

An Overview of the NEGF Formalism

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Fundamental equations governing NEGF



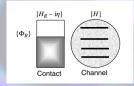


Figure: Channel contains no electrons and is disconnected from the contact where the electrons occupy the states described by $\{\Phi_R\}$

Fundamental equations governing NEGF



$$\begin{array}{ccc} \text{contact} & \text{device} \\ \text{contact} & \begin{pmatrix} EI_R - H_R + i\eta & -\tau^\dagger \\ -\tau & EI - H \end{pmatrix} \begin{pmatrix} \Phi_R + \chi \\ \psi \end{pmatrix} = \begin{pmatrix} S_R \\ 0 \end{pmatrix}$$

$$G_R = [EI_R - H_R + i\eta]^{-1}$$
$$\{\chi\} = G_R \ \tau^{\dagger}\{\psi\}$$
$$\{\psi\} = G \ \{S\}$$

where

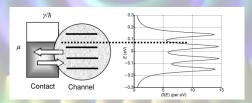
$$G = [EI - H - \Sigma]^{-1}$$

$$\Sigma = \tau G_R \tau^{\dagger}, \quad S = \tau \Phi_R$$

Level broadening and LDOS



The contact/reservoir has a near-continuous distribution of energy levels, whereas the channel has a discrete set of energy levels. When the channel is coupled to a contact, one would expect an intermediate state of such energy levels- something in between a continuous distribution and discrete levels. This also implies that the energy levels within the channel are no longer eigenvalues of the Hamiltonian, [H], but act as independent variables.



Level broadening and LDOS



LDOS:

$$D(\vec{r}; E) = \sum_{\alpha} |\Phi_{\alpha}(\vec{r})|^2 \delta(E - \epsilon_{\alpha}),$$

which is the diagonal element of the spectral matrix, [A(E)]:

$$A(\vec{r}, \vec{r'}; E) = 2\pi \sum_{\alpha} \phi_{\alpha}(\vec{r}) \ \delta(E - \epsilon_{\alpha}) \ \phi_{\alpha}^{*}(\vec{r'})$$

In matrix notation, $[A(E)] = 2\pi\delta(E[I] - [H])$, $D(E) = \frac{1}{2\pi}\text{Tr}[A(E)]$

Level broadening and LDOS



$$2\pi\delta(E - \epsilon_{\alpha}) = \left[\frac{2\eta}{(E - \epsilon_{\alpha})^2 + \eta^2}\right]_{\eta \to 0^+}$$

the spectral matrix can be expressed as:

$$[A(E)] = i[G(E) - G^{\dagger}(E)]$$

where

Retarded Green's function :
$$G(E) = [(E + i0^+)I - H]^{-1}$$

Advanced Green's function : $G^{\dagger}(E) = [(E - i0^+)^{-1}I - H]^{-1}$

Note that the above expressions hold true for the reservoir/contact Green functions. The equivalent device/channel Green functions are given by:

Retarded Green's function :
$$G(E) = [(E + i0^+)I - H - \tau G_R \tau^{\dagger}]^{-1}$$

Advanced Green's function : $G^{\dagger}(E) = [(E - i0^+)^{-1}I - H - \tau G_R \tau^{\dagger}]^{-1}$

Green's functions in time domain



 $[\tilde{G}^R(t)]$: Fourier transform of $[G^R(E)]$

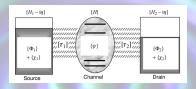
$$[\tilde{G}^R(t)] = -rac{i}{\hbar}
u(t)e^{-i0^+t} egin{bmatrix} e^{-i\epsilon_1t/\hbar} & 0 & 0 & \dots \ 0 & e^{-i\epsilon_2t/\hbar} & 0 & \dots \ 0 & 0 & e^{-i\epsilon_3t/\hbar} & \dots \ \dots & \dots & \dots & \dots \end{pmatrix}$$

Similarly,

$$[ilde{G}^{A}(t)]=rac{i}{\hbar}
u(-t)e^{i0^+t}egin{bmatrix} e^{-i\epsilon_1t/\hbar} & 0 & 0 & \dots \ 0 & e^{-i\epsilon_2t/\hbar} & 0 & \dots \ 0 & 0 & e^{-i\epsilon_3t/\hbar} & \dots \ \dots & \dots & \dots \end{bmatrix}$$

