

Unveiling the Surface Structure of Amorphous Solid Water via Selective Infrared Irradiation of OH Stretching Modes

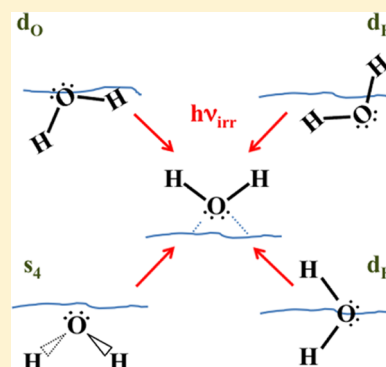
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ABSTRACT: In the quest to understand the formation of the building blocks of life, amorphous solid water (ASW) is one of the most widely studied molecular systems. Indeed, ASW is ubiquitous in the cold interstellar medium (ISM), where ASW-coated dust grains provide a catalytic surface for solid phase chemistry, and is believed to be present in the Earth's atmosphere at high altitudes. It has been shown that the ice surface adsorbs small molecules such as CO, N₂, or CH₄, most likely at OH groups dangling from the surface. Our study presents completely new insights concerning the behavior of ASW upon selective infrared (IR) irradiation of its dangling modes. When irradiated, these surface H₂O molecules reorganize, predominantly forming a stabilized monomer-like water mode on the ice surface. We show that we systematically provoke “hole-burning” effects (or net loss of oscillators) at the wavelength of irradiation and reproduce the same absorbed water monomer on the ASW surface. Our study suggests that all dangling modes share one common channel of vibrational relaxation; the ice remains amorphous but with a reduced range of binding sites, and thus an altered catalytic capacity.

SECTION: Spectroscopy, Photochemistry, and Excited States



Amorphous solid water (ASW) is a molecular system that has long provoked interest due, in part, to its role in the formation of molecules key to the origins of life.^{1–5} ASW has long been known to accrete small molecules such as CO, H₂O, N₂, or CH₄,^{2,6,7} initiating chemical and photochemical surface reactivity.^{2,8,9} In the interstellar medium (ISM), water in the form of ASW is the most abundant solid phase molecular species.^{10,11} The production of molecules, from the most simple, H₂,^{12,13} to the more complex CH₃OH,⁸ and even precursors to the simplest amino acid, glycine,¹⁴ is catalyzed by the ASW surface;¹⁵ both the outer surface and surfaces within its porous structure are involved. The selective IR irradiation of crystalline ice^{16,17} and water clusters¹⁸ has already been studied. In the former case, the desorption of H₂O molecules, and in the latter, the dissociation of clusters, was stimulated. We are interested in the behavior of ASW upon selective IR irradiation and have studied the irradiation of the four surface modes of this ice, assigned in the literature^{1,2} and illustrated in Figure 1. Theoretical calculations, supported by experimental studies, suggest that water molecules in the dH mode are bi- or tricoordinated, presenting one free OH bond dangling at the surface; dO molecules present a free oxygen electronic doublet; and s4 molecules have a tetrahedral structure at the surface, which is distorted compared to the tetrahedra of bulk ASW.

In these experiments we prepared a pure ASW sample (Figure 1) as follows: deionized water was subjected to multiple freeze–pump–thaw cycles under vacuum to remove dissolved gases. Mixtures of purified H₂O and helium (Air Liquide, ≥

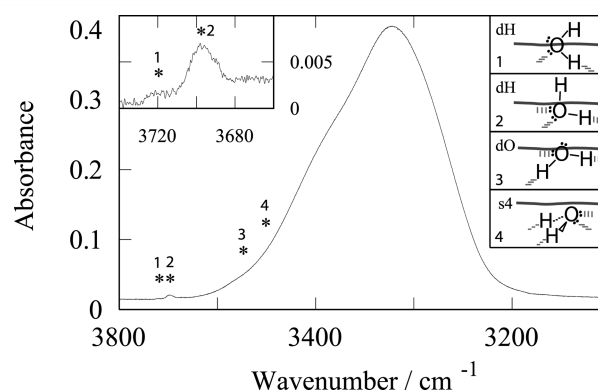


Figure 1. The four surface modes of water ice—dH (3720 and 3698 cm^{−1}), dO (3549 cm^{−1}), and s4 (3503 cm^{−1})—are illustrated on a representative ASW sample.

