Fluid Simulation for Video Games (part 19)

By Dr. Michael J. Gourlay

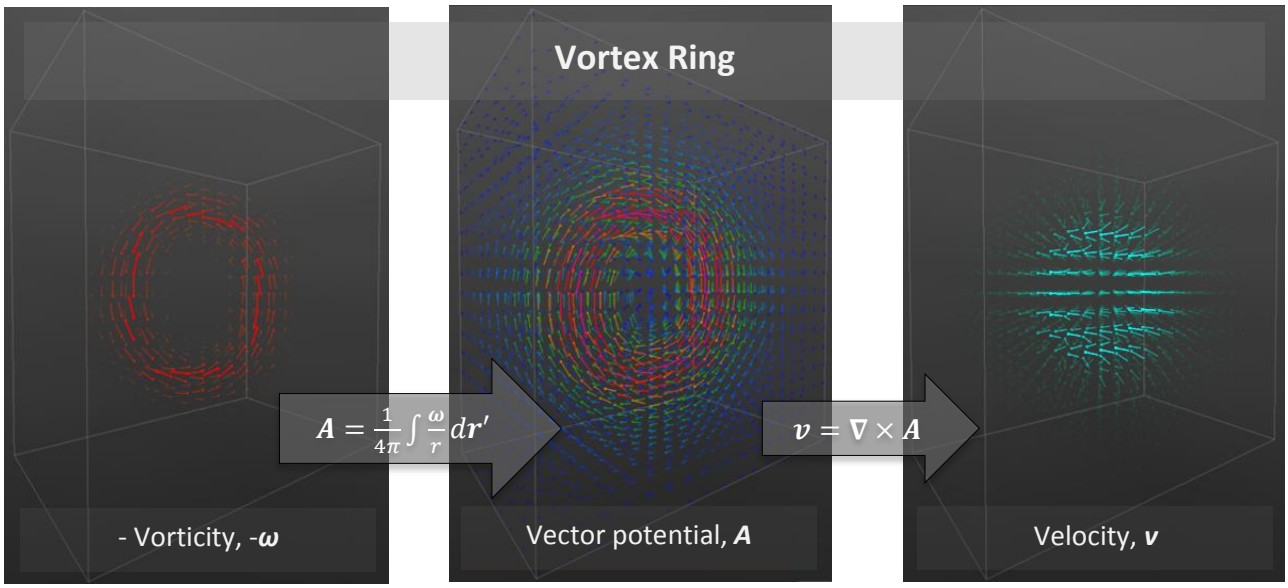

Figure 1. *Vector field visualizations of negative vorticity, vector potential and velocity for a vortex ring.*

From Vorticity through Vector Potential to Velocity

This article plus the next one describe how to improve fidelity while reducing computational cost, by combining techniques described in earlier articles in this series.

Fluid simulation entails computing flow velocity everywhere in the fluid domain. Most articles in this series advocate tracking *vorticity* (the *curl* of velocity) which describes how a flow rotates. So the simulation needs to map vorticity into velocity. Part 3 and part 4 described how to do that using an integral technique and Part 6 described a differential technique. The integral technique has computational complexity O(N log N) and memory complexity O(N), whereas the differential technique has computational complexity O(N) and memory complexity O(N log N). When computers tend to have more abundant memory than compute, the latter seems more attractive.

But the differential technique (solving a vector Poisson equation) had another problem: Boundary conditions. To solve a partial differential equation, you effectively already need a solution at the boundaries. This seems to have a circular dependency: To solve the equation, you need a solution.

In Part 6, I dodged this problem by assuming some solution at the boundaries, then keeping the boundaries far from the region of interest. But this greatly expanded the domain size. Done the obvious way, that would consume more compute and memory: If the domain increased by 3x in each direction, memory and compute would increase by 3^3=27 times. Instead, I decreased spatial resolution and held memory and compute constant. So the overall result was that the differential method had lower fidelity.

