

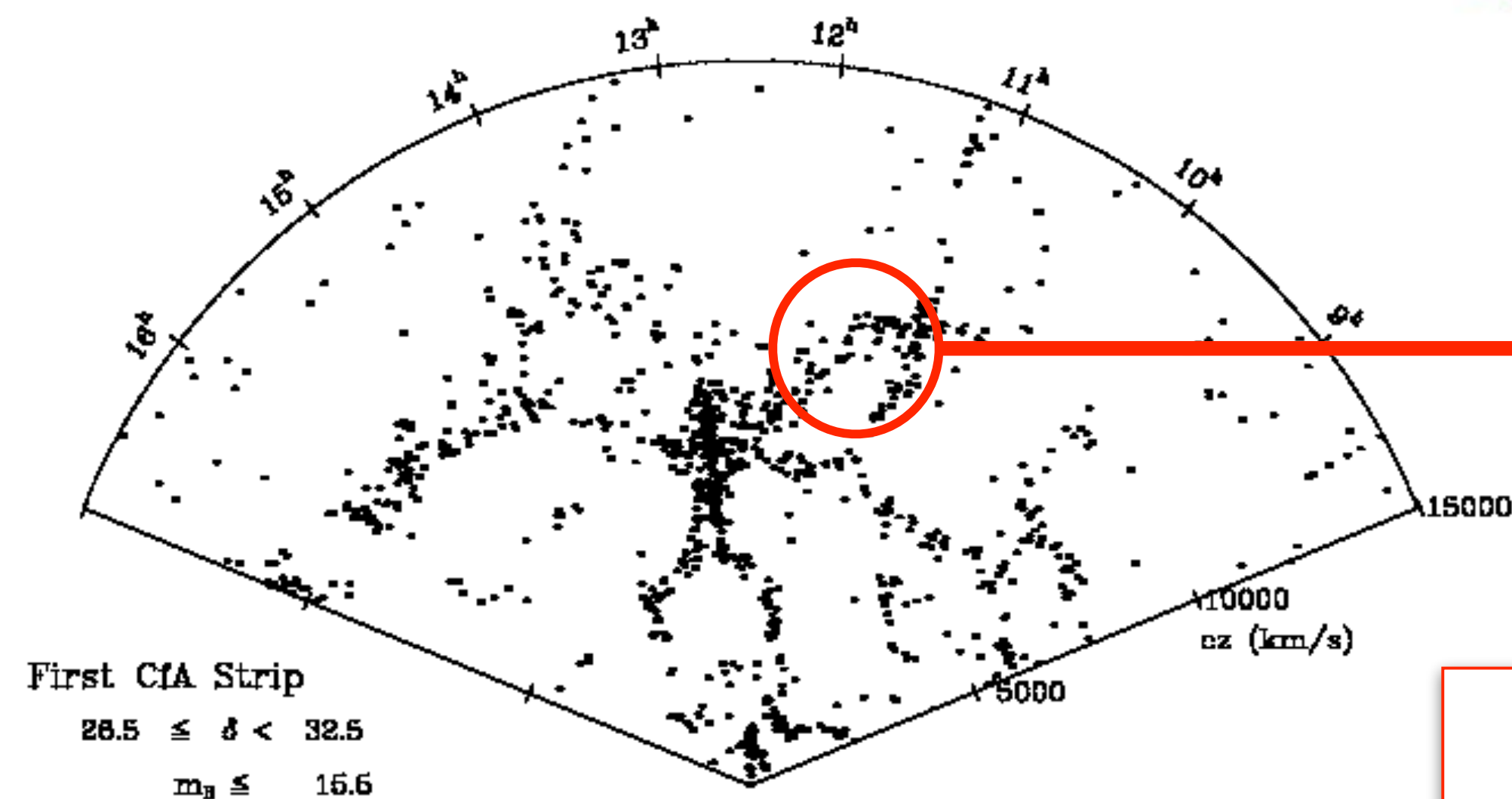
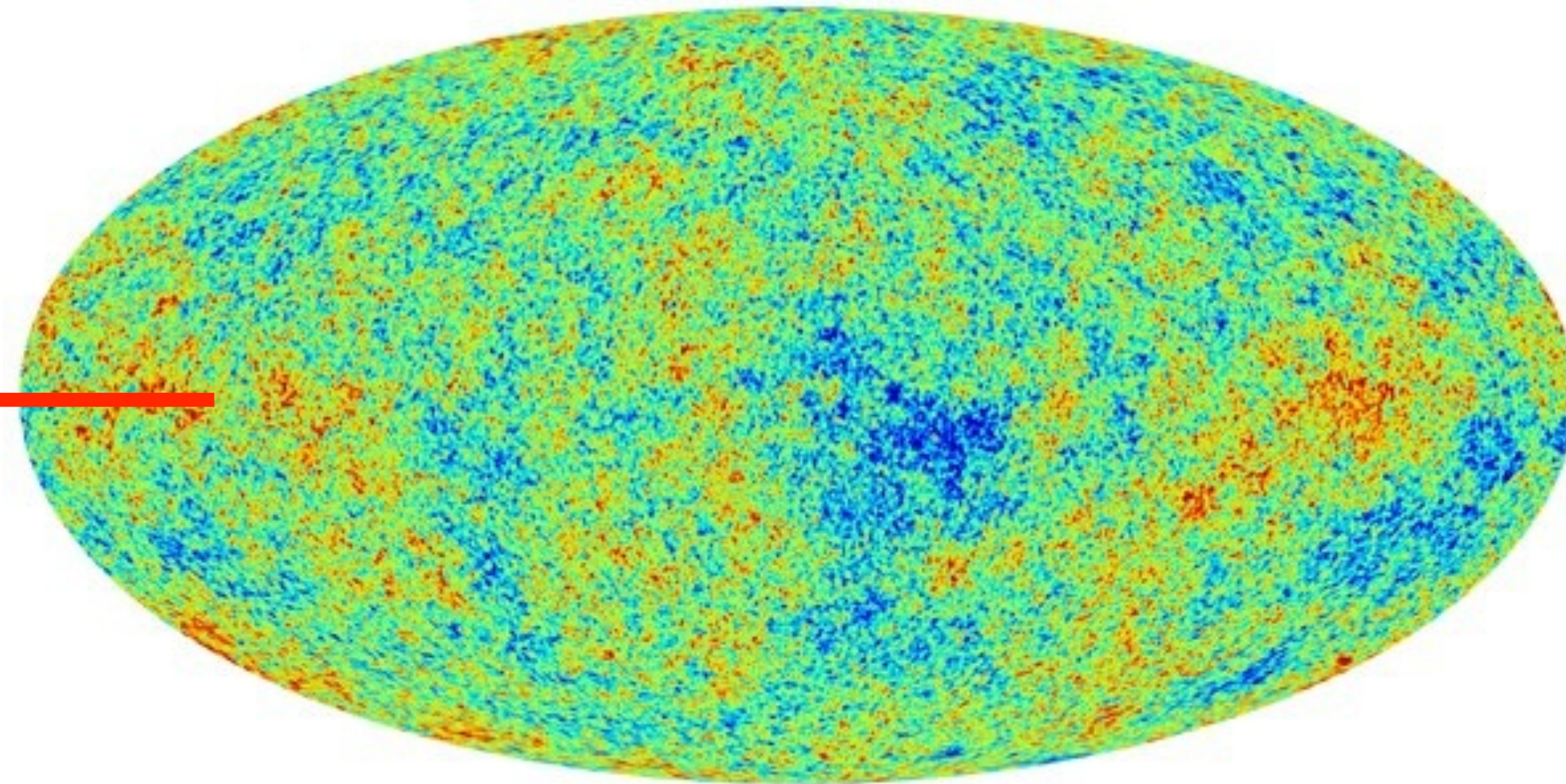
# 2PCF