99.9999%) gas were prepared in a stainless steel dosing line with base pressure 10^{−4} mbar in ratios of H₂O: He = 1: 25. Ices were produced by depositing the gas mixture directly onto a gold-plated copper surface held at 50 K (to avoid trapping of the vector gas or nitrogen) then cooled to 3.7 K (the cooling typically takes around five minutes due to the high cryogenic power of 0.5 W at 4 K). The cooled surface is located in a high

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vacuum chamber with a base pressure of 10^{-8} mbar at 3.7 K. IR spectra were recorded in reflection mode using a Bruker 66/S FTIR spectrometer equipped with a MCT detector ($4000\text{--}800\text{ cm}^{-1}$). Full details of the experimental setup are given in the work by Coussan et al.¹⁹ The ices grown in our study were characterized as purely amorphous in nature due to the position of the bulk OH stretch and the characteristic dangling modes at 3720 and 3698 cm^{-1} (see Figure 1).^{1,2} Some previous studies using a helium carrier gas have reported the formation of ice nanocrystals or clusters,²⁰ but this can be ruled out due to the absence of an absorption feature centered at 3692 cm^{-1} . This band was never observed in any deposition of a $\text{H}_2\text{O}:\text{He}$ mixture during the work detailed here. After deposition, ices were selectively irradiated using a tunable IR OPO Laserspec ($1.5\text{--}4\text{ }\mu\text{m}$), pumped at 10 Hz by a pulsed Nd:YAG Quantel Brilliant B laser (1064 nm, pulse duration 6 ns). The average laser power is $\approx 35\text{ mW}$ in the ν_{OH} domain, except in the range $3520\text{--}3500\text{ cm}^{-1}$, where it is $\approx 10\text{ mW}$, with a fwhm $\geq 1.5\text{ cm}^{-1}$. Each irradiation was performed for 1 h to ensure saturation of the effects.

Figure 2 shows the results of irradiations carried out on the four dangling mode bands. The irradiations provoke permanent “hole-burnings” in the irradiated band, which are slightly shifted (by up to 3 cm^{-1}) compared to the irradiation frequency. The “hole-burning” is clearest for the irradiation of dH at 3698 cm^{-1} but is also easily visible for the irradiations at 3720 and 3549 cm^{-1} . Interestingly, upon each irradiation, we observe the growth of a new band centered at $\sim 3725\text{ cm}^{-1}$, with fwhm $\approx 5\text{ cm}^{-1}$. Both the shifts in frequency upon irradiation, and the narrowness of the “hole-burning” effects and the newly created bands at 3725 cm^{-1} illustrate the inhomogeneity of the bands; each band contains a distribution of oscillators, but only one class of oscillator isomerizes upon irradiation at a given frequency, producing one new oscillator class. After irradiation at 3698 cm^{-1} , a second, smaller peak at 3638 cm^{-1} is also observed (see Figure 3a). Unirradiated ASW samples and their IR spectra were stable and remained unchanged over the time scale of an irradiation study.

The common feature of all the irradiations in Figure 2 is the growth of the 3725 cm^{-1} band. What is the source of this new band, previously unidentified in ASW spectra? Considering energetics only, irradiating between 3720 and 3503 cm^{-1} could potentially break one or two hydrogen bonds, as the average weak H-bond strength is around 1800 cm^{-1} . Is the band, therefore, due to H_2O surface molecules in a different conformation than those in dH, dO and s4 modes, or is it due to water molecules that have desorbed then readsorbed at the surface? The latter response can be immediately discarded based on the results of molecular adsorption studies, which have shown a red-shift of the OH dangling frequency upon adsorption of multiple molecular species at the dangling modes.^{2,4,21} For example, in the case of nitrogen adsorption onto ASW, Manca et al.²¹ observed a red shift of 22 cm^{-1} of the dH mode.^{22,23} In our experiments, the new band at 3725 cm^{-1} is blue-shifted with respect to the dH modes. Moreover, the dynamic vacuum of 10^{-8} mbar rapidly evacuates any desorbing molecules, such that the residual pressure is too low to allow redeposition. Thermal effects are discounted based upon annealing of ASW samples, which provoked a global decrease of the dangling bonds and no production of narrow peaks at 3725 cm^{-1} , in agreement with previous studies.² Upon irradiation, the newly produced band is narrow, indicating a single, homogeneous vibrational mode. Thus, the 3725 cm^{-1}

