

2º Examen Entrenamiento

1)
$$\begin{cases} v_t = (1 + \ln(1+t)) v_x + v \\ v(x, 0) = e^{-x^2} \end{cases}$$
 respecto a x

Transformamos v en $\hat{v} = \mathcal{F}_x[v]$:

$$\begin{aligned} \hat{v}_t &= (1 + \ln(1+t)) \hat{v} \cdot (i\omega)^2 + \hat{v} \Rightarrow \hat{v}(\omega, 0) = \mathcal{F}[e^{-x^2}] = \sqrt{\pi} e^{-\omega^2/4} \\ \Rightarrow \frac{d}{dt} (\ln(\hat{v})) &= 1 - \omega^2 (1 + \ln(1+t)) \Rightarrow \\ \Rightarrow \ln(\hat{v}(\omega, t)) - \ln(\hat{v}(\omega, 0)) &= t - \omega^2 t - \omega^2 \int_0^t \ln(1+s) ds \Rightarrow \\ \Rightarrow \hat{v}(\omega, t) &= \sqrt{\pi} e^{-\omega^2/4 + t - \omega^2(t+1)\ln(t+1)} \end{aligned}$$

Transformada inversa:

$$\mathcal{F}^{-1}[e^{-b x^2}] = \sqrt{\frac{b}{\pi}} e^{-b x^2}, \text{ luego escribimos:}$$

$$\hat{v} = \underbrace{\sqrt{\pi} e^t}_{\text{constante respecto a } x} e^{-\omega^2 \cdot \left[\frac{1}{4} + (t+1)\ln(t+1) \right]} \Rightarrow \frac{1}{4b} \Rightarrow b = \frac{1}{1+4(t+1)\ln(t+1)}$$

$$\Rightarrow v(x, t) = \sqrt{\pi} e^t \cdot \sqrt{\frac{b}{\pi}} e^{-b x^2} = e^t \sqrt{\frac{1}{1+4(t+1)\ln(t+1)}} e^{-\frac{1}{1+4(t+1)\ln(t+1)} x^2}$$

$$v(2, e^3 - 1) = \dots = \sqrt{\frac{1}{1+12e^3}} e^{e^3 - \frac{4}{1+12e^3}} \quad \boxed{E_3 \ln(2)}$$

Usamos que $(1+t)\ln(1+t) = \ln((1+t)^{1+t})$

2)
$$\begin{cases} w_{tt} + 2w_t + 8w = f(t), & w(0) = 0, \quad w'(0) = 1 \\ f = \begin{cases} \pi - t & \text{en } (0, \pi) \\ \text{sat} & \text{en } (\pi, \infty) \end{cases} \end{cases}$$

Aplicamos \mathcal{L} a la ecuación:

$$\begin{aligned} \hat{w} z^2 - \cancel{w(0)} \cdot z - \cancel{w'(0)} + 2(z\hat{w} - \cancel{w(0)}) + 8\hat{w} &= \hat{f} \Rightarrow \\ \Rightarrow \hat{w} &= (\hat{f} + 1) \cdot \frac{1}{z^2 + 2z + 8} \end{aligned}$$

Calculamos \hat{g} :

$$g(t) = (\pi - t) \cdot [H(t) - H(t - \pi)] + \sin t \cdot H(t - \pi) =$$

$$= (\pi - t) H(t) + (t - \pi) H(t - \pi) - \sin(t - \pi) H(t - \pi) \Rightarrow$$

$$\Rightarrow \hat{g} = \frac{\pi}{2} - \frac{1}{z} + e^{-\pi z} \cdot \frac{1}{z^2} - e^{-\pi z} \cdot \frac{1}{1+z^2}$$

Luego $\hat{\omega}(4) = \frac{1}{32} \left(\frac{\pi}{4} - \frac{1}{16} + e^{-\pi 4} \left(\frac{1}{16} - \frac{1}{17} \right) \right) = \dots$

$$= \frac{1}{2^9} \left(15 + 4\pi + e^{-4\pi} \cdot \frac{1}{17} \right). \quad \boxed{\text{Opción (5)}}$$

(3) $\sqrt{z} = \sqrt{|z|} \left(\cos\left(\frac{\text{Arg}(z)}{2}\right) + i \sin\left(\frac{\text{Arg}(z)}{2}\right) \right) \rightarrow$ Ramo principal de \sqrt{z} .

