

# Aerodynamics?

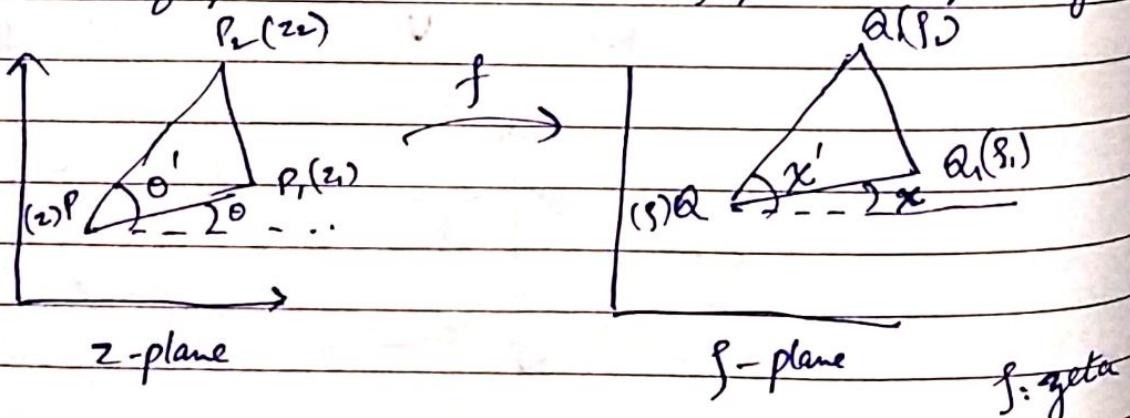
## 1. Aerofoils & Kutta-Joukowski Theorem :-

**Conformal mapping :** A mapping that preserves angle in both magnitude & orientation.

$$w = \phi + i\psi = f(z)$$

$\zeta = \xi + i\eta$   
 $g(s) = \alpha + i\beta$   
 simpler form  
 conformal mapping

Let  $f(z)$  be a single valued diff. complex function within a closed contour  $C$  in  $z$ -plane. Let  $\zeta = \xi + i\eta$  be another complex variable such that  $\zeta = f(z)$ . Then corresponding function to each point in  $z$ -plane or in  $C$ , there will be a point in  $\zeta$ -plane or in  $C'$ . The necessary condition for such map to exist is  $f'(z) \neq 0$  at any pt. in  $z$ -plane or in  $C$ . This also means that  $\frac{d\zeta}{dz}$  must exist independently of the direction of  $dz$ . Let  $P, P_1, P_2$  &  $Q, Q_1, Q_2$  be the '2' sets of points in  $z$ -plane &  $\zeta$ -plane, respectively



$$\text{Now, } \frac{\zeta_1 - \zeta}{z_1 - z} = \frac{f(z_1) - f(z)}{z_1 - z}$$

$$\frac{\zeta_2 - \zeta}{z_2 - z} = \frac{f(z_2) - f(z)}{z_2 - z}$$

When  $P_1 \rightarrow P$ ,  $P_L \rightarrow P$

$$\frac{f_1 - f}{z_1 - z} = f'(z) ; \quad \frac{f_L - f}{z_L - z} = f'(z)$$

$$\Rightarrow \frac{f_1 - f}{z_1 - z} = \frac{f_L - f}{z_L - z} = f'(z) + \frac{df}{dz}$$
①

Next,

$$\frac{QQ^* e^{ix}}{PP_1 e^{i\theta}} = \frac{QQ_L e^{ix'}}{PP_L e^{i\theta'}}$$

$$\Rightarrow \frac{QQ_1}{PP_1} e^{i(x-\theta)} = \frac{QQ_L}{PP_L} e^{i(x'-\theta')}$$

When  $P_1 \rightarrow P$  &  $P_L \rightarrow P$

$$\Rightarrow e^{i(x-\theta)} = e^{i(x'-\theta')}$$

$$\Rightarrow x - \theta = x' - \theta' \Rightarrow x' - x = \theta' - \theta$$

i.e.  $\angle Q, QQ_1 = \angle P, PP_L$

$$\text{Also, } \frac{QQ_1}{PP_1} = \frac{QQ_L}{PP_L} = |f'(z)| = \frac{df}{dz}$$

$\Rightarrow \Delta P, PP_L$  &  $\Delta Q, QQ_1$  are similar

This establishes the similarity of the corresponding infinitesimal element of the 2 planes. Such a relation between the 2 planes is called conformal representation of either one of the other.

$$\begin{aligned}f &= f(z), \quad z = x+iy \\&= \xi + i\eta\end{aligned}$$

$$\begin{aligned}\frac{df}{dz} &= f'(z) = \frac{d}{dz}(\xi + i\eta) \\&= \frac{\partial \xi}{\partial x}(x+iy) \frac{dx}{dz} \\&= \frac{\partial \xi}{\partial x} + i \frac{\partial \eta}{\partial x} \quad \text{--- (1)} \\ \text{Also, } \frac{df}{dz} &= \frac{d}{dy}(\xi + i\eta) \frac{dy}{dz} \\&= \frac{\partial \xi}{\partial y} + i \frac{\partial \eta}{\partial y} \quad \text{--- (2)}\end{aligned}$$

