Theta & \cos\Theta & 0 \\ -i\sin\Theta & 0 & 0 & \cos\Theta \end{pmatrix}$$
 (2)

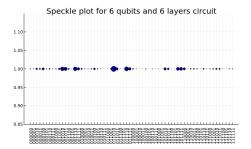
where $R(\theta, \varphi)$ and $M(\theta)$ are gates that can be build with trapped ions.

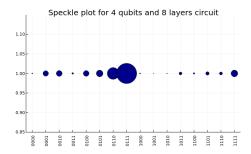
The probability distribution for different number N=4,6 of qubits and depths 3,6,8 is shown in Fig.1 and exhibits a speckled pattern as expected. Each point represents a state in Hilbert space, i.e. one of the possible outcomes of the measurement (bit-string x). There are 2^N states and the "interference" pattern changes every time we do the sampling.











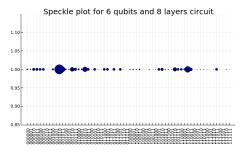


Figure 1: Probability distribution $P(x)=|\langle x|\Psi\rangle|^2$ for a 4-qubit (left column) and a 6-qubit (right column) quantum random-circuit of depths =3, 6, 8. The probability is computed from S=1024 measurements (shots)

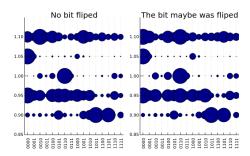


Figure 2: Probability distributions from 4-qubit depth = 4 random-circuit over S=25 measurements. The left plot shows speckled patterns for the random circuits without the bit flipped, sampled in different pseudo-experiments, while the right plot displays the speckled patterns using the same parameters as on the left, however now the circuit has a 50 % chance of one spin being flipped at a random location.

3 Task 2: Effects of a one spin-flip error at random location

A large sensitivity to perturbations (i.e. to errors in the application of the gates) is an indication of the emergence of chaos. Here we explore the effects of introducing a perturbation in the form of a spin-flip error, which we model using,

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \tag{3}$$

4 Task 3: Convergence to Potter-Thomas distribution

As the depth of the circuit increases we expect the probability of a given bit-string x to converge to the exponential distribution,

$$p \to 2^N e^{-2Np} \tag{4}$$

In Fig. 3 we study the convergence of the cumulative distribution p as we increase the depth of the circuit. In the limit of large depth the distribution should converge to a Porter-Thomas distribution [2].

5 Task 4: Cross-entropy benchmarking fidelity

An efficient way to assess the quality of the probability distribution is to compute its cross-entropy benchmarking fidelity \mathcal{F}_{XEB} ,

$$\mathcal{F}_{XEB} = 2^N \langle P \rangle - 1 = \frac{2^N}{S} \sum_{i=1}^S P(x_i) - 1$$
 (5)

A uniform distribution where every bit-string x occurs with the same frequency yields $\mathcal{F}_{XEB} = 0$ whereas an error-free quantum random-circuit distribution yields $\mathcal{F}_{XEB} = 1$. A sampling performed on a NISQ quantum computer is therefore expected to yield $0 > \mathcal{F}_{XEB} < 1$. Here we investigate the effect of a sistematic error in the circuit, which we model by perturbing the angle of each 2-qubit gate by a fixed amount $\Delta\Theta$.

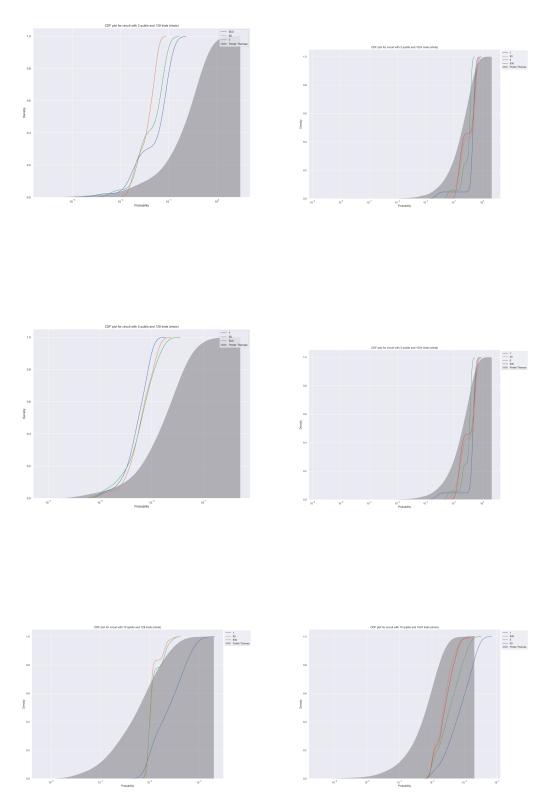


Figure 3: Cumulative distribution for a 2-qubit (upper panel) quantum random-circuit of depths =1,5,50,500 and S=128,1024 (left and right columns respectively). Middle panel: same for a 5-qubit circuit and lower panel for a 10-qubit circuit.

6 Supplementary material

Find in the repo the codes we used to run tasks 1-4.

References

- [1] Arute et al, Quantum supremacy using a programmable superconducting processor, Nature 574 (2019).
- $[2] \ \ {\rm Sean \ Mullane}, \ \textit{Sampling \ random \ quantum \ circuits: \ a \ pedestrian's \ guide \ (2020)}.$