But there is another solution: Compute the solution at the boundaries using an integral technique. Remember, part 3 and part 4 explain how to integrate vorticity to get *velocity*, so you're already familiar with an integral technique. In this

case, however, we want to integrate vorticity to get the *vector potential*, not velocity. Fortunately the technique for computing vector potential from vorticity also applies to computing velocity from vorticity, simply by using a different formula for the integrand.

This article explains how to use O(N^2) and O(N log N) integral techniques to compute vector potential, then reuses the same approach as in part 6 to compute velocity from vector potential. The subsequent article will apply the technique from this article, to impose boundary conditions to solve the vector Poisson equation using the O(N) multi-grid algorithm, yielding a faster algorithm, with better fidelity, than any of the algorithms presented so far in this series.

Integral 1 (part 3)

Velocity from vorticity

$$\boldsymbol{v} = \frac{1}{4\pi}\int \frac{\boldsymbol{\omega} \times \boldsymbol{r}}{r^3} d\boldsymbol{r}'$$

Differential with generic, simplistic BC (part 6)

Velocity from vector potential from vorticity

$$\nabla^2 \mathbf{A} = -\boldsymbol{\omega} \quad \text{in } \Omega \text{ (interior)}$$

where $$a\boldsymbol{A} + b\frac{\partial \mathbf{A}}{\partial \boldsymbol{n}} = 0 \quad \text{on } \partial\Omega \text{ (boundary)}$$

$$\boldsymbol{v} = \boldsymbol{\nabla} \times \boldsymbol{A}$$

Integral 2 (part 19, this article)

Velocity from vector potential from vorticity

$$\boldsymbol{v} = \boldsymbol{\nabla} \times \boldsymbol{A}$$

$$\boldsymbol{A} = \frac{1}{4\pi}\int \frac{\boldsymbol{\omega}}{r} d\boldsymbol{r}'$$

Differential with integral BC (part 20, next article)

Velocity from vector potential from vorticity

$$\nabla^2 \mathbf{A} = -\boldsymbol{\omega} \quad \text{in } \Omega \text{ (interior)}$$

where $$\boldsymbol{A} = \frac{1}{4\pi}\int \frac{\boldsymbol{\omega}}{r} d\boldsymbol{r}' \quad \text{on } \partial\Omega \text{ (boundary)}$$

$$\boldsymbol{v} = \boldsymbol{\nabla} \times \boldsymbol{A}$$

Figure 2: *Comparison of fluid simulation techniques.*

Parts 1 and 2 summarized fluid dynamics and simulation techniques. Parts 3 and 4 presented a vortex-particle fluid simulation with two-way fluid-body interactions that run in real time. Part 5 profiled and optimized that simulation code. Part 6 described a differential method for computing velocity from vorticity. Figure 1 shows the relationships between the various techniques and articles. Part 7 showed how to integrate a fluid simulation into a typical particle system. Parts 8, 9, 10 and 11 explained how to simulate density, buoyancy, heat and combustion in a vortex-based fluid simulation. Part 12 explained how improper sampling caused unwanted jerky motion and described how to mitigate it. Part 13 added convex polytopes and lift-like forces. Parts 14, 15, 16, 17 and 18 added containers, smoothed particle hydrodynamics (SPH), liquids and fluid surfaces.

Vector Potential

Recall from Part 6 that the Helmholtz theorem states that a vector function $\boldsymbol{v}$ can be decomposed into an irrotational part ($\boldsymbol{\nabla}\varphi$) that has only divergence and a solenoidal part ($\boldsymbol{\nabla} \times \boldsymbol{A}$) that has only curl:

$$\boldsymbol{v} = \boldsymbol{\nabla} \times \boldsymbol{A} - \boldsymbol{\nabla}\varphi \qquad (1)$$

where:

$$\boldsymbol{A}(\boldsymbol{r}) = \frac{1}{4\pi}\int \frac{\boldsymbol{\nabla}' \times \boldsymbol{v}(\boldsymbol{r}')}{|\boldsymbol{r} - \boldsymbol{r}'|} d^3\boldsymbol{r}' \quad (2)$$