Channel connected to two contacts





$$\begin{pmatrix} EI - H_1 + i\eta & -\tau_1^{\dagger} & 0 \\ -\tau_1 & EI - H & -\tau_2 \\ 0 & -\tau_2^{\dagger} & EI - H_2 + i\eta \end{pmatrix} \begin{pmatrix} \Phi_1 + \chi_1 \\ \psi \\ \Phi_2 + \chi_2 \end{pmatrix} = \begin{pmatrix} S_1 \\ 0 \\ S_2 \end{pmatrix}$$

Channel Green's function, $G = [EI - H - \Sigma_1 - \Sigma_2]^{-1}$

Source term,
$$S = S_1 + S_2 = \tau_1 \{ \Phi_1 \} + \tau_2 \{ \Phi_2 \}$$

Spectral function,
$$A = A_1 + A_2 = G(\Gamma_1 + \Gamma_2)G^{\dagger}$$

Channel connected to two contacts



$$[\rho] = \int_{-\infty}^{+\infty} \frac{dE}{2\pi} [G^n(E)]$$

$$I_i = -\frac{q}{h} \int_{-\infty}^{+\infty} dE \left\{ \underbrace{\mathsf{Tr}[\Gamma_i A] f_i}_{\mathsf{Inflow}} - \underbrace{\mathsf{Tr}[\Gamma_i G^n]}_{\mathsf{Outflow}} \right\}$$

$$\overline{T}(E) = \mathsf{Tr}[\Gamma_2 G \Gamma_1 G^{\dagger}]$$

where $[G^n] = [A_1]f_1 + [A_2]f_2 = G\Sigma_{in}G^{\dagger} = G\{[\Gamma_1]f_1 + [\Gamma_2]f_2\}G^{\dagger}$



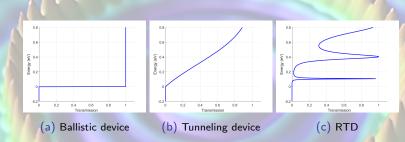


Figure: Transmission function for different devices



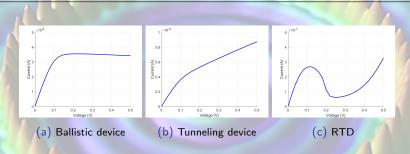


Figure: I-V characteristics for different devices

2-D conductors





$$\begin{array}{c|c}
\varepsilon_1 \\
\hline
\varepsilon_2 & t \\
\hline
\varepsilon_3 \\
\hline
\beta^+ & \overline{\beta}
\end{array}$$

where

$$\alpha = \begin{bmatrix} \epsilon & t & 0 \\ t & \epsilon & t \\ 0 & t & \epsilon \end{bmatrix}, \ \beta = \begin{bmatrix} t & 0 & 0 \\ 0 & t & 0 \\ 0 & 0 & t \end{bmatrix}$$

A basis transformation can be done such that α is diagonalized:

$$\tilde{\alpha} = V^{\dagger} \alpha V, \quad \tilde{\beta} = \beta$$

In the transformed basis, the 2D conductors can be visualized as a set of independent 1D conductors, with diagonal elements $\epsilon_n = \epsilon - 2t \cos(k_n a)$, where $k_n a = \frac{n\pi}{N+1}$.

2-D conductors



Previously shown transformations convert the lattice basis to an abstract mode basis.

$$\overset{\tilde{X}}{\underbrace{X}} = V^{\dagger} & \overset{X}{\underbrace{X}} & V$$
Mode basis
$$\overset{X}{\underbrace{X}} = V & \overset{X}{\underbrace{X}} & V^{\dagger}$$
Lattice basis
Mode basis

$$ilde{\Sigma}_1 = egin{bmatrix} te^{ik_1 a} & 0 & 0 & \dots \\ 0 & te^{ik_2 a} & 0 & \dots \\ 0 & 0 & te^{ik_3 a} & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

Graphene



$$\begin{array}{c|c} \Sigma_1 & [H] & \Sigma_2 \\ \hline & \overline{a} \\ \hline & \varepsilon \\ \hline & \overline{\epsilon} \\ \hline & \underline{\beta}^+ \alpha & \underline{\beta} \\ \end{array}$$

Recursive relation:

$$g_2^{-1} = (E + i0^+)I - \alpha - \beta^{\dagger}g_2\beta$$

Graphene







