

# FORMATION OF COMPACT AMORPHOUS H<sub>2</sub>O ICE BY CODEPOSITION OF HYDROGEN ATOMS WITH OXYGEN MOLECULES ON GRAIN SURFACES

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## ABSTRACT

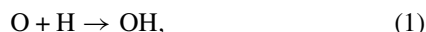
Formation of H<sub>2</sub>O molecules through the codeposition of oxygen molecules and hydrogen atoms is examined in situ using IR spectroscopy at 10–40 K under various O<sub>2</sub> and H fluxes. It is found that H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub> are continuously formed by reaction, even at 40 K. The H<sub>2</sub>O ice formed is amorphous, but has a compact (not microporous) structure compared to vapor-deposited amorphous H<sub>2</sub>O ice, because dangling OH bonds are not observed in the IR spectrum. This is consistent with astronomical observations in molecular clouds and theoretical predictions, which suggest that hydrogenation of O<sub>2</sub> is one of the potential routes for reproducing these IR spectral characteristics. The composition of the ice formed by codeposition varies with the O<sub>2</sub>/H ratio and temperature. Although no data are available at present for the H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> ratio of ice in molecular clouds, this study suggests that hydrogenation of O<sub>2</sub> has a potential to yield a H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> ratio of 5 or more in molecular clouds.

*Key words:* atomic processes – ISM: clouds

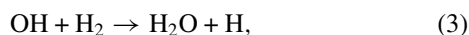
## 1. INTRODUCTION

Interstellar icy grain mantles consist of various molecules such as H<sub>2</sub>O, CO, CO<sub>2</sub>, H<sub>2</sub>CO, CH<sub>3</sub>OH, and NH<sub>3</sub>. Among these components, H<sub>2</sub>O is predominant in dense molecular clouds (e.g., Gibb et al. 2004). The elucidation of H<sub>2</sub>O formation in space is one of the most challenging objectives because of the ubiquity of H<sub>2</sub>O in space and its significant importance for chemical evolution and the origin of life. Theoretical studies predict that the formation of H<sub>2</sub>O in the gas phase is difficult to explain given the observed abundance in molecular clouds; therefore, it has been suggested that grain surface reactions are essential for the formation of H<sub>2</sub>O in such environments (d’Hendecourt et al. 1985; Hasegawa et al. 1992; Hasegawa & Herbst 1993; Cuppen & Herbst 2007).

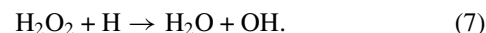
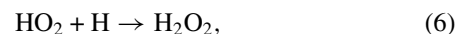
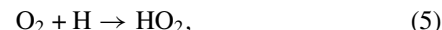
The simplest process for H<sub>2</sub>O formation on grain surfaces is the sequential hydrogenation of O atoms:



Because Reactions (1) and (2) are radical–radical reactions, these reactions have no activation barriers. Hiraoka et al. (1998) examined the above reaction pathways using D atoms instead of H atoms to irradiate O atoms in a N<sub>2</sub>O matrix at 12 K, and reported the detection of D<sub>2</sub>O by temperature programmed desorption (TPD) spectra through reactions of O atoms with D atoms. However, in their experimental procedures, it is not clear whether D<sub>2</sub>O was produced through D addition to O atoms at 12 K or during heating. Subsequent to Reaction (1), H<sub>2</sub>O is also produced through the reaction of OH with H<sub>2</sub>:



with an activation barrier of 2600 K (Schiff 1973). Tielens & Hagen (1982) have proposed an alternative route for the formation of H<sub>2</sub>O from an oxygen atom:



Reaction (5) has essentially no activation barrier (Walch et al. 1988), and Reaction (7) has an activation energy of approximately 2000 K (Koussa et al. 2006).

Miyauchi et al. (2008) have experimentally demonstrated that H<sub>2</sub>O is formed via the addition of H atoms to O<sub>2</sub> at 10 K, where solid O<sub>2</sub> was deposited at 10 K and H or D atoms were then irradiated onto the solid O<sub>2</sub> (hereafter, the experiment is referred to as the H irradiation experiment onto solid O<sub>2</sub>). H<sub>2</sub>O<sub>2</sub> formation was observed during the experiments, indicating that H<sub>2</sub>O is formed through Reactions (5)–(7). In addition, D<sub>2</sub>O is also formed through the successive addition of D to solid O<sub>2</sub>. They also reported that the ratios of rate constants were  $k_1\text{H}/k_1\text{D} = 1$ , while  $k_2\text{H}/k_2\text{D} = 8$ , where  $k_1\text{H}$  and  $k_2\text{H}$  are the rate constants for H atom addition in Reactions (5) and (7), respectively, and  $k_1\text{D}$  and  $k_2\text{D}$  are those for D atom addition, respectively. The former rate constant ratio is consistent with the fact that Reaction (5) has essentially no activation barrier (Walch et al. 1988). The latter result, where the ratio of  $k_2\text{H}/k_2\text{D} = 8$ , is reasonably explained as the isotope effect of a tunneling reaction with a reasonable activation energy of ~2000 K (Koussa et al. 2006).

After publication of Miyauchi et al. (2008), Ioppolo et al. (2008) studied the same reactions at 12–28 K. They confirmed our results qualitatively, but not quantitatively. The experiments performed by Ioppolo et al. (2008) were basically similar to those by Miyauchi et al. (2008) except for the type of atomic source, and flux and temperature of H atoms; a thermal-cracking-type H-atom source was used with a nose-shaped quartz pipe to cool down the H atoms to room temperature. The flux of H atoms was not measured, but was estimated to be  $5 \times 10^{13}$  atoms cm<sup>−2</sup> s<sup>−1</sup>. Different results for the rate constants of Reactions (5) and (7) were obtained:  $k_1\text{H} = k_2\text{H}$  and  $k_1\text{D} = k_2\text{D}$ , which are inconsistent with the barrier height estimations reported by Koussa et al. (2006) and Walch et al. (1988). Furthermore,  $k_1\text{H}/k_2\text{D} = 0.4$  and  $k_2\text{H}/k_2\text{D} = 0.7$  ratios

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were obtained, which indicate that the reaction of D atoms is faster than that of H atoms. These different results may be due to the method used to determine the reaction rates; different fitting models were used in the initial and latter stages of the reactions, of which the physicochemical meaning is unclear.