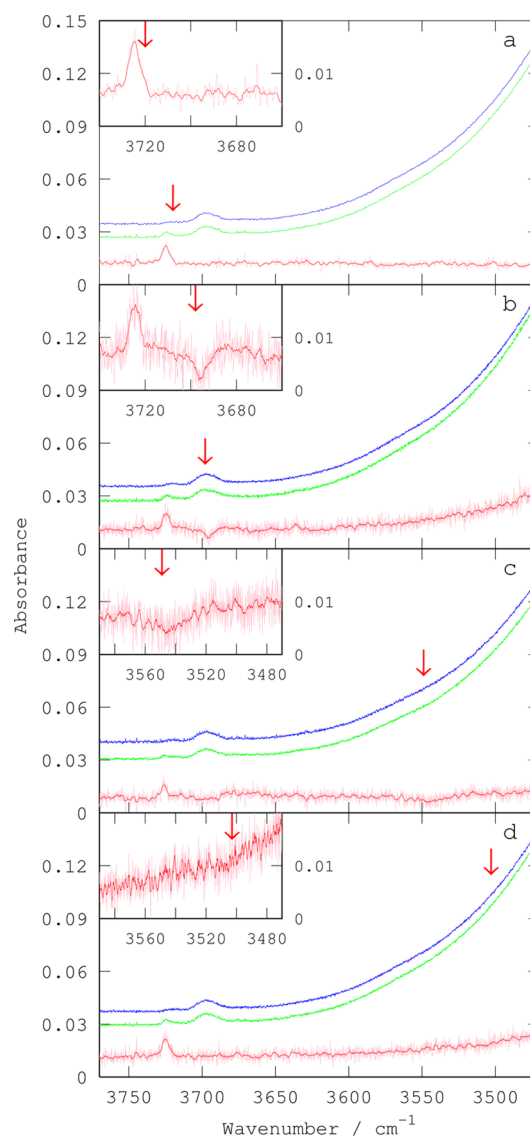


Figure 2. Effect of selective irradiation. Panels show the effects of selective irradiation at (a) 3720 cm^{-1} , (b) 3698 cm^{-1} , (c) 3549 cm^{-1} , and (d) 3503 cm^{-1} . The spectra are colored as follows: before irradiation, blue; after irradiation, green; difference spectrum, red. The insets magnify the irradiated range.

band is clearly not due to a perturbed dH mode, but it rather has dangling OH character, and we propose that it samples a bicoordinated H_2O molecule, with two dangling OH and two coordinated electron pairs (d2H). This structure explains the blue shift of the peak with respect to the dH mode at 3720 cm^{-1} as, because neither of the two OH oscillators is directly hydrogen bonded, the free OH oscillators are less perturbed than the surface dH modes.

We performed further experiments to verify the source of the two new peaks at 3725 and 3638 cm^{-1} , as illustrated in Figure 3. The background deposition of water on a pure ASW sample does not result in a new band, but rather the global growth of the dH and bulk OH bands (Figure 3b). However, subsequent background deposition of pure nitrogen, a molecule present as a low-level source of pollution in the chamber, provokes the appearance of a band at 3726 cm^{-1} (Figure 3c), revealed thanks to the magnifying effect of nitrogen.²⁴ The background deposition of N_2 upon a new ASW sample does not produce

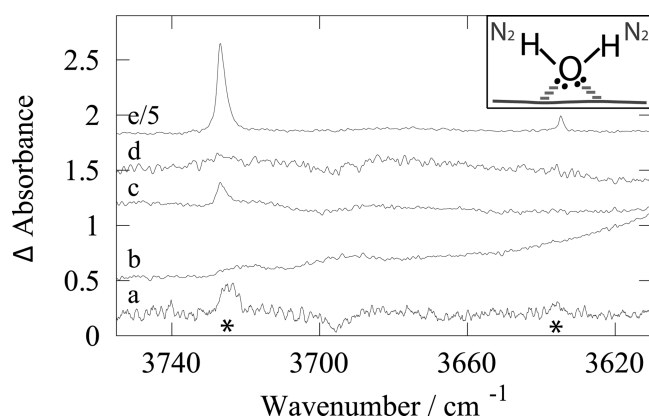


Figure 3. Background deposition experiments. We present the difference spectra of (a) ASW irradiated at 3698 cm^{-1} , reproduced from Figure 2; (b) an ASW sample plus background-deposited water; (c) spectrum b plus background-deposited nitrogen; (d) an ASW sample plus background-deposited nitrogen; and (e) spectrum d plus background-deposited $\text{H}_2\text{O}:\text{N}_2$ mixture.

a clear peak (Figure 3d), but the background deposition of a $\text{H}_2\text{O}:\text{N}_2$ (1:10) mixture provokes the appearance of two narrow bands at 3725 cm^{-1} and 3638 cm^{-1} (Figure 3), with a peak area ratio of approximately 10:1, compared to 6:1 seen after irradiation at 3698 cm^{-1} . When this deposition spectrum is compared to the spectrum of a gas-phase water monomer (see Table 1), the two peaks show red shifts of nearly 200 cm^{-1} with respect to the ν_3/ν_1 modes of the monomer (one always compares gas-phase experimental results with those of a “classical” theoretical calculation, as classical calculations are carried out for isolated species. It also allows us to approximate the perturbation induced by the medium). These red shifts are almost precisely those observed in the solid phase for water in a nitrogen matrix, where the $\nu_3:\nu_1$ ratio is 8:1, and a $\text{H}_2\text{O}-\text{N}_2$ complex.^{25,27} Considering our experimental results and the literature, the band at 3725 cm^{-1} is positively attributed to a water monomer interacting with the surface via its two electronic doublets; its large intensity compared to the other dH bands is explained by the magnifying effect of nitrogen, as extensively investigated by Hujo et al.²⁴ We suggest that nitrogen molecules, present as a low-level pollutant in the chamber, serendipitously complex the water molecule, as illustrated in Figure 3, stabilizing the molecule, preventing any further adsorption, and magnifying the OH stretching bands. However, the nitrogen must only be present as a trace pollutant on the surface, as no shifting of the dH peak positions of the ASW is seen, either during cooling of the water ice sample from 50 to 3.7 K, or during the irradiation period. If nitrogen were present at multilayer concentrations, we would expect to see a red shift^{21–23} of up to 22 cm^{-1} . The ν_1 peak at 3638 cm^{-1} is not observed after irradiation of the other three