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Conformal Mapping :-

$$f(z) = f(x+iy) = \xi + i\eta$$

From complex analysis :-

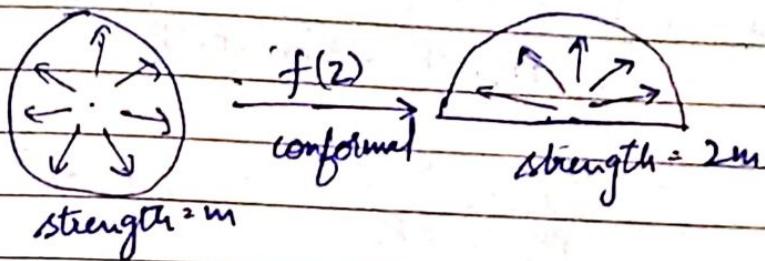
$$\frac{df}{dz} = \xi_x + i\eta_x$$

From CR-equation :-

$$\frac{\partial \xi}{\partial x} = \frac{\partial \eta}{\partial y}, \quad \frac{\partial \xi}{\partial y} = -\frac{\partial \eta}{\partial x}$$

$$\frac{df}{dz} = \frac{\partial \eta}{\partial y} + i \cancel{\frac{\partial \eta}{\partial x}} = i \left( \frac{\partial \eta}{\partial x} - i \frac{\partial \eta}{\partial y} \right)$$

$$\rightarrow f'(z) = \sqrt{\left( \frac{\partial \eta}{\partial x} \right)^2 + \left( \frac{\partial \eta}{\partial y} \right)^2}$$



$\omega = F_1(z) = \phi + i\psi$ , where  $\phi$  &  $\psi$  are velocity & current functions of any motion within the contour  $C'$  in  $\{z\}$ -plane.

Then within  $C'$ ,

$$\phi + i\psi = F_1(\xi + i\eta) \Rightarrow f_1(\xi, \eta) + i\bar{f}_1(\xi, \eta)$$

&amp. c' given by,

$$\psi = \bar{f}_1(\xi, \eta)$$

Similarly,

$$\omega = F_1(z)$$

$$= F_1(\xi + i\eta) = F_1(\xi(x, y) + i\eta(x, y))$$

$$\Rightarrow \phi + i\psi = F_2(x, y) = f_2(x, y) + i\bar{f}_2(x, y)$$

In this case the contour is :-

$$\psi(x, y) = \bar{f}_2(x, y) = \text{constant}$$

This shows that  $\psi$  &  $\phi$  are same in both the planes.

$$q_1 = \frac{dw}{dz}, \quad q_2 = \frac{dw}{ds}$$

$$|q_1|^2 = \left| \frac{dw}{dz} \right|^2, \quad |q_2|^2 = \left| \frac{dw}{ds} \right|^2$$

$$\Rightarrow |q_1|^2 = |q_2|^2$$

$$\left| \frac{dw}{dz} \right|^2 = \left| \frac{ds}{dz} \right|^2 \quad \text{--- (1)}$$

$$\frac{\Delta Q_1, Q_2 Q_2}{\Delta P_1, P_2 P_2} = \frac{\frac{1}{2} Q_1 Q_2 \cdot Q_2 Q_2 \cdot \sin \angle Q_1 Q_2 Q_2}{\frac{1}{2} P_1 P_2 \cdot P_2 P_2 \cdot \sin \angle P_1 P_2 P_2}$$

$$= \frac{Q_1 Q_2 \cdot Q_2 Q_2}{P_1 P_2 \cdot P_2 P_2} = |f'(z)|^2$$

$$= \left( \frac{dx}{dz} \right)^2 + \left( \frac{dy}{dz} \right)^2$$

$$=: h^2$$

$$\Rightarrow \Delta Q_1, Q_2 Q_2 = h^2 \Delta P_1, P_2 P_2$$

Area of the  $\Delta^{de}$  in  $\xi$  &  $Z$  planes are of ratio  $h^2$  to 1  
 $\Rightarrow d\xi d\eta = h^2 dx dy$

$$\text{From } ①, |q_2|^2 = \left| \frac{d\omega}{d\xi} \right|^2 = \left| \frac{d\omega}{dz} \right|^2 \left| \frac{dz}{d\xi} \right|^2$$

