Flow Rates of Mass, Momentum, and Energy:

Using the conservative form of the shear layer equation:

$$\frac{\partial}{\partial x} ((\bar{u})^2) + \frac{\partial}{\partial y} (\bar{u}\bar{v}) = -\frac{\partial}{\partial y} \langle u'v' \rangle \quad (ignoring \ viscous \ terms)$$

We can make some fundamental observations about flow rates. Integrate the above from  $y=-\infty$  to  $y=+\infty$  to obtain:

$$\frac{d}{dx} \int_{-\infty}^{\infty} (\bar{u})^2 dy = -\left[\bar{u}\bar{v} + \sqrt{u}v'\right]_{-\infty}^{\infty}$$

for a planar jet or wake, both  $\bar{u}\bar{v}\Big|_{\infty}^{\infty}$  and  $(u'v')\Big|_{\infty}^{\infty}$  to zero (as  $y\to\pm\infty$ ). Thus:

$$\frac{d}{dx} \int_{0}^{\infty} (\bar{u})^{2} dy = 0 \quad \text{jet or wake}$$

This implies that the momentum flow rate of the mean flow:

$$\dot{M}(x) = \int_{-\infty}^{\infty} (\bar{u}(x,y))^2 dy$$

is conserved in that the rate is independent of x for jets and wakes. For a mixing layer,  $\bar{u}\bar{v}|_{\infty}^{\infty} \pm 0$ , so the rate is not conserved.

We will show later that we expect:

$$\bar{u}(x,y) = u_{\text{offset}} + u_{s}(x) \bar{f}(y)$$
  $\gamma = f(x)$ 

for a fully developed free shear flow that is planar. If unffset = 0 (as for a non-coflowing jet), then the flow rates of mass and energy are , in the mean:

Flow Rate of Energy of the Mean Flow: 
$$\dot{p}(x) \equiv \int_{-\infty}^{\infty} g \, u_s(x) \, \vec{f}(\vec{s}(x)) \, dy$$

$$= g \, u_s(x) \, \delta(x) \int_{-\infty}^{\infty} \vec{f}(\eta) \, d\eta$$

$$= \int_{-\infty}^{\infty} \frac{1}{2} g \, u_s(x) \, \vec{f}(s(x)) \, dy$$

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For a non-coffewing jet, we will later show that Jux and Us ~ x - 1/2 so:

$$\dot{m}(x) \sim x^{\frac{1}{2}}$$
 For a non-conflowing jet!  $\dot{E}(x) \sim x^{-\frac{1}{2}}$ 

Therefore, the mass flow rate increases and the energy flow rate decreases as x increases.