Numerical methods

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Consider the govering equations
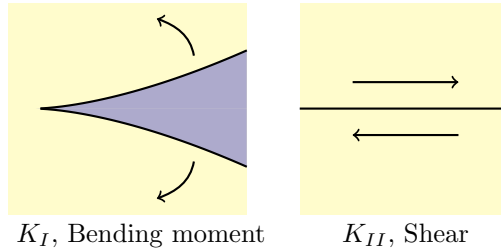
$$\left(\begin{array}{c} p(z) \\ 0 \end{array} \right) = \left(\begin{array}{c} \sigma_y \\ \tau_{xy} \end{array} \right) = \int_0^\infty \left(\begin{array}{cc} K_{11}(x-z) & K_{12}(x-z) \\ K_{21}(x-z) & K_{22}(x-z) \end{array} \right) \left(\begin{array}{c} g'(x) \\ h'(x) \end{array} \right) dx \quad (1)$$

$$h^2 p' = \lambda \quad (2)$$

Have the "input" parameters as

- BC's P, M (or equivalently g', h'' at $x \to \infty$)

- λ , the speed

Want to solve for the toughness K_I and K_{II}. In this project, we have so far focused on K_I.

K_I, Bending moment $\qquad$ K_{II}, Shear

Goal: Find λ such that $K_I(\lambda) = 0$, "Zero toughness solution". Given this we then want to investigate the behaviour for small $K_I \approx 0$. To do this, take some given value of λ and then solve equations 1, 2.

Discretization of problem

The method chosen to discretize the problem is to take a vector $(x_1, \ldots, x_n)$ of n points at which we measure g', h' and have a vector $(z_1, \ldots, z_{n-1})$ of $n - 1$ intermediate points at which p is measured. The spacing chosen is a $\tan^2$ spacing.

The "obvious" way to interpolate h' in between the x_i's is simple linear interpolation. But both h', g' become singular near 0. However, we expect a $x^{-1/2}$

x, the n points at which h', g' are measured

z, the $n-1$ points at which p is measured

singularity, which allows us to "remove" said singularity. The interpolation used is

$$g'(x) = \begin{cases} \frac{1}{\sqrt{x}}(a_i x + b_i) & i < t \\ a_i x + b_i & i \geq t \end{cases}$$

for x in the spline $x \in [x_i, x_{i+1}]$. Choose $1 < t < n$, typically $t = n/2$. Similarly

$$h'(x) = \begin{cases} \frac{1}{\sqrt{x}}(c_i x + d_i) & i < t \\ c_i x + d_i & i \geq t \end{cases}$$

With the same t used. We will interpolate beyond x_n via $g'(x) = a_n x + b_n$, for $x > x_n$, and similar for h'.

The values of g', h' are stored via

$$\boldsymbol{\theta} = \begin{pmatrix} a_1 x_1 + b_1 \\ \vdots \\ a_n x_n + b_n \\ c_1 x_1 + d_1 \\ \vdots \\ c_n x_n + d_n \end{pmatrix}$$

Once one has $\boldsymbol{\theta}$, it is trivial to recover, say $g'(x_i)$, since either $g'(x_i) = \boldsymbol{\theta}_i$ or $g'(x_i) = \boldsymbol{\theta}_i/\sqrt{x_i}$. Similarly, given $g'(x_i)$; $\boldsymbol{\theta}_i$ can be calculated.

Recovering the a_i's

Suppose we know $\boldsymbol{\theta}$, (and always assume we know the x_i). Can we recover a_i, b_i, c_i, d_i? The answer is yes, once we add in the boundary conditions at ∞.

2

Further we have that

$$\gamma = \begin{pmatrix} a_1 \\ \vdots \\ a_n \\ b_1 \\ \vdots \\ b_n \\ c_1 \\ \vdots \\ c_n \\ d_1 \\ \vdots \\ d_n \end{pmatrix} = T\boldsymbol{\theta}$$

Where T is a $4n \times n$ interpolation matrix. A quick check reveals we have $4n$ unknowns, in $\boldsymbol{\gamma}$. Knowing θ provides $2n$ equations. Demanding continuity of the interpolated g', h' provides another $2(n-1)$ equations, (match at $x_2, \ldots, x_n$). Finally boundary conditions on the spline at ∞ provide another 2 equations.

