

# Viscous control of shallow elastic fracture

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This paper considers the problem of a semi-infinite crack parallel to the boundary of a half plane, with the crack filled by an incompressible viscous fluid. The dynamics are driven by a bending moment applied to the arm of the crack, and we look for travelling wave solutions. We examine two models of fracture; fracture with a single tip, and fracture with a wet tip preceded by a region of dry fracture.

**Key words:** Authors should not enter keywords on the manuscript, as these must be chosen by the author during the online submission process and will then be added during the typesetting process (see <http://journals.cambridge.org/data/relatedlink/jfm-keywords.pdf> for the full list)

## 1. Introduction

Here we review the literature as well as describe the problem in more detail. We have the vertical displacement  $h$ , the horizontal displacement  $g$ , the thickness of the arm  $l$ , and the pressure  $p$ . We look for a travelling wave solution (propagating left), with speed  $c$ .

## 2. Formulation of problem

### 2.1. Single tip

From lubrication, we have Poiseuille flow in the crack. We obtain the flux, and conservation of mass as

$$q = -\frac{1}{12\mu} \frac{dp}{dx} h^3, \quad \frac{\partial q}{\partial x} + \frac{\partial h}{\partial t} = 0, \quad (2.1)$$

which combined gives

$$\frac{dp}{dx} = 12\mu c/h^2. \quad (2.2)$$

Setting  $p \rightarrow 0$  at  $x \rightarrow \infty$ , we can write this in integral form,

$$p(x) = - \int_x^\infty 12\mu c/h(\tilde{x})^2 d\tilde{x}. \quad (2.3)$$

From the linear theory of elasticity, due to others who have studied this problem, we have

$$\begin{bmatrix} -\sigma_y \\ -\tau_{xy} \end{bmatrix} = \begin{bmatrix} p(x) \\ 0 \end{bmatrix} = \frac{E}{4\pi l(1-\nu^2)} \int_0^\infty \mathbf{K} \left( \frac{\tilde{x}-x}{l} \right) \begin{bmatrix} g'(\tilde{x}) \\ h'(\tilde{x}) \end{bmatrix} d\tilde{x}, \quad (2.4)$$

where the integral kernel is

$$\mathbf{K}(\xi) = \begin{bmatrix} K_{11} & K_{12} \\ K_{21} & K_{22} \end{bmatrix} = \begin{bmatrix} \frac{(32-24\xi^2)}{(\xi^2+4)^3} & \frac{(48\xi^2-64)}{\xi(\xi^2+4)^3} \\ -\frac{(16\xi^4+16\xi^2+4)}{\xi(\xi^2+4)^3} & -\frac{(32-24\xi^2)}{(\xi^2+4)^3} \end{bmatrix}. \quad (2.5)$$

The boundary conditions near  $x = 0$  are governed by fracture mechanics

$$K_I = \lim_{x \rightarrow 0} \frac{E}{1-\nu^2} \sqrt{\frac{\pi}{8}} \sqrt{x} h'(x), \quad K_{II} = \lim_{x \rightarrow 0} \frac{E}{1-\nu^2} \sqrt{\frac{\pi}{8}} \sqrt{x} g'(x). \quad (2.6a, b)$$

As we go to  $x \gg l$ , we are looking at the problem of peeling off a thin strip from an elastic half space. We can then use beam theory approximations, which give

$$M(x) = \frac{El^3}{12(1-\nu^2)} \frac{d^2 h}{dx^2} = \frac{El^2}{6(1-\nu^2)} \frac{dg}{dx}, \quad p = \frac{El^3}{12(1-\nu^2)} h^{(4)}(x) \quad (2.7a, b)$$

As  $x \rightarrow \infty$ ,  $M(x) \rightarrow M$ , the applied bending moment, so this gives us boundary conditions on  $h''$ ,  $g'$ .

## 2.2. Double tip

The equations are mostly unchanged from the single tip problem. From the geometry of the double tip problem, a second crack tip is situated at  $x = -lL$ , and  $h(x), h'(x) = 0$  for  $x < 0$ . The boundary conditions at  $\infty$  are unaffected, and the boundary conditions at the crack tip become

$$\lim_{x \rightarrow 0} \sqrt{x} h'(x) = L, \quad \lim_{x \rightarrow -Ll} \frac{E}{1-\nu^2} \sqrt{\frac{\pi}{8}} \sqrt{x} g'(x) = K_{II}. \quad (2.8a, b)$$

The limits in the elasticity integral become  $(-Ll, \infty)$ .

## 2.3. Rescaling

We can define the following dimensionless variables

$$x = l\xi, \quad h(x) = \frac{12M(1-\nu^2)}{El} H(\xi), \quad g(x) = \frac{12M(1-\nu^2)}{El} G(\xi), \quad (2.9)$$

$$p = \frac{3M}{\pi l^2} \Pi(\xi), \quad K_I = Ml^{-3/2} \kappa_I, \quad K_{II} = Ml^{-3/2} \kappa_{II}, \quad \lambda = \frac{4\pi\mu p^* l^3}{M^2}. \quad (2.10)$$

With these scalings, the equations become

$$\begin{bmatrix} \Pi \\ 0 \end{bmatrix} = \int_0^\infty \mathbf{K}(\tilde{\xi} - \xi) \begin{bmatrix} G'(\tilde{\xi}) \\ H'(\tilde{\xi}) \end{bmatrix} d\tilde{\xi} \quad (2.11)$$

$$H^2 \frac{d\Pi}{d\xi} = \lambda \quad \text{or} \quad \Pi(\xi) = - \int_\xi^\infty \lambda / H(\tilde{\xi})^2 d\tilde{\xi} \quad (2.12a, b)$$

$$\lim_{\xi \rightarrow \infty} H'' = 1, \quad \lim_{\xi \rightarrow \infty} G' = \frac{1}{2}, \quad \lim_{\xi \rightarrow 0} 3\sqrt{2\pi\xi} H' = \kappa_I, \quad \lim_{\xi \rightarrow 0} 3\sqrt{2\pi\xi} G' = \kappa_{II}, \quad (2.13)$$

These shall be the governing equations for the rest of this paper.