Compare (2) with the formula from Part 2 for obtaining velocity $\boldsymbol{v}$ from vorticity $\boldsymbol{\omega}$, where $\boldsymbol{\omega} = \boldsymbol{\nabla} \times \boldsymbol{v}$, using the Biot-Savart law:

$$\boldsymbol{v}(\boldsymbol{r}) = \frac{\boldsymbol{1}}{\boldsymbol{4\pi}}\int \frac{\boldsymbol{\omega}(\boldsymbol{r}') \times (\boldsymbol{r} - \boldsymbol{r}')}{|\boldsymbol{r} - \boldsymbol{r}'|^3} \boldsymbol{d^3 r}' \quad (3)$$

Equation (3) formed the basis for the algorithm presented in Part 3, which was used throughout most of the articles in this series.

Notice that the integrals in equations (2) and (3) vaguely resemble each other. Numerical integration techniques used to compute velocity (***v***, eqn. 3) also apply to vector potential (***A***, eqn. 2). So you can numerically integrate the vector potential ***A***, then take the curl of ***A*** to get velocity, ***v***=∇×***A***.

This might seem odd; why bother going through the extra step of computing vector potential from vorticity, then taking its curl, if you can directly compute velocity from vorticity? Because computing ***A*** by solving a Poisson equation ($\boldsymbol{\nabla^2 A} = -\boldsymbol{\omega}$) can be much faster within a domain interior, but you need to know ***A*** on the domain boundaries. So you can integrate vorticity to compute ***A*** on the domain boundaries, then solve a Poisson equation to compute ***A*** in the domain interior. That's a lot of steps. So this article will only focus on computing ***A*** everywhere using the integral (2), then the next article will combine integral and differential techniques into a coherent solution.

What is Vector Potential?

I find it useful to develop an intuition for quantities like velocity, vorticity and vector potential. Figure 1 shows those as vector fields. Most people have a sense for what velocity means: The direction and speed something moves. Vorticity is less familiar but still accessible: it's like the axis about which some fluid rotates, where the axis points away from the clockwise face and its length indicates how fast the fluid spins. Vector potential is a little more abstract, but in Figure 1, vector potential looks like a diffused or blurred version of vorticity. That's a clue for how to think about it.

The interpretation of vector potential as "diffused vorticity" dovetails well with Poisson's equation. Consider an analogy to a block of metal heated or cooled at a collection of points. For example, it could be held at a higher temperature at one end and at a lower temperature at the opposite end. Given enough time, the heat will conduct through the block and its temperature distribution will reach a steady state. The scalar version of Poisson's equation describes the distribution of heat throughout the block.

With that mental model in mind, look at Figure 1 again. Imagine that vorticity is a heat source inside of a block of metal suspended within a cooler environment. (In Figure 1, that heat source would be a ring, like the outer ring of a stove burner.) The vector potential looks something like what the temperature would be throughout that block: warm where red and cool where blue. One obvious difference here is that vorticity and vector potential are vectors, not scalars (like heat and temperature), but if you consider each component (x, y and z) separately, each are scalar fields. (The x-component of vorticity is the "heat source" for the x-component of the vector potential "temperature", and so on for y and z.) So the vector Poisson equation is almost like it describes 3 "flavors" of heat that propagate separately.

Computing Vector Potential by Integrating Vorticity

Analogous to the approach in Part 3, you could replace the integral in equation (2) with a summation:

$$\boldsymbol{A}(\boldsymbol{x}) = \frac{1}{4\pi}\sum_{i=1}^{N} \frac{\boldsymbol{\omega}_i}{r_i} \quad (4)$$

The summand in (4) has the problem that as r_i gets small (that is, as the query point x gets close to a vorton center), the $1/r_i$ term gets very large. Using the same approach as in Part 3, mollify the vortex by giving it a finite size. The summation then becomes

$$\boldsymbol{A}(\boldsymbol{x}) = \frac{1}{4\pi}\sum_{i=1}^{N}\frac{\boldsymbol{\omega}_i}{r_i}\Delta V_i \begin{cases} 1 & : r \geq \sigma \\ (r_i/\sigma_i)^3 & : r \leq \sigma \end{cases} \qquad (5)$$