The continuity conditions are

$$a_1 x_2 + b_1 = a_2 x_2 + b_2$$

$$\vdots$$

$$a_{t-2} x_{t-1} + b_{t-2} = a_{t-1} x_{t-1} + b_{t-1}$$
$$(a_{t-1} x_t + b_{t-1})/\sqrt{x_t} = a_t x_t + b_t$$
$$a_t x_{t+1} + b_t = a_{t+1} x_{t+1} + b_{t+1}$$

$$\vdots$$

$$a_{n-1} x_n + b_{n-1} = a_n x_n + b_n$$

and similar for c, d. This means that with the exceptions of $i = t - 1, n$, have that

$$\frac{\theta_{i+1} - \theta_i}{x_{i+1} - x_i} = a_i$$

$$\frac{\theta_i x_{i+1} - \theta_{i+1} x_i}{x_{i+1} - x_i} = b_i$$

Same idea with x_{t-1}, just have to be a little careful about the switch in the continuity condition,

$$\frac{\sqrt{x_t}\theta_t - \theta_{t-1}}{x_t - x_{t-1}} = a_{t-1}$$

$$\frac{\theta_{t-1} x_t - \theta_t x_{t-1}\sqrt{x_t}}{x_t - x_{t-1}} = b_{t-1}$$

3

So we are almost done, just missing 4 rows in our matrix. Have the nth row as all zeros, i.e. $a_n = 0$ due to boundary conditions, and so trivially the $2n^{th}$ row is 0 except $T_{2n,n} = 1$ Now, B.C. for h' implies $h''(x_n) \approx h''(x_{n-1})$ i.e. $c_n = c_{n-1}$ and thus from continuity $d_n = d_{n-1}$. This completes our interpolation matrix T. We still have some extra boundary conditions to impose, which we will do later. Naively, these are $g'(x_n) = b_n = 1/2$, $c_n = 1$. These are approximately true, but we need to do a bit better. (I will add this into the explanation once I understand it...)

Analytic expressions

Now that we have made the piecewise analytic approximation[1] we can avoid making any more approximations. Recall K_{ij} has an analytic expression (even better, it's a rational function). If we take h', g' to be piecewise analytic, then the integrand becomes an analytic expression that can be exactly integrated. For example,

$$p(z) = \int_0^\infty K_{11}(x-z)g'(x) + K_{12}(x-z)h'(x) \, dx$$

$$= \sum_{i=1}^{t-1} \int_0^\infty K_{11}(x-z)(a_i x + b_i)/\sqrt{x} + K_{12}(x-z)(c_i x + d_i)/\sqrt{x} \, dx$$

$$+ \sum_{i=t}^{n} \int_0^\infty K_{11}(x-z)(a_i x + b_i) + K_{12}(x-z)(c_i x + d_i) \, dx$$

The right hand side of this equation may look ghastly, (and we haven't even expanded the K_{ij}'s yet ...) but it is an analytic expression in z. Further, we can do the integration before knowing any of the values of a, b, c, d. It is also clear that $p(z)$ is linear in a, b, c, d. Therefore, once we know the spacing of the z_i's we see that

$$\begin{pmatrix} p(z_1) \\ \vdots \\ p(z_{n-1}) \\ 0 \\ \vdots \\ 0 \end{pmatrix} = \begin{pmatrix} B_{1,1} & \cdots & B_{1,2n} \\ \vdots & \ddots & \vdots \\ B_{2(n-1),1} & \cdots & B_{2(n-1),2n} \end{pmatrix} \boldsymbol{\gamma} = BT\boldsymbol{\theta}$$

Where the matrix B depends on the choice of spacings (x, z), but does not depend on $\boldsymbol{\gamma}$.