## 2.4. Beam theory asymptotics

In the dimensionless variables, the outer asymptotics are of the form

$$\frac{d^2 H}{d\xi^2} = \frac{1}{2} \frac{dG}{d\xi}, \quad H^{(4)}(\xi) = \frac{3}{\pi} \Pi(\xi), \quad \frac{d^2 H}{d\xi^2} \rightarrow 1 \quad (2.13a, b, c)$$

From integration by parts, we can write

$$H''(\xi) = 1 - \frac{1}{2} \int_{\xi}^{\infty} (\tilde{\xi} - \xi)^2 H^{(5)}(\tilde{\xi}) d\tilde{\xi}, \quad (2.14)$$

provided  $\lim_{\xi \rightarrow \infty} \xi H^{(3)}(\xi) = \lim_{\xi \rightarrow \infty} \xi^2 H^{(4)}(\xi) = 0$ . Then using equation 2.12a, we have that

$$H''(\xi) = 1 - \frac{3\lambda}{2\pi} \int_{\xi}^{\infty} \frac{(\tilde{\xi} - \xi)^2}{H(\tilde{\xi})^2} d\tilde{\xi}. \quad (2.15)$$

Since  $H(\xi) = \frac{1}{2}\xi^2 + o(\xi^2)$ , as  $\xi \rightarrow \infty$ , we can use this to get a better estimate of  $H''$ ;

$$H''(\xi) = 1 - \frac{2\lambda}{\pi} \frac{1}{\xi} + o(1/\xi). \quad (2.16)$$

This new expression can be used to refine the error estimate from  $o(1/\xi)$ , to  $O(\log(\xi)/\xi^2)$ .

### 2.5. Linear perturbation problem

*This section is probably better placed elsewhere...* The equations of the linear perturbation problem:

$$\Pi = \Pi_0 + \mathcal{E}\Pi_1 + O(\mathcal{E}), \quad H = H_0 + \mathcal{E}H_1 + O(\mathcal{E}) \quad (2.17)$$

$$\begin{bmatrix} \Pi_1 \\ 0 \end{bmatrix} = \int_0^{\infty} \mathbf{K}(\xi - \tilde{\xi}) \begin{bmatrix} G'_1(\tilde{\xi}) \\ H'_1(\tilde{\xi}) \end{bmatrix} d\tilde{\xi}, \quad H_0^2 \Pi'_1 + 2H_0 H_1 \Pi'_0 = \lambda_1 \quad (2.18a, b)$$

$$H'_1 \rightarrow 0 \text{ as } \xi \rightarrow \infty, \quad H_1 \sim \xi^s + \frac{\tilde{A}\lambda_1}{3\lambda_0^{2/3}} \xi^{2/3} + \dots \text{ as } \xi \rightarrow 0 \quad (2.19a, b)$$

But these can be made into a more convenient form, by considering instead  $\tilde{\Pi} = \Pi_0 - 3\lambda_0/\lambda_1 \Pi_1$ , and similar for  $\tilde{H}$ ,  $\tilde{G}$ . The equations become

$$\begin{bmatrix} \tilde{\Pi} \\ 0 \end{bmatrix} = \int_0^{\infty} \mathbf{K}(\xi - \tilde{\xi}) \begin{bmatrix} \tilde{G}'(\tilde{\xi}) \\ \tilde{H}'(\tilde{\xi}) \end{bmatrix} d\tilde{\xi}, \quad H_0^2 \tilde{\Pi}' + 2H_0 \tilde{H} \Pi'_0 = 0 \quad (2.20a, b)$$

$$\tilde{H}'' \rightarrow 1 \text{ as } \xi \rightarrow \infty, \quad \tilde{H} \sim -\frac{3\lambda_0}{\lambda_1} \xi^s + \dots \text{ as } \xi \rightarrow 0 \quad (2.21a, b)$$

## 3. Numerical scheme

### 3.1. Single Tip

We discretize the problem by taking  $n+1$  points  $\boldsymbol{\xi} = (\xi_0 = 0, \xi_1, \dots, \xi_n)$  at which we measure  $H'$ ,  $G'$ , and  $n$  intermediate points  $\boldsymbol{\zeta} = (\zeta_0, \dots, \zeta_{n-1})$  at which to measure  $\Pi$ , so that  $\xi_0 < \zeta_0 < \dots < \zeta_{n-1} < \xi_n$ . We work with  $\sqrt{\xi} G'(\xi)$ ,  $\sqrt{\xi} H'(\xi)$  near the tip to avoid singularities. We define  $\boldsymbol{\theta}_G = [\sqrt{\xi_0} G'(\xi_0), \dots, \sqrt{\xi_{t-1}} G'(\xi_{t-1}), G'(\xi_t), \dots, G'(\xi_n)]$ , and  $\boldsymbol{\theta}_H$  similarly, as well as  $\boldsymbol{\theta} = [\boldsymbol{\theta}_G, \boldsymbol{\theta}_H]$ . Typically  $t \approx n/2$  was used. From the linearity of the elasticity integral (and the discretized integral) we may write

$$[\Pi(\zeta_1), \dots, \Pi(\zeta_{n-1}), \underbrace{0, \dots, 0}_{n-1}] = \mathbf{J}\boldsymbol{\theta}, \quad (3.1)$$

for some matrix  $\mathbf{J}$ . One can recover  $H(\xi_i)$  from  $\boldsymbol{\theta}_H$ . Therefore, a discretized lubrication integral, yields an expression for  $\Pi(\xi_i)$  as a function of  $\boldsymbol{\theta}_H$ . So we can write

$$[\Pi(\xi_1), \dots, \Pi(\xi_{n-1}), \underbrace{0, \dots, 0}_{n-1}] = \mathbf{J}\boldsymbol{\theta} = \mathbf{f}(\boldsymbol{\theta}_H), \quad (3.2)$$

for some function  $\mathbf{f}$ .

Both  $G'(\xi_n)$ , and  $H''(\xi_n)$  are known from our beam theory asymptotic expansion. But these are linear in  $\boldsymbol{\theta}$ , as  $G'(\xi_n) = \theta_n$ , and  $H''(\xi_n) \approx (\theta_{2n} - \theta_{2n-1})/(\xi_n - \xi_{n-1})$ . Therefore we can add another two rows to  $\mathbf{J}$ , so that

$$\mathbf{A}\boldsymbol{\theta} = [\mathbf{f}(\boldsymbol{\theta}), G'(\xi_n), H''(\xi_n)]. \quad (3.3)$$

Where the  $\mathbf{A}$  is the enlarged matrix. This can be solved by Newton's method from quite arbitrary initial guesses.

