

EPFL

SEMESTER PROJECT

Minimal Rational Interpolation for Time-Harmonic Maxwell's Equations

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1 ABSTRACT

2 INTRODUCTION

3 FINITE ELEMENT DISCRETIZATION OF THE TIME-HARMONIC MAXWELL'S EQUATIONS

3.1 VECTOR POTENTIAL FORMULATION OF THE TIME-HARMONIC MAXWELL'S EQUATIONS

Let \mathbf{E} denote an electric field, \mathbf{B} a magnetic field strength, ρ an electric charge density, and \mathbf{j} an electric current density. Maxwell's equations are stated in ? as

$$\nabla \cdot (\epsilon \mathbf{E}) = \rho \quad (3.1)$$

$$\nabla \cdot \mathbf{B} = 0 \quad (3.2)$$

$$\nabla \times \mathbf{E} = -\partial_t \mathbf{B} \quad (3.3)$$

$$\nabla \times (\mu^{-1} \mathbf{B}) = \partial_t (\epsilon \mathbf{E}) + \mathbf{j} \quad (3.4)$$

with ϵ being the permittivity and μ the permeability.

Equation (3.2) allows for an expression of the magnetic field $\mathbf{B} = \nabla \times \mathbf{u}$ in terms of a vector valued function \mathbf{u} , the vector potential (in literature commonly denoted with \mathbf{A}). Similarly, (3.3) suggests rewriting the electric field $\mathbf{E} = -\nabla \phi - \partial_t \mathbf{u}$ using a scalar function ϕ , referred to as the scalar potential.

The physical quantities \mathbf{E} and \mathbf{B} remain unchanged if we transform $\mathbf{u} \rightarrow \mathbf{u}' = \mathbf{u} + \nabla \psi$ or $\phi \rightarrow \phi' = \phi - \partial_t \psi$ for arbitrary functions ψ . A convenient choice of ψ is suggested in ? to be

$$\psi = \int_0^t \phi dt' \quad (3.5)$$

which transforms $\phi \rightarrow \phi' = 0$ and $\mathbf{u} \rightarrow \mathbf{u}' = \mathbf{u} + \nabla \int_0^t \phi dt'$. Thus, the expressions for the electrical and magnetic field become

$$\mathbf{E} = -\partial_t \mathbf{u} \quad (3.6)$$

$$\mathbf{B} = \nabla \times \mathbf{u} \quad (3.7)$$

where I renamed the variable \mathbf{u}' to \mathbf{u} for simplicity.

Plugging the identities (3.6) and (3.7) into (3.4) yields

$$\nabla \times (\mu^{-1} \nabla \times \mathbf{u}) = \epsilon \partial_t^2 \mathbf{u} + \mathbf{j} \quad (3.8)$$

For the rest of this report, I restrict myself to vector potentials \mathbf{u} that exhibit a harmonic dependence on time t , i.e. may be factorized into a term solely depending on the position \mathbf{x} and a complex exponential: $\mathbf{u}(\mathbf{x}, t) = \mathbf{u}(\mathbf{x}) \exp(i\omega t)$. Substituting this expression into (3.8) results in the

Time-harmonic potential equation: $\nabla \times (\mu^{-1} \nabla \times \mathbf{u}) - \epsilon \omega^2 \mathbf{u} = \mathbf{j}$

(3.9)

3.2 WEAK FORMULATION FOR THE TIME-HARMONIC POTENTIAL EQUATION

Equation (3.9) may be multiplied by a vector-valued function $\mathbf{v} \in H_{\text{curl}}(\Omega)$, where

$$H_{\text{curl}}(\Omega) = \{\mathbf{u} : \Omega \rightarrow \mathbb{C}, \text{ such that } \mathbf{u} \in L^2(\mathbb{C})^3, \nabla \times \mathbf{u} \in L^2(\mathbb{C})^3\} \quad (3.10)$$

and then integrated over all of Ω to obtain

$$\int_{\Omega} (\nabla \times (\mu^{-1} \nabla \times \mathbf{u})) \cdot \mathbf{v} - \omega^2 \int_{\Omega} \epsilon \mathbf{u} \cdot \mathbf{v} = \int_{\Omega} \mathbf{j} \cdot \mathbf{v} \quad (3.11)$$