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[1] I think I have made up some terminology here, piecewise linear isn't quite right due to the $x^{-1/2}$ parts, and "sometimes piecewise linear sometimes piecewise $x^{-1/2} \times$(linear function)" is not ideal either

We can go further and incorporate the boundary conditions into this equation. The discretized versions of the boundary conditions, become $g'(x_n) = 1/2$ and $\dfrac{h'(x_n) - h'(x_{n-1})}{x_n - x_{n-1}} = 1$. [2] These conditions are linear in terms of g', h', and so by adding another two rows onto the matrix BT, get that

$$
\begin{pmatrix} p(z_1) \\ \vdots \\ p(z_{n-1}) \\ 0 \\ \vdots \\ 0 \\ g'(\infty) \\ h''(\infty) \end{pmatrix} = \begin{pmatrix} A_{1,1} & \cdots & A_{1,2n} \\ \vdots & \ddots & \vdots \\ A_{2n,1} & \cdots & A_{2n,2n} \end{pmatrix} \boldsymbol{\theta}
$$

Where $g'(\infty), h''(\infty)$ are the (constant) boundary conditions.

Now we use the second equation for p, namely $p = \int_z^\infty \lambda/h^2 dx$. We can integrate our piecewise analytic expression for h' to recover h, imposing both $h(x_1) = 0$ as well as continuity at the x_i for $i = 2, \ldots n$. The result is

$$
h(x) = \begin{cases} \sqrt{x}(w_i x + e_i) + r_i & i < t \\ w_i x^2 + e_i x + r_i & i \geq t \end{cases}
$$

for some constants w, e, r. These constants are related to $\boldsymbol{\gamma}$, and so $\boldsymbol{\theta}$, linearly. Given this piecewise analytic expression for h, we can find an analytic expression for $p(z)$. Since we know the spacings, we have

$$
f(\boldsymbol{\theta}) = \begin{vmatrix} p(z_1) \\ \vdots \\ p(z_{n-1}) \\ 0 \\ \vdots \\ 0 \\ g'(\infty) \\ h''(\infty) \end{vmatrix} = A\boldsymbol{\theta}
$$

Where we now just need to solve for $\boldsymbol{\theta}$.

Newton's method

Suppose $\boldsymbol{\theta}$ is iterate 1. To get the next iterate you need to solve (to first order)

$$
f(\boldsymbol{\theta} + \delta\boldsymbol{\theta}) = A(\boldsymbol{\theta} + \delta\boldsymbol{\theta})
$$

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[2] Again, we need to do a bit better than this. For now, this illustrates the point.

$$f(\boldsymbol{\theta}) + (Df|_{\boldsymbol{\theta}})(\delta\boldsymbol{\theta}) = A\boldsymbol{\theta} + A\delta\boldsymbol{\theta}$$

Where $Df|_{\boldsymbol{\theta}}$ is a matrix of partial derivatives. Therefore, get to first order that

$$\delta\boldsymbol{\theta} = (A - Df|_{\boldsymbol{\theta}})^{-1}(f(\boldsymbol{\theta}) - A\boldsymbol{\theta})$$

Ingredients:

- Matrix A itself (of which the $2(n-1) \times 2n$ part is the integral kernel)

- The function $f(\boldsymbol{\theta})$. I.e. given $\boldsymbol{\theta}$ you need to calculate $\int_z^\infty \lambda/h^2 dx$ (Key functions "hprime_to_h" and "hprime_to_p").

- Need to calculate Df which involves calculating $\dfrac{\partial}{\partial\boldsymbol{\theta}} \displaystyle\int_z^\infty \dfrac{\lambda}{h(x)^2} dx$

So we have worked out numerically $K(\lambda)$, now we want to solve $K(\lambda_0) = 0$ for λ_0. We do a "march". Sublety in that $K < 0$ is unphysical, so a guess of $\lambda > \lambda_0$ where $K(\lambda_0) = 0$ does not make any physical sense (& will get bad numerical results). To get around this difficulty, take the next iterate of λ as smaller than predicted.