$$= \frac{|q_1|^2}{h^2}$$

$$\Rightarrow |q_2|^2 d\xi d\eta = \frac{|q_1|^2}{h^2} \cdot h^2 dx dy$$

$$\Rightarrow |q_2|^2 d\xi d\eta = |q_1|^2 dx dy$$

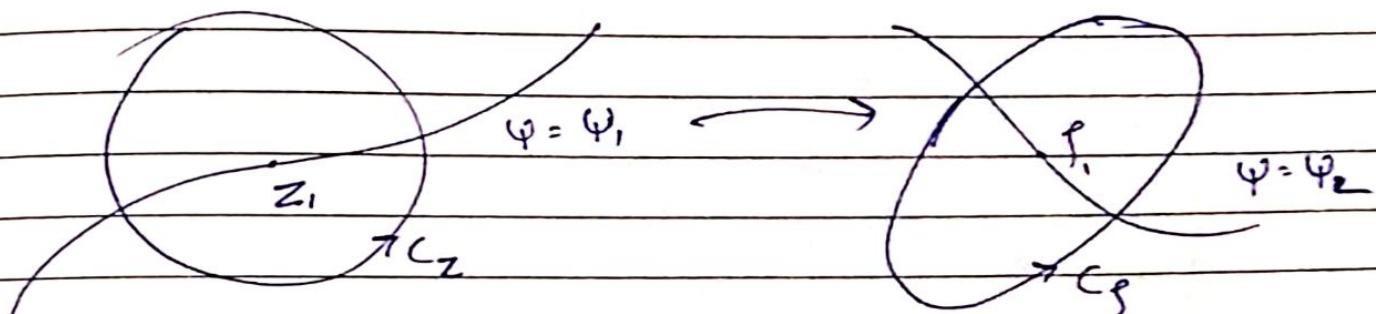
$$\Rightarrow \frac{1}{L} \int p |q_2|^2 d\xi d\eta = \frac{1}{L} \int p |q_1|^2 dx dy$$

$$\Rightarrow T_\xi = T_Z$$

K.E is same in both planes.

## \* Transformation of a source :-

S-plane



Let there be a source of strength  $m$  at  $z_1$ , &  $s_1$  be the corresponding pt in S-plane. Let these be regular points of transformation. Then a small curve  $C_{z_1}$  may be drawn to exclude  $z_1$  & similarly for  $C_{s_1}$  to exclude  $s_1$ . Since we know the value of stream function is independent of the domain considered, we have

$$\int_{C_{z_1}} d\psi = \int_{s_1} d\psi$$

$$\text{But. } \int_{C_{z_1}} d\psi = \int_{C_{z_1}} \left( \frac{\partial \psi}{\partial x} dx + \frac{\partial \psi}{\partial y} dy \right) = \int_{C_{z_1}} (-v dx + u dy)$$

= the total flow across the contour  $C_{z_1}$   
= sum of sources " of "  $C_{z_1}$ .

4.

representation

Theorem: Under conformal definition a uniform line source maps into another uniform line source of the same strength.

Proof:-

Let there be a uniform line source of strength 'm' per unit length through the pt.  $z = z_0$  & suppose the conformal transformation  $\xi = f(z)$ . Let  $C_{z_0}$  be the curve around  $z_0$  in  $z$ -plane &  $C_{z_0}$  is mapped into  $C_\xi$  in  $\xi$ -plane. Then  $\xi = \xi_0$  lies in  $C_\xi$ .

The complex potential in  $z$ -plane &  $\xi$ -plane is same and has the form

$$\omega = \phi + i\psi \text{ in } z\text{-plane}$$

$$\omega = \phi' + i\psi, \text{ in } \xi\text{-plane}$$

$$\Rightarrow \phi = \phi' \quad \psi = \psi'$$

Since ' $\psi$ ' is the same at the corresponding pts.  $z = z_0$  &  $\xi = \xi_0$  we have  $\int_{C_{z_0}} d\psi = \int_{C_\xi} d\psi' \quad \text{--- (1)}$

But in  $z$ -plane,

$$\omega = -m \log(z - z_0)$$

$$\Rightarrow d\omega = -m dz / (z - z_0)$$

$$\Rightarrow \int_{C_{z_0}} d\omega = - \int_{C_{z_0}} m dz / (z - z_0)$$

$$\Rightarrow \int_{C_{z_0}} (d\phi + d\psi) = -m \cdot 2\pi i \quad \text{by Cauchy's residue theorem}$$

$$\Rightarrow \int_{C_{z_0}} d\psi = -2\pi m \quad \text{--- (2)}$$

(2) indicates the value of fluid crossing unit thickness of  $C_{z_0}$  per unit time

$$\rightarrow \int_{C_S} d\psi' = -2\pi m$$

← indicates

$\Rightarrow$  the volume of the fluid crossing unit thickness of  $C_S$  per unit time

$\therefore$  The conformal mapping preserves the strength of the simple source.

### 5. Kutta - Joukowski theorem :-

When a cylinder of any shape is placed in a uniform stream of speed  $V$ , the resultant ~~thrust~~ thrust on the cylinder is a lift of magnitude  $KDU$  per unit length  $d$  at right angles to the stream.  $K$  is the circulation around cylinder.

