Repetitive experiments: probability and frequency

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Formulas 9.19 to 9.22

We are trying to solve a recurrence relation 9.19

$$M(n,G) = \sum_{j=1}^{m} M(n-1, G - g_j)$$

subject to the boundary condition

$$M(0,G) = \delta(G,0).$$

Divinely inspired (but see below for some comments) we pretend that there exists some $h(\lambda)$ such that M(0, -G) is a Laplace transform of $h(\lambda)$, i.e.

$$\delta(G,0) = M(0,G) = \int d\lambda \exp{\{\lambda G\}} h(\lambda)$$

Of course this prompts the question of whether such an h actually exists. We ignore this for now (but see below).

Then, we compute M(n,G) by induction on n. Let's start with n=1 from n=0:

$$M(1,G) = \sum_{j=1}^{m} M(0,G-g_j) = \sum_{j=1}^{m} \int d\lambda \, \exp\{\lambda(G-g_i)\}h(\lambda) = \int d\lambda \, \left(\sum_{j=1}^{m} \exp\{-\lambda g_j\}\right) \exp\{\lambda(G-g_i)\}h(\lambda)$$

where $Z(\lambda) = \sum_{j=1}^{m} \exp\{-\lambda g_j\}.$

Of course we then have in the same way

$$M(n,G) = \int d\lambda \ Z^n(\lambda) \exp{\{\lambda G\}} h(\lambda)$$

aka 9.22.

Remark 1, inspiration: This is an idea as old as Fourier. IF the initial condition was $M(0,G) = \exp\{\lambda_0 G\}$ for some fixed λ then the fact that for such exponential functions shifting input G to $G - g_j$ simply multiplies the result by $\exp\{-\lambda_0 g_j\}$, we would have by induction $M(n,G) = Z^n(\lambda_0)M(0,G)$. In other words, the function $\exp\{\lambda_0 G\}$ is an eigenfunction of the operator $M(\cdot) \to \sum_{j=1}^m M(\cdot - g_j)$ with eigenvalue $Z(\lambda_0)$, so repeated application of this

operator to increase n just keeps multiplying the function by $Z(\lambda_0)$. This is the content of the formulas 9.20 and 9.21: setting n=0 in 9.20 gives the boundary value $M(0,G)=\exp\{\lambda_0 G\}$, and then the eigenvalue $\exp\{\alpha\}$ has to be equal to $Z(\lambda_0)$.

Now the equation 9.19 is linear. So if we can make the initial condition M(0, n) a sum of exponentials, the solution is the sum of corresponding solutions. This is what we did above.

This is analogous to how any exponential function $\exp\{\lambda_0 x\}$ being boundary condition to heat equation (at t=0) leads to $F(t,x)=\exp\{\lambda_0^2 t\}\exp\{\lambda_0 x\}$ solution, and so if one can write F(0,x) as sum of exponentials then the solution follows. The heat equation is second order, so one can use sines and cosines as boundary conditions instead, and then decompose the actual boundary values into those (aka Fouries analysis). We have a first order difference equation, so use exponentials.

Remark 2, finding h: I don't think there is such an $h(\lambda)$ (but have not proved this). If one iterprets $\delta(G,0)$ as a delta function (aka a distribution), then one can get arbitrarrily close by taking h to be zero-mean Gaussian with ever increasing variance (the two-sided Laplace transformis then a zero-mean Gaussian with ever decreasing vriance, which converges to the delta function).

Remark 3: I don't think any of this matters. As far as I can tell, we are only going to be using 9.27, and not 9.22. That said, I'd Jaynes's derivation of 9.27 is suspect (and the talk of "inverse Laplace transform" seems doubly suspect; this would be right if we were doing Fourier transforms, but not for Laplace.) So we better derive 9.27. We do this after an unnecessary detour through the 1D linear recurrences.

Bonus: 1D case for no good reason (maybe for motivation). The 9.19 is a 2-variable version of recurrence relation, more particularly of a constant-recursive sequence,

$$M(n) = \sum_{k=1}^{d} c_k M(n-k)$$

or putting $c_0 = -1$:

$$\sum_{k=0}^{d} c_k M(n-k) = 0$$

The answer is that if $P(x) = \sum_{k=0}^{d} c_k x^k$ has **distinct** roots μ_j , then M(n) has the form $\sum_{j=0}^{d} k_j \mu_j^n$. This is perhaps reminiscent of 9.20 (bearing in mind $\mu_j^n = \exp\{\alpha_j n\}$ where $\alpha_j = \log \mu_j$, with the caveat that μ_j are complex and could be 0).