For  $\xi_i < \xi < \xi_{i+1}$ , we interpolate as

$$G'(\xi) = \begin{cases} \xi^{-1/2}(a_i\xi + b_i) \\ a_i\xi + b_i \end{cases}, \quad H'(\xi) = \begin{cases} \xi^{-1/2}(c_i\xi^{1/2} + d_i) \\ c_i\xi + d_i \end{cases}, \quad \text{for } \begin{cases} i < t \\ i \geq t \end{cases} \quad (3.4)$$

The choice of interpolating function was based on the appearance of the relevant functions. We will also define  $a_n, b_n, c_n, d_n$  for interpolation beyond  $\xi_n$ . With this choice of interpolation, there exist exact closed form expressions for both the lubrication integral, and the elasticity integral, in terms of the  $a_i - d_i$  coefficients.

It therefore remains to determine  $a_i - d_i$  in terms of  $\boldsymbol{\theta}$ . Continuity of  $G', H'$  imposes  $2(n-1)$  linear equations. We also have the  $2n$  equations following from the definition of  $\boldsymbol{\theta}$ , (such as  $a_i\xi_i + b_i = \theta_i$  for  $t \leq i \leq n$ ).

From our asymptotic expansion (via beam theory) we know  $\theta_n = G'(\xi_n)$  and  $a_n = G''(\xi_n)$ . Therefore we can write

$$a_n = \frac{G''(\xi_n)}{G'(\xi_n)}\theta_n, \quad b_n = \theta_n - a_n\xi_n = \left(1 - \frac{G''(\xi_n)}{G'(\xi_n)}\xi_n\right)\theta_n \quad (3.5)$$

With  $H$ , we know that  $c_n = H''(\xi_n)$ ,  $c_{n-1} = H''(\xi_{n-1})$ , and so we have that

$$c_n = \frac{H''(\xi_n)}{H''(\xi_{n-1})}c_{n-1}, \quad d_n = -c_n\xi_n + c_{n-1}\xi_n + d_{n-1} \quad (3.6)$$

Therefore, we have enough equations to know the  $a_i - d_i$  in terms of  $\boldsymbol{\theta}$ .

Note that we choose a value of  $\lambda$ , fix the boundary conditions at  $\xi \rightarrow \infty$ , then solve the problem and subsequently recover the boundary conditions at  $\xi = 0$  ( $\kappa_I, \kappa_{II}$ ). This can then be inverted, so that we think of  $\lambda = \lambda(\kappa_I)$ . Physically, we know  $\kappa_I$ , and want to find  $\lambda$ , but in numerically solving the problem, it makes more sense to choose  $\lambda$  and recover  $\kappa_I$ .

The spacing of the points should reflect that the important part of the problem is happening near the tip, and this is where the points should be concentrated. The spacing that was typically used in numerical calculations was

$$\xi_i = \tan^2(\chi i/m), \quad i = 1, \dots, m < n \quad (3.7)$$

where  $\chi$  is chosen so that  $\tan^2(\chi) = O(10)$ , and the remaining points are added in a geometric progression, so that

$$\xi_{i+1} = (\xi_m/\xi_{m-1})\xi_i, \quad i = m, \dots, n-1 \quad (3.8)$$

### 3.2. Linear Perturbation Problem

From equation 2.21b, we anticipate a singularity of the form  $\xi^{s-1}$  in  $\tilde{H}'$ , (we still expect a  $\xi^{-1/2}$  singularity in  $\tilde{G}'$ ). Therefore, the interpolation was changed to reflect this. Some of the integrals no longer have exact expressions. In this case, they are calculated by a numerical integration routine.

The lubrication equation for the linear perturbation problem (2.20b), is linear in  $\tilde{H}$ . Therefore, we can obtain two expressions for  $\tilde{H}(\zeta_i)$  that are linear in  $\tilde{G}'(\xi_j)$ ,  $\tilde{H}'(\xi_j)$ . Together with the boundary conditions and beam theory asymptotics, (we haven't changed the integral kernel, so the asymptotics remain the same) there are enough equations to numerically solve the linear perturbation problem. There is no need to use Newton's method, as we can simply solve the linear set of equations.

### 3.3. Double Tip

In solving the problem of two tips situated at  $-L$  and  $0$ , an additional  $r$  points are taken to cover  $-L \leq \xi < 0$ . The spacing of points for  $\xi < 0$  was chosen so that there was a concentration of points near  $-L$  and near  $0$ .

We interpolate  $G'$  expecting a  $\xi^{-1/2}$  singularity at  $\xi = -L$ , and  $H'$  expecting a  $\xi^{-1/2}$  singularity at  $\xi = 0$ . We do not calculate  $\Pi$  for  $\xi < 0$  (although it is easily done), but just require that  $\sigma_{xy} = 0$  for  $\xi < 0$ . This provides enough equations for the problem to be solved as before, with Newton's method.

Note that we input  $-L$  and  $\lambda$  and recover  $\kappa_I$ ,  $\kappa_{II}$ , where  $\kappa_I$  is measured at  $0$ . Physically, for  $L > 0$ , we must have  $\kappa_I = 0$ . Numerically we solve for some  $\lambda$ ,  $L$ , find  $\kappa_I > 0$  and extrapolate to  $\kappa_I = 0$ .