Slides stolen from Emiliano Sefussati

# Cosmological perturbations

CMB temperature fluctuations

$$T(\hat{n})$$



$$n_g(\vec{x})$$

number density of galaxies

We can only study the **statistical properties** of cosmological perturbations

Mathematically, these are **random fields**

# Random fields

If  $\phi$  is a **random variable** with Probability Distribution Function (PDF)  $\mathcal{P}(\phi)$  we can compute:

$$\langle \phi \rangle = \int d\phi \mathcal{P}(\phi) \phi$$

mean

$$\langle \phi^2 \rangle = \int d\phi \mathcal{P}(\phi) \phi^2$$

2-nd-order moment

$$\langle \phi^n \rangle = \int d\phi \mathcal{P}(\phi) \phi^n$$

n-th-order moment

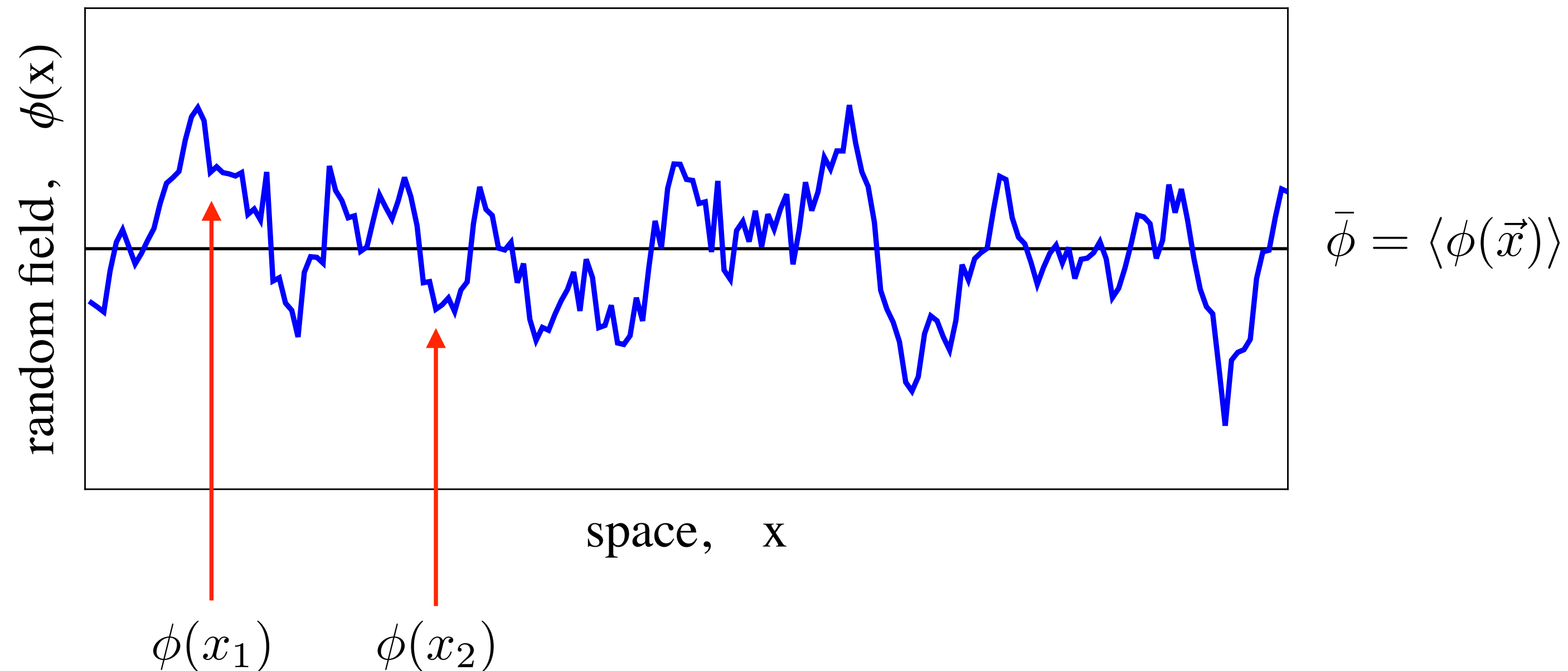
$$\sigma_\phi^2 = \langle \phi^2 \rangle - \langle \phi \rangle^2$$