One common approach to such equations is to study a generating function of the M(n), i.e.

$$f_0(x) := \sum_{n \ge 0} M(n) x^n$$

One sees that – modulo the edge effects – shifting the sequence to M(n-k) just multiplies f_0 by x^k :

$$f_k(x) := \sum_{n \ge k} M(n-k)x^n = \sum_{n-k \ge 0} M(n-k)x^{n-k}x^k = x^k f_0(x)$$

Then multiplying this by c_k and summing over $k = 1, \dots d$ one sees that all the terms of degree $\geq d$ cancel, so that

$$\sum_{k=0}^{d} c_k x^k f_0(x) = Q(x)$$

where Q(x) is a polynomial of degree of at most d-1 consisting of all the terms that don't cancel. Note that the Q(x) will vary depending on the initial values $M(0), M(1), \ldots M(d-1)$. Then

$$f_0(x) = \frac{Q(x)}{\sum_{k=0}^{d} c_k x^k}$$

Then one can factorize the polynomial $P(x) = \sum_{k=0}^{d} c_k x^k$ and use partial fraction expansion of $\frac{Q(x)}{P(x)}$ to get a formula fo the coefficients M(n).

We want to make the following remarks about this procedure:

- 1) Upon substitution $x=\frac{1}{z}$ the generating function becomes Z-transform, a discrete time analog of Laplace transform (think $z=\exp(s)$ and $\sum_{n\geq 0} M(n)z^{-n}$ as $\int_{t\geq 0} M(t)(\exp s)^{-t}$). The "shifting the signal multiplies the transform" property that we used is the analogue of the corresponding property of Laplace transform. In fact the very recurrence we are solving is a discrete time analogue of an ODE. So we are using the "Laplace transform method for ODEs", just in discrete time.
- 2) An advantage of working with generating functions is that they are *formal* power series, meaning we don't worry about convergence.
- 3) From this perspective what we did is instead of solving for M(n) directly, we solved for it's Z/Laplace transform and then deduced what M(n) must be.

A version of Exercise 9.1

For more on the "missing species" problem see Wikipedia, the paper of Fisher and Section 6.2 of the Computer Age Statistical Inference by Efron and Hastie.

We follow the above references answer a different, albeit related, question: namely, how many new species we may hope to find if we sample more i.e. collect more specimens.

According to Fisher,

$$S = -\alpha \ln(1-x), N = \alpha \frac{x}{1-x}$$

where α is some kind of "average Poisson rate" parameter and is independent of size of sample, while $x=\frac{\sigma}{\sigma+1}$ where $\sigma=\frac{x}{1-x}$ the scale parameter in the prior Gamma distribution over rates of Poissons, and grows linearly with sample size/sample time.

Then

$$S = \alpha \ln(1+\sigma), N = \alpha\sigma, \frac{S}{N} = \frac{\ln(1+\sigma)}{\sigma}.$$

For
$$N=122,\,S=19$$
 this gives $\frac{\ln(1+t)}{t}=\frac{19}{122}$

and

$$\sigma = 19.34583762707454276246963164708576325676864924349274639666$$

This can be expected to be reasonably accurate up to about doubling the sample size σ .

After double the sample size i.e. doubling σ , the expected number of species S becomes

So after collecting 122 more samples we expect to discover about 4 new species.

Exercise 9.2

Presumably we are to rederive formula 3.89 (and no 3.77).

Given k linear functions $\mathcal{G}_1, \ldots, \mathcal{G}_k$ of n_j with $\mathcal{G}_l(n_1, \ldots, n_m) = \sum_j g_j^l n_j$ we introduce $M(n, G_1, \ldots, G_k)$ as the number of tuples $\vec{n} = (n_1, \ldots, n_m)$ such that $\mathcal{G}_l(\vec{n}) = G_l$.