This macro implements the formula of equation 5:

```
#define VORTON_ACCUMULATE_VECTOR_POTENTIAL_private( vVecPot , vPosQuery , mPosition , mAngularVelocity , mSize ,
mSpreadingRangeFactor , mSpreadingCirculationFactor )    \
{                                                                                                                  \
    const Vec3          vOtherToSelf = vPosQuery - mPosition ;                                                     \
    const float         radius       = mSize * 0.5f * mSpreadingRangeFactor ;                                      \
    const float         radius2      = radius * radius ;                                                           \
    const float         dist2        = vOtherToSelf.Mag2() ;                                                       \
    const float         distLaw      = ( dist2 < radius2 ) ? ( dist2 / ( radius2 * radius ) ) : ( finvsqrtf( dist2 ) ) ; \
    const Vec3  vecPotContribution = TwoThirds * radius2 * radius * mAngularVelocity * distLaw ;                   \
    vVecPot +=  vecPotContribution * mSpreadingCirculationFactor ;                                                 \
}
```

Note: See the archive file associated with this article. It contains the rest of the code. Also see Part 3 for elaboration of these code snippets.

Direct Summation

You can use equation 5 by directly summing it, for all vortons, for all query points. Assuming you want to query at each vorton, that makes direct summation have asymptotic time complexity $O(N^2)$. The treecode presented below runs faster, but makes some approximations. You can use the direct summation algorithm to compare other methods:

```
Vec3 VortonSim::ComputeVectorPotential_Direct( const Vec3 & vPosition )
{
    const size_t    numVortons      = mVortons->Size() ;
    Vec3            vecPotAccumulator( 0.0f , 0.0f , 0.0f ) ;
    for( unsigned iVorton = 0 ; iVorton < numVortons ; ++ iVorton )
    {   // For each vorton...
        const Vorton &  rVorton = (*mVortons)[ iVorton ] ;
        VORTON_ACCUMULATE_VECTOR_POTENTIAL( vecPotAccumulator , vPosition , rVorton ) ;
    }
    return vecPotAccumulator ;
}
```