surface modes, but this is likely due to its low intensity compared to that of ν_3 .

One explanation for the observed “hole-burning” is that amorphous ice is unable to relax all of the vibrational energy injected at the surface through bulk relaxation channels. As a result, some fraction of this energy is accumulated at the surface and in the immediate sublayers, where it induces reconstruction of the surface. We observe saturation of the “hole-burning” events within the time scale of the irradiations performed, suggesting that the rearrangement of surface molecules is not an efficient relaxation channel, but is a minority effect compared to the main relaxation channels via the bulk ice. Thus, the production of the 3725 cm^{-1} band is not the major result of irradiation, but is due to the inability of the ice to fully dissipate the injected energy. It has previously been suggested that ASW is a disordered material that has no long-range organization,¹ which could help to explain the lack of efficiency in the bulk relaxation channels. Although it is possible that some H_2O molecules desorb upon irradiation, as in the studies of Focsa et al.,^{16,17} we consider this effect to be minor because no increase in pressure is observed in the chamber, and the energy injected is not enough to break more than two H-bonds.

It is curious that the absorption band at 3725 cm^{-1} is produced upon irradiation of all dangling bonds. The irradiation effects at 3720 and 3698 cm^{-1} , in particular, are very similar, except that the “hole-burning” of the doubly coordinated dH is less pronounced than that of the triply coordinated dH. Such molecular rearrangement requires only a reorientation of the water molecule and thus is not “energy consuming” compared to the energy injected into the system. These results also provide evidence of a local ordering to the ASW structure. Upon irradiation of the dH modes, both the “hole-burning” in absorption bands and the newly produced monomer band at 3725 cm^{-1} are relatively narrow (fwhm $\approx 5\text{ cm}^{-1}$). This suggests that each surface molecule is surrounded by a locally ordered oscillator network, resulting reproducibly in the production of one oscillator class upon selective irradiation. If we consider irradiation of the s4 band, centered at 3503 cm^{-1} , we were unable to observe a definitive “hole-burning” event, likely due to the width of the band. However, we observed an increase at 3725 cm^{-1} , as for the dH modes. It is unlikely that the tetra-coordinated s4 molecules are themselves ejected from within the surface layer, as this would be highly energetically unfavorable. However, as we see an increase at 3725 cm^{-1} , it is likely that the relaxation channel involves the breaking of H-bonds, with two breaks being enough to “transform” a s4 molecule into a monomer-like water molecule. The case of dO is likely intermediate between those of dH and s4. During each of these irradiations, we see no interconversion between the modes, suggesting that there is only a single “surface” channel for the release of excess vibrational energy.

In this work we have provided new insights into the bonding and structure of the surface molecules in ASW. The ensemble

Table 1. Comparison between Calculated and Observed ν_3 and ν_1 Water Monomer Frequencies^a

vibrational mode	calculated ^b (gas phase)	gas phase ^c	observed		present work
			N_2 matrix ^d	$\text{H}_2\text{O}-\text{N}_2$ complex ^b	
ν_3	3924	3943	3728	3730	3725
ν_1	3822	3832	3635	3640	3638
$\Delta_\nu = \nu_3 - \nu_1$	102	111	93	90	87

^aFrequencies are given in cm^{-1} . ^bReference 25. ^cReference 26. ^dReference 27.

of our results show that surface modes, in particular the dH dangling bonds, are sensitive to photoinduced rearrangement due to competition between surface reorganization and the main relaxation channels in the bulk water ice. Rather than desorbing from the surface, molecules embedded in the surface layer become loosely associated with the surface in the form of monomer-like structures interacting through their two free electron pairs. The fortuitous presence of nitrogen in the chamber both promoted the magnification of the OH stretching mode of the monomer-like molecule and stabilized it on the ice surface. Inducing such conformational changes in an ASW surface potentially alters its physicochemical properties, most notably its catalytic potential.

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Notes

The authors declare no competing financial interest.

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