

The central, Gaussian or normal distribution

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The Herschel–Maxwell derivation and kinetic energy.

Observe that if we postulate that kinetic E energy of a point particle is a function of its velocity $\vec{v} = (v_x, v_y)$ then we can write postulates analogous to the ones in Section 7.2:

- (P1) The total energy is sum of the energies due to the x and y motion:
 $E(x, y) = E(x, 0) + E(0, y)$
- (P2) This energy should be independent of the angle: $E(r, \theta) = E(r)$

These are log versions of the conditions in 7.2, so the function that satisfies them $E(x, y) = c(x^2 + y^2)$ is (minus) log of Gaussian density. This fact is of use in Hamiltonian Monte Carlo.

Exercise 7.1

TO DO

Exercise 7.2

Consider the family of distributions $p_{\mu, \sigma}(v)$.

We want to express fact that convolution of $p_{\mu, \sigma}(v)$ with $q(v)$ still belongs to the same family. I will prefer to work with parameter $\nu = \sigma^2$ instead. To avoid confusion the variable v will be replaced by x . So we work with the family $p_{\mu, \nu}(x)$.

Let $\mu_q = \langle \epsilon \rangle_q$ be the mean and $\nu_q = \langle \epsilon^2 \rangle_q - \langle \epsilon \rangle_q^2$ the variance of q . Then the new distribution must be $p_{\mu + \mu_q, \nu + \nu_q}(x)$. At the same time the expansion 7.20 becomes

$$p_{\mu + \mu_q, \nu + \nu_q}(x) = p_{\mu, \nu}(x) - \mu_q \frac{\partial}{\partial x} p_{\mu, \nu}(x) + \frac{1}{2}(\nu_q + \mu_q^2) \frac{\partial^2}{\partial x^2} p_{\mu, \nu}(x) + \dots$$

Taylor expanding around μ, ν

$$p_{\mu + \mu_q, \nu + \nu_q}(x) = p_{\mu, \nu}(x) + \mu_q \frac{\partial}{\partial \mu} p_{\mu, \nu}(x) + \nu_q \frac{\partial}{\partial \nu} p_{\mu, \nu}(x) +$$

$$\frac{1}{2}\mu_q^2 \frac{\partial^2}{\partial^2 \mu} p_{\mu,\nu}(x) + \frac{1}{2}\nu_q^2 \frac{\partial^2}{\partial^2 \nu} p_{\mu,\nu}(x) + \mu_q \nu_q \frac{\partial^2}{\partial \mu \partial \nu} p_{\mu,\nu}(x)$$

Now **if** we wanted this to be true for arbitrary (small) μ_q, ν_q we would have equality of Taylor coefficients:

$$-\frac{\partial}{\partial x} p_{\mu,\nu}(x) = \frac{\partial}{\partial \mu} p_{\mu,\nu}(x)$$

$$\frac{1}{2} \frac{\partial^2}{\partial^2 x} p_{\mu,\nu}(x) = \frac{\partial}{\partial \nu} p_{\mu,\nu}(x) = \frac{1}{2} \frac{\partial^2}{\partial^2 \mu} p_{\mu,\nu}(x)$$

where the first equation says that $p_{\mu,\nu}(x)$ is a function of $x - \mu$ and not of μ and x separately, $p_{\mu,\nu}(x) = f_\nu(x - \mu)$. From this $\frac{\partial^2}{\partial^2 x} p_{\mu,\nu}(x) = \frac{\partial^2}{\partial^2 \mu} p_{\mu,\nu}(x)$ follows, and we simply recover the more general Gaussina family $p_{\mu,\nu}(x) = \frac{1}{\sqrt{2\pi\nu}} \exp\{-\frac{(x-\mu)^2}{2\nu}\}$ as in 7.23.

However, **if** we instead think of μ and ν as $\mu(t)$ and $\nu(t)$ so that the family $p_t(x) = p_{\mu(t),\nu(t)}(x)$ is a single-parameter family, then the expansions become expansions in terms of t : with $\mu_q(t) = \mu'_q(0)t + o(t^2)$, $\nu_q(t) = \nu'_q(0)t + o(t^2)$

$$p_{\mu+\mu_q(t),\nu+\nu_q(t)}(x) = p_{\mu,\nu}(x) + [-\mu'_q(0) \frac{\partial}{\partial x} p_{\mu,\nu}(x) + \frac{1}{2} \nu'_q(0) \frac{\partial^2}{\partial^2 x} p_{\mu,\nu}(x)]t +$$

$$p_{\mu+\mu_q(t),\nu+\nu_q(t)}(x) = p_{\mu,\nu}(x) + \frac{\partial}{\partial t} p_{\mu,\nu}(x)t + o(t^2)$$

Equating Taylor coefficients:

$$\frac{\partial}{\partial t} p_t(x) = -\mu'_q \frac{\partial}{\partial x} p_t(x) + \frac{1}{2} \nu'_q \frac{\partial^2}{\partial^2 x} p_t(x)$$

This is a Fokker-Plank equation, albeit a very special one, with $\mu(x, t) = \mu'(0)$, $\sigma^2(x, t) = \nu'(0)$, corresponding to the stochastic process where the drift μ and diffusion coefficient $\nu/2$ are both constant. Denote $\mu'(0) = m$ and $\nu'(0) = v$.

