

The metric (??) in real variables reads

$$\begin{aligned} ds^2 &= (v_t + u_{xx})(dt^2 + dx^2) + \Delta(u)(dy^2 + dz^2) \\ &\quad - 2b(dt\,dz - dx\,dy) + 2c(dt\,dy + dx\,dz). \end{aligned} \quad (2.5)$$

The coframe of one-forms becomes

$$\begin{aligned} \Omega_1 &= \frac{1}{2\sqrt{v_t + u_{xx}}}[(c + ib)(dy + idz) + (v_t + u_{xx})(dt + idx)] \\ \Omega_2 &= \bar{\Omega}_1, \quad \Omega_3 = \frac{dy + idz}{2\sqrt{v_t + u_{xx}}}, \quad \Omega_4 = \bar{\Omega}_3 \end{aligned} \quad (2.6)$$

with the metric

$$ds^2 = \Omega_1 \otimes \bar{\Omega}_1 + \varepsilon \Omega_2 \otimes \bar{\Omega}_2. \quad (2.7)$$

### 3 Bi-Hamiltonian representations of CMA system

The *CMA* system (??) can be put in the Hamiltonian form

$$\begin{pmatrix} u_t \\ v_t \end{pmatrix} = J_0 \begin{pmatrix} \delta_u H_1 \\ \delta_v H_1 \end{pmatrix} \quad (3.1)$$

where  $\delta_u$  and  $\delta_v$  are Euler-Lagrange operators [?] with respect to  $u$  and  $v$ . Here  $J_0$  is the Hamiltonian operator

$$J_0 = \begin{pmatrix} 0 & \frac{1}{a} \\ -\frac{1}{a} & \frac{1}{a}(cD_y + D_y c - bD_z - D_z b)\frac{1}{a} \end{pmatrix} \quad (3.2)$$

determining the structure of Poisson bracket and  $H_1$  is the corresponding Hamiltonian density

$$H_1 = \frac{1}{2}[v^2\Delta(u) - u_{xx}(u_y^2 + u_z^2)] - \varepsilon u. \quad (3.3)$$

The first real recursion operator has the form

$$\begin{aligned} R_1 &= \begin{pmatrix} 0 & 0 \\ QD_z - cD_x & b \end{pmatrix} + \\ &\Delta^{-1} \begin{pmatrix} D_y(-aD_x + bD_y + cD_z) + D_z(cD_y - bD_z) & -D_z a \\ D_x[D_y(cD_y - bD_z) + D_z(aD_x - bD_y - cD_z)] & -D_x D_y a \end{pmatrix} \end{aligned} \quad (3.4)$$