variance



# Random fields

If  $\phi(\vec{x})$  is a **random field** we can also compute **correlation functions**



**two-point function**  $\langle \phi(x_1)\phi(x_2) \rangle = \langle \phi(x_1) \rangle \langle \phi(x_2) \rangle + \langle \phi(x_1)\phi(x_2) \rangle_c$

**three-point function**  $\langle \phi(x_1)\phi(x_2)\phi(x_3) \rangle = \langle \phi(x_1) \rangle \langle \phi(x_2) \rangle \langle \phi(x_3) \rangle +$   
 $+ \langle \phi(x_1)\phi(x_2) \rangle_c \langle \phi(x_3) \rangle + \text{perm.} +$   
 $+ \langle \phi(x_1)\phi(x_2)\phi(x_3) \rangle_c$

...

**n-point function**  $\langle \phi(x_1)\phi(x_2) \dots \phi(x_n) \rangle$

# The distribution of galaxies

The galaxy number density and its perturbations as random fields

$$n_g(\vec{x}) \equiv \bar{n}_g [1 + \delta_g(\vec{x})]$$

**galaxy number density**

 **mean** galaxy number

$$\delta_g(\vec{x}) \equiv \frac{n_g(\vec{x}) - \bar{n}_g}{\bar{n}_g}$$

**galaxy overdensity**  
or density contrast

N.B.  $\langle \delta_g(\vec{x}) \rangle \equiv 0$

$$\delta_g(\vec{x}) \geq -1$$



# The galaxy two-point correlation function

What is the probability of finding two galaxies in the volume elements  $dV_1$  and  $dV_2$ ?

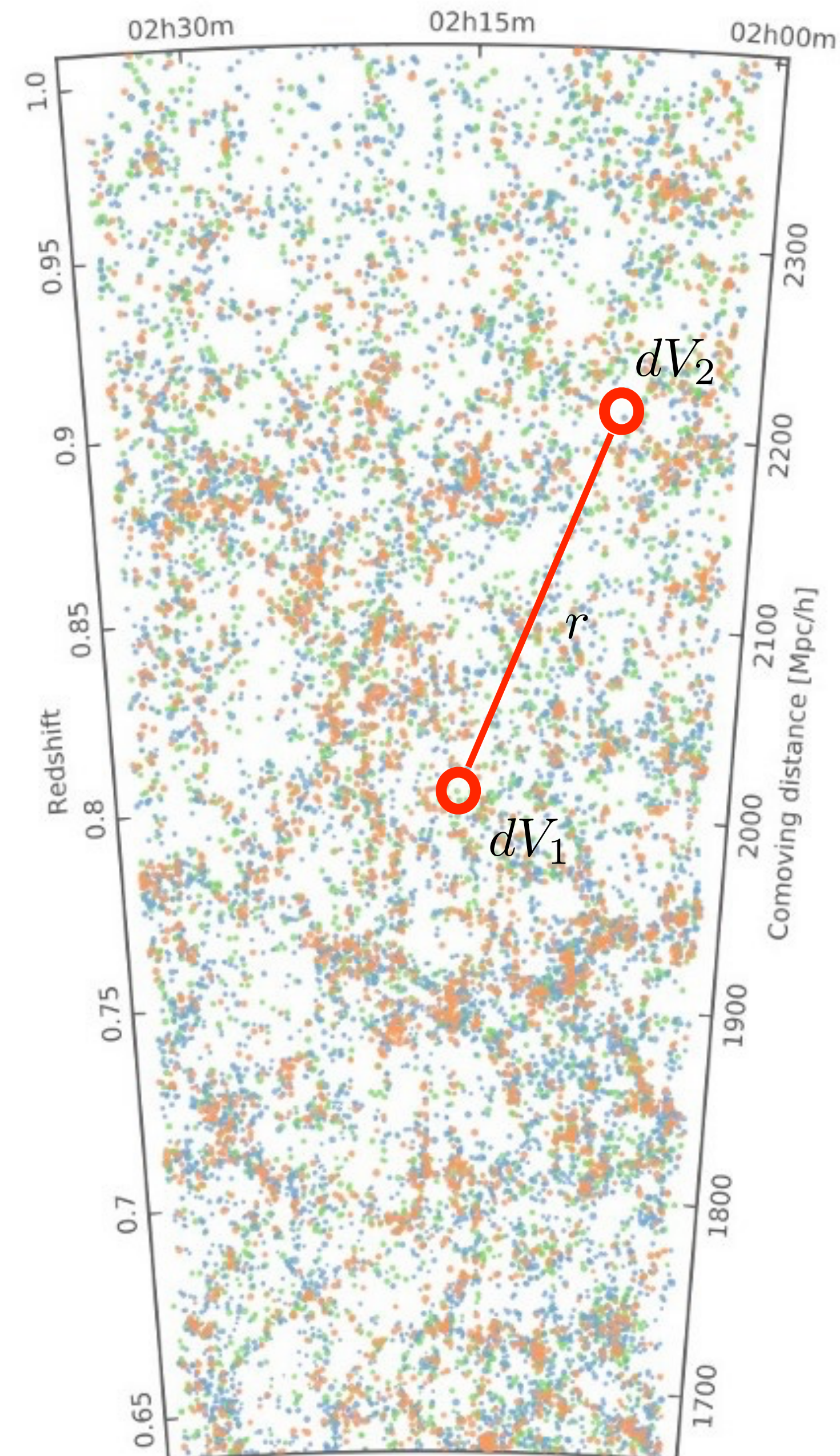
$$\begin{aligned} dP &= dV_1 dV_2 \langle n_g(\vec{x}_1) n_g(\vec{x}_2) \rangle \\ &= dV_1 dV_2 \bar{n}_g^2 [1 + \langle \delta_g(\vec{x}_1) \delta_g(\vec{x}_2) \rangle] \end{aligned}$$