Changing coordinates to $y(x, t) = x - mt$ aka $x(y, t) = y + mt$, we have

$$p_t(x(y, t)) = p_t(y + mt) =: q_t(y)$$

and compute by chain rule

$$\frac{\partial}{\partial t} q_t(y) = \frac{\partial}{\partial t} p_t(x(y, t)) = \frac{\partial}{\partial t} p_t(y + mt) + m \frac{\partial}{\partial x} p_t(y + mt)$$

while

$$\frac{1}{2} v \frac{\partial^2}{\partial^2 y} q_t(y) = \frac{1}{2} v \frac{\partial^2}{\partial^2 y} p_t(x(y, t)) = \frac{1}{2} v \frac{\partial^2}{\partial^2 x} p_t(y + mt)$$

So the substitution we made reduces the Fokker-Plank equation we have (with drift) to the diffusion equation (without drift) i.e. 7.22 (with $\sigma^2 = t$), which by 7.23 has solution $q_t(y) = \frac{1}{\sqrt{2\pi t}} \exp\{-\frac{y^2}{2t}\}$, or, after substitution

$$p_t(x) = \frac{1}{\sqrt{2\pi t}} \exp\{-\frac{(x - mt)^2}{2t}\}$$

This has variance $\sigma^2 = t$ so we can rewrite it as $p(x) = \frac{1}{\sqrt{2\pi\sigma^2}} \exp\{-\frac{(x - m\sigma^2)^2}{2\sigma^2}\}$, as in the formulation of the exercise.

Exercise 7.3 preliminary remarks.

I interpret this as follows. We are now considering evolution of beliefs about noise-generating process, and want to update this based on data and see that as the amount of data grows our posterior beliefs about the noise-generating process will imply beliefs about frequencies of noise that will converge to those produced by true frequencies $f(e)$.

In order to talk about the above meaningfully in mathematical sense, one needs to define some space of noise-generating processes in such a way that formulating a prior probability over it is possible, and also define in what sense the posterior beliefs about frequencies “converge”.

Since we only observe real-valued noise data (and no other type of information about the noise-generating processes) the space of noise-generating processes can be taken to be the set of real-valued stochastic processes. Now this is still very large. Now one could either 1) take only noise-generating processes that generate noise in i.i.d. way 2) focus only on the beliefs about “frequencies” i.e. how often various values (close to) various e s are observed - or both.

The intuition is that those processes that produce incorrect frequency of e s will be suppressed in the update (due to a mechanism like that in the Borel’s law of large numbers), leaving only the process that do have the correct frequencies in the “posterior distribution over processes”, whatever that means. To illustrate some of the difficulties, consider the simple case of discrete time and processes that just i.i.d. sample from some probability distribution. Consider a distribution which moreover has a density p . Then a particular dataset e_1, \dots, e_n has likelihood $\prod_{i=1}^N p(e_i)$. At each finite N the closer our distribution p is to the empirical distribution (of $\frac{1}{N} \sum_i \delta e_i$) the higher the likelihood; this seems problematic.

Perhaps one has to discretize the set of possible es and then take limit, or simply use some more sophisticated analysis.

Footnote 12

General distributions

Sum of weights.

We are doing MLE for the location parameter μ , meaning $p(\mu, y) = f(\mu - y)$. The likelihood $L_{\vec{y}}(\mu) = \prod_i f(\mu - y_i)$.

Suppose we have, for all \vec{y} near \vec{y}_0 , an isolated local minimum of $L_{\vec{y}}$ at $\hat{\mu}(\vec{y})$. Suppose further we can write $\hat{\mu}(\vec{y}) = \sum y_i w_i(\vec{y})$ with w_i (continuous) functions of the differences $y_k - y_l$. Shift each y_i by a small δ , to $\tilde{y}_i = y_i + \delta$. Then, since μ is a location parameter, likelihood is also shifted: $L_{\tilde{\vec{y}}}(\mu + \delta) = L_{\vec{y}}(\mu)$ so

$$\hat{\mu}(\tilde{\vec{y}}) = \hat{\mu}(\vec{y}) + \delta.$$

On other hand we have $\hat{\mu}(\vec{y}) = \sum y_i w_i(\vec{y})$ and since all $w_i(\vec{y})$ are unchanged by the shift plugging in $\tilde{\vec{y}}$ we get

$$\hat{\mu}(\tilde{\vec{y}}) = \sum (y_i + \delta) w_i = \hat{\mu}(\vec{y}) + \delta \left(\sum w_i(\vec{y}) \right).$$

We conclude $\sum_i w_i(\vec{y}) = 1$.