In the H irradiation experiment onto solid O<sub>2</sub>, the formation of H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub> stops at some stage, although some O<sub>2</sub> molecules remain unreacted (Miyauchi et al. 2008; Ioppolo et al. 2008). A similar phenomenon was observed in experiments on the hydrogenation of CO (Watanabe et al. 2004). In order to explain the above observations Miyauchi et al. (2008) proposed a layer model, in which the reaction occurs only at the surface of the deposited O<sub>2</sub>, resulting in an onion-like structure of molecular layers with H<sub>2</sub>O on the top, H<sub>2</sub>O<sub>2</sub> in the middle, and O<sub>2</sub> at the bottom, as suggested by Watanabe et al. (2004). Probably due to the layer structure, abundant H<sub>2</sub>O in molecular clouds was not reproduced in the laboratory experiments (Miyauchi et al. 2008; Ioppolo et al. 2008), although the reaction kinetics were revealed by the H irradiation experiments onto solid O<sub>2</sub>. In molecular clouds, reactions of H atoms with solid O<sub>2</sub> having a thickness of several monolayers are almost not expected to occur, because O<sub>2</sub> molecules could be continuously formed on the surface through Reaction (4), and then react with H atoms, which eventually results in the formation of H<sub>2</sub>O<sub>2</sub> and H<sub>2</sub>O on the grain surface (Tielens & Hagen 1982). In order to simulate the reactions in molecular clouds, it is more appropriate that O<sub>2</sub> molecules and H atoms are supplied simultaneously onto a substrate (hereafter, the experiment is referred to as the O<sub>2</sub> and H codeposition experiment), with in situ monitoring of the products. Miyauchi et al. (2008) have demonstrated that Reactions (5)–(7) occur very quickly (<1 minute), even at 10 K, despite the significant activation barrier of Reaction (7) (Koussa et al. 2006). Therefore, it is expected that codeposition of H atoms with O<sub>2</sub> molecules leads to the continuous formation of H<sub>2</sub>O<sub>2</sub> and H<sub>2</sub>O under conditions closer to those of dense molecular clouds.

In addition to the mechanism for the formation of H<sub>2</sub>O, the IR spectral features for H<sub>2</sub>O ice in molecular clouds have been of significant interest; the OH-stretching mode of H<sub>2</sub>O, the so-called 3  $\mu$ m band, reflects the structure, composition, and temperature of H<sub>2</sub>O-dominated ice (Hagen et al. 1981, 1983). Therefore, the shape of the 3  $\mu$ m band will be useful for investigating the characteristics of H<sub>2</sub>O ice. It has been shown, both observationally (Smith et al. 1989) and theoretically (Kouchi et al. 1994), that the structure of water ice in molecular clouds is amorphous. However, in the H irradiation experiments onto solid O<sub>2</sub>, it was very difficult to discuss the structure of H<sub>2</sub>O ice, because the shape of the 3  $\mu$ m band assigned to H<sub>2</sub>O was obscured, due to the low levels of product overlapping with H<sub>2</sub>O<sub>2</sub> (Miyauchi et al. 2008; Ioppolo et al. 2008). In contrast, it is expected that more H<sub>2</sub>O could be formed during continuous codeposition of O<sub>2</sub> with H, which may provide more suitable IR spectral data for examination of the 3  $\mu$ m band.

In this study, the formation of H<sub>2</sub>O from O<sub>2</sub> molecules is re-examined at 10–40 K by the codeposition of O<sub>2</sub> with H atoms onto a substrate. We present new information regarding the temperature dependence of water formation and the structure of amorphous H<sub>2</sub>O ice produced by surface reactions, which have not been reported previously. The present study enables us to compare the experimental results with astronomical observations more appropriately than previous experiments.