Proof :- let there be a fixed cylinder of some form in the finite region of the plane, its cross-section containing origin. The disturbance of the stream caused by the cylinder can be represented at a great distance in the form :

$$\omega = Az + B/z + \dots, \text{ where } A, B \text{ depend on } U \text{ & } K.$$

Let the direction of the stream makes an angle  $\alpha$  with  $x$ -axis. Then the complex potential due to uniform stream velocity  $U$  is

$$\omega_1 = Ue^{-i\alpha z} - \textcircled{2}$$

And due to circulation,

$$\omega_3 = \frac{ik}{2\pi} \log z - \textcircled{3}$$

∴ The complete complex potential is

$$\begin{aligned} \omega &= \omega_1 + \omega_L + \omega_3 \\ &= U e^{-i\alpha z} + \frac{iR}{2\pi} \log z + AZ + \frac{B}{z} + \dots \end{aligned}$$

By Blasius Theorem :-

- (4)

The force exerted  
on the cylinder is

$$x - iy = \frac{1}{2} i P \oint_C \left( \frac{d\omega}{dz} \right)^L$$

$$= \frac{1}{2} i P \oint_C \left[ -i\alpha U e^{-i\alpha z} + \frac{iR}{2\pi z} + A - \frac{B}{z^2} - \frac{C}{z^3} \right] dz$$

$$= \frac{1}{2} i P \oint_C \left[ i\alpha U e^{-i\alpha z} dz + \cancel{\int_{C'} \left( \frac{iR}{2\pi z} + A - \frac{B}{z^2} - \frac{C}{z^3} \dots \right) dz} \right]$$

$$= \frac{1}{2} i P \oint_C \left[ (-i\alpha U e^{-i\alpha z})^L dz + \left( \frac{iR}{2\pi z} + A - \frac{B}{z^2} - \frac{C}{z^3} \dots \right)^L dz + 2(-i\alpha U e^{-i\alpha z}) \left( \frac{iR}{2\pi z} + A - \frac{B}{z^2} \dots \right) dz \right]$$

$$\boxed{\frac{dw}{dz} = U e^{-i\alpha z} + \frac{iR}{2\pi z} - \dots}$$

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$$x - iy = \frac{1}{2} i P \cdot 2\pi i \quad (\text{Sum of the residues at } z=0)$$

Since  $z=0$  is the pole of  $\frac{dw}{dz}$  inside  $C$ , therefore  
the residue at  $z=0 = \frac{iR}{\pi} U e^{-i\alpha}$

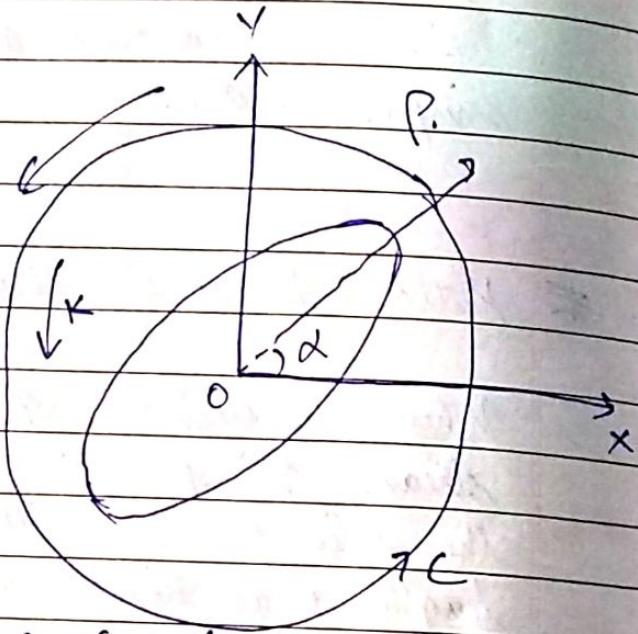
$$x - iy = \frac{1}{2} i P \cdot 2\pi i \cdot \frac{iR}{\pi} U e^{-i\alpha}$$

$$= -i P K U e^{-i\alpha}$$

$$\Rightarrow x - iy = -i P K U (\cos\alpha - i\sin\alpha)$$

$$x = -K U p \sin\alpha$$

$$y = P K U \cos\alpha$$



$$\text{Resultant lift} = \sqrt{x^2 + y^2}$$

$$= \sqrt{(\rho K U)^2 (\cos^2 \alpha + \sin^2 \alpha)} \\ = \rho K U$$

which always acts right angles to cylinder.

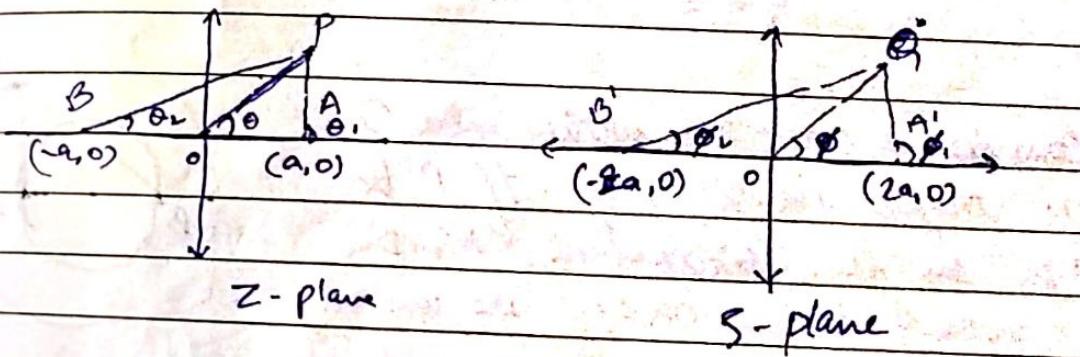
Joukowskii Transformation, Joukowskii aerofoils & Hypothesis