Treecode

The treecode algorithm for computing vector potential has identical structure to the treecode that computes velocity. It runs in O(N log N) time:

```
Vec3 VortonSim::ComputeVectorPotential_Tree( const Vec3 & vPosition , const unsigned indices[3] , size_t iLayer , const
UniformGrid< VECTOR< unsigned > > & vortonIndicesGrid , const NestedGrid< Vorton > & influenceTree )
{
ComputeVectorPotential_Direct.
```

```
const UniformGrid< Vorton > &   rChildLayer             = influenceTree[ iLayer - 1 ] ;
unsigned                        clusterMinIndices[3] ;
const unsigned *                pClusterDims            = influenceTree.GetDecimations( iLayer ) ;
influenceTree.GetChildClusterMinCornerIndex( clusterMinIndices , pClusterDims , indices ) ;

const Vec3 &            vGridMinCorner          = rChildLayer.GetMinCorner() ;
const Vec3              vSpacing                = rChildLayer.GetCellSpacing() ;
unsigned                increment[3]            ;
const unsigned &        numXchild               = rChildLayer.GetNumPoints( 0 ) ;
const unsigned          numXYchild              = numXchild * rChildLayer.GetNumPoints( 1 ) ;
Vec3                    vecPotAccumulator( 0.0f , 0.0f , 0.0f ) ;
const float             vortonRadius = (*mVortons)[ 0 ].GetRadius() ;
static const float      marginFactor    = 0.0001f ;

// When domain is 2D in XY plane, min.z==max.z so vPos.z test below would fail unless margin.z!=0.
const Vec3          margin          = marginFactor * vSpacing + ( 0.0f == vSpacing.z ? Vec3(0,0,FLT_MIN) : Vec3(0,0,0) );

// For each cell of child layer in this grid cluster...
for( increment[2] = 0 ; increment[2] < pClusterDims[2] ; ++ increment[2] )
{
    unsigned idxChild[3] ;
    idxChild[2] = clusterMinIndices[2] + increment[2] ;
    Vec3 vCellMinCorner , vCellMaxCorner ;
    vCellMinCorner.z = vGridMinCorner.z + float( idxChild[2]     ) * vSpacing.z ;
    vCellMaxCorner.z = vGridMinCorner.z + float( idxChild[2] + 1 ) * vSpacing.z ;
    const unsigned offsetZ = idxChild[2] * numXYchild ;
    for( increment[1] = 0 ; increment[1] < pClusterDims[1] ; ++ increment[1] )
    {
        idxChild[1] = clusterMinIndices[1] + increment[1] ;
        vCellMinCorner.y = vGridMinCorner.y + float( idxChild[1]     ) * vSpacing.y ;
        vCellMaxCorner.y = vGridMinCorner.y + float( idxChild[1] + 1 ) * vSpacing.y ;
        const unsigned offsetYZ = idxChild[1] * numXchild + offsetZ ;
        for( increment[0] = 0 ; increment[0] < pClusterDims[0] ; ++ increment[0] )
        {
            idxChild[0] = clusterMinIndices[0] + increment[0] ;
            vCellMinCorner.x = vGridMinCorner.x + float( idxChild[0]     ) * vSpacing.x ;
            vCellMaxCorner.x = vGridMinCorner.x + float( idxChild[0] + 1 ) * vSpacing.x ;
            if(
                    ( iLayer > 1 )
                &&  ( vPosition.x >= vCellMinCorner.x - margin.x )
                &&  ( vPosition.y >= vCellMinCorner.y - margin.y )
                &&  ( vPosition.z >= vCellMinCorner.z - margin.z )
                &&  ( vPosition.x <  vCellMaxCorner.x + margin.x )
                &&  ( vPosition.y <  vCellMaxCorner.y + margin.y )
                &&  ( vPosition.z <  vCellMaxCorner.z + margin.z )
              )
            {   // Test position is inside childCell and currentLayer > 0...
                // Recurse child layer.
  vecPotAccumulator += ComputeVectorPotential_Tree( vPosition , idxChild , iLayer - 1 , vortonIndicesGrid , influenceTree ) ;
            }
            else
            {   // Test position is outside childCell, or reached leaf node.
                //    Compute velocity induced by cell at corner point x.
                //    Accumulate influence, storing in velocityAccumulator.
                const unsigned  offsetXYZ       = idxChild[0] + offsetYZ ;
                const Vorton &  rVortonChild    = rChildLayer[ offsetXYZ ] ;

                if( 1 == iLayer )
                {   // Reached base layer.
                    // Instead of using supervorton, use direct summation of original vortons.
                    // This only makes a difference when a leaf-layer grid cell contains multiple vortons.
                    const unsigned numVortonsInCell = vortonIndicesGrid[ offsetXYZ ].Size() ;
                    for( unsigned ivHere = 0 ; ivHere < numVortonsInCell ; ++ ivHere )
                    {   // For each vorton in this gridcell...
                        const unsigned &    rVortIdxHere    = vortonIndicesGrid[ offsetXYZ ][ ivHere ] ;
                        Vorton &            rVortonHere     = (*mVortons)[ rVortIdxHere ] ;
                        VORTON_ACCUMULATE_VECTOR_POTENTIAL( vecPotAccumulator , vPosition , rVortonHere ) ;
                    }
                }
                else
                {   // Current layer is either...
                    // not a base layer, i.e. this layer is an "aggregation" layer of supervortons, ...or...
                    // this is the base layer and USE_ORIGINAL_VORTONS_IN_BASE_LAYER is disabled.
                    VORTON_ACCUMULATE_VECTOR_POTENTIAL( vecPotAccumulator , vPosition , rVortonChild ) ;
```

```
                }
            }
        }
    }
    return vecPotAccumulator ;
}
```

Parallelize Vector Potential Algorithms with Intel® Threading Building Blocks

You can use Intel® Threading Building Blocks to parallelize either the direct summation or treecode algorithms. The approach follows that taken in Part 3:

- Write a worker routine that operates on a slice of the problem.
- Write a wrapper functor class that wraps the worker routine.
- Write a wrapper routine that directs TBB to call a functor.