## 2. EXPERIMENTAL DETAILS

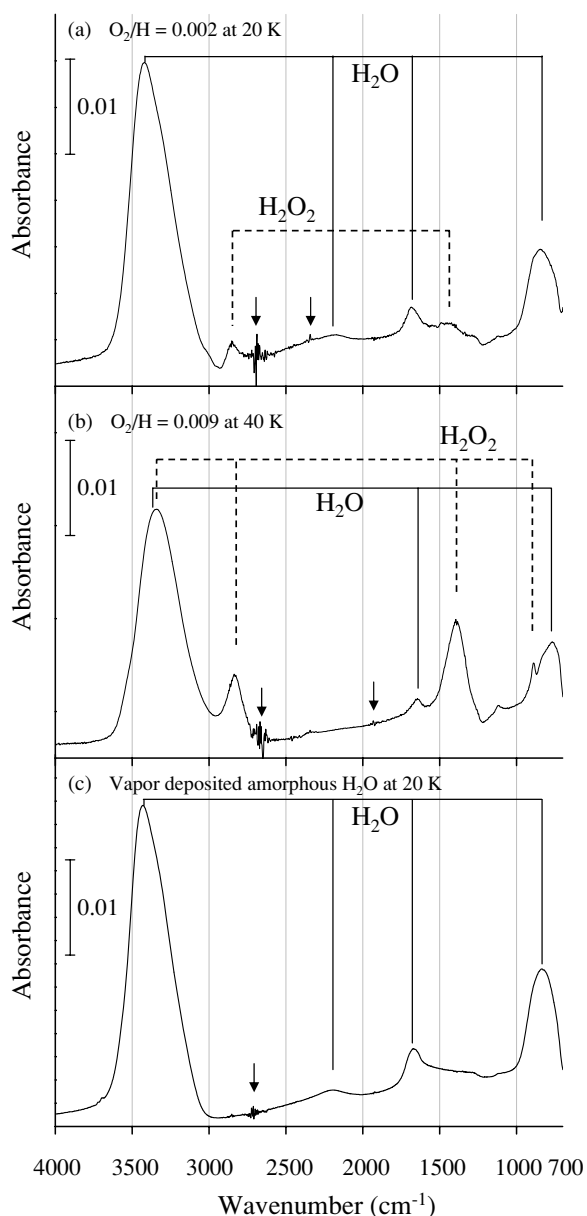
Experiments were performed using the Apparatus for Surface Reaction in Astrophysics (ASURA) system. ASURA consists of a main chamber, an atomic source, and a Fourier transform infrared (FTIR) spectrometer. Details of this apparatus have been described previously (Watanabe et al. 2006; Nagaoka et al. 2007). At the center of the main chamber, a mirror-finished aluminum (Al) substrate is mounted on the cold head of a He refrigerator. The base pressure of the main chamber is  $\sim 10^{-10}$  torr, but it reaches  $(1\text{--}2) \times 10^{-7}$  torr during operation of the atomic source. Atomic hydrogen (H) was produced by microwave-induced plasma in a Pyrex tube and transferred via a series of polytetrafluoroethylene (PTFE) and aluminum tubes to the substrate. The atomic beam was cooled to 100 K in the aluminum tube that was connected to another He refrigerator. The flux of H atoms in the present experiment was measured, using the method reported by Hidaka et al. (2007), to be  $2 \times 10^{14}$  atoms cm<sup>-2</sup> s<sup>-1</sup>. The O<sub>2</sub> molecules were introduced into the main chamber through a capillary plate located approximately 5 cm from the substrate, with an incident angle of 30° to the normal of the substrate surface. Although O<sub>2</sub> molecules were introduced at room temperature, they were immediately cooled down to the temperature of the substrate once they were absorbed. Because O<sub>2</sub> molecules are IR inactive, the deposition rate of O<sub>2</sub> cannot directly be determined by an IR spectrometer. Instead, for the preparatory experiment, CO molecules were introduced into the main chamber under a condition and the deposition rate of CO was first determined. Then, O<sub>2</sub> molecules were introduced into the chamber under the same deposition (gas flow) condition as that of the CO. Since the conductance of the gas line is considered to be almost the same for O<sub>2</sub> and CO gases, the O<sub>2</sub> deposition rate is nearly equal to that of CO. This assumption is based on the fact that the binding energy of CO on an H<sub>2</sub>O surface (1900 K) is very close to that of O<sub>2</sub> (1600 K) (Tielens 2005), and CO and O<sub>2</sub> have a close mass. Deposition rates of O<sub>2</sub> were varied from 7.5 to  $380 \times 10^{10}$  molecules cm<sup>-2</sup> s<sup>-1</sup>, which results in O<sub>2</sub>/H flux ratios of  $3.8 \times 10^{-4}$  to  $1.9 \times 10^{-2}$ . The O<sub>2</sub>/H ratio of  $3.8 \times 10^{-4}$  is the minimum value attainable with this experimental setup, while values larger than  $1.9 \times 10^{-2}$  are possible, but not utilized in this study. The reaction products were monitored in situ by infrared reflection-absorption spectroscopy with a resolution of 4 cm<sup>-1</sup> through the spectral range of 700–4000 cm<sup>-1</sup>. The temperature of the Al substrate was maintained at 10–40 K during each experiment. For the experiments with O<sub>2</sub>/H ratios of 0.019, 0.009, 0.002, and  $3.8 \times 10^{-4}$ , H and O<sub>2</sub> codeposition was conducted for 90, 120, 360, and 1020 minutes, respectively.

To investigate the effect of amorphous ice on the reactions, similar experiments were performed on a predeposited amorphous D<sub>2</sub>O ice with a thickness of  $\sim 30$  monolayer (ML), where 1 ML indicates the amount of D<sub>2</sub>O to form  $10^{15}$  molecules cm<sup>-2</sup>. Amorphous D<sub>2</sub>O ice was produced by vapor deposition at 10 K on the Al substrate with a rate of  $\sim 1$  ML minute<sup>-1</sup>.

## 3. RESULTS AND DISCUSSION

### 3.1. Overview of the Reaction

Under all experimental conditions, both H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub> were formed by codeposition of H atoms with O<sub>2</sub> molecules on the Al substrate and amorphous D<sub>2</sub>O ice (Figure 1). In the same manner as utilized in the previous experiment (Miyauchi et al. 2008), we confirmed that H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub> formed are not contaminants. No



**Figure 1.** IR spectra of products after codeposition of O<sub>2</sub> with H under conditions of (a) O<sub>2</sub>/H = 0.002 at 20 K after 360 minutes and (b) O<sub>2</sub>/H = 0.009 at 40 K after 120 minutes. (c) For comparison, the IR spectra of pure amorphous H<sub>2</sub>O ice produced by vapor deposition at 20 K. Solid and dashed lines indicate peaks assigned to H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub>, respectively. Arrows indicate noise caused by vibration of the He refrigerator.

intermediate radicals such as HO<sub>2</sub> (1142 cm<sup>-1</sup>, Cooper et al. 2006) were observed in this study, which is consistent with the H irradiation experiment onto solid O<sub>2</sub> (Miyachi et al. 2008). Ioppolo et al. (2008) showed that, for the H irradiation experiment onto solid O<sub>2</sub>, the reactions proceed at temperatures lower than 28 K. This is consistent with O<sub>2</sub> desorption from a Au substrate at temperatures of 29.5–31.3 K in the typical experimental timescale (Acharyya et al. 2007). However, in the codeposition experiments conducted in this study, it is clear that the formation of H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub> occurs even at 40 K, which could be above the desorption temperature of O<sub>2</sub>.