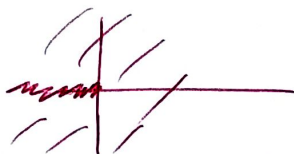
$$z \left(\cosh(\sqrt{z}) - 1 \right) w'' + \frac{\sinh(2z)}{z} w' - \sinh(z) w = 0.$$

Para ver el tipo de singularidad del punto debemos, en primer lugar, decidir si los coeficientes son analíticos en un entorno de 0, y luego el tipo de polo que surge al dividir por el coeficiente de w'' .

- $\frac{\sinh(2z)}{z}$ y $\sinh(z)$ son analíticos.

- ¿Es $\cosh(\sqrt{z}) - 1$ analítico?

Sabemos que \sqrt{z} es analítico en:



En principio, $\cosh(\sqrt{z}) - 1$ no está definida en el semieje real negativo.

Sin embargo, $\cosh(w) - 1 = \sum_{k=1}^{\infty} \frac{w^{2k}}{(2k)!}$, y esto es válido $\forall w$.

En particular, si $w = \sqrt{z}$ (donde \sqrt{z} está definida), se tiene:

$$\cosh(\sqrt{z}) - 1 = \sum_{k=1}^{\infty} \frac{z^k}{(2k)!} \rightarrow \text{Esta serie converge a todo } \mathbb{C}. \text{ Luego}$$

podemos extender la definición de $\cosh(\sqrt{z}) - 1$ a todo \mathbb{C} de manera analítica.

Escribamos la ecuación en forma estándar:

$$w'' = - \underbrace{\frac{\operatorname{sen}(2z)}{8z(\cosh(\sqrt{z})-1)}}_{b(z)/z^2} w' + \underbrace{\frac{\operatorname{sen} z}{z(\cosh(\sqrt{z})-1)}}_{a(z)/z^2} w$$

$$b(z) = - \frac{\operatorname{sen}(2z)}{8(\cosh(\sqrt{z})-1)} \Rightarrow b(0) = \lim_{z \rightarrow 0} - \frac{\operatorname{sen}(2z)}{8(\cosh(\sqrt{z})-1)} = \lim_{z \rightarrow 0} - \frac{2z}{8 \cdot (\frac{z}{2})} = -\frac{1}{2}$$

$$a(z) = \frac{z \operatorname{sen} z}{\cosh(\sqrt{z})-1} \Rightarrow a(0) = \lim_{z \rightarrow 0} \frac{z \operatorname{sen} z}{\cosh(\sqrt{z})-1} = 0$$

El sistema queda determinado por $C = \begin{bmatrix} 0 & 1 \\ 0 & 1/2 \end{bmatrix}$.

Autovalores: $\lambda_1 = \frac{1}{2}, \lambda_2 = 0$.

Sol. general: $w(z) = C_1 z^{1/2} p_1(z) + C_2 p_2(z)$, $p_1(0) = p_2(0) = 1$.

La solución es la (14), porque $z^{1/2} p_1(z)$ es $o(z^{1/4})$, y es que $\frac{z^{1/2} p_1(z)}{z^{1/4}} \xrightarrow{z \rightarrow 0} 0$.

④ $(1+t^2)w'' + \frac{t^2}{1+t^2}w = 0$, $w(0)=0$, $w'(0)=1$.