This may further be simplified (3.11) to (see Section 1 for details)

$$\boxed{\int_{\Omega} (\mu^{-1} \nabla \times \mathbf{u}) \cdot (\nabla \times \mathbf{v}) - \epsilon \omega^2 \int_{\Omega} \mathbf{u} \cdot \mathbf{v} = \int_{\Omega} \mathbf{j} \cdot \mathbf{v} + \int_{\partial\Omega} ((\mu^{-1} \nabla \times \mathbf{u}) \times \mathbf{n}) \cdot \mathbf{v}} \quad (3.12)$$

where \mathbf{n} denotes the surface normal to the boundary $\partial\Omega$.

3.3 TWO-DIMENSIONAL RESONANT CAVITY

A resonant cavity is a region Ω enclosed by a boundary $\partial\Omega$. The boundary is subdivided into one (or more) inlets Γ_N and a perfect electrically conducting wall $\Gamma_D = \partial\Omega \setminus \Gamma_N$.



FIGURE 3.1 – Example of a two-dimensional resonant cavity: The rectangular cavity.

Suppose the current density $\mathbf{j} \equiv 0$ and orient the coordinate system in such a way that $\mathbf{u} = u_z \mathbf{e}_z$ and $\mathbf{v} = v_z \mathbf{e}_z$. Consequently,

$$(\mu^{-1} \nabla \times \mathbf{u}) \cdot (\nabla \times \mathbf{v}) = (\mu^{-1} \nabla u_z) \cdot (\nabla v_z) \quad (3.13)$$

Define $g_z = ((\mu^{-1} \nabla \times \mathbf{u}) \times \mathbf{n})_z$ along the boundary Γ_N , to convert (3.12) into the weak formulation for a two-dimensional resonant cavity

$$\int_{\Omega} (\mu^{-1} \nabla u_z) \cdot (\nabla v_z) - \omega^2 \int_{\Omega} \epsilon u_z v_z = \int_{\partial\Omega} g_z v_z \quad (3.14)$$

3.4 WAVEGUIDE

3.5 IMPERFECT CONDUCTOR

4 MINIMAL RATIONAL INTERPOLATION FOR THE TIME-HARMONIC MAXWELL'S EQUATIONS

5 EXAMPLES

5.1 TWO-DIMENSIONAL RECTANGULAR CAVITY

5.2 DUAL MODE CIRCULAR WAVEGUIDE FILTER



FIGURE 5.1 – Dual-mode circular waveguide filter.

5.3 IMPERFECTLY CONDUCTING BOUNDARIES

6 CONCLUSION AND OUTLOOK

7 APPENDIX

7.1 DETAILED DERIVATION FOR THE WEAK FORMULATION OF THE TIME-HARMONIC POTENTIAL EQUATION

The goal is to rewrite the curl-integral on the left-hand side of (3.11):

$$\int_{\Omega} (\nabla \times (\mu^{-1} \nabla \times \mathbf{u})) \cdot \mathbf{v} \quad (7.1)$$