The reasons for the formation of H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub>, as observed in the codeposition experiments at temperatures higher than the O<sub>2</sub> desorption temperature, will be discussed. The residence time of a species ( $\tau_r$ ) on a substrate is given by the following

equation:

$$\tau_r = \nu_0^{-1} \exp(E_{i-j}/kT), \quad (8)$$

where  $\nu_0$  and  $E_{i-j}$  are the vibrational frequency and the binding energy of the adsorbed species  $i$  on the substrate  $j$ , respectively,  $k$  is the Boltzmann constant, and  $T$  is temperature. The value of  $E_{O_2-Al}$  is unknown. Instead, by substituting  $E_{O_2-Au} = 912$  K (Acharyya et al. 2007), and  $\nu_0$  as  $1.4 \times 10^{12}$  s<sup>-1</sup> (Tielens 2005), the  $\tau_r$  of O<sub>2</sub> molecules on the Au substrate are calculated to be  $\sim 11$  s and  $5.7 \times 10^{-3}$  s at 30 and 40 K, respectively. In contrast, because no data are available for  $E_{H-Al}$ , and also for  $E_{H-Au}$ ,  $\tau_r$  of H atoms cannot be calculated from equation (8). Then, considering that  $E_{H-H_2O}$  (650 K, Al-Halabi & van Dishoeck 2007) is much smaller than  $E_{O_2-H_2O}$  (1600 K, Tielens 2005),  $E_{H-Au}$  is expected to be much smaller than  $E_{O_2-Au}$ , leading to an assumption that  $\tau_r$  of H atoms is much shorter than that of O<sub>2</sub> molecules on the Au substrate shown above. If these relationships can be applied to the Al substrate, then the results indicate that Reactions (5)–(7) proceed quickly enough, so that H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub> are produced within the short residence time of H and O<sub>2</sub> before the H atoms and O<sub>2</sub> molecules desorb from the Al substrate. Once H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub> are formed, these products can remain even at 40 K because of their high desorption temperatures.

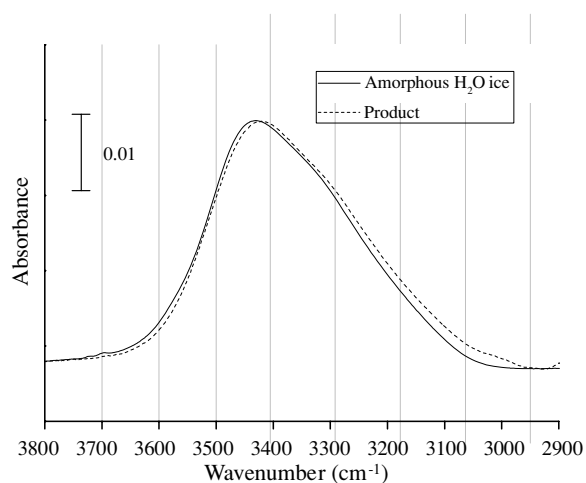
### 3.2. Formation of Compact Amorphous Ice

Figures 1(a) and (b) show examples of IR spectra after codeposition of H with O<sub>2</sub> obtained with an O<sub>2</sub>/H ratio of 0.002 at 20 K and with an O<sub>2</sub>/H ratio of 0.009 at 40 K, respectively. For comparison, Figure 1(c) shows an IR spectrum of the amorphous H<sub>2</sub>O ice produced by vapor deposition at 20 K with a rate of  $\sim 0.4$  ML minute<sup>-1</sup>. Some peaks in Figure 1 are assigned to H<sub>2</sub>O (3414, 2173, 1683, and 843 cm<sup>-1</sup>) and H<sub>2</sub>O<sub>2</sub> (3340, 2837, 1394, and 888 cm<sup>-1</sup>). The OH-stretching mode (3414 cm<sup>-1</sup>) in Figure 1(a) is almost the same as that for the vapor-deposited amorphous H<sub>2</sub>O ice (Figure 1(c)), which indicates that the H<sub>2</sub>O ice formed in this experiment is amorphous (Bergren et al. 1978; Hagen et al. 1981). The spectral feature is consistent with the observations of H<sub>2</sub>O ice in molecular clouds (e.g., Hagen et al. 1981; van de Bult et al. 1985; Smith et al. 1989) and theoretical studies (Kouchi et al. 1994). However, the width of the 3  $\mu$ m band of ice produced after the codeposition is slightly wider in the small-wavenumber region than that of vapor-deposited amorphous H<sub>2</sub>O ice at 20 K (Figure 2). Hagen et al. (1981) showed that the width of the 3  $\mu$ m band becomes irreversibly narrower with temperature; however, the results of the present study do not exhibit such a tendency. It is therefore concluded that the wider 3  $\mu$ m band from the codeposition experiment is not due to local heating by the heat of reactions, but is probably due to the overlapping of the H<sub>2</sub>O<sub>2</sub> band. Although the heats of reaction (2.1, 3.7, and 3.9 eV for Reactions (5), (6), and (7), respectively;<sup>2</sup> 4.5 eV for H–H recombination) are considerably large, it is suggested that the heats of reactions do not affect the microscopic structure of amorphous H<sub>2</sub>O ice. This suggestion is consistent with the theoretical discussion by Kouchi et al. (1994).

