

The solution to (??) is determined by the elementary quadratures

$$\begin{aligned}\sigma_0(t) = & -\varepsilon c_1 c_0 c_3^2 \left[\sin c_1 t \int \cos c_1 t \mu dt - \cos c_1 t \int \sin c_1 t \mu dt \right] \\ & + A \cos c_1 t + B \sin c_1 t\end{aligned}\quad (7.12)$$

where $\mu = \mu_1 \cos 3c_1 t + \mu_2 \sin 3c_1 t$.

8 The metric

All our equations being completely solved, we can explicitly construct the corresponding metric (??)

$$\begin{aligned}ds^2 = & (v_t + u_{xx})(dt^2 + dx^2) \mp \Lambda^{-1/2}(dy^2 + dz^2) \\ & - 2b(dt dz - dx dy) + 2c(dt dy + dx dz).\end{aligned}\quad (8.1)$$

The metric coefficients are defined as follows

$$u_{xx} = -\frac{1}{c_0 \mu} e^{-c_1 x} (\pm \Lambda^{1/2} + \nu) \mp e^{-2c_1 x} \Lambda^{-1/2} \quad (8.2)$$

$$\begin{aligned}-b = v_z - u_{xy} = & \frac{1}{c_0 c_1 c_3 \mu} \left[-\frac{\mu'}{\mu} (\nu \pm \Lambda^{1/2}) + \nu' \pm \frac{1}{2} \Lambda^{-1/2} \Lambda_t \right] \cos \theta \\ & + \left[\frac{1}{c_0 c_3 \mu} (\pm \Lambda^{1/2} + \nu) \pm c_3 e^{-c_1 x} \Lambda^{-1/2} \right] \sin \theta\end{aligned}\quad (8.3)$$

$$\begin{aligned}c = v_y + u_{xz} = & \frac{1}{c_0 c_1 c_3 \mu} \left[-\frac{\mu'}{\mu} (\nu \pm \Lambda^{1/2}) + \nu' \pm \frac{1}{2} \Lambda^{-1/2} \Lambda_t \right] \sin \theta \\ & - \left[\frac{1}{c_0 c_3 \mu} (\pm \Lambda^{1/2} + \nu) \pm c_3 e^{-c_1 x} \Lambda^{-1/2} \right] \cos \theta\end{aligned}\quad (8.4)$$

while for v_t we can use the r.h.s. of equation (??). Here Λ is defined in (??) whence it follows

$$\Lambda_t = 2c_0 \mu' (c_3^2 e^{-c_1 x} - \sigma) - 2c_0 \mu \sigma_t + 2\nu \nu'$$

where σ is defined in (??), $\theta = -c_1 t + \theta_0$ and the expression for σ_t is given right after the equation (??).

Since $\sigma(t)$ and $\sigma'(t)$ involve two different combinations of y and z and there is obviously no reduction in either x or t , there is no symmetry reduction of the metric (??) in the number of independent variables and hence