Buscamos una solución $w = \sum C_n t^n$. Nota: la ecuación está definida en $t \geq 0$. Pero si extendemos w a todo \mathbb{R} usando la misma serie se podrá ~~extender~~ hablar del desarrollo de Taylor de w .

$$(1+t^2)w'' + \frac{t^2}{1+t^2}w = 0 \Rightarrow (1+t^2)^2 w'' = -t^2 w \Rightarrow$$

$$\Rightarrow (1+t^4+2t^2)w'' = -t^2 w \Rightarrow$$

$$\Rightarrow (1+t^4+2t^2) \sum_{n=2}^{\infty} n(n-1) t^{n-2} C_n = -t^2 \sum_{n=0}^{\infty} t^n C_n$$

$$t^0 \rightarrow 2C_2 = 0$$

$$t^1 \rightarrow 6C_3 = 0$$

$$t^2 \rightarrow 4C_2 + 12C_4 = -C_0$$

$$t^3 \rightarrow 12C_3 + 20C_5 = -C_1$$

$$t^4 \rightarrow 2C_2 + 24C_4 + 30C_6 = -C_2$$

\vdots

$$t^n \rightarrow (n+2)(n+1)C_{n+2} + 2n(n-1)C_n + (n-2)(n-3)C_{n-2} = -C_{n-2}$$

$$C_{k+2} = - \frac{C_{k-2}(1 + (k-2)(k-3)) + 2k(k-1)C_k}{(k+2)(k+1)}.$$

La recurrence borne a la (9) $\times C_5 = -1/20$. luego (9)

5)
$$u_t = \left(1 + \frac{t}{\sqrt{2+2t+t^2}}\right) u_{xx}$$

$$u(x,0) = e^{-x^2}$$

$\hat{u}(\omega, t) = \hat{\mathcal{F}}_x(u(x, t))$. Aplicamos $\hat{\mathcal{F}}$ a la ecuación?

$$\hat{u}_t = \left(1 + \frac{t}{\sqrt{2+2t+t^2}}\right) \cdot (-\omega^2 \hat{u}) \Rightarrow$$

$$\Rightarrow \frac{\partial}{\partial t} (\ln(\hat{u})) = \left(1 + \frac{t}{\sqrt{2+2t+t^2}}\right) \cdot (-\omega^2) \Rightarrow$$

$$\begin{aligned} \Rightarrow \ln(\hat{u}(\omega, t)) - \ln(\hat{u}(\omega, 0)) &= -\omega^2 \int_0^t \left(1 + \frac{s}{\sqrt{2+2s+s^2}}\right) ds = \\ &= -\omega^2 t - \omega^2 \int_0^t \frac{s}{\sqrt{2+2s+s^2}} ds \end{aligned}$$

$$\begin{aligned} \int \frac{x}{\sqrt{2+2x+x^2}} dx &= \int \frac{x}{\sqrt{1+(x+1)^2}} dx = \frac{1}{2} \int \frac{2x}{\sqrt{1+(x+1)^2}} = \\ &= \frac{1}{2} \left(\int \frac{2(x+1)}{\sqrt{1+(x+1)^2}} dx - \int \frac{2}{\sqrt{1+(x+1)^2}} \right) = \\ &= \sqrt{1+(x+1)^2} - \operatorname{argsh}(x+1) \end{aligned}$$

$\int_0^t \frac{s}{\sqrt{s^2+2s+2}} ds = \sqrt{1+(t+1)^2} - \sqrt{2} - [\operatorname{argsh}(t+1) - \operatorname{argsh}(1)]$

Luego
$$\hat{u}(\omega, t) = \hat{u}(\omega, 0) \cdot e^{-\omega^2 \left(t + \sqrt{1+(t+1)^2} - \operatorname{argsh}(t+1) - \sqrt{2} + \operatorname{argsh}(1) \right)}$$