↑  
excess probability

We now make the assumption of  
**statistical homogeneity and isotropy**

$$\xi(|\vec{x}_1 - \vec{x}_2|) \equiv \langle \delta_g(\vec{x}_1) \delta_g(\vec{x}_2) \rangle$$

*the two-point correlation function  $\xi(r)$   
only depends on the distance  $r = |\vec{x}_1 - \vec{x}_2|$   
between the two points*



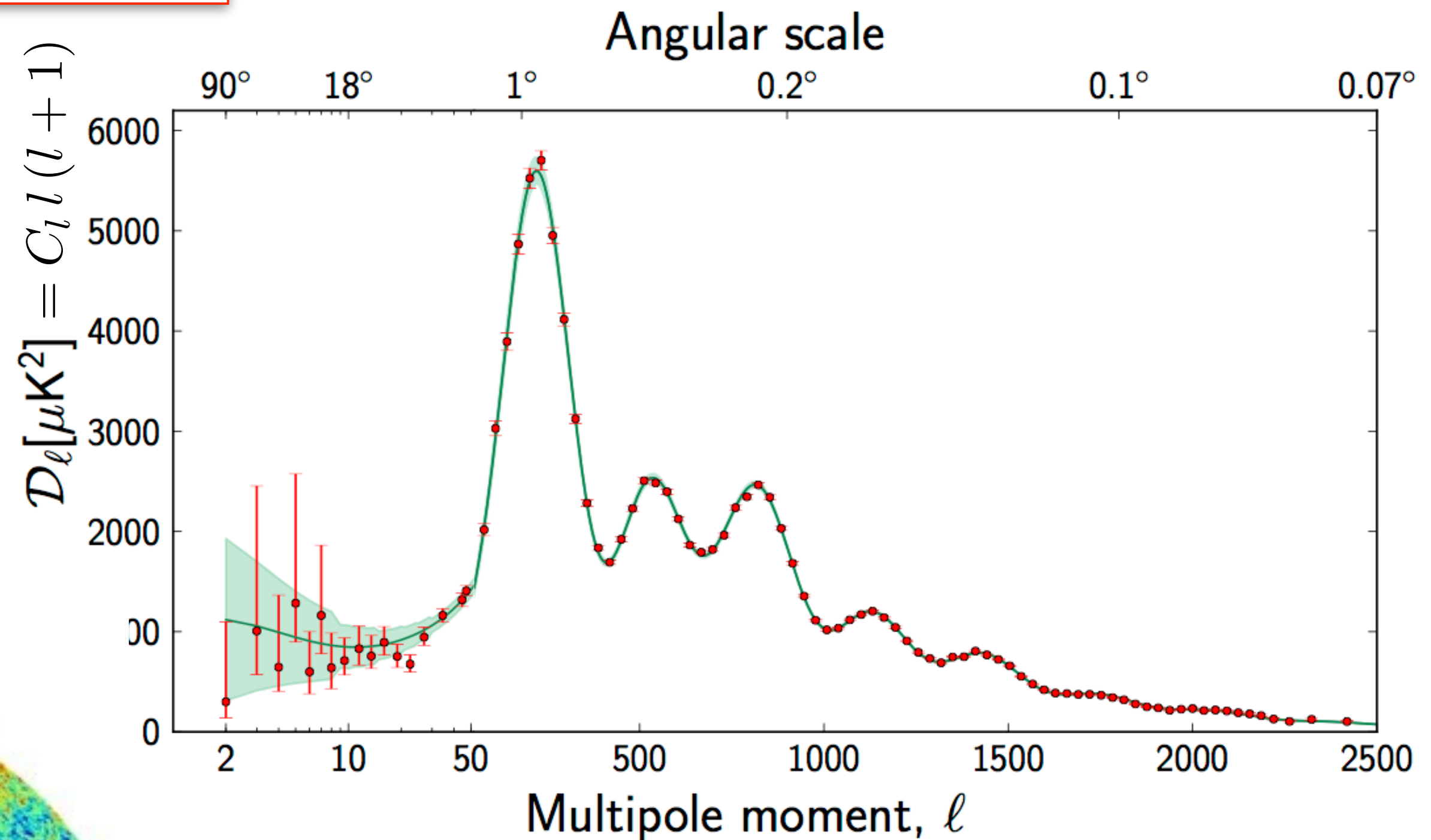
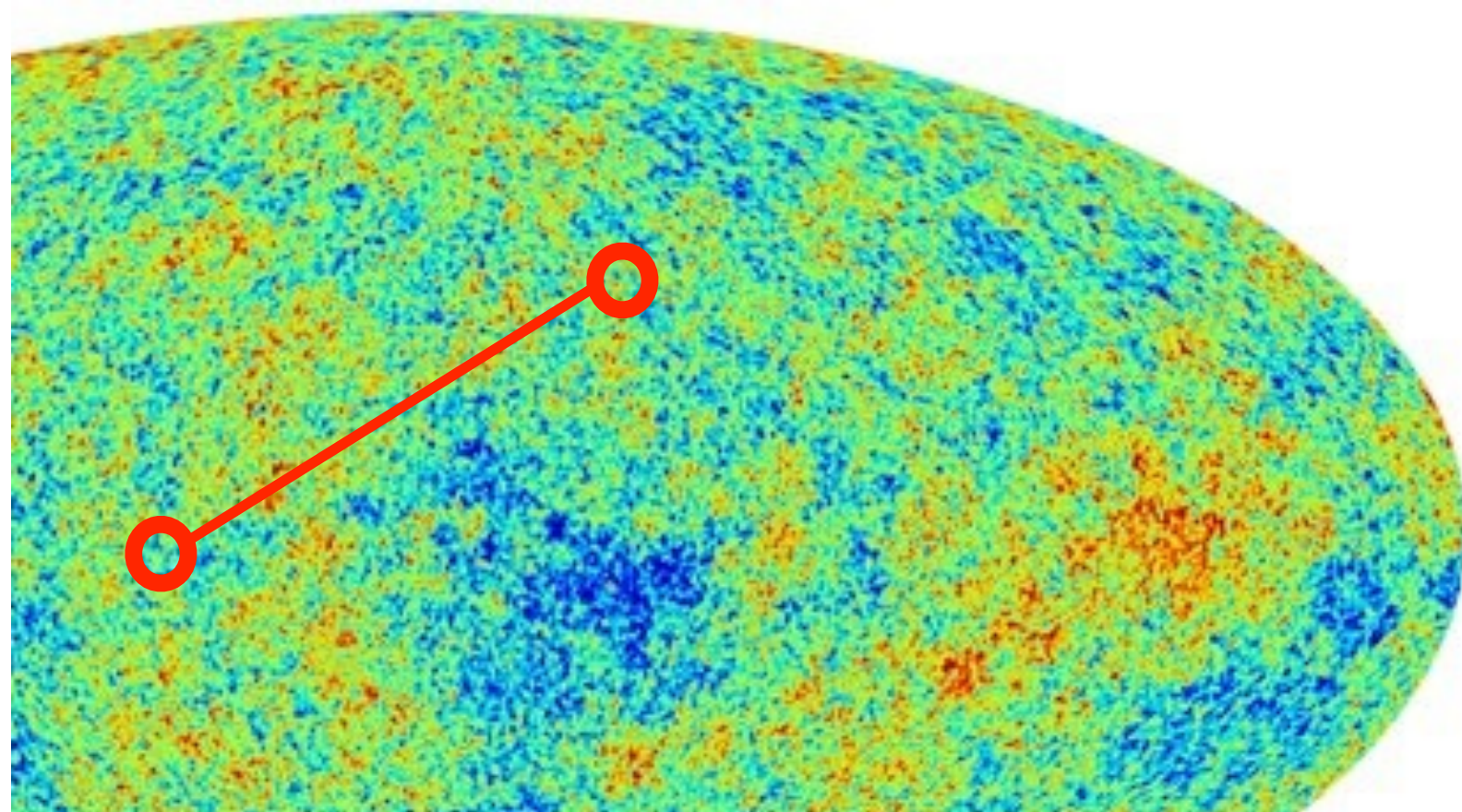


# Gaussian and non-Gaussian random fields

The statistical properties of a Gaussian random field are completely characterised by its 2-point correlation function. All higher-order, *connected* correlation functions are vanishing

$$\delta_T(\hat{n}) \equiv \frac{T(\hat{n}) - \bar{T}}{\bar{T}}$$

$$\mathcal{P}[\delta_T(\hat{n})] = \frac{1}{\sqrt{2\pi\sigma_T^2}} e^{-\frac{1}{2} \frac{\delta_T^2}{\sigma_T^2}}$$