Two small peaks assigned to the dangling OH bonds are typically observed in the spectrum for vapor-deposited amorphous H<sub>2</sub>O ice at 3721 and 3698 cm<sup>-1</sup> (Figure 3(a)), and these are assigned to two- and three-coordinated surface H<sub>2</sub>O molecules,

<sup>2</sup> NIST CCCBDB Web site; <http://cccbdb.nist.gov/hf0k.asp>.



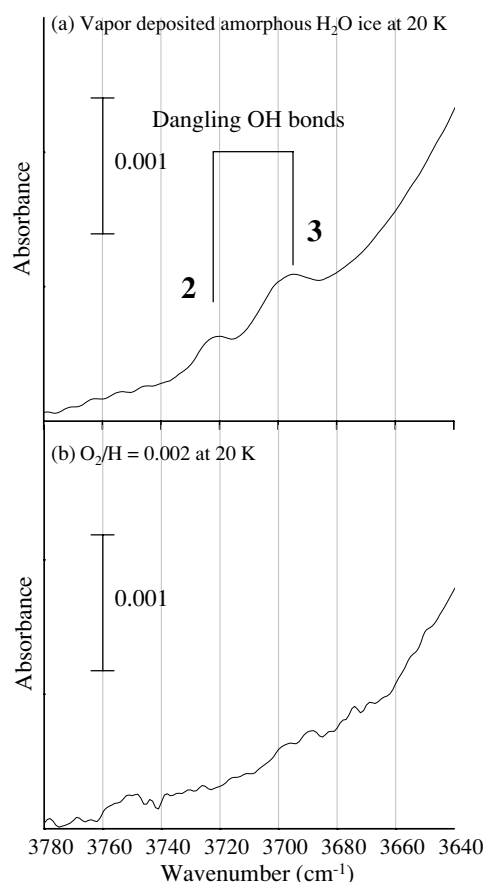


**Figure 2.** Comparison of IR spectra designated as the OH-stretching mode for the product of O<sub>2</sub>-H codeposition and pure amorphous H<sub>2</sub>O ice produced by vapor deposition at 20 K.

respectively (Buch & Devlin 1991). However, these peaks are not observed in the amorphous H<sub>2</sub>O ice formed by codeposition (Figure 3(b)). The dangling OH bonds appear in the IR spectrum for amorphous H<sub>2</sub>O ice due to the microporous structure (Buch & Devlin 1991; Rowland & Devlin 1991; Zondlo et al. 1997). On the other hand, it has been reported that the dangling OH bonds disappear when the amorphous ice is annealed to 120 K (Rowland & Devlin 1991; Zondlo et al. 1997), when the ice is irradiated with 200 keV protons at 15 K (Palumbo 2006), or when the ice is irradiated with 100 keV Ar<sup>+</sup> ions at 40 K (Raut et al. 2007), all of which are due to collapse of the microporous structure. Although the peak position of the dangling OH bonds of H<sub>2</sub>O<sub>2</sub>-H<sub>2</sub>O mixed ice (3664 cm<sup>-1</sup>, Palumbo 2006) is different from that of vapor-deposited pure amorphous H<sub>2</sub>O ice (3698 and 3721 cm<sup>-1</sup>), such a feature does not appear in the IR spectrum of amorphous H<sub>2</sub>O ice formed by codeposition (Figure 3(b)). The lack of dangling OH bonds in H<sub>2</sub>O ice formed by codeposition is consistent with the feature for H<sub>2</sub>O ice in molecular clouds (Keane et al. 2001). It was confirmed that codeposition of O<sub>2</sub> and H<sub>2</sub>O molecules at various ratios also produces spectral bands for dangling bonds, which indicates that the bands do not disappear due to the coexistence of O<sub>2</sub> and H<sub>2</sub>O. Therefore, the less clear signature of dangling OH bonds in Figure 3(b) implies that the structure of the amorphous H<sub>2</sub>O ice formed by codeposition is more compact compared with vapor-deposited amorphous H<sub>2</sub>O ice. This might be due to the heats of reactions. Small amounts of H<sub>2</sub>O<sub>2</sub> present with H<sub>2</sub>O may also contribute to the lack of dangling OH bonds.

### 3.3. Conditions for Abundant H<sub>2</sub>O Formation

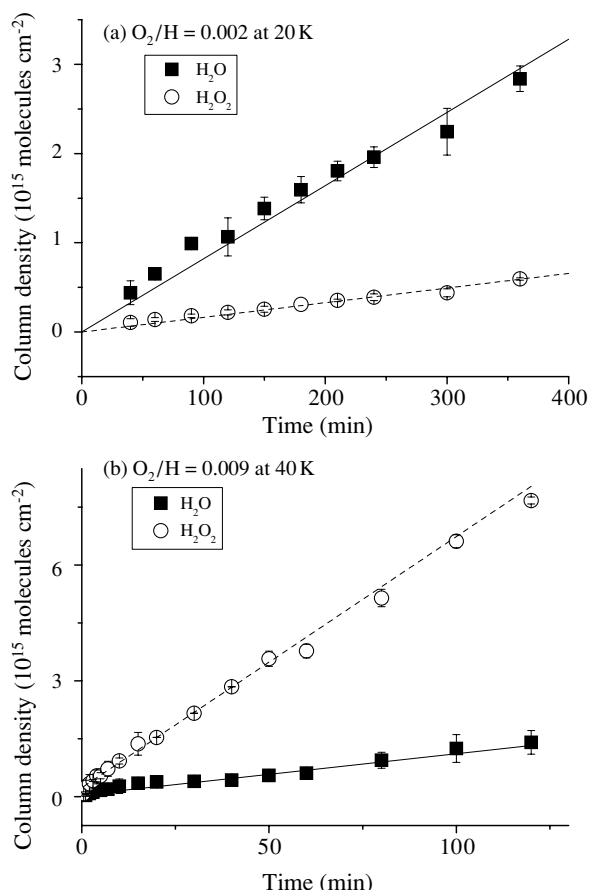
It is known that the peak positions assigned to the OH-stretching mode for both H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub> molecules (3000–3600 cm<sup>-1</sup>) overlap with each other in the IR spectrum of H<sub>2</sub>O-H<sub>2</sub>O<sub>2</sub> mixed ice (Boudin et al. 1998); therefore, the peak is not suitable for quantification. Instead, the peaks assigned to the OH-bending mode for H<sub>2</sub>O (1590–1740 cm<sup>-1</sup>) and H<sub>2</sub>O<sub>2</sub> (1270–1530 cm<sup>-1</sup>) were used to obtain their abundances in this study. The column densities are calculated from our obtained peak areas and the previously reported integrated band strengths, as described in Hidaka et al. (2007) in which correction for the light path with an IR incident angle of about 83° was made. The band strengths used are  $1.2 \times 10^{-17}$  and