Here is the worker routine:

```
void VortonSim::ComputeVectorPotentialAtGridpoints_Slice( size_t izStart , size_t izEnd
                  , const UniformGrid< VECTOR< unsigned > > & vortonIndicesGrid , const NestedGrid< Vorton > & influenceTree )
{
    const size_t        numLayers   = influenceTree.GetDepth() ;

    UniformGrid< Vec3 > &   vectorPotentialGrid = mVectorPotentialMultiGrid[ 0 ] ;

    const Vec3 &        vMinCorner  = mVelGrid.GetMinCorner() ;
    static const float  nudge       = 1.0f - 2.0f * FLT_EPSILON ;
    const Vec3          vSpacing    = mVelGrid.GetCellSpacing() * nudge ;
    const unsigned      dims[3]     =   { mVelGrid.GetNumPoints( 0 )
                                        , mVelGrid.GetNumPoints( 1 )
                                        , mVelGrid.GetNumPoints( 2 ) } ;
    const unsigned      numXY       = dims[0] * dims[1] ;
    unsigned            idx[ 3 ] ;

    for( idx[2] = static_cast< unsigned >( izStart ) ; idx[2] < izEnd ; ++ idx[2] )
    {   // For subset of z index values...
        Vec3 vPosition ;
        // Compute the z-coordinate of the world-space position of this gridpoint.
        vPosition.z = vMinCorner.z + float( idx[2] ) * vSpacing.z ;
        // Precompute the z contribution to the offset into the velocity grid.
        const unsigned offsetZ = idx[2] * numXY ;
        for( idx[1] = 0 ; idx[1] < dims[1] ; ++ idx[1] )
        {   // For every gridpoint along the y-axis...
            // Compute the y-coordinate of the world-space position of this gridpoint.
            vPosition.y = vMinCorner.y + float( idx[1] ) * vSpacing.y ;
            // Precompute the y contribution to the offset into the velocity grid.
            const unsigned offsetYZ = idx[1] * dims[0] + offsetZ ;
            for( idx[0] = 0 ; idx[0] < dims[0] ; ++ idx[0] )
            {   // For every gridpoint along the x-axis...

                // Compute the x-coordinate of the world-space position of this gridpoint.
                vPosition.x = vMinCorner.x + float( idx[0] ) * vSpacing.x ;
                // Compute the offset into the velocity grid.
                const unsigned offsetXYZ = idx[0] + offsetYZ ;

                // Compute the fluid flow velocity at this gridpoint, due to all vortons.
                static const unsigned zeros[3] = { 0 , 0 , 0 } ; // Starter indices for recursive algorithm
                if( numLayers > 1 )
                {
                    vectorPotentialGrid[ offsetXYZ ] = ComputeVectorPotential_Tree( vPosition , zeros , numLayers - 1
                                                                                  , vortonIndicesGrid , influenceTree ) ;
                }
                else
                {
```

```
                    vectorPotentialGrid[ offsetXYZ ] = ComputeVectorPotential_Direct( vPosition ) ;
                }
            }
        }
    }
}
```