Con $\hat{u}(\omega, 0) = \hat{\mathcal{F}}[e^{-x^2}] = \sqrt{\pi} e^{-\omega^2/4} \Rightarrow$

$$\Rightarrow \hat{u}(\omega, t) = \sqrt{\pi} e^{-\omega^2 \left(\frac{1}{4} + t + \sqrt{1+(t+1)^2} - \operatorname{argsh}(t+1) - \sqrt{2} + \operatorname{argsh}(1) \right)}$$

Reemplazando $\frac{1}{4b} = \frac{1}{4} + t + \sqrt{1+(t+1)^2} - \operatorname{argsh}(t+1) - \sqrt{2} + \operatorname{argsh}(1)$, podemos hallar \hat{u}^{-1}

○ mejor, $\mathcal{F}^{-1}[\sqrt{\pi} e^{-w^2 c}] = (?) :$

$$\mathcal{F}^{-1}\left[\sqrt{\frac{\pi}{b}} e^{-w^2/4b}\right] = e^{-bx^2} \Rightarrow c = 1/4b \Rightarrow b = \frac{1}{4c}$$

$$\Rightarrow \frac{1}{\sqrt{4c}} \mathcal{F}^{-1}\left[\sqrt{\pi} e^{-w^2 c}\right] = e^{-\frac{1}{4c}x^2} \Rightarrow$$

$$\Rightarrow \boxed{\mathcal{F}^{-1}\left[\sqrt{\pi} e^{-w^2 c}\right] = \frac{1}{2\sqrt{c}} e^{-x^2/4c}} \Rightarrow$$

$$\Rightarrow v(x,t) = \frac{e^{-x^2/4 \cdot [\frac{1}{4} + t + \sqrt{1+(t+1)^2} - \alpha y \operatorname{sh}(t+1) - \sqrt{2} + \alpha y \operatorname{sh}(1)]}}{2 \sqrt{\frac{1}{4} + t + \sqrt{1+(t+1)^2} - \alpha y \operatorname{sh}(t+1) - \sqrt{2} + \alpha y \operatorname{sh}(1)}}$$

Option (2)

⑥ $v(x,\alpha) = \frac{2}{\pi} \int_{-\infty}^{\infty} \frac{f(t)}{(x-t)^2 + \alpha^2} dt = \frac{2}{\pi} \int_0^1 \frac{t(1-t)}{(x-t)^2 + \alpha^2} dt \rightarrow$

$$\Rightarrow v(1,\alpha) = \frac{2}{\pi} \int_0^1 \frac{t(1-t)}{(1-t)^2 + \alpha^2} dt$$

Integral indefinida

$$\int \frac{t(1-t)}{(1-t)^2 + \alpha^2} dt$$

1) División

$$\begin{array}{r} -t^2 + t + 0 \\ + t^2 - 2t + 1 + \alpha^2 \\ \hline -t + 1 + \alpha^2 \end{array} \quad \begin{array}{r} t^2 - 2t + 1 + \alpha^2 \\ -1 \end{array}$$

$$\int \frac{t(1-t)}{(1-t)^2 + \alpha^2} dt = -t + \int \frac{-t + 1 + \alpha^2}{(1-t)^2 + \alpha^2} dt =$$

$$= -t + \int \frac{-t + 1}{(1-t)^2 + \alpha^2} dt + \int \frac{\alpha^2}{(1-t)^2 + \alpha^2} dt = -t - \frac{\ln(\alpha^2 + (1-t)^2)}{2} +$$

$$+ \boxed{\int \frac{\alpha^2}{(1-t)^2 + \alpha^2} dt} \rightarrow \int \frac{\alpha^2}{\alpha^2 \left(\left(\frac{1-t}{\alpha} \right)^2 + 1 \right)} dt = \int \frac{1}{\left(\frac{1-t}{\alpha} \right)^2 + 1} dt = \alpha \tan\left(\frac{1-t}{\alpha}\right) \cdot (-\alpha)$$