Consider the transformation, :-

$$\zeta = z + a^2 z^{-1} \quad \dots \quad (1)$$

Now, let A & B be the points at  $z = -a, a$ . These points are getting mapped to  $A' \& B'$  resp. i.e.

$$\zeta = 2a \quad \& \quad \zeta = -2a$$



$$\zeta - 2a = z - 2a + a^2 z^{-1}$$

$$= \frac{(z-a)^2}{z} \quad \dots \quad (2)$$

Similarly,

$$\zeta + 2a = \frac{(z+a)^2}{z} \quad \dots \quad (3)$$

So from (3) :-

$$A'Q e^{i\phi} = (AP e^{i\theta_1})^2$$

$$OP e^{i\theta_0}$$

$$\Rightarrow A'Q = \frac{AP^2}{OP} \quad \& \quad e^{i\theta_1} = e^{i(2\theta_1 - \theta_0)} \Rightarrow \theta_1 = 2\theta_1 - \theta_0$$

- (4)

Again from (2) :-

$$B'Q = BP^2/OP \quad \& \quad \phi_2 = \theta_2 - \theta \quad -(3)$$

$$\therefore A'QB' = \phi_1 - \phi_2 = 2(\theta_1 - \theta_2) = 2\angle APB \quad -(4)$$

From (4), (5) :-

$$A'Q + B'Q = \frac{AP^2 + BP^2}{OP} = \frac{2(OP^2 + OA^2)}{OP}$$

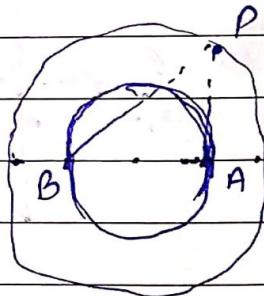
Since  $OA$  is  
median of  $\triangle APB$ ,

$$\text{i.e. } AP^2 + BP^2 = 2(OP^2 + OA^2)$$

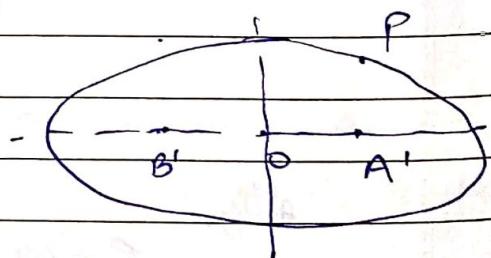
When  $|z|$  is very large, then

$$f = \bar{z}$$

Consider a circle  $C$  with centre  $O$   
at origin in  $z$ -plane. If  $P$  be  
a pt on this circle then,  $A'Q + B'Q$   
= constant because  $OA$  &  $OP$  are constant.



from (6). It follows that  $Q$   
will describe an ellipse with  
foci  $A'$  &  $B'$ . Hence Joukowski  
transformation (i) transforms a  
circle in  $z$ -plane to an ellipse  
in  $f$ -plane. Also the angle is getting doubled via the  
transformation. i.e if  $P$  be a point on the smaller circle  
with  $\angle APB = \pi/2$  then  $\angle A'QB' = \pi$ . i.e.  $Q$  lies on  $A'B'$ .



Let  $P$  be the pt. "z" on the bigger circle &  $P'$  be its inverse point w.r.t. circle  $AB$  &  $P''$  be the reflection of  $P'$  on  $x$ -axis, then  $P''$  be the pt.  $a^2 z$ . Then we draw a parallelogram with  $OP, OP'$

as adjacent sides &  $OQ$  be the diagonal st.  $Q$  be the pt.  $Q = z + a^2 z$ . The locus of the point  $Q$  is a fish shaped contour which touches the line  $BA$  on both sides & such ~~contour~~ is known as Joukowski's aerofoils where  $B$  is the trailing edge &  $R$  is the leading edge.

Also,

$$\frac{ds}{dz} = 1 - a^2 z^2$$

$$\text{So, } \frac{ds}{dz} = 0 \Rightarrow z = \pm a$$

And  $z = a$  is mapped  $s = za$  &  
 $s = -za$

But we need to know

$$\frac{ds}{dz} \text{ at } B \text{ so either } \frac{ds}{dz} = 0 \text{ or } \frac{ds}{dz} = \infty$$

If  $\vec{q}$  be the velocity at  $B$  for the circle &  $\vec{q}'$  be the velocity at  $B'$  of the aerofoil, then