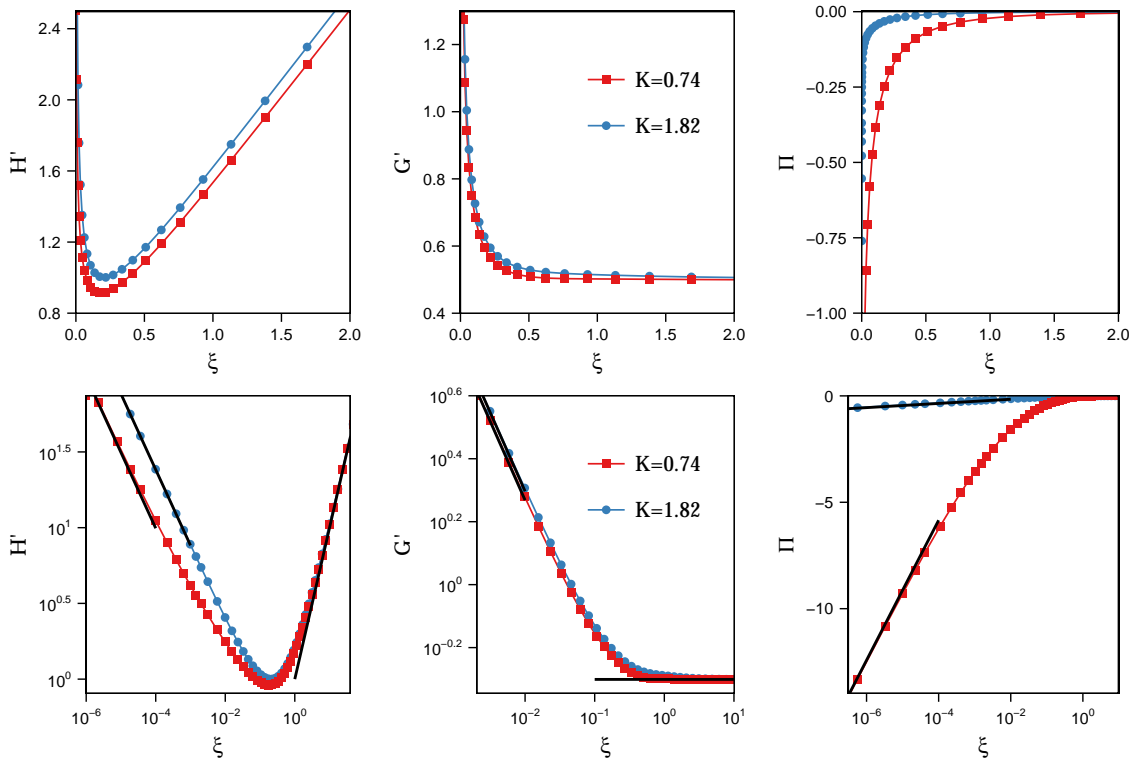


FIGURE 1. Numerical solutions for two typical values of  $\kappa_I$ . Logarithmic scales are shown, with solid lines indicating the predicted asymptotics;  $H', G' \propto \xi^{-1/2}$ ,  $\Pi \propto \log(\xi)$ , near  $\xi = 0$ ,  $H' \rightarrow \xi$ ,  $G' \rightarrow 1/2$  as  $\xi \rightarrow \infty$ . Figure produced with  $n = 465$ ,  $\xi_n = 819$ .

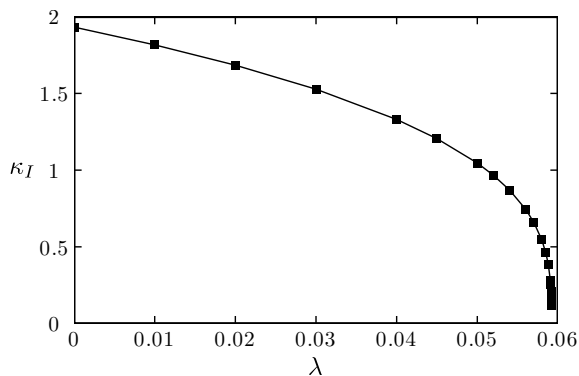


FIGURE 2. Here we vary the parameter  $\lambda$  and plot the change in  $\kappa_I$ . Figure produced with  $n = 465$ ,  $\xi_n = 819$ .

## 4. Results

### 4.1. Single tip

Start off with some of the basic graphs showing  $H', G'$ , and  $\Pi$  against  $\xi$ .

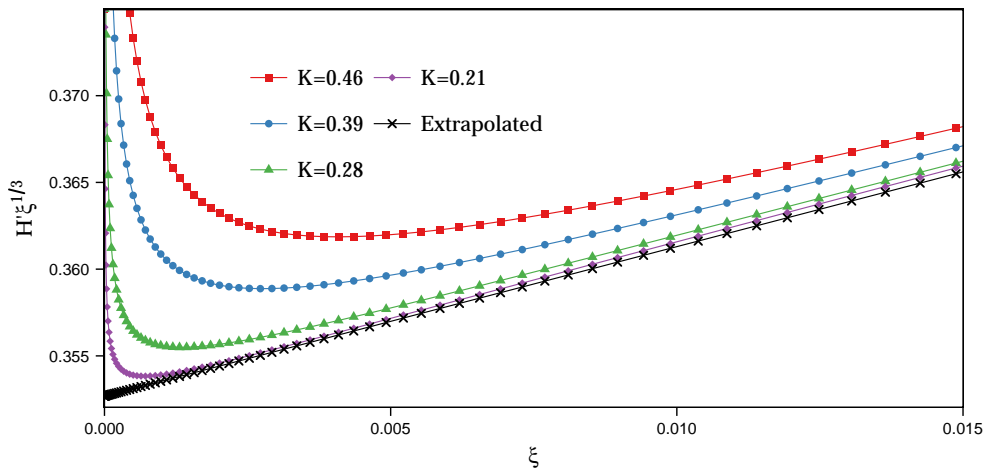


FIGURE 3. As  $\kappa_I \rightarrow 0$ ,  $H'$  moves from a  $\xi^{-1/2}$  singularity to a  $\xi^{-1/3}$  singularity. We can not calculate  $\kappa_I = 0$ , but the extrapolation to it is shown. Figure produced with  $n = 465$ ,  $\xi_n = 819$ .

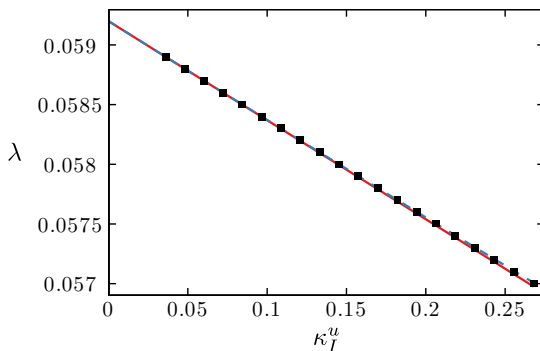


FIGURE 4. The numerical values of  $\kappa_I^u$  are plotted as points against the values of  $\lambda$ . A linear fit from the two smallest  $\kappa_I$  values is plotted as a solid line, a quadratic fit from the three smallest  $\kappa_I$  values is plotted as a dashed line. They are almost indistinguishable at this scale. The difference between the two extrapolations to  $\kappa_I = 0$ , provides an estimate of the error in calculating  $\lambda_0$ , (not accounting for the error due to  $n$ ), which in this instance is  $\approx 0.002\%$ . This figure was made with  $n = 524$ ,  $\xi_n = 846$ .