In order to simplify the curls and apply the Gauss theorem, I first show prove the vector calculus identity

$$(\nabla \times \mathbf{a}) \cdot \mathbf{b} = \nabla \cdot (\mathbf{a} \times \mathbf{b}) + \mathbf{a} \cdot (\nabla \times \mathbf{b}) \quad (7.2)$$

where \mathbf{a}, \mathbf{b} are vector-value functions. The completely antisymmetric tensor ε_{ijk} , frequently referred to as the Levi-Civita tensor, may be employed to rewrite the components of the curl of a vector-function \mathbf{a} as the sum

$$(\nabla \times \mathbf{a})_k = \sum_i \sum_j \varepsilon_{ijk} \partial_i u_j \quad (7.3)$$

where ∂_i denotes the partial derivative with respect to the i -th coordinate direction. This yields

$$\begin{aligned} (\nabla \times \mathbf{a}) \cdot \mathbf{b} &= \sum_k (\nabla \times \mathbf{a})_k b_k \\ &= \sum_k \left(\sum_i \sum_j \varepsilon_{ijk} \partial_i a_j \right) b_k \\ &= \sum_k \sum_i \sum_j \partial_i (\varepsilon_{ijk} a_j b_k) - \sum_k \sum_i \sum_j a_j (\varepsilon_{ijk} \partial_i b_k) \\ &= \sum_k \sum_i \sum_j \partial_i (\varepsilon_{jki} a_j b_k) - \sum_k \sum_i \sum_j a_j ((-\varepsilon_{ikj}) \partial_i b_k) \\ &= \sum_i \partial_i (\mathbf{a} \times \mathbf{b})_i + \sum_j u_j (\nabla \times \mathbf{b})_j \\ &= \nabla \cdot (\mathbf{a} \times \mathbf{b}) + \mathbf{a} \cdot (\nabla \times \mathbf{b}) \end{aligned} \quad (7.4)$$

by expressing the scalar product as a component-sum, using the product rule and applying the symmetry and anti-symmetry properties of the Levi-Civita tensor. Now the identity (7.2) to (7.1) together with Gauss' theorem gives

$$\begin{aligned} \int_{\Omega} (\nabla \times (\mu^{-1} \nabla \times \mathbf{u})) \cdot \mathbf{v} &= \int_{\Omega} \nabla \cdot ((\mu^{-1} \nabla \times \mathbf{u}) \times \mathbf{v}) + \int_{\Omega} (\mu^{-1} \nabla \times \mathbf{u}) \cdot (\nabla \times \mathbf{v}) \\ &= \int_{\partial\Omega} ((\mu^{-1} \nabla \times \mathbf{u}) \times \mathbf{v}) \cdot \mathbf{n} + \int_{\Omega} (\mu^{-1} \nabla \times \mathbf{u}) \cdot (\nabla \times \mathbf{v}) \end{aligned} \quad (7.5)$$

For later convenience, the boundary integral can further be simplified using

$$(\mathbf{a} \times \mathbf{b}) \cdot \mathbf{c} = -(\mathbf{a} \times \mathbf{c}) \cdot \mathbf{b} \quad (7.6)$$

This identity follows immediately from a small manipulation with the Levi-Civita tensor:

$$\begin{aligned} (\mathbf{a} \times \mathbf{b}) \cdot \mathbf{c} &= \sum_k \left(\sum_i \sum_j \varepsilon_{ijk} a_i b_j \right) c_k \\ &= \sum_j \left(\sum_i \sum_k (-\varepsilon_{ikj}) a_i c_k \right) b_j \\ &= -(\mathbf{a} \times \mathbf{c}) \cdot \mathbf{b} \end{aligned} \quad (7.7)$$

The boundary integral becomes

$$\int_{\partial\Omega} ((\mu^{-1} \nabla \times \mathbf{u}) \times \mathbf{v}) \cdot \mathbf{n} = - \int_{\partial\Omega} ((\mu^{-1} \nabla \times \mathbf{u}) \times \mathbf{n}) \cdot \mathbf{v} \quad (7.8)$$

This concludes the short derivation, because now (7.1) may be rewritten as

$$- \int_{\partial\Omega} ((\mu^{-1} \nabla \times \mathbf{u}) \times \mathbf{v}) \cdot \mathbf{n} + \int_{\Omega} (\mu^{-1} \nabla \times \mathbf{u}) \cdot (\nabla \times \mathbf{v}) \quad (7.9)$$