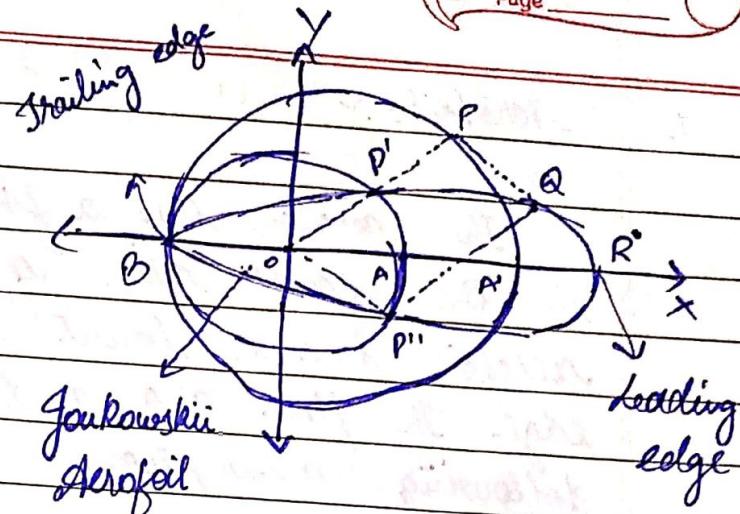
$$|\vec{q}'| = \left| \frac{dw/ds}{dz} \right| = \left| \frac{dw/dz}{1 - a^2 z^2} \right| \cdot \left| \frac{dz}{ds} \right| = z \left| \frac{dz}{ds} \right|$$

$$\text{If } \frac{ds}{dz} = 0 \text{ then } \frac{dz}{ds} = \infty \Rightarrow \vec{q}' = \infty \text{ at } B.$$

To avoid infinite velocity at trailing edge.  $\frac{ds}{dz} = 0$  at trailing edge  $B$ .  $\rightarrow$  not suitable

The velocity at  $B$  is taken as 0, i.e. taken as stagnation pt.

of flow in  $z$ -plane. This is Joukowski Hypothesis.



## 7. Aerofoils :-

The aerofoil has a fish type profile. It is employed in the construction of modern airplanes. Such an aerofoil has a blunt leading edge & a sharp trailing edge. The flow around the aerofoil depends on the following assumptions.

- (i) The air behaves as an incompressible fluid
- (ii) The aerofoil is a cylinder whose cross-section is a curve of fish type.
- (iii) The flow is 2D irrotational cyclic motion

## Flow past a circle :-

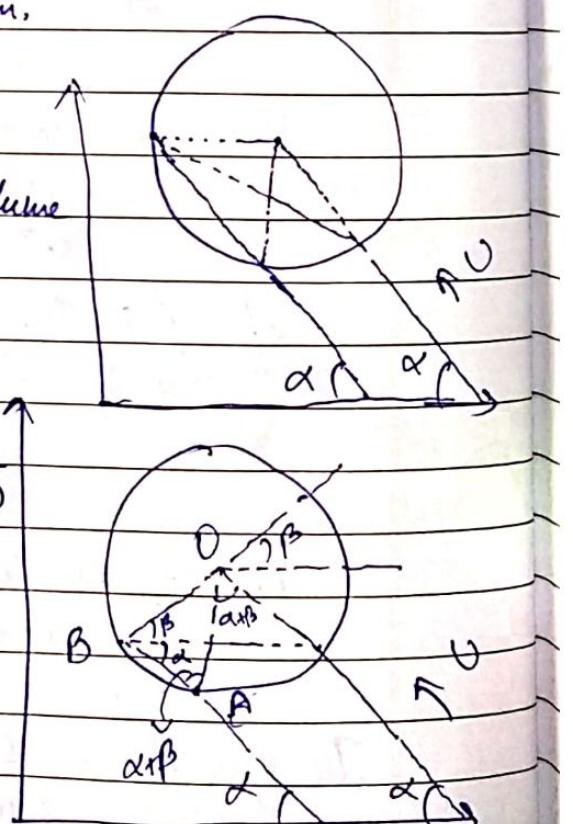
Let  $U$  be the velocity of the stream at infinity. Its direction making an angle  $\alpha$  with the  $-ve$   ~~$x$ -axis~~,  $K$  be the circulation around the ~~cylinder~~ circle whose center is at  $Z_0$  & radius  $b$ . Then,

$$\omega = U e^{i\alpha} (z - z_0) + \frac{Ub^2 e^{-i\alpha}}{z - z_0} + \frac{iK}{2\pi} \log(z - z_0), \quad b \text{ is the volume}$$

To find the stagnation point

$$\frac{d\omega}{dz} = 0 \Rightarrow U e^{i\alpha}$$

$$-\frac{Ub^2 e^{-i\alpha}}{(z - z_0)^2} + \frac{iK}{2\pi(z - z_0)} = 0 \quad \text{--- (1)}$$



Taking the stagnation pt. as  $z = z_0 + be^{i(\alpha+\beta)}$   
 $= z_0 - be^{i\beta}$ .