Figure 1: March to find λ_0
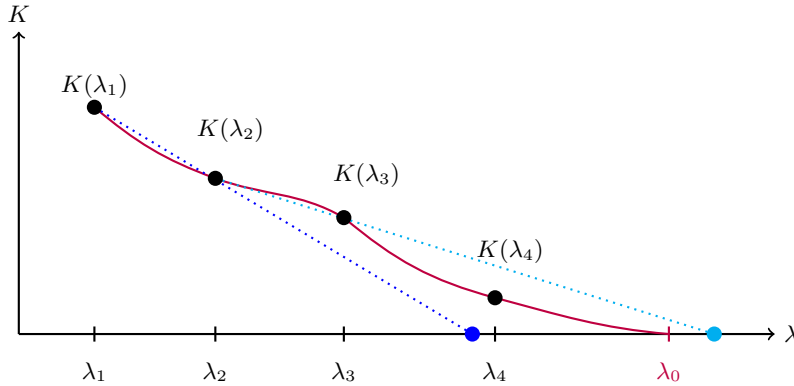
For example in Figure 1, the obvious choice for λ_4 (the light blue circle) is larger than the true value of λ_0, and therefore the naive extrapolation method won't quite work.

Guide to programs

Important note: What I have called $\boldsymbol{\theta}$ in this document is (confusingly) called hprime in the code. From now on I will call it $\boldsymbol{h}'$, in bold to try to avoid confusion.

K_of_c_march

First the program sets up the spacing as $\tan^2$. It also sets the initial $\boldsymbol{h}' = (\underbrace{1, \ldots, 1}_{g'} \underbrace{x_1 + 1, \ldots, x_n + 1}_{h'})$ (This first guess is based on the the the boundary conditions at ∞.)

Most of the work is then done by `fixed_lambda_M_iteration` which then solves for K_I and $\boldsymbol{h}'$.

N.B. $\boldsymbol{h}'$ is updated via $\boldsymbol{h}'_i = \dfrac{\boldsymbol{h}'_{i-1} - \boldsymbol{h}'_{i-2}}{\lambda_{i-1} - \lambda_{i-2}} \lambda_i + \dfrac{\lambda_{i-1} \boldsymbol{h}'_{i-2} - \lambda_{i-1} \boldsymbol{h}'_{i-1}}{\lambda_{i-1} - \lambda_{i-2}}$ which is just linear extrapolation. In the absence of any better ideas this is the sensible choice.

After iterating for a few values, get near λ_0. Hear we suspect that something like $K^3 \sim \lambda - \lambda_0$ near $\lambda = \lambda_0$, $K = 0$. So given two prior guesses, extrapolate via $\lambda_i = \dfrac{K_{i-1}^3 \lambda_{i-2} - K_{i-2}^3 \lambda_{i-1}}{K_{i-1}^3 - K_{i-2}^3}$ But as noted earlier, must be careful to not extrapolate further than λ_0. So an idea is to take $(\lambda_i + \lambda_{i-1})/2$ as the next guess, i.e.

$$\lambda_i = \frac{\lambda_{i-1} - \lambda_{i-2}}{K_{i-1}^3 - K_{i-2}^3} \frac{K_{i-1}^3}{2} + \frac{K_{i-1}^3 \lambda_{i-2} - K_{i-2}^3 \lambda_{i-1}}{K_{i-1}^3 - K_{i-2}^3}$$

Then the program just iterates. If it doesn't converge, it simply tries a smaller value of λ.

fixed_lambda_M_iteration

Takes a value of λ and returns the corresponding K value. Now using the "scaled" version, so the values of P, M are implicitly assumed to be $0, 1$ respectively.

Somewhat concerningly, $\boldsymbol{h}'$ is assumed to already have this spacing, which could potentially cause issues. If you wanted to change the spacing you would have to do it in two different places.

Subroutines then return the kernel matrix & the interpolate matrix. The kernel matrix is in lieu of $\begin{pmatrix} p \\ 0 \end{pmatrix} = \int \underline{\underline{K}} \begin{pmatrix} g' \\ h' \end{pmatrix}$. The interpolate matrix actually only appears here to add in boundary conditions to the matrix A.

The matrix A is set up, which is part kernel, part interpolate matrix, same matrix as described earlier.

