Evidently the Reynolds stresses  $\langle u', u', u', v' \rangle$  play a critical role in the evolution of the mean velocity field  $\overline{u}$ ; and mean pressure field  $\overline{p}$ . Consequently, it is important to understand the properties of the Reynolds stresses.

First and foremost, the Reynolds stresses are indeed stresses. The Reynolds stresses are the Components of a second-order tensor, which is obviously symmetric. The diagonal components are normal stresses while the off-diagonal components are stresses.

The turbulent kinetic energy  $k(\vec{x},t)$  is defined to be half the trace of the Reynolds stress tensor:  $k = \frac{1}{2} < \vec{u} \cdot \vec{u} > = \frac{1}{2} < u_i \cdot u_i >$ 

It is the mean kinetic energy per unit mass in the fluctuating velocity field. Note that:

$$\frac{1}{2} \langle \vec{u} \cdot \vec{u} \rangle = \frac{1}{2} \langle \vec{u} \cdot \vec{u} \rangle + \frac{1}{2} \langle \vec{u} \cdot \vec{u}' \rangle$$

$$= \frac{1}{2} \quad \vec{u} \cdot \vec{u} + \frac{1}{2} \langle \vec{u}' \cdot \vec{u}' \rangle$$

$$= \frac{1}{2} \quad \vec{u} \cdot \vec{u} + \frac{1}{2} \langle \vec{u}' \cdot \vec{u}' \rangle$$

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$$= \frac{1}{2} \quad \vec{u} \cdot \vec{u} \cdot \vec{u} + \frac{1}{2} \langle \vec{u}' \cdot \vec{u}' \rangle$$

$$= \frac{1}{2} \quad \vec{u} \cdot \vec{u} \cdot \vec{u} \cdot \vec{u} + \frac{1}{2} \langle \vec{u}' \cdot \vec{u}' \rangle$$

$$= \frac{1}{2} \quad \vec{u} \cdot \vec{$$

The Reynolds stress tensor is also positive - definite. That is:

$$\underline{x}^{T}\underline{\Gamma}^{x}\underline{x}$$
 =  $x_{i} < u_{i}'u_{j}' > x_{j} = < x_{i}u_{i}' \times_{j}u_{j}' > = < (x_{i}u_{i}')^{2} > \ge 0$ 

Reynolds Stresses

As a consequence, the eigenvalues of the Reynolds stress tensor are all positive or zero! Moreover, the normal stresses are always compressive. In general, all eigenvalues of the tensor are positive.

The Reynolds stress tensor may be decomposed as follows:

Does not affect mean flow!

We denote the anisotropic stress components by aij, and define the normalized anisotropy tensor as:

$$b_{ij} = \frac{a_{ij}}{2k} = \frac{\langle u'_i u'_j \rangle}{\langle u'_i u'_j \rangle} - \frac{1}{3} \delta_{ij}$$

In the event that the flow is irrotational, then:

$$\frac{\partial u_i}{\partial x_j} - \frac{\partial u_j}{\partial x_i} = 0$$

As a consequence:

$$\frac{\partial}{\partial x_{j}}\left(\frac{1}{2}\langle u_{i}'u_{i}'\rangle\right) - \frac{\partial}{\partial x_{i}}\left(\langle u_{i}'u_{j}'\rangle\right) = \left\langle u_{i}'\left(\frac{\partial u_{i}'}{\partial x_{j}} - \frac{\partial u_{j}'}{\partial x_{i}}\right)\right\rangle = 0$$

This gives rise to the Corrsin-Kistler equation:

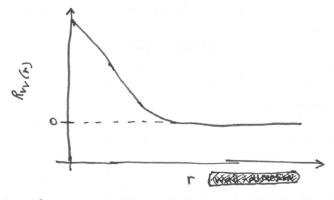
$$\frac{\partial}{\partial x_i}(\langle u_i'u_j'\rangle) = \frac{\partial k}{\partial x_j}$$

The above implies that the Reynolds stress is a pressure-like term for irrotational flow and has absolutely no effect on the mean velocity field. Therefore, a flow is turbulent only if it is notational!

The Reynolds Stress tensor is the Zero-Separation value of the two-point correlation:

$$R_{ij}(\vec{r}_s\vec{x}_st) = \langle u_i'(\vec{x}_st) u_j'(\vec{x}_t+\vec{r}_st) \rangle$$

The two-point correlation is a commonly studied statistical quantity, both because of its relation to the Reynolds stress tensor and because of its association with turbulent structures. For turbulent flows, two-point correlations generally decay smoothy to zero for large separation, with finite correlations for small separation.



Representative plot of two-point correlation of the well-normal velocity component for turbulent channel flow for separations in the streamvise direction.

Finally, for some flows, symmetries in the flow geometry, determine properties of the Reynolds stress tensor. See page 89 of Pope for more details.