Then, ① reduces to,

$$U(e^{i\alpha} - e^{-i(\alpha+\beta)}) - \frac{ik}{2\pi b} e^{-i\beta} = 0$$

$$\Rightarrow K = 4\pi b U \sin(\alpha+\beta)$$

$$\Rightarrow K \leq 4\pi b U \text{ since } \sin(\alpha+\beta) < 1 \quad -(2)$$

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\* Flow around a circle :-

$$\omega = Ue^{i\alpha}(z-z_0) + Ub^2 e^{-i\alpha} / z-z_0 + iR/2\pi \log(z-z_0)$$

$$\Rightarrow \frac{d\omega}{dz} = Ue^{i\alpha} - Ub^2 e^{-i\alpha} / (z-z_0)^2 + iR/2\pi \log(z-z_0)$$

For stagnation points,  $\frac{d\omega}{dz} = 0$ , let us take the stagnation point as :-

$$z = z_0 + be^{i(\alpha+\beta)} = z_0 - be^{i\beta}$$

Eq. 1 reduces to,

$$U[e^{i\alpha} - e^{-i(\alpha+2\beta)}] - iR/2\pi b e^{-i\beta} = 0$$

$$\Rightarrow 2\pi b U [e^{i(\alpha+\beta)} - e^{-i(\alpha+\beta)}] = iR$$

$$\Rightarrow 2\pi b U 2i \sin(\alpha+\beta) = iR$$

$$\Rightarrow R = 4\pi b U \sin(\alpha+\beta)$$

$$\Rightarrow R < 4\pi b U, \text{ since } b, U > 0$$

The stagnation points are given by  $z - z_0 = be^{i(\alpha+\beta)} = t$

From ①,

$$Ue^{i\alpha} - \frac{1}{t^2} Ub^2 e^{-i\alpha} + \frac{iR}{2\pi t} = 0$$

$$\Rightarrow t = be^{-i\alpha} [-i \sin(\alpha+\beta) \pm \cos(\alpha+\beta)], \text{ putting } R = 4\pi b U$$

$$= be^{-i\alpha} e^{i(\alpha+\beta)} \quad \text{or, } -be^{i\beta}$$

$$= be^{-i(2\alpha+\beta)} \quad \text{or, } -be^{i\beta}$$

$$\therefore \text{the points } t_1 = be^{-i(2\alpha+\beta)}$$

$$= be^{-i[2\alpha-(2\alpha+\beta)]}$$

$$\& t_2 = -be^{i\beta} = be^{i(\pi+\beta)}$$

are points B & A, respectively where

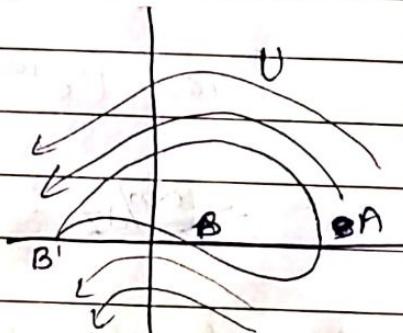
$\frac{dw}{dz} = 0$ , i.e., A & B are the stagnation pts.

~~Following~~

### 9. Flow past an aerofoil :-

From Joukowski transformation,

①  $\zeta = z + a\gamma z$  transforms a circle into an aerofoil. Thus the stagnation point P will be transformed into trailing edge B' and the other stagnation point A will be transformed to A''. B' will lie on the -ve side of  $\zeta$ -axis s.t.  $OB' = -2a$ . Let  $U$  be the velocity of the stream at infinity. its direction making an angle  $\alpha$  with  $\zeta$ -axis. We take the origins of both  $z$  &  $\zeta$  planes to coincide then B & B' will lie on the left side of origin O.



By Blasius Theorem,

$$x - iy = \gamma_2 i P_C \int_{\zeta} \left( \frac{dw}{ds} \right)^2 ds - (2)$$

From ① :-

$$\frac{dw}{ds} = \frac{dw/dz}{ds/dz} = \frac{dw/dz}{1-a\gamma_z}$$

$$\left( \frac{dw}{ds} \right)^2 = \frac{dw}{ds} \cdot \frac{dw}{ds} = \frac{(dw/dz)^2}{(ds/dz)^2} \quad -(3)$$

Cauchy Residual  
Theorem  
Exam point of view

discuss

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Now,

$$w = U e^{i\alpha} (z - z_0) + U b^L e^{-i\alpha} / z + iR / 2\pi \log(z - z_0)$$

$$\frac{dw}{dz} = U e^{i\alpha} - U b^L e^{-i\alpha} / (z - z_0) + iR / 2\pi (z - z_0)$$

$$\begin{aligned} \frac{1}{(1 - a^L/z^L)} \cdot \left( \frac{dw}{dz} \right) &= (1 - a^L z^L)^{-1} \left[ U e^{i\alpha} - \frac{U b^L e^{-i\alpha}}{z^L} (1 - z^L z)^{-1} + \frac{iR}{2\pi z} (1 - z^L z)^{-1} \right] \\ &= (1 + 2a^L z^L + \dots) \left[ U e^{-i\alpha} - \frac{U b^L e^{i\alpha}}{z^L} (1 + 2z^L z + \dots) + \frac{iR}{2\pi z} (1 + 2a^L z^L + \dots) \right] \end{aligned}$$

- (4)