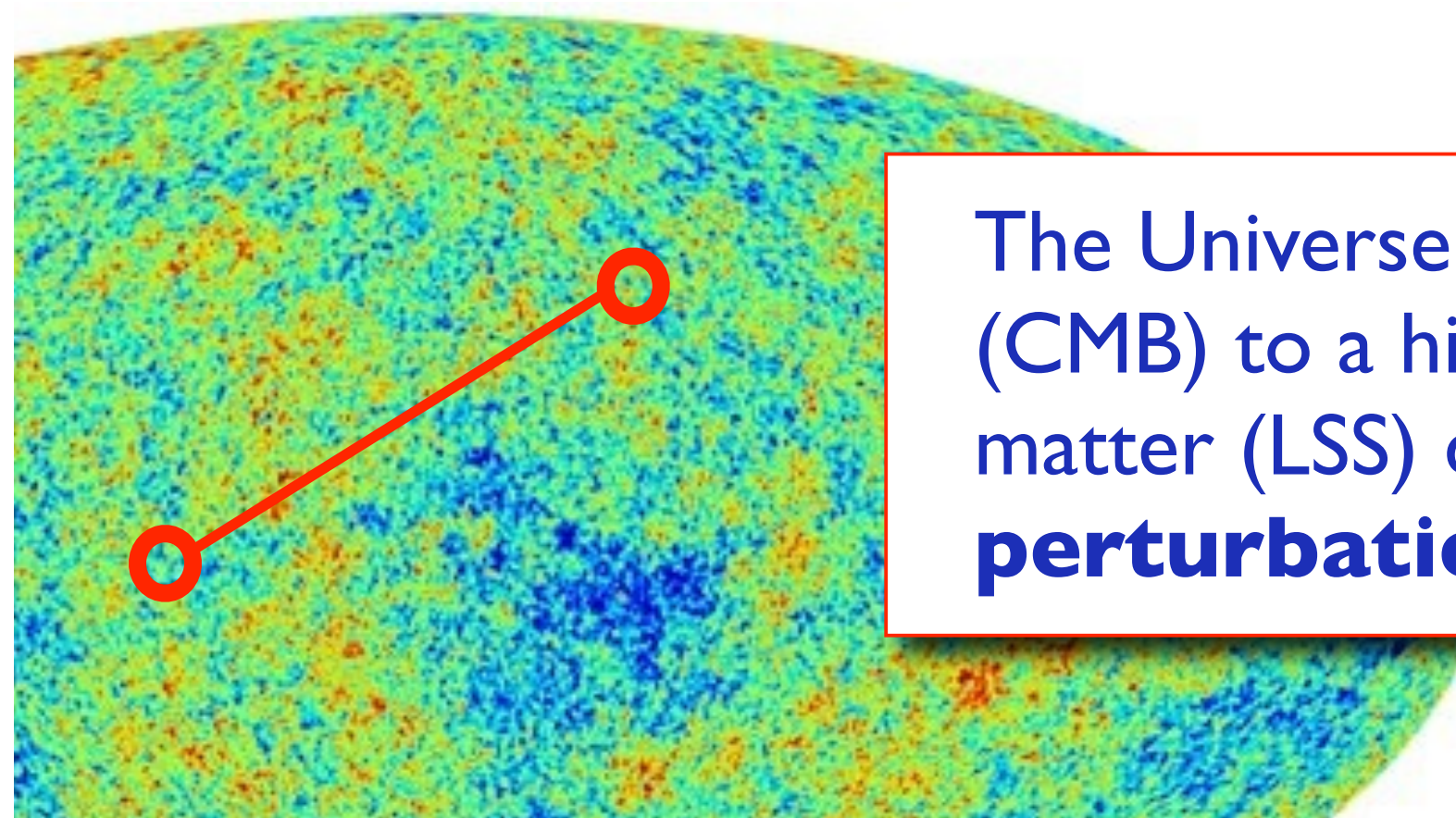
Perturbations in the CMB are one of the best examples of Gaussian random field



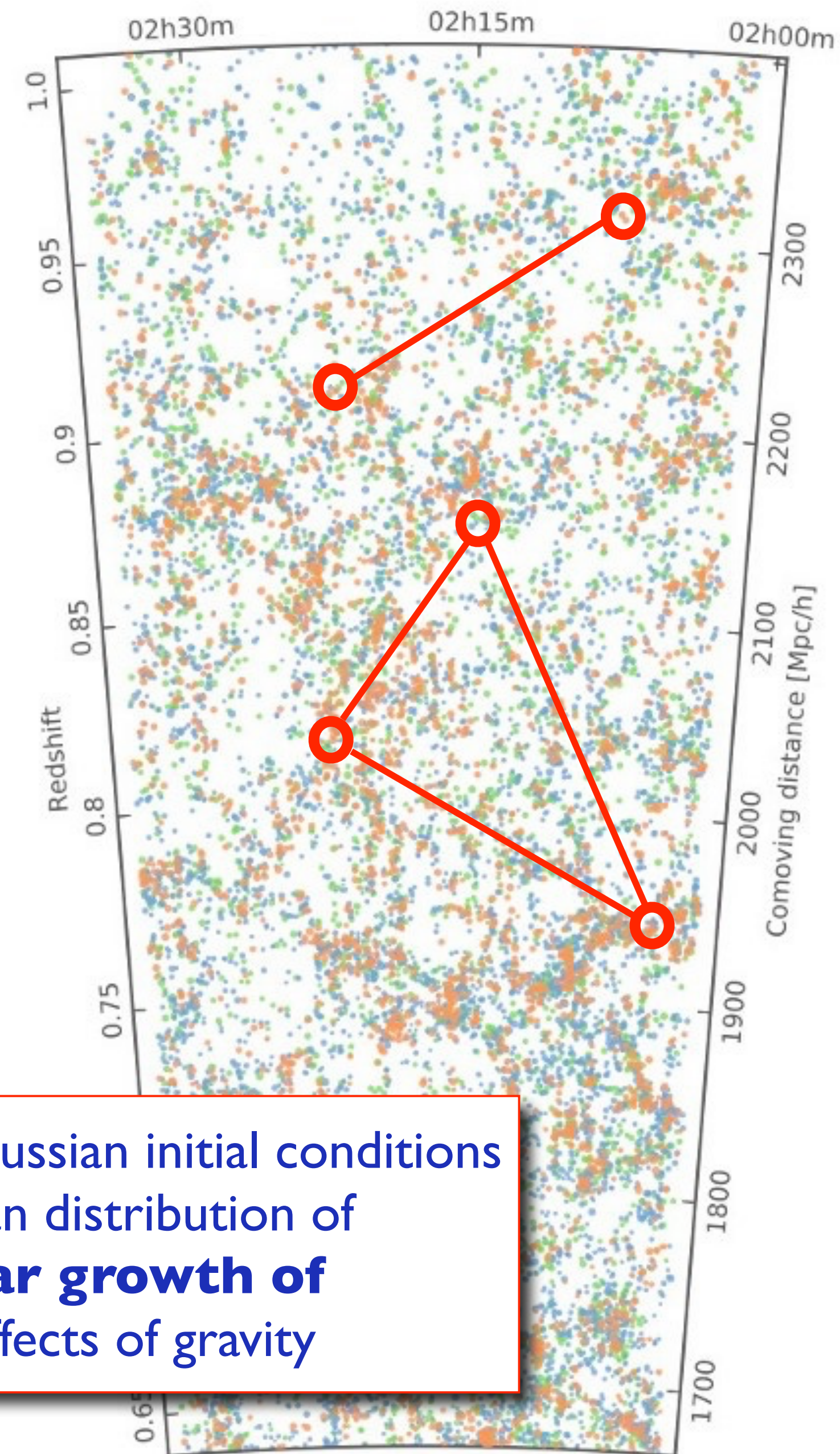
# Gaussian and non-Gaussian random fields