#### 4.2. Linear perturbation problem

We solve the linear perturbation problem. All that we really want to know is that we see the  $\xi^{s-1}$  behaviour that we expect, and we ask what the intercept of  $\tilde{H}_1$  is. It is perhaps worth mentioning the difficulties in measuring the intercept and perhaps a notion of the sensitivity of the result on the estimate provided for  $H_0$ . Illustrating that is the next figure

#### 4.3. Two tips

After the linear perturbation problem, we move on to the two tip problem. Perhaps some graphs that show an outline of the full numerical problem with non-zero  $\kappa_I$  and  $\kappa_{II}$ , although these are not physical.

We now move on to the  $\kappa_I = 0$  set of relations.

Then we could move on to talk about the decoupling between the fluid problem and the dry fracture problem. Relevant graphs to include would show that  $H$  really doesn't

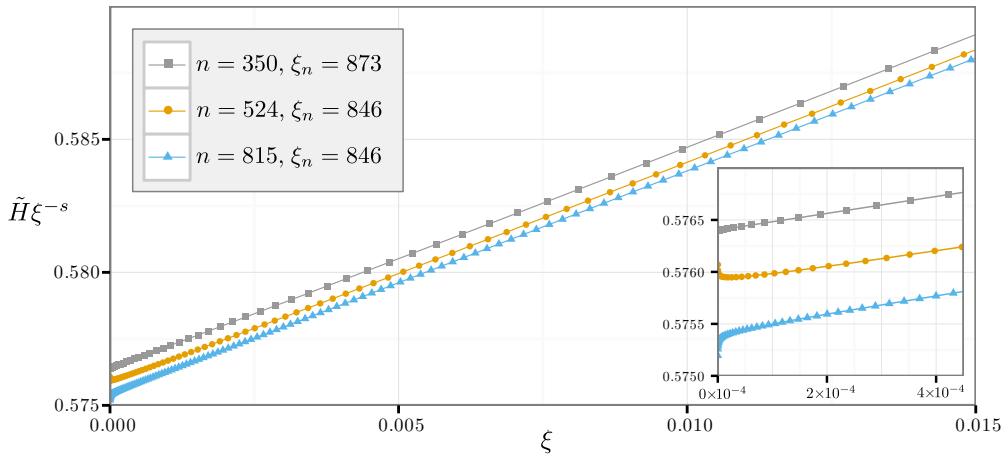


FIGURE 5. The numerical solution of the linear perturbation problem for a selection of resolutions. Of interest is the value of the intercept, which as shown is dependent on the resolution. Also shown is the numerical instability near the tip, due to our difficulty in calculating  $H_0$  for  $\xi \ll 1$ .

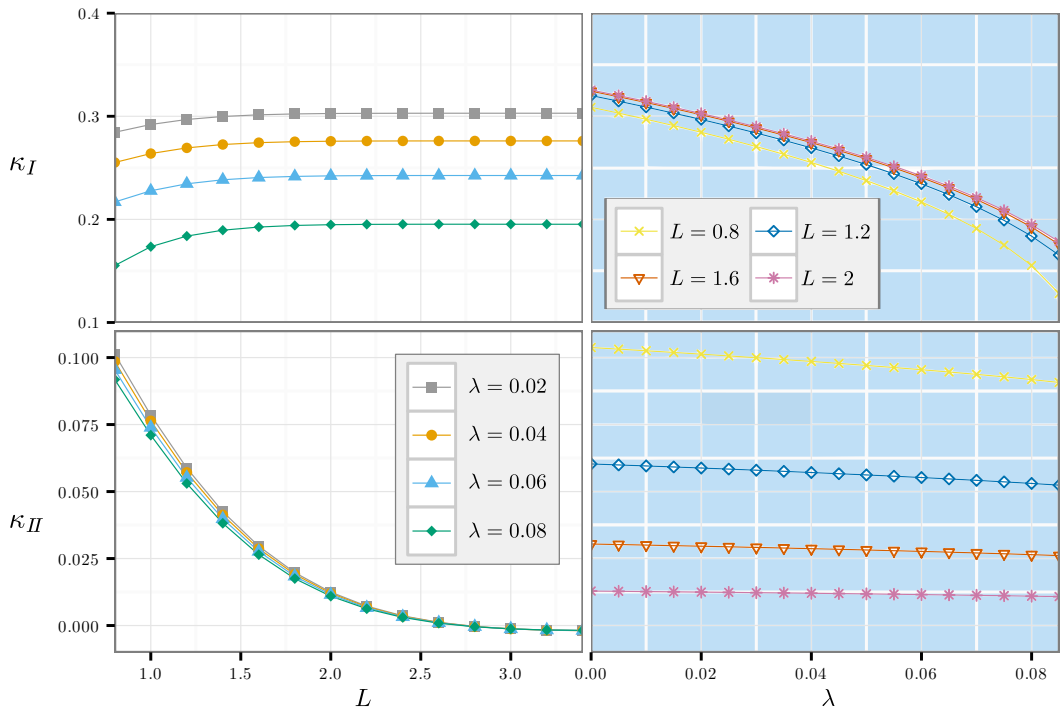


FIGURE 6. Some of the numerical results for the two tip problem. Having  $\kappa_I \neq 0$  at  $\xi = 0$  and  $L \neq 0$  is unphysical, but is what is found numerically. We can recover the physical solution by increasing  $\lambda$  for fixed  $L$  until  $\kappa_I = 0$ . Figure made with  $n = 995$ ,  $\xi_n = 846$ .

vary much with  $\lambda_0$ , and that given a reference  $H'$ , one can construct  $G'$  with relative ease.



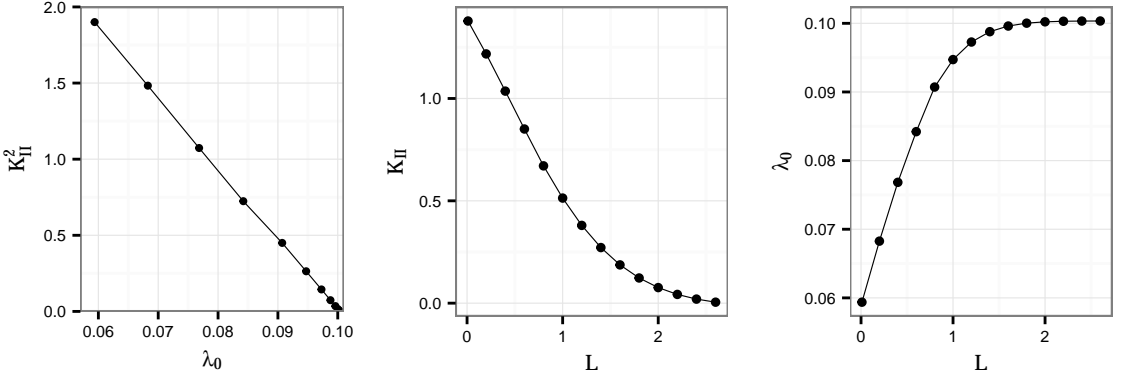


FIGURE 7. The results of extrapolating to  $\kappa_I = 0$ . Figure made with  $n = 995$ ,  $\xi_n = 846$ .