Here is the functor class:

```
class VortonSim_ComputeVectorPotentialAtGridpoints_TBB
{
        VortonSim *                                 mVortonSim          ;
        const UniformGrid< VECTOR< unsigned > > &   mVortonIndicesGrid  ;
        const NestedGrid< Vorton > &                mInfluenceTree      ;
    public:
        void operator() ( const tbb::blocked_range<size_t> & r ) const
        {   // Compute subset of vector potential grid.
            SetFloatingPointControlWord( mMasterThreadFloatingPointControlWord ) ;
            SetMmxControlStatusRegister( mMasterThreadMmxControlStatusRegister ) ;
            mVortonSim->ComputeVectorPotentialAtGridpoints_Slice( r.begin() , r.end() , mVortonIndicesGrid , mInfluenceTree ) ;
        }
        VortonSim_ComputeVectorPotentialAtGridpoints_TBB( VortonSim * pVortonSim
                                                        , const UniformGrid< VECTOR< unsigned > > & vortonIndicesGrid
                                                        , const NestedGrid< Vorton > & influenceTree )
            : mVortonSim( pVortonSim )
            , mVortonIndicesGrid( vortonIndicesGrid )
            , mInfluenceTree( influenceTree )
        {
            mMasterThreadFloatingPointControlWord = GetFloatingPointControlWord() ;
            mMasterThreadMmxControlStatusRegister = GetMmxControlStatusRegister() ;
        }
    private:
        WORD        mMasterThreadFloatingPointControlWord   ;
        unsigned    mMasterThreadMmxControlStatusRegister   ;
} ;
```

Here is the wrapper routine:

```
void VortonSim::ComputeVectorPotentialFromVorticity_Integral( const UniformGrid< VECTOR< unsigned > > & vortonIndicesGrid
                                                            , const NestedGrid< Vorton > & influenceTree )
{
    const unsigned numZ = mVelGrid.GetNumPoints( 2 ) ;
    // Estimate grain size based on size of problem and number of processors.
    const size_t grainSize =  Max2( size_t( 1 ) , numZ / gNumberOfProcessors ) ;
    // Compute velocity at gridpoints using multiple threads.
    parallel_for( tbb::blocked_range<size_t>( 0 , numZ , grainSize )
                , VortonSim_ComputeVectorPotentialAtGridpoints_TBB( this , vortonIndicesGrid , influenceTree ) ) ;
}
```

Performance

The table shows the duration (in milliseconds per frame) of various routines, run on a 3.50 GHz Intel Core i7-3770K that has 4 physical cores and 2 local cores per physical core.

# threads	Frame	Vorton Sim	Vector Potential	Render
1	31.7	7.80	2.78	13.8
2	18.6	5.47	1.30	7.33
3	13.8	4.89	0.982	4.94
4	13.6	4.79	0.845	4.93
8	11.6	4.55	0.683	3.84

Summary

This article applied the integral techniques presented in Part 3 to compute the vector potential described in Part 6. This is meant as a stepping stone to a hybrid integral-differential technique: The next article (Part 20) in this series will describe how to use the differential technique (also presented in Part 6) to compute for vector potential by solving a vector Poisson equation, where the boundary conditions will be specified using integral techniques. The hybrid solution will run faster and with better results than either the purely integral or purely differential techniques.

Further Reading

Normally I list references that I think elaborate on the topic, but in this unusual case, I could find very little on the subject. I found this in a web search and I admit that I have neither obtained nor read it. But I found so few relevant publications that I chose to list this one here in case it helps any readers.

- Tutty (1986): On vector potential-vorticity methods for incompressible flow problems. Journal of Computational Physics, volume 64, Issue 2, June 1986, Pages 368-379.

About the Author

Dr. Michael J. Gourlay works at Microsoft as a principal development lead on HoloLens in the Environment Understanding group.

He previously worked at Electronic Arts (EA Sports) as the software architect for the Football Sports Business Unit, as a senior lead engineer on *Madden NFL* and original architect of FranTk (the engine behind Connected Careers mode), on character physics and ANT (the procedural animation system used by EA), on *Mixed Martial Arts*, and as a lead programmer on *NASCAR.* He wrote Lynx (the visual effects system used in EA games worldwide) and patented algorithms for interactive, high-bandwidth online applications.

He also developed curricula for and taught at the University of Central Florida, Florida Interactive Entertainment Academy, an interdisciplinary graduate program that teaches programmers, producers, and artists how to make video games and training simulations.

Prior to joining EA, he performed scientific research using computational fluid dynamics and the world's largest massively parallel supercomputers. His previous research also includes nonlinear dynamics in quantum mechanical

systems and atomic, molecular, and optical physics. Michael received his degrees in physics and philosophy from Georgia Tech and the University of Colorado at Boulder.