The statistical properties of a Gaussian random field are completely characterised by its 2-point correlation function. All higher-order, *connected* correlation functions are vanishing

all other random fields are non-Gaussian!



The Universe evolves from Gaussian initial conditions (CMB) to a highly non-Gaussian distribution of matter (LSS) due to **nonlinear growth of perturbations** under the effects of gravity





# Ergodic hypothesis

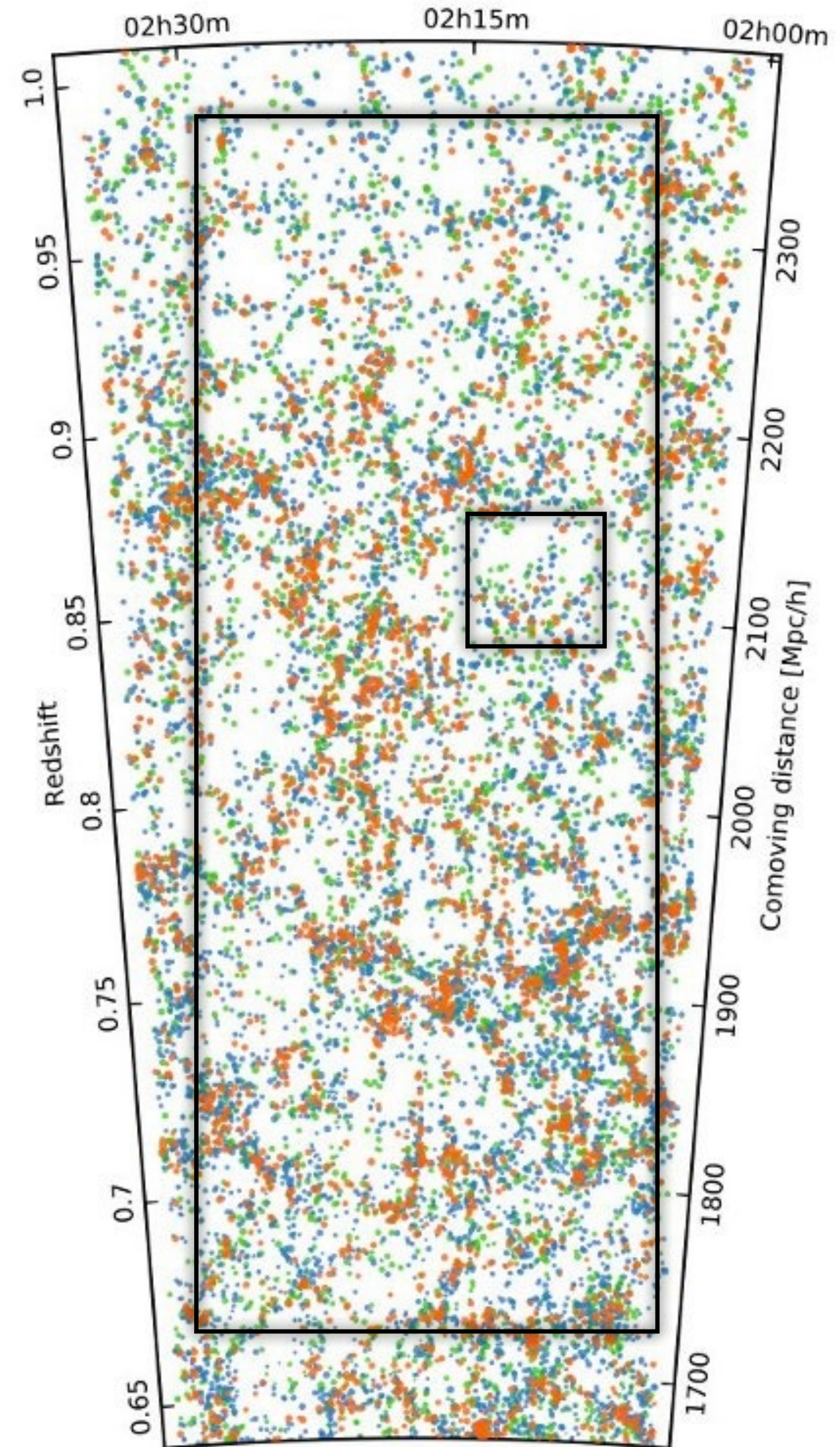
Expectation values, in principle, are to be intended as *ensemble averages*, i.e. averages over many “realisations of the Universe” ...