By, (2)

$$x - iy = \frac{iP}{2} \int_{C_2} \frac{\left( \frac{dw}{dz} \right)^L}{(1 - a^L z^L)^L} dz$$

$$= \frac{iP}{2} \cdot 2\pi i \left[ \text{sum of the residues of } \frac{\left( \frac{dw}{dz} \right)^L}{(1 - a^L z^L)^L} \right]$$

$$= -\pi P \text{ coeff. of } \frac{1}{z} \text{ in RHS of (4)}$$

$$= -\pi P \cdot 2U e^{i\alpha} iK / 2\pi = -iPK e^{i\alpha} U$$

$$x = PKU \sin \alpha \quad Y = PKU \cos \alpha \Rightarrow F = PKU$$

Lift acting on the aerofoil =  $\rho K U = 4\pi b U^2 \sin(\alpha + \beta)$

### 10. Moment of the force acting on aerofoil

$$\left(\frac{dw}{ds}\right)^2 \times s ds = \frac{z + a^L z}{1 - a^L z} \times \left(\frac{dw}{dz}\right)^2 dz$$

$$\frac{z + a^L z}{1 - a^L z} \left(\frac{dw}{dz}\right)^2 = (z + a^L z) \left(1 + a^L z + \dots\right) \int \frac{U e^{iz} - U b^L e^{-iz}}{z^2} dz$$

$$= \left(z + \frac{2a^L}{z} + \dots\right) (\dots)$$

$$\Rightarrow \text{coeff. of } \frac{1}{z} = -2U^2 b^L \frac{-K^L}{4\pi z} + \frac{iU e^{iz}}{\pi} z_0 + 2a^L U^L e^{2iz}$$

$$M = \text{Real part of } \left\{ -\frac{P}{2} \times \text{coeff. of } \frac{1}{z} \times 2\pi i \right\}$$

$$= \text{Re} \left\{ -\frac{P}{2} \left[ -2b^L U^L - \frac{K^L}{4\pi z} + \frac{iU e^{iz} z_0}{\pi} + 2a^L U^L e^{2iz} \right] \right\}$$

$$\text{Taking } z_0 = c e^{i\alpha}, c > 0 \text{ then, } M = \text{Re} \left\{ -\frac{P}{2} \left[ -2b^L U^L - \frac{K^L}{4\pi z} + \frac{iU e^{iz} c e^{i\alpha}}{\pi} + 2a^L U^L e^{2iz} \right] \right\}$$

$$M = 2\pi P U^L [ 2bc \sin(\alpha + \beta) \cos(\alpha + \alpha) + a^L \sin 2\alpha ]$$

This is the moment at  $z_0$ . The moment at the trailing edge is  $aY + M$

Required moment at B

$$= \alpha PKU \cos \alpha + 2\pi PV^2 [2bc \sin(\alpha+\beta) \cos(\alpha+\gamma) + a^2 \sin 2\alpha]$$

$$= 4a\pi bV^2 P \cos \alpha + 2\pi PV^2 [2bc \sin(\alpha+\beta) \cos(\alpha+\gamma) + a^2 \sin 2\alpha]$$

### Navier-Stokes Equation.

With  $P(x, y, z)$  as center and edges of length  $S_x, S_y, S_z$  parallel to coordinate axes, we construct an elementary rectangular parallelopiped.

Let us consider the fluid motion is viscous. We assume that the fluid element is moving & its mass is  $\rho S_x S_y S_z$ . Let the coordinate

of  $P_2$  &  $P_1$  be  $(x + \frac{S_x}{2}, y, z)$  &  $(x - \frac{S_x}{2}, y, z)$  respectively. At P the force components // to  $Ox, Oy, Oz$  on the rectangular surface ABCD of area  $S_y S_z$  through P having  $\hat{n}$  as the unit normal are  $= (-\sigma_{xx} S_y S_z, -\sigma_{yy} S_y S_z, -\sigma_{zz} S_y S_z)$  — ①

At the point  $P_2 (x + S_x/2, y, z)$  the components of the force are  $A_2 B_2 C_2 D_2$  is

$$= \left[ \left( -\sigma_{xx} + S_y/2 \frac{\partial \sigma_{yy}}{\partial x} \right) S_y S_z, \left( -\sigma_{yy} + S_y/2 \frac{\partial \sigma_{yy}}{\partial x} \right) S_y S_z, \right. \\ \left. \left( -\sigma_{zz} + S_y/2 \frac{\partial \sigma_{zz}}{\partial x} \right) S_y S_z \right] \quad \text{— ②}$$

Only force on  $P_1$  in A, B, C, D,

$$= \left[ \mathbf{F} - \left( -\sigma_{yy} - S_y/2 \frac{\partial \sigma_{yy}}{\partial x} \right) S_y S_z, - \left( \sigma_{yy} - S_y/2 \frac{\partial \sigma_{yy}}{\partial x} \right) S_y S_z, \right. \\ \left. - \left( \sigma_{zz} - S_y/2 \frac{\partial \sigma_{zz}}{\partial x} \right) S_y S_z \right] \quad \text{— ③}$$