**Figure 3.** IR spectra of (a) pure amorphous H<sub>2</sub>O ice produced by vapor deposition at 20 K and (b) products of the codeposition experiments under the conditions of O<sub>2</sub>/H = 0.002 at 20 K. Peaks numbered with 2 and 3 in (a) indicate two- and three-coordinated dangling OH bonds, respectively.

$2.1 \times 10^{-17}$  cm molecules<sup>-1</sup> for H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub>, respectively, which are obtained by the *transmission* absorption spectroscopy of pure amorphous H<sub>2</sub>O ice (Gerakins et al. 1995) and the *reflection* absorption spectroscopy of H<sub>2</sub>O<sub>2</sub> in solid mixture (Loeffler et al. 2006), respectively. Our estimates of column densities by the reflection absorption spectroscopy can include some errors especially in H<sub>2</sub>O due to the use of band strength obtained by the transmission spectroscopy. In order to check the difference in column density estimates between the transmission and reflection methods, we also performed the transmission spectroscopy for pure amorphous H<sub>2</sub>O ice and pure solid CO and found that using the same reported band strength the difference in the estimated column densities between the transmission and reflection methods is within approximately a factor of 2. It should be noted that even in the transmission spectroscopy, it is very difficult to derive the exact column densities of molecules in solid mixture. Furthermore, all of our discussion in the present study is based on the relative column densities. Therefore, we conclude that our estimated column densities are reasonably acceptable.

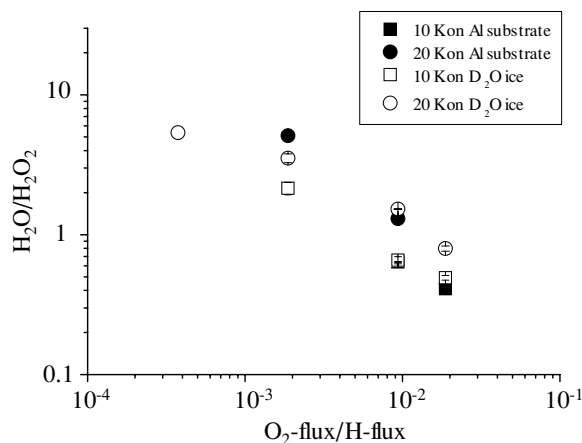
Figure 4 presents typical variations in the column densities of H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub> as a function of time. These molecules display a linear increase, because the source materials (O<sub>2</sub> + H) that can yield H<sub>2</sub>O<sub>2</sub> and H<sub>2</sub>O are continuously supplied. The results clearly show that the relative abundance of H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub> (H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub>) varies with the experimental conditions. For example, under conditions with an O<sub>2</sub>/H ratio of 0.002 at 20 K (Figure 4(a)), the H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> ratio is approximately 5.



**Figure 4.** Variations in the column density of  $\text{H}_2\text{O}$  and  $\text{H}_2\text{O}_2$  after codeposition of  $\text{O}_2$  with H under the conditions of (a)  $\text{O}_2/\text{H} = 0.002$  at 20 K and (b)  $\text{O}_2/\text{H} = 0.009$  at 40 K.

However, under conditions with an  $\text{O}_2/\text{H}$  ratio of 0.009 at 40 K (Figure 4(b)), the  $\text{H}_2\text{O}/\text{H}_2\text{O}_2$  ratio is approximately 0.15. Figure 5 shows that as the  $\text{O}_2/\text{H}$  flux ratio decreases, the  $\text{H}_2\text{O}/\text{H}_2\text{O}_2$  ratio increases up to approximately 5, which is qualitatively understandable when considering the onion-like structure described in section 1. When the  $\text{O}_2$  flux is large,  $\text{O}_2$  molecules on the surface, as well as  $\text{H}_2\text{O}_2$  formed by Reaction (6), are covered with additional  $\text{O}_2$  prior to subsequent Reaction (7) that produces  $\text{H}_2\text{O}$ . Since the H atoms cannot react with molecules buried in the bulk, part of the buried  $\text{O}_2$  and  $\text{H}_2\text{O}_2$  remains unreacted. Therefore, when the  $\text{O}_2$  flux is large, most H atoms are used for the reaction with  $\text{O}_2$  rather than  $\text{H}_2\text{O}_2$ , which leads to ice with a smaller  $\text{H}_2\text{O}/\text{H}_2\text{O}_2$  ratio. On the other hand, when the  $\text{O}_2$  flux is small,  $\text{O}_2$  and  $\text{H}_2\text{O}_2$  have more opportunity to react with H atoms to form  $\text{H}_2\text{O}$  before being covered by additional  $\text{O}_2$ , which leads to ice with a larger  $\text{H}_2\text{O}/\text{H}_2\text{O}_2$  ratio. However, it is interesting to note that although the H flux is quantitatively 2–4 orders of magnitude larger than the  $\text{O}_2$  flux, the  $\text{H}_2\text{O}/\text{H}_2\text{O}_2$  ratio is still sensitive to the change in the  $\text{O}_2/\text{H}$  ratio. In the gas phase, when the number density ratio of two reactant partners is significant, the final yields exhibit little dependence on the ratio. That is, reactions will eventually proceed to completion until the reactants are consumed.