... but we only have one Universe!

We assume the **ergodic hypothesis**:  
**ensemble averages are equal to spatial averages**

$$\langle \phi(\vec{x}) \rangle \equiv \int d\phi \phi \mathcal{P}(\phi) \overset{\text{red arrow}}{=} \frac{1}{V} \int_V d^3x \phi(\vec{x})$$

We should make sure, however, that the observed volume correspond to a “fair sample” of the Universe

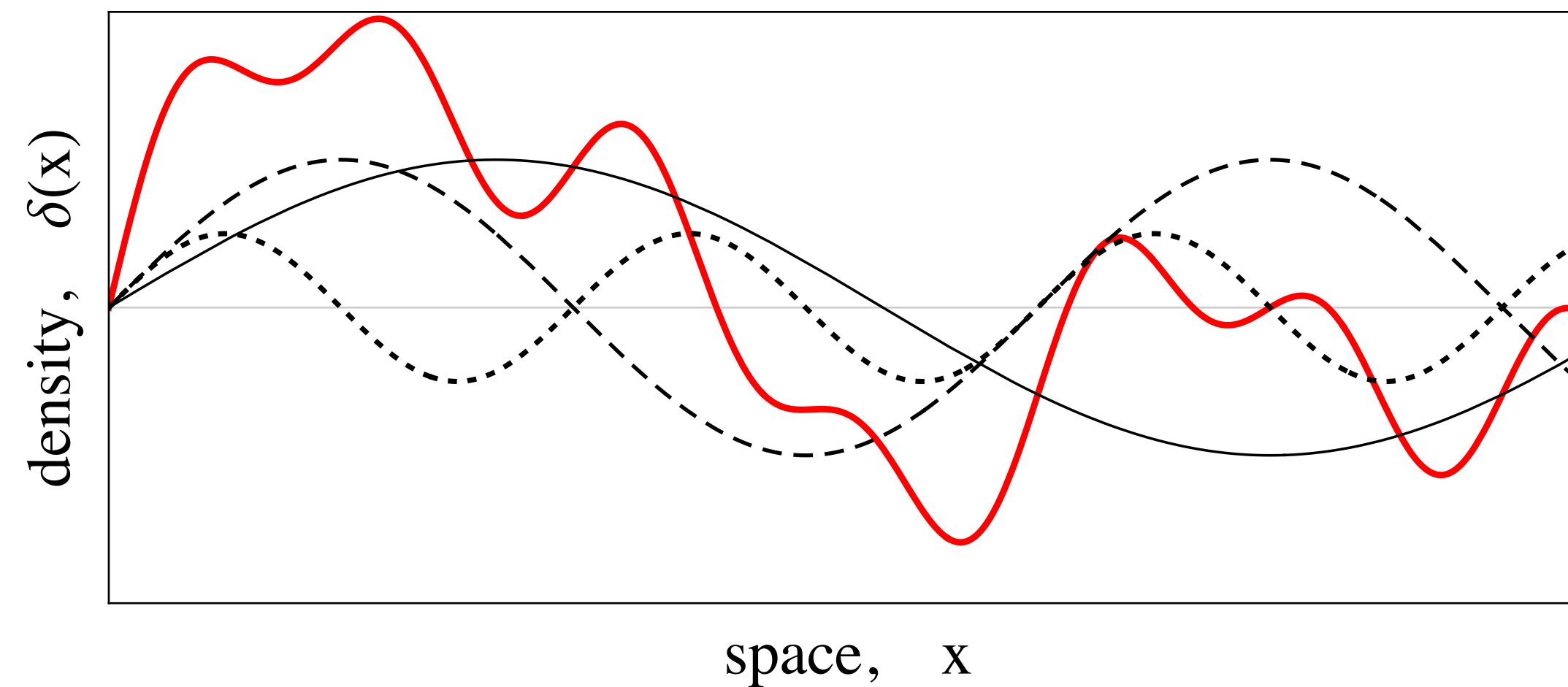




# Fourier space

Theoretical predictions for the matter correlation functions are performed in **Fourier space**

Fourier analysis naturally separates perturbations at different scales:



$$\delta_{\vec{k}} = \int \frac{d^3x}{(2\pi)^3} e^{-i\vec{k}\cdot\vec{x}} \delta(\vec{x})$$

$$\delta(\vec{x}) = \int d^3k e^{i\vec{k}\cdot\vec{x}} \delta_{\vec{k}}$$

- Since  $\delta(\vec{x})$  is a random field  $\delta_{\vec{k}}$  is also a random field
- Since  $\delta(\vec{x})$  is real  $\delta_{\vec{k}}^* = \delta_{-\vec{k}}$



# Fourier space: correlation functions

The 2-point function in Fourier space: the **power spectrum**

$$\langle \delta_{\vec{k}_1} \delta_{\vec{k}_2} \rangle = \delta_D(\vec{k}_1 + \vec{k}_2) P(k_1) \qquad P(k) = \int \frac{d^3x}{(2\pi)^3} e^{i\vec{k} \cdot \vec{x}} \xi(x)$$

  
homogeneity & isotropy

The power spectrum is the *Fourier Transform*  
of the 2-point correlation function

**The power spectrum is a measure of the  
amplitude of perturbations as a function of scale**

$$\Delta(k) \equiv 4\pi k^3 P(k) \qquad \text{adimensional power spectrum}$$

$$\sigma_\delta^2 \equiv \langle \delta^2(\vec{x}) \rangle = 4\pi \int dk k^2 P(k) = \int \frac{dk}{k} \Delta(k)$$