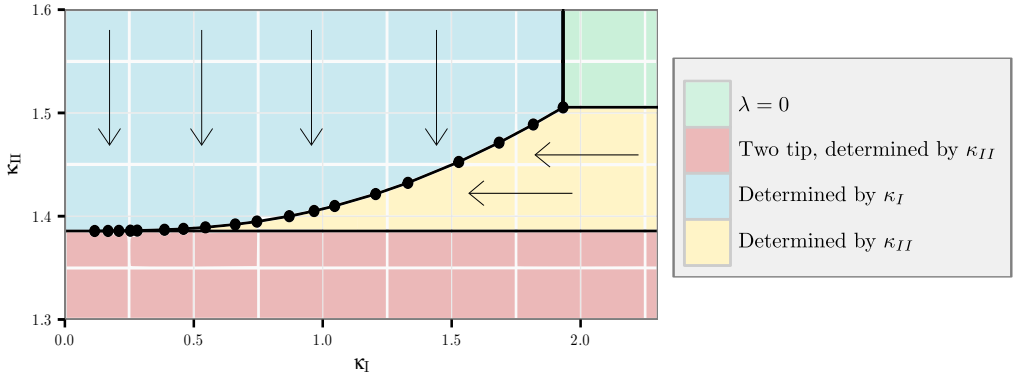


FIGURE 8. Given values  $(\kappa_I, \kappa_{II})$ , this graph determines which fracture regime occurs and so how  $\lambda$  and/or  $L$  should be calculated. Figure made with  $n = 465$ ,  $\xi_n = 819$ .

## 5. Discussion

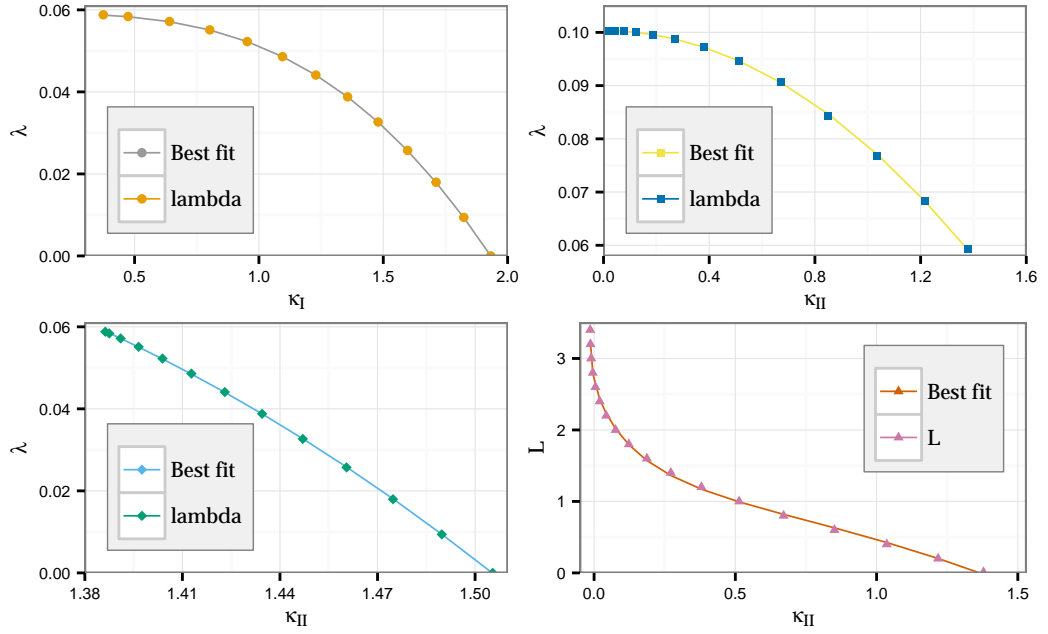


FIGURE 9. The formula valid for all  $\kappa_I$ . For the single tip calculations,  $n = 815$  was used, for the double tip  $n = 995$ .  $\xi_n = 846$  in both cases.

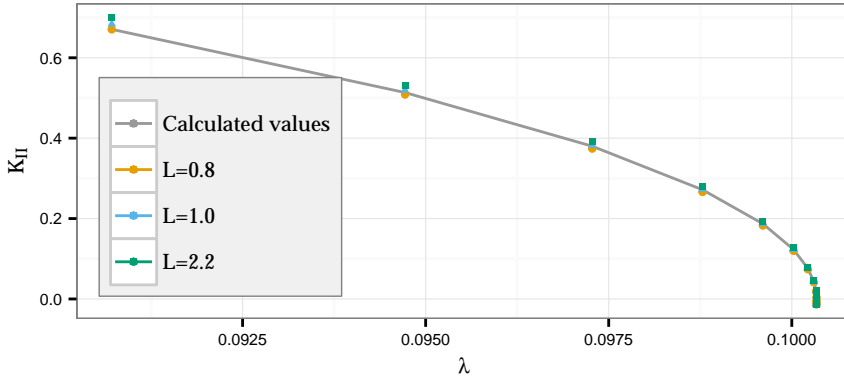


FIGURE 10. Reconstructing the full solution given a reference  $H'$ . Figure made with  $n = 995$ ,  $\xi_n = 846$ .

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$a/d$	$M = 4$	$M = 8$	Callan <i>et al.</i>
0.1	1.56905	1.56	1.56904
0.3	1.50484	1.504	1.50484
0.55	1.39128	1.391	1.39131
0.7	1.32281	10.322	1.32288
0.913	1.34479	100.351	1.35185

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TABLE 1. Values of  $kd$  at which trapped modes occur when  $\rho(\theta) = a$ 

## 6. Citations and references

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Hwang & Tuck (1970, pp. 2-4)

`\citep[p. 6]{Worster92}`:

(Worster 1992, p. 6)

`\citep[see][]{Koch83, Lee71, Linton92}`:

(see Koch 1983; Lee 1971; Linton & Evans 1992)

`\citep[see][p. 18]{Martin80}`:

(see Martin 1980, p. 18)

`\citep{Brownell104, Brownell107, Ursell150, Wijngaarden68, Miller91}`:

(Brownell & Su 2004, 2007; Ursell 1950; van Wijngaarden 1968; Miller 1991)

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Where there are up to ten authors, all authors’ names should be given in the reference list. Where there are more than ten authors, only the first name should appear, followed by *et al.*

Acknowledgements should be included at the end of the paper, before the References section or any appendices, and should be a separate paragraph without a heading. Several anonymous individuals are thanked for contributions to these instructions.

