# Electromagnetic Theory Handbook for Objective-First Optimization

Jesse Lu, jesselu@stanford.edu

May 20, 2011

#### Contents

| 1 | Maxwell's equations                  | 1 |
|---|--------------------------------------|---|
| 2 | Maxwell's equations for a waveguide. | 1 |
| 3 | Perfectly matched layers             | 2 |

### 1 Maxwell's equations

According to Eqs. 3.7 and 3.8 in [1], Maxwell's time-harmonic equations (E, H, J, and  $M \propto e^{-i\omega t}$ ) are

$$-i\omega H = -\frac{1}{\mu}\nabla \times E - \frac{1}{\mu}M\tag{1}$$

$$-i\omega E = \frac{1}{\epsilon}\nabla \times H - \frac{1}{\epsilon}J \tag{2}$$

where M and J are the magnetic and electric current densities, respectively. The wave equations are then,

$$\nabla \times \frac{1}{\mu} \nabla \times E - \omega^2 \epsilon E = i\omega J - \nabla \times \frac{1}{\mu} M \tag{3}$$

and

$$\nabla \times \frac{1}{\epsilon} \nabla \times H - \omega^2 \mu H = i\omega M + \nabla \times \frac{1}{\epsilon} J. \tag{4}$$

## 2 Maxwell's equations for a waveguide.

Assume that we have a uniform waveguide (not periodic, since periodic waveguides such as photonic crystal waveguides must be dealt with using Bloch's theorem). We want to find solutions of the form,

$$E(x, y, z, t) = E(x, y)e^{i\beta z - i\omega t},$$
(5)

where  $\beta$  is the wave-vector in the direction of propagation (z).

The solution for a two-dimensional waveguide of non-magnetic material ( $\mu = \mu_0$  everywhere) is messy[2], but the end result is

$$\left(\nabla \frac{1}{\epsilon_z} \nabla \cdot \epsilon - \nabla \times \nabla \times + \mu_0 \omega^2 \epsilon - \beta^2\right) E_t = 0, \tag{6}$$

where the transverse E-field components are  $E_t = \hat{x}E_x + \hat{y}E_y$ . We can then back-out the longitudinal component  $E_z$  using

$$\nabla \cdot \epsilon E = 0, \tag{7}$$

resulting in

$$E_z = \frac{i}{\beta \epsilon_z} \nabla \cdot \epsilon E_t. \tag{8}$$

If there is no variation in y (slab or one-dimensional waveguide), then we obtain

$$\left(\frac{\partial}{\partial x}\frac{1}{\epsilon_z}\frac{\partial}{\partial x}\epsilon_x + \mu_0\omega^2\epsilon_x - \beta^2\right)E_x = 0.$$
(9)

### 3 Perfectly matched layers

The upshot of ref. [3] is that a PML can be implemented by simply substituting partial derivatives in the following manner,

$$\frac{\delta}{\delta x} \to \frac{1}{1 + i \frac{\sigma_x(x)}{\delta x}} \frac{\delta}{\delta x},\tag{10}$$

where  $\sigma_x(x) > 0$  in the PML and  $\sigma_x = 0$  outside of it.

Further considerations include complex  $\sigma$ , Im  $\sigma < 0$ , to attenuate evanescent waves. Quadratic or cubic growth of  $\sigma$  to reduce numerical reflections arising from discretization error.

Generally, a half-wavelength thick PML layer is sufficient for acceptable attentuation.

### References

- [1] Allen Taflove, Susan C. Hagness, Computational Electrodynamics, Third Edition (Artech House, 2005).
- [2] Jesse Lu, 2.5D Waveguide Equations.pdf and 2.5D Waveguide Equations (simplified).pdf, https://github.com/JesseLu/misc/tree/master/scribbling.
- [3] Steven G. Johnson, Notes on Perfectly Matched Layers (PMLs) (2007).