2) Es decir:

$$\int \frac{t(1-t)}{(1-t)^2 + \alpha^2} dt = -t \frac{\ln(\alpha^2 + (1-t)^2)}{2} - \alpha \operatorname{atan}\left(\frac{1-t}{\alpha}\right) \Rightarrow$$

$$\begin{aligned} \Rightarrow U(1, \alpha) &= \frac{1}{\pi} \left(-\alpha + \frac{\alpha}{2} \left(-\ln(\alpha^2) + \ln(\alpha^2 + 1) \right) - \alpha^2 \left(\operatorname{atan}(0) - \operatorname{atan}\left(\frac{1}{\alpha}\right) \right) \right) = \\ &= \frac{1}{\pi} \left(-\alpha + \frac{\alpha}{2} \left(\ln\left(\frac{\alpha^2 + 1}{\alpha^2}\right) \right) + \alpha^2 \operatorname{atan}\left(\frac{1}{\alpha}\right) \right) = \\ &= \frac{1}{\pi} \left(-\alpha + \frac{\alpha}{2} \ln\left(1 + \frac{1}{\alpha^2}\right) + \alpha^2 \operatorname{atan}\left(\frac{1}{\alpha}\right) \right). \quad \boxed{\text{Solución: (18)}} \end{aligned}$$

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$$\begin{cases} U_t = U_{rr} + \frac{1}{r} U_r + \frac{1}{r^2} U_{\theta\theta} \\ U(1, \theta, t) = 0 \end{cases}$$

$$U(r, \theta, 0) = J_1(\alpha r) \cos \theta + J_2(\beta r) \sin(2\theta)$$

Con $\alpha \neq \beta$, α cero de J_1 y β cero de J_2 .

$$U(r, \theta, t) = \cos \theta \sum w_k(t) J_1(\lambda_{k1} r) + \sin(2\theta) \sum h_k(t) J_2(\lambda_{k2} r)$$

- Con esta definición, $U(1, \theta, t) = 0$ siempre.

$$\begin{aligned} - U(r, \theta, 0) &= \cos \theta \sum w_k(0) J_1(\lambda_{k1} r) + \sin(2\theta) \sum h_k(0) J_2(\lambda_{k2} r) = \\ &= J_1(\alpha r) \cos \theta + J_2(\beta r) \sin(2\theta). \end{aligned}$$

Por un lado, $\cos \theta$ y $\sin(2\theta)$ son funciones linealmente independientes (de hecho, ortogonales). Luego esta igualdad implica que:

$$\sum w_k(0) J_1(\lambda_{k1} r) = J_1(\alpha r)$$

$$\sum h_k(0) J_2(\lambda_{k2} r) = J_2(\beta r)$$

Ahora, aplicando unicidad de serie de Bessel:

$$\begin{cases} w_k(0) = 0 & \text{si } \lambda_{k1} \neq \alpha; \quad w_k(0) = 1 & \text{si } \lambda_{k1} = \alpha \\ \beta_k(0) = 0 & \text{si } \lambda_{k2} \neq \beta; \quad \beta_k(0) = 1 & \text{si } \lambda_{k2} = \beta. \end{cases}$$

Veamos ahora los EDOs que ~~resolven~~ cumplen w_k o β_k .

$$\begin{aligned}
 & \cos \theta \sum w_k'(t) J_1(\lambda_{k_1} r) + \sin(2\theta) \sum \beta_k'(t) J_2(\lambda_{k_2} r) = \\
 & = \cos \theta \sum w_k(t) \lambda_{k_1}^2 J_1''(\lambda_{k_1} r) + \sin(2\theta) \sum \beta_k(t) \lambda_{k_2}^2 J_2''(\lambda_{k_2} r) + \\
 & + \frac{1}{r} \left[\cos \theta \sum w_k(t) \lambda_{k_1} J_1'(\lambda_{k_1} r) + \sin(2\theta) \sum \beta_k(t) \lambda_{k_2} J_2'(\lambda_{k_2} r) \right] + \\
 & + \frac{1}{r^2} \left[-\cos \theta \sum w_k(t) J_1(\lambda_{k_1} r) - \sin(2\theta) \sum \beta_k(t) J_2(\lambda_{k_2} r) \right].
 \end{aligned}$$