## Appendix A

This appendix contains sample equations in the JFM style. Please refer to the L<sup>A</sup>T<sub>E</sub>X source file for examples of how to display such equations in your manuscript.

$$(\nabla^2 + k^2)G_s = (\nabla^2 + k^2)G_a = 0 \quad (\text{A } 1)$$

$$\nabla \cdot \mathbf{v} = 0, \quad \nabla^2 P = \nabla \cdot (\mathbf{v} \times \mathbf{w}). \quad (\text{A } 2)$$

$$G_s, G_a \sim 1/(2\pi) \ln r \quad \text{as} \quad r \equiv |P - Q| \rightarrow 0, \quad (\text{A } 3)$$

$$\left. \begin{aligned} \frac{\partial G_s}{\partial y} &= 0 \quad \text{on} \quad y = 0, \\ G_a &= 0 \quad \text{on} \quad y = 0, \end{aligned} \right\} \quad (\text{A } 4)$$

$$-\frac{1}{2\pi} \int_0^\infty \gamma^{-1} [\exp(-k\gamma|y-\eta|) + \exp(-k\gamma(2d-y-\eta))] \cos k(x-\xi)t \, dt, \quad 0 < y, \quad \eta < d, \quad (\text{A } 5)$$

$$\gamma(t) = \begin{cases} -i(1-t^2)^{1/2}, & t \leq 1 \\ (t^2-1)^{1/2}, & t > 1. \end{cases} \quad (\text{A } 6)$$

$$-\frac{1}{2\pi} \int_0^\infty B(t) \frac{\cosh k\gamma(d-y)}{\gamma \sinh k\gamma d} \cos k(x-\xi)t \, dt$$

$$G = -\frac{1}{4}i(H_0(kr) + H_0(kr_1)) - \frac{1}{\pi} \int_0^\infty \frac{e^{-k\gamma d}}{\gamma \sinh k\gamma d} \cosh k\gamma(d-y) \cosh k\gamma(d-\eta) \quad (\text{A } 7)$$

Note that when equations are included in definitions, it may be suitable to render them in line, rather than in the equation environment:  $\mathbf{n}_q = (-y'(\theta), x'(\theta))/w(\theta)$ . Now  $G_a = \frac{1}{4}Y_0(kr) + \widetilde{G}_a$  where  $r = \{[x(\theta) - x(\psi)]^2 + [y(\theta) - y(\psi)]^2\}^{1/2}$  and  $\widetilde{G}_a$  is regular as  $kr \rightarrow 0$ . However, any fractions displayed like this, other than  $\frac{1}{2}$  or  $\frac{1}{4}$ , must be written on the line, and not stacked (ie 1/3).

$$\begin{aligned} \frac{\partial}{\partial n_q} \left( \frac{1}{4}Y_0(kr) \right) &\sim \frac{1}{4\pi w^3(\theta)} [x''(\theta)y'(\theta) - y''(\theta)x'(\theta)] \\ &= \frac{1}{4\pi w^3(\theta)} [\rho'(\theta)\rho''(\theta) - \rho^2(\theta) - 2\rho'^2(\theta)] \quad \text{as} \quad kr \rightarrow 0. \end{aligned} \quad (\text{A } 8)$$

$$\frac{1}{2}\phi_i = \frac{\pi}{M} \sum_{j=1}^M \phi_j K_{ij}^a w_j, \quad i = 1, \dots, M, \quad (\text{A } 9)$$

where

$$K_{ij}^a = \begin{cases} \partial G_a(\theta_i, \theta_j) / \partial n_q, & i \neq j \\ \partial \widetilde{G}_a(\theta_i, \theta_i) / \partial n_q + [\rho'_i \rho''_i - \rho_i^2 - 2\rho_i'^2] / 4\pi w_i^3, & i = j. \end{cases} \quad (\text{A } 10)$$

$$\rho_l = \lim_{\zeta \rightarrow Z_l^-(x)} \rho(x, \zeta), \quad \rho_u = \lim_{\zeta \rightarrow Z_u^+(x)} \rho(x, \zeta) \quad (\text{A } 11a, b)$$

$$(\rho(x, \zeta), \phi_{\zeta\zeta}(x, \zeta)) = (\rho_0, N_0) \quad \text{for} \quad Z_l(x) < \zeta < Z_u(x). \quad (\text{A } 12)$$

$$\tau_{ij} = (\overline{u_i u_j} - \overline{u_i} \overline{u_j}) + (\overline{u_i u_j^{SGS}} + \overline{u_i^{SGS} u_j}) + \overline{u_i^{SGS} u_j^{SGS}}, \quad (\text{A } 13a)$$

$$\tau_j^\theta = (\overline{u_j \theta} - \overline{u_j} \overline{\theta}) + (\overline{u_j \theta^{SGS}} + \overline{u_j^{SGS} \theta}) + \overline{u_j^{SGS} \theta^{SGS}}. \quad (\text{A } 13b)$$

$$\mathbf{Q}_C = \begin{bmatrix} -\omega^{-2}V'_w & -(\alpha^t\omega)^{-1} & 0 & 0 & 0 \\ \frac{\beta}{\alpha\omega^2}V'_w & 0 & 0 & 0 & i\omega^{-1} \\ i\omega^{-1} & 0 & 0 & 0 & 0 \\ iR_\delta^{-1}(\alpha^t + \omega^{-1}V''_w) & 0 & -(i\alpha^t R_\delta)^{-1} & 0 & 0 \\ \frac{i\beta}{\alpha\omega}R_\delta^{-1}V''_w & 0 & 0 & 0 & 0 \\ (i\alpha^t)^{-1}V'_w & (3R_\delta^{-1} + c^t(i\alpha^t)^{-1}) & 0 & -(\alpha^t)^{-2}R_\delta^{-1} & 0 \end{bmatrix}. \quad (\text{A } 14)$$

$$\boldsymbol{\eta}^t = \hat{\boldsymbol{\eta}}^t \exp[i(\alpha^t x_1^t - \omega t)], \quad (\text{A } 15)$$

where  $\hat{\boldsymbol{\eta}}^t = \mathbf{b} \exp(i\gamma x_3^t)$ .