Recursion 9.19 becomes

$$M(n, G_1, \dots, G_k) = \sum_{j=1}^m M(n-1, G_1 - g_j^1, \dots, G_k - g_j^k)$$

Ansatz 9.20 becomes

$$M(n, G_1, \dots, G_k) = \exp{\{\alpha n + \sum_{l} \lambda_l G_l\}}$$

and condition 9.21

$$\exp\{\alpha\} = Z(\lambda_1, \dots, \lambda_k) = \sum_{j=1}^m \exp\{\sum_{l=1}^k -\lambda_l g_j^l\}$$

To shorten notation we write $\vec{G} = (G_1, \dots, G_k)$ and $\vec{\lambda} = (\lambda_1, \dots, \lambda_k)$. Then general solution 9.22 becomes

$$H(n, \vec{G}) = \int d\vec{\lambda} \ Z^n(\vec{\lambda}) \exp{\{\vec{\lambda} \cdot \vec{G}\} h(\vec{\lambda})}$$

Then with initial condition $M(0, \vec{G}) = \delta(\vec{G}, \vec{0})$ we have as before

$$Z^n(\vec{\lambda}) = \sum_{\vec{G}} M(n, \vec{G}) \exp\{-\vec{\lambda} \cdot \vec{g}\}$$

Now we apply this to derive multinomial distribution. Partition outcomes n_j into k sets S_1, \ldots, S_k of size s_1, \ldots, s_k and let $\mathcal{G}_l(n_j) = \begin{cases} 1 \text{ if } n_j \in S_l \\ 0 \text{ else} \end{cases}$

Then $M(n, c_1, ..., c_k)$ is the number of ways to get a tuple \vec{n} with $\sum_{j \in S_l} n_j = c_l$ and we can compute

$$Z(\vec{\lambda}) = \sum_{l} s_l \exp\{-\lambda_l\} = \sum_{l} s_l x_l$$

$$Z^{n}(\vec{\lambda}) = \sum_{(c_1, \dots, c_k)} \binom{n}{c_1, \dots c_k} \prod_{l} (s_l x_l)^{c_l}$$

and so

$$M(n, c_1, \dots, c_k) = \binom{n}{c_1, \dots c_k} \prod_l s_l^{c_l}$$

Probability of getting sequence of this type is $\frac{M(n,c_1,\ldots,c_k)}{n^m}$ and denoting $\frac{s_l}{n}=f_l$ we have

$$P(c_1, \dots, c_k) = \binom{n}{c_1, \dots c_k} \prod_{l} f_l^{c_l}$$

in agreement with 3.89.

Comments on Section 9.7

"If A is linear in the n_j , then it is the same as our G in (9.17)" – I think Jaynes means that "for a given linear function \mathcal{G} of n_j s, let A be the proposition ' \mathcal{G} takes value G'". In that case, indeed, M(n,G) = M(n,A) by definition.

"the notion of entropy inherent in probability theory independently of the work of Shannon" – in the work of Shannon, the notion of entropy is derived from combinatorics of sequences, much like here.

Exercise 9.3 (the first one)

$$f_k'' = \frac{n_k - \delta_{jk} - \delta_{tk}}{n - 2}$$

$$\delta f_k = f_k'' - f_k = f_k'' - \frac{n_k}{n} = \frac{2f_k - \delta_{jk} - \delta_{tk}}{n - 2}$$

so

$$\delta H = \sum_{k} (-1 - \ln(f_k)) \frac{2f_k - \delta_{jk} - \delta_{tk}}{n - 2}$$
$$\frac{2H + \ln(f_j) + \ln(f_t)}{n - 2}$$

and

$$H'' = \frac{nH + \ln(f_j) + \ln(f_t)}{n-2}$$

$$M(n-2, G - g_i - g_t) = \exp\{(n-2)H''\} = f_i f_t \exp\{nH\}$$

so

$$p(r_i = j, r_s = t | GnI_0) = f_i f_t$$

The trials are still pairwise independent (and by similar analysis l-wise independent as long as $k \ll n$).

Formula 9.78

Assuming not all g_j are the same, strict convexity of $\ln Z$ implies that the derivative $-\frac{\partial \ln Z}{\partial \lambda}$ attains each value in its range once. However, it only attains values $\min g_j < \lambda < \max g_j$. Averages \bar{G} outside this range - with exception of the endpoints $\min g_j$ and $\max g_j$ - are impossible (the endpoints correspond to $\lambda = \pm \infty$ and would require those u_i with non-extremal g_i to be set to zero, while not restricting the others; this is a rather special case, not directly covered by Jaynes' analysis in this section). We better condition on an achievable average.