

The application of the principle of virtual work is considerably simplified when the potential energy of a system is known. In the case of a virtual displacement, formula (10.19) becomes $\delta U = -\delta V$. Moreover, if the position of the system is defined by a single independent variable θ , we can write $\delta V = (dV/d\theta) \delta\theta$. Since $\delta\theta$ must be different from zero, the condition $\delta U = 0$ for the equilibrium of the system becomes

$$\frac{dV}{d\theta} = 0 \quad (10.21)$$

In terms of potential energy, therefore, the principle of virtual work states that *if a system is in equilibrium, the derivative of its total potential energy is zero*. If the position of the system depends upon several independent variables (the system is then said to possess *several degrees of freedom*), the partial derivatives of V with respect to each of the independent variables must be zero.

Consider, for example, a structure made of two members AC and CB and carrying a load W at C . The structure is supported by a pin at A and a roller at B , and a spring BD connects B to a fixed point D (Fig. 10.13a). The constant of the spring is k , and it is assumed that the natural length of the spring is equal to AD and thus that the spring is undeformed when B coincides with A . Neglecting the friction forces and the weight of the members, we find that the only forces which do work during a displacement of the structure are the weight \mathbf{W} and the force \mathbf{F} exerted by the spring at point B (Fig. 10.13b). The total potential energy of the system will thus be obtained by adding the potential energy V_g corresponding to the gravity force \mathbf{W} and the potential energy V_e corresponding to the elastic force \mathbf{F} .

Choosing a coordinate system with origin at A and noting that the deflection of the spring, measured from its undeformed position, is $AB = x_B$, we write

$$V_e = \frac{1}{2} k x_B^2 \quad V_g = W y_C$$

Expressing the coordinates x_B and y_C in terms of the angle θ , we have

$$\begin{aligned} x_B &= 2l \sin \theta & y_C &= l \cos \theta \\ V_e &= \frac{1}{2} k (2l \sin \theta)^2 & V_g &= W(l \cos \theta) \\ V &= V_e + V_g = 2kl^2 \sin^2 \theta + Wl \cos \theta \end{aligned} \quad (10.22)$$

The positions of equilibrium of the system are obtained by equating to zero the derivative of the potential energy V . We write

$$\frac{dV}{d\theta} = 4kl^2 \sin \theta \cos \theta - Wl \sin \theta = 0$$

or, factoring $l \sin \theta$,

$$\frac{dV}{d\theta} = l \sin \theta (4kl \cos \theta - W) = 0$$

There are therefore two positions of equilibrium, corresponding to the values $\theta = 0$ and $\theta = \cos^{-1} (W/4kl)$, respectively.†

†The second position does not exist if $W > 4kl$.

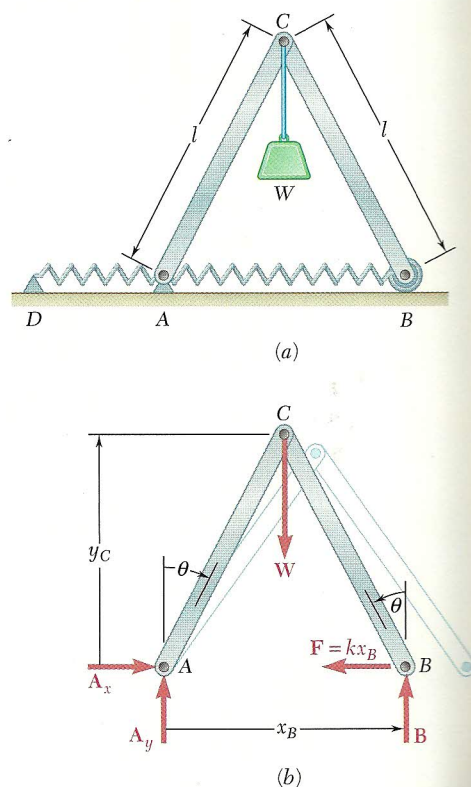


Fig. 10.13