Como $\cos \theta$ y $\sin(2\theta)$ son independientes, esto implicó:

$$(*) \quad \begin{cases} \cdot \sum w_k'(t) J_1(\lambda_{k_1} r) = \sum w_k(t) \lambda_{k_1}^2 J_1''(\lambda_{k_1} r) + \frac{1}{r} \sum w_k(t) \lambda_{k_1} J_1'(\lambda_{k_1} r) - \\ - \frac{1}{r^2} \sum w_k(t) J_1(\lambda_{k_1} r) \\ \cdot \sum \beta_k'(t) J_2(\lambda_{k_2} r) = \sum \beta_k(t) \lambda_{k_2}^2 J_2''(\lambda_{k_2} r) + \frac{1}{r} \sum \beta_k(t) \lambda_{k_2} J_2'(\lambda_{k_2} r) - \\ - \frac{1}{r^2} \sum \beta_k(t) J_2(\lambda_{k_2} r) \end{cases}$$

$$\begin{aligned}
 & J_1 \text{ cuaple: } x^2 J_1'' + x J_1' + (x^2 - 1) J_1 = 0 \quad \Rightarrow \text{Evaluamos en } \lambda_{k_1} r \\
 & J_2 \text{ cuaple: } x^2 J_2'' + x J_2' + (x^2 - 4) J_2 = 0
 \end{aligned}$$

$$\Rightarrow \begin{cases} \lambda_{k_1}^2 r^2 J_1''(\lambda_{k_1} r) + \lambda_{k_1} r J_1'(\lambda_{k_1} r) + (\lambda_{k_1}^2 r^2 - 1) J_1(\lambda_{k_1} r) = 0 \\ \lambda_{k_2}^2 r^2 J_2''(\lambda_{k_2} r) + \lambda_{k_2} r J_2'(\lambda_{k_2} r) + (\lambda_{k_2}^2 r^2 - 4) J_2(\lambda_{k_2} r) = 0 \end{cases}$$

Luego el sistema (*) queda:

$$\begin{cases} \sum w_k'(t) J_1(\lambda_{k_1} r) = - \sum \lambda_{k_1}^2 w_k(t) J_1(\lambda_{k_1} r) \\ \sum \beta_k'(t) J_2(\lambda_{k_2} r) = - \sum \lambda_{k_2}^2 \beta_k(t) J_2(\lambda_{k_2} r) \end{cases}$$

Por unicidad de Bessel:

$$\begin{cases} w_k' = -\lambda_{k_1}^2 w_k(t) \\ \beta_k' = -\lambda_{k_2}^2 \beta_k(t) \end{cases} \quad \begin{aligned} & w_k(0) = 0 \text{ salvo en } k_1 = k_\alpha, \text{ que vale } 1 \\ & \beta_k(0) = 0 \text{ salvo en } k_2 = k_\beta, \text{ que vale } 1 \end{aligned}$$

$$\text{La solución es } w_k(t) = \begin{cases} e^{-\alpha^2 t}, & k = k_\alpha \\ 0, & k \neq k_\alpha \end{cases} \quad \beta_k(t) = \begin{cases} e^{-\beta^2 t}, & k = k_\beta \\ 0, & k \neq k_\beta \end{cases}$$

$$\text{Luego } u(r, \theta, t) = e^{-\alpha^2 t} \cos \theta J_1(\alpha r) + e^{-\beta^2 t} \sin(2\theta) J_2(\beta r). \quad \boxed{\text{Es b (17)}}$$