$$\text{Det}[\rho\omega^2\delta_{ps} - C_{pqrs}^t k_q^t k_r^t] = 0, \quad (\text{A } 16)$$

$$\langle k_1^t, k_2^t, k_3^t \rangle = \langle \alpha^t, 0, \gamma \rangle \quad (\text{A } 17)$$

$$\mathbf{f}(\theta, \psi) = (g(\psi) \cos \theta, g(\psi) \sin \theta, f(\psi)). \quad (\text{A } 18)$$

$$f(\psi_1) = \frac{3b}{\pi[2(a+b\cos\psi_1)]^{3/2}} \int_0^{2\pi} \frac{(\sin\psi_1 - \sin\psi)(a+b\cos\psi)^{1/2}}{[1-\cos(\psi_1-\psi)](2+\alpha)^{1/2}} dx, \quad (\text{A } 19)$$

$$\begin{aligned} g(\psi_1) = & \frac{3}{\pi[2(a+b\cos\psi_1)]^{3/2}} \int_0^{2\pi} \left( \frac{a+b\cos\psi}{2+\alpha} \right)^{1/2} \left\{ f(\psi)[(\cos\psi_1 - b\beta_1)S + \beta_1 P] \right. \\ & \times \frac{\sin\psi_1 - \sin\psi}{1-\cos(\psi_1-\psi)} + g(\psi) \left[ \left( 2+\alpha - \frac{(\sin\psi_1 - \sin\psi)^2}{1-\cos(\psi-\psi_1)} - b^2\gamma \right) S \right. \\ & \left. \left. + \left( b^2 \cos\psi_1 \gamma - \frac{a}{b} \alpha \right) F\left(\frac{1}{2}\pi, \delta\right) - (2+\alpha) \cos\psi_1 E\left(\frac{1}{2}\pi, \delta\right) \right] \right\} d\psi, \end{aligned} \quad (\text{A } 20)$$

$$\alpha = \alpha(\psi, \psi_1) = \frac{b^2[1-\cos(\psi-\psi_1)]}{(a+b\cos\psi)(a+b\cos\psi_1)}, \quad \beta - \beta(\psi, \psi_1) = \frac{1-\cos(\psi-\psi_1)}{a+b\cos\psi}. \quad (\text{A } 21)$$

$$\left. \begin{aligned} H(0) &= \frac{\epsilon \bar{C}_v}{\tilde{v}_T^{1/2}(1-\beta)}, \quad H'(0) = -1 + \epsilon^{2/3} \bar{C}_u + \epsilon \hat{C}'_u; \\ H''(0) &= \frac{\epsilon u_*^2}{\tilde{v}_T^{1/2} u_P^2}, \quad H'(\infty) = 0. \end{aligned} \right\} \quad (\text{A } 22)$$

LEMMA 1. Let  $f(z)$  be a trial Batchelor (1971, pp. 231–232) function defined on  $[0, 1]$ . Let  $\Lambda_1$  denote the ground-state eigenvalue for  $-d^2g/dz^2 = \Lambda g$ , where  $g$  must satisfy  $\pm dg/dz + \alpha g = 0$  at  $z = 0, 1$  for some non-negative constant  $\alpha$ . Then for any  $f$  that is not identically zero we have

$$\frac{\alpha(f^2(0) + f^2(1)) + \int_0^1 \left( \frac{df}{dz} \right)^2 dz}{\int_0^1 f^2 dz} \geq \Lambda_1 \geq \left( \frac{-\alpha + (\alpha^2 + 8\pi^2\alpha)^{1/2}}{4\pi} \right)^2. \quad (\text{A } 23)$$

COROLLARY 1. Any non-zero trial function  $f$  which satisfies the boundary condition  $f(0) = f(1) = 0$  always satisfies

$$\int_0^1 \left( \frac{df}{dz} \right)^2 dz. \quad (\text{A } 24)$$

## REFERENCES

- BATCHELOR, G. K. 1971 Small-scale variation of convected quantities like temperature in turbulent fluid. part 1. general discussion and the case of small conductivity. *J. Fluid Mech.* **5**, 113–133.
- BROWNELL, C. J. & SU, L. K. 2004 Planar measurements of differential diffusion in turbulent jets. *AIAA Paper 2004-2335*.
- BROWNELL, C. J. & SU, L. K. 2007 Scale relations and spatial spectra in a differentially diffusing jet. *AIAA Paper 2007-1314*.
- DENNIS, S. C. R. 1985 Compact explicit finite difference approximations to the Navier–Stokes equation. In *Ninth Intl Conf. on Numerical Methods in Fluid Dynamics* (ed. Soubbaramayer & J. P. Boujot), *Lecture Notes in Physics*, vol. 218, pp. 23–51. Springer.
- HWANG, L.-S. & TUCK, E. O. 1970 On the oscillations of harbours of arbitrary shape. *J. Fluid Mech.* **42**, 447–464.
- KOCH, W. 1983 Resonant acoustic frequencies of flat plate cascades. *J. Sound Vib.* **88**, 233–242.
- LEE, J.-J. 1971 Wave-induced oscillations in harbours of arbitrary geometry. *J. Fluid Mech.* **45**, 375–394.
- LINTON, C. M. & EVANS, D. V. 1992 The radiation and scattering of surface waves by a vertical circular cylinder in a channel. *Phil. Trans. R. Soc. Lond.* **338**, 325–357.
- MARTIN, P. A. 1980 On the null-field equations for the exterior problems of acoustics. *Q. J. Mech. Appl. Maths* **33**, 385–396.
- MILLER, P. L. 1991 Mixing in high schmidt number turbulent jets. PhD thesis, California Institute of Technology.
- ROGALLO, R. S. 1981 Numerical experiments in homogeneous turbulence. *Tech. Rep.* 81835. NASA Tech. Mem.
- URSELL, F. 1950 Surface waves on deep water in the presence of a submerged cylinder i. *Proc. Camb. Phil. Soc.* **46**, 141–152.
- VAN WIJNGAARDEN, L. 1968 On the oscillations near and at resonance in open pipes. *J. Engng Maths* **2**, 225–240.
- WORSTER, M. G. 1992 The dynamics of mushy layers. In *Interactive dynamics of convection and solidification* (ed. S. H. Davis, H. E. Huppert, W. Muller & M. G. Worster), pp. 113–138. Kluwer.