The rcond statement is testing how conditioned the matrix A is, or how ameniable it is to being numerically inverted. There does not seem to be a huge amount of point in adding this step in.

Then the iteration loop begins. Follows Newton's method for the equation $f(\boldsymbol{h}') = A\boldsymbol{h}'$ and iterates via $\boldsymbol{h}'_{new} = \boldsymbol{h}'_{old} + (A - Df|_{\boldsymbol{h}'_{old}})^{-1}(f(\boldsymbol{h}'_{old}) - A\boldsymbol{h}'_{old})$ Where we already know A. f, Df are provided via `hprime_to_p` and f, Df are called p, dp respectively in the program.

pprime_to_p

This function is given h, x, z (h in the form of the h_coeff vector) and returns

$$\frac{1}{\lambda}\begin{pmatrix} p(z_1) \\ \vdots \\ p(z_{n-1}) \end{pmatrix} \text{ via the equation } p = -\lambda \int_z^\infty \frac{1}{h^2}dx$$

Again, once one has made the approximation of h as piecewise analytic, can actually work out the exact integrals. So all this function really has to do is to evaluate some known functions, plugging in the values of h_coeff.

pressure_map_derivative

This function calculates Df or $\frac{\partial}{\partial\boldsymbol{\theta}}f$. Most of f is either zeros, or the boundary conditions which we assume to be constant. (actually, our more refined guess of the boundary conditions makes them non-constant, but `bending_p_adjust` will do these corrections.

The easiest way to think of p is as a function of the coefficients of h, i.e. the w, e, r. So $p(z) = p(z; w, e, r)$, and actually we can find analytic expressions for things like $\frac{\partial}{\partial e_i}p$. Now there is a linear relationship between (w, e, r) and $\boldsymbol{\theta}$, say $(w, e, r) = H\boldsymbol{\theta}$ where H is the h_coeff matrix, provided by `h_coefficient_matrix.m`

We use this, together with the chain rule, to find an expression for Df.

hprime_to_p

finds p and dp given $\boldsymbol{h}'$ The pressure is found via `pprime_to_p`. dp is found via `pressure_map_derivative`.

Then `bending_p_adjust` corrects the shoddy boundary conditions in both p and dp.

hprime_to_h

This function takes h' and returns H the coefficient matrix for h. Starting point, suppose we know the c, d coefficients for h'. Then have that

$$h'(x) = \begin{cases} \frac{1}{\sqrt{x}}(c_i x + d_i) & i < t \\ c_i x + d_i & i \geq t \end{cases} \quad x \in [x_i, x_{i+1}]$$

Then integrating gives that

$$h'(x) = \begin{cases} w_i x^{3/2} + e_i x^{1/2} + r_i & i < t \\ w_i x^2 + e_i x + r_i & i \geq t \end{cases} \quad x \in [x_i, x_{i+1}]$$

8

where

$$w_i = \begin{cases} \frac{2}{3}c_i & i < t \\ \frac{1}{2}c_i & i \geq t \end{cases}$$

$$e_i = \begin{cases} 2d_i & i < t \\ d_i & i \geq t \end{cases}$$

We recover the r_i's through imposing continuity and $c_1 = 0$. Get that

$$r_i = \begin{cases} r_i + \frac{2}{3}(c_i - c_{i+1})x_{i+1}^{3/2} + 2(d_i - d_{i+1})x_{i+1}^{2/3} & i < t \\ r_{i+1} + \frac{1}{2}(c_i - c_{i+1})x_{i+1}^2 + (d_i - d_{i+1})x_{i+1} & i \geq t \end{cases}$$

Which looks complicated, but can see that r is related to c, d linearly.

So all this function does is return a matrix H. Important note: The coefficients of h are stored $(w_1, e_1, r_1, \ldots, w_n, e_n, r_n)$, rather than what you might expect, $(w_1, \ldots, w_n, e_1, \ldots, e_n, r_1, \ldots, r_n)$. (For reasons unknown)

pressure_shear_matrix

This function finds the matrix which replaces the elastic integral. Recall that once we make the piecewise analytic approximation, we can do all the integrals exactly. The resulting functions are stored in this program.