The present results are readily understandable when the following factors are considered. In the present case, comparison of the sum of column densities of products with H atom fluence shows that, at most, only  $\sim 2\%$  of H atoms are used for the hydrogenation of  $\text{O}_2$ . This is consistent with a previous study in

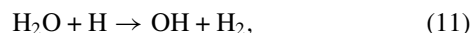
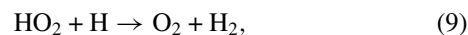


**Figure 5.**  $\text{H}_2\text{O}/\text{H}_2\text{O}_2$  ratio vs.  $\text{O}_2/\text{H}$  at 10 and 20 K. Open and solid symbols represent the experimental results for codeposition with and without amorphous  $\text{D}_2\text{O}$  ice, respectively.

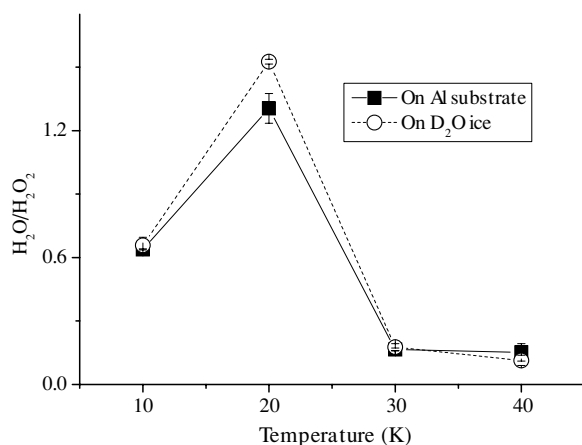
which only a small portion of H atoms ( $\sim 0.1\%$ ) was consumed for hydrogenation of CO (Watanabe et al. 2004). Therefore, most of the H atoms are used for the formation of  $\text{H}_2$  by recombination. Recombination is a no-barrier radical–radical reaction and has a rate proportional to the square of the H atom density at the surface. Therefore, under the condition of a very high H-atom flux, recombination proceeds preferentially compared with the hydrogenation of  $\text{O}_2$ . As a result, the number of H atoms used for hydrogenation reactions in the present experiment is much smaller than that irradiated onto the substrate.

At an  $\text{O}_2/\text{H}$  ratio of  $\sim 0.009$ , the  $\text{H}_2\text{O}/\text{H}_2\text{O}_2$  ratio varies with temperature (Figure 6); at 10 K, the  $\text{H}_2\text{O}/\text{H}_2\text{O}_2$  ratio is approximately 0.6, while it becomes approximately 1.3 at 20 K. At 30 and 40 K, the  $\text{H}_2\text{O}/\text{H}_2\text{O}_2$  ratio is significantly smaller than that at 10 and 20 K, both at approximately 0.15. The temperature dependence of  $\text{H}_2\text{O}/\text{H}_2\text{O}_2$  is related to the residence time and the H-atom surface diffusion distance. At 10 K, both H and  $\text{O}_2$  can reside on the surface for sufficient time to react by a tunneling process (Miyauchi et al. 2008). At 20 K, the H atoms are assumed to be more mobile on the surface than at 10 K; therefore, more H atoms can meet and react with  $\text{O}_2$  and  $\text{H}_2\text{O}_2$  which results in the formation of more  $\text{H}_2\text{O}$ . On the other hand, at 30 and 40 K, the tunneling Reaction (7) becomes less effective, due to the very short residence time of H atoms on the surface. However, the formation of  $\text{H}_2\text{O}_2$  occurs regardless of the very short residence times of  $\text{O}_2$  and H because Reaction (5) has essentially no barrier (Walch et al. 1988). These characteristics result in smaller values of  $\text{H}_2\text{O}/\text{H}_2\text{O}_2$  (Figure 4).

In addition to reactions that produce  $\text{H}_2\text{O}_2$  and  $\text{H}_2\text{O}$ , the following H-subtraction reactions,



must be considered to discuss the  $\text{H}_2\text{O}/\text{H}_2\text{O}_2$  ratio of the products. Reaction (11) is an endothermic reaction (Brouard et al. 2002); therefore, H subtraction from  $\text{H}_2\text{O}$  can be neglected under the present experimental conditions. As for Reactions (9) and (10) on the surface, no data are available for the activation



**Figure 6.** H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> ratio vs. temperature under constant O<sub>2</sub>/H (0.009) conditions. Open and solid symbols represent the experimental results for codeposition with and without amorphous D<sub>2</sub>O ice, respectively.

energy. In the gas phase, it is reported that the barrier heights for Reactions (9) and (10) are 1070 K (Lee & Hochgreb 1998) and 1890–4780 K (Lee & Hochgreb 1998; Tarchouna et al. 2006), respectively. Despite the relatively low barrier height for Reaction (9), the progress of the reaction should be negligible, because Reaction (5) proceeds with essentially no barrier (Walch et al. 1988) and Reaction (6) proceeds much faster than Reaction (5) (Miyauchi et al. 2008). The barrier height of Reaction (10) is almost the same or slightly higher than that for Reaction (7) (~2000 K, Koussa et al. 2006), which suggests that H<sub>2</sub>O<sub>2</sub> formed by Reaction (6) could be used for both Reactions (7) and (10). However, the high H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> ratio of 5 obtained for a small O<sub>2</sub>/H flux ratio of 0.002 demonstrates that even if H subtraction from H<sub>2</sub>O<sub>2</sub> proceeds to some extent, Reaction (10) would have little contribution to the H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> ratio of the products, because the product of Reaction (10), HO<sub>2</sub>, is quickly hydrogenated again by additional H atoms to yield H<sub>2</sub>O<sub>2</sub> by Reaction (6).

The hydrogenation of O<sub>2</sub> to form H<sub>2</sub>O proceeds even at 40 K in the codeposition experiments, regardless of the presence of amorphous D<sub>2</sub>O ice as the substrate. The variations in H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> with temperature, both with and without amorphous D<sub>2</sub>O ice, are quite similar (Figure 6). However, previous studies reported that amorphous ice has a catalytic effect to enhance the reaction rate for reactions of H atoms with CO molecules (e.g., Hidaka et al. 2007) and H–D substitution of solid methanol (Watanabe & Kouchi 2008). In addition, it has been theoretically demonstrated that H<sub>2</sub>O molecules have a catalytic effect to reduce the reaction barrier height of some surface reactions (Woon 2002; Xie et al. 2006). Although such a catalytic effect is not clearly observed for the value of H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> in the present study, we suggest that the effect could occur in the formation of H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub> from O<sub>2</sub>. This is because when H atoms are irradiated onto 30 ML amorphous D<sub>2</sub>O ice covered with 3 ML O<sub>2</sub>, H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub> are formed even at 40 K, while the formation of these molecules is not observed above 23 K without amorphous ice (Y. Oba et al. 2009, in preparation).

#### 4. ASTROPHYSICAL IMPLICATIONS

It is widely recognized that H<sub>2</sub>O ice in dense molecular clouds is amorphous (e.g., Smith et al. 1989). Amorphous H<sub>2</sub>O ice produced by vapor deposition in the laboratory contains micropores (Mayer & Pletzer 1986), and these pores contain dangling OH

bonds. Therefore, it has been implied that amorphous ice in dense molecular clouds may have a large number of dangling bonds. However, there have been no reports on the detection of dangling bonds in interstellar ices (Keane et al. 2001). To explain this, two possibilities have been suggested: (1) there are no significant amounts of micropores formed with the formation of amorphous H<sub>2</sub>O ice; (2) dangling bonds are present at the formation of amorphous H<sub>2</sub>O ice, but disappear due to cosmic ray irradiation (Palumbo 2006; Raut et al. 2007). However, the former possibility has not been experimentally investigated. The present study clearly demonstrates that amorphous H<sub>2</sub>O ice does not exhibit bands for dangling bonds when it is formed by surface atomic reactions.

In addition to the lack of dangling OH bonds, the IR spectra of amorphous H<sub>2</sub>O ices in molecular clouds have a characteristic feature at 3.2–3.5 μm, which is referred to as the long-wavelength wing (Smith et al. 1989) or low-frequency wing (Hagen et al. 1983). It has been considered that the wing is derived from the absorption of ammonia (Merrill et al. 1976), hydrated silicates (Knacke & Krätschmer 1980), hydrocarbons (Smith et al. 1989), or methanol (Grim et al. 1991). The contributions from ammonia and hydrous silicates to the wing were ruled out by Smith et al. (1989). However, the source of the long-wavelength wing is still unclear to date. In addition to these candidate species, another candidate is proposed from the present study, in that small amounts of H<sub>2</sub>O<sub>2</sub> may explain the long-wavelength wing, as shown in Figure 2. Although the presence of H<sub>2</sub>O<sub>2</sub> in molecular clouds has not been confirmed (see below), the weak absorption at 2900–3300 cm<sup>−1</sup> (Figure 2) is similar to previously reported astronomical observations (e.g., Smith et al. 1989).

In contrast to the ubiquitous presence of water, hydrogen peroxide has not been observed in molecular clouds. Blake et al. (1987) theoretically estimated that the H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> ratio is approximately 30 in the core of the Orion molecular cloud. In addition, the upper limit of H<sub>2</sub>O<sub>2</sub> abundances in the NGC7538:IRS9 protostar has been estimated to be <0.5 compared to H<sub>2</sub>O abundances (Boudin et al. 1998). The H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> ratio in these environments is considerably high when compared to that obtained in the present study (up to ~5). If the O<sub>2</sub>/H ratio is extrapolated to a lower value (~1 × 10<sup>−4</sup>), then the expected H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> ratio becomes 15 to 30 (Figure 5). However, in molecular clouds, O<sub>2</sub> is formed by the reaction of two oxygen atoms (Reaction (4)) absorbed on a grain surface, followed by reactions with H atoms to produce H<sub>2</sub>O<sub>2</sub> and H<sub>2</sub>O (Reactions (5)–(7)). Furthermore, there remains some uncertainty regarding the absolute abundance of O and H atoms in molecular clouds. Therefore, an O<sub>2</sub>/H ratio of ~1 × 10<sup>−4</sup> with an estimated abundance of H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> (15 to 30) may not correspond directly with the actual H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> ratio in molecular clouds. The probability of H–H recombination is proportional to the square of the number of H atoms on a surface (Vidali et al. 2007); therefore, the fraction of H atoms consumed by recombination becomes smaller, and that consumed by hydrogenation of O<sub>2</sub> becomes larger under the condition of low H flux in molecular clouds. Thus, in molecular clouds, the ratio of H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> tends to be larger than that in the present experiments. Regardless, the experimental results cannot be applied directly for discussion of the actual H<sub>2</sub>O/H<sub>2</sub>O<sub>2</sub> versus O/H ratios in molecular clouds at present. For future work, the kinetics of Reactions (2) and (3) and astronomical observations of H<sub>2</sub>O<sub>2</sub> are required for the further elucidation of H<sub>2</sub>O formation mechanisms and the relative abundance of H<sub>2</sub>O<sub>2</sub>.

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