Is MLE for location parameter a weighted average?

Cases $n = 1$ and $n = 2$.

Now we look at whether $\hat{\mu}(\vec{y}) = \sum y_i w_i(\vec{y})$ is indeed true.

It is instructive to examine the case $n = 1$. Then we are just looking at local optima of $f(\mu - y_1)$. If f is unimodal with unique maximum at 0 then $\hat{\mu} = y$ and of course w is - as prescribed - function of no inputs that “sums” to 1, i.e. is the constant 1. For general f we have various local maxima m_j and MLE $\hat{\mu}_j = y + m_j$ which are not in the required form. It may be prudent to assume unimodality and peak at 0 for this problem.

For $n = 2$ such a form $\hat{\mu}(y_1, y_2) = y_1 w_1(y_1 - y_2) + y_2 w_2(y_1 - y_2)$ is achievable as follows. Denote $\frac{y_1 - y_2}{2}$ by y . If $y = 0$ the two data points coincide and $L_{\vec{y}}(\mu) = f(\mu - y_1)^2$ which has the same minima as $f(\mu - y_1)$ - we effectively have only one data point, a case which we have examined above. So below we assume $y \neq 0$.

Then by translating by $\delta = \frac{y_1 + y_2}{2}$ we have $\hat{\mu}(y_1, y_2) = \frac{y_1 + y_2}{2} + \hat{\mu}(-y, y)$. Write $\hat{\mu}(-y, y) = y w(y)$. Then

$$\hat{\mu}(y_1, y_2) = \frac{y_1 + y_2}{2} + \hat{\mu}(-y, y)$$

$$= \frac{y_1 + y_2}{2} + \frac{y_1 - y_2}{2} w(y_1 - y_2) =$$

$$y_1(w_1(y_1 - y_2)) + y_2 w_2(y_1 - y_2)$$

where

$$w_1(y_1 - y_2) = \frac{1}{2} + \frac{1}{2} w(y_1 - y_2)$$

$$w_2(y_1 - y_2) = \frac{1}{2} - \frac{1}{2} w(y_1 - y_2).$$

We note that if f is unimodal with maximum at 0 then $|\hat{\mu}(-y, y)| \leq |y|$ so $|w(y)| \leq 1$ and so $w_1, w_2 \geq 0$.

MLE as weighted average, $n > 2$.

For larger n we can still look at $\delta = \frac{1}{n} \sum y_i$, take $\tilde{y} = (y_1 - \delta, \dots, y_n - \delta)$. Observe that \tilde{y} is a vector that depends only on \vec{d} where \vec{d} is the vector of all pairwise differences $y_k - y_l$.

Now

$$\hat{\mu}(\vec{y}) = \delta + \hat{\mu}(\tilde{y}) = \delta + \mu(\vec{d})$$

Pick ANY $\vec{\alpha} = (\alpha_1, \dots, \alpha_n)$ with $\sum \alpha_i = 0$. Then $\vec{\alpha} \cdot \vec{y} = \sum \alpha_i y_i$ is a function of \vec{d} .

We can then define $w(\vec{d})$ via

$$\mu(\vec{d}) = (\vec{\alpha} \cdot \vec{y}) w(\vec{d})$$

and recover

$$\hat{\mu}(\vec{y}) = \sum y_i w_i(\vec{d})$$

with

$$w_i = \frac{1}{n} + \alpha_i w(\vec{d})$$

for all i .

This is non-unique if $n > 2$ (ther are many $\vec{\alpha}$ leding to different $w_i(\vec{d})$).

Locally near a specific \vec{y} we can pick $\alpha = \tilde{y}$.

Then $w(\vec{d}) = \frac{\hat{\mu}(\tilde{y})}{|\tilde{y}|^2}$, so that $|\tilde{y}|^2 \geq \tilde{y}_{min}^2 + \tilde{y}_{max}^2 \geq 2|\tilde{y}_{min}\tilde{y}_{max}|$. For unimodal f peaked at zero we still have $\tilde{y}_{min} \leq \hat{\mu}(\tilde{y}) \leq \tilde{y}_{max}$. Putting these together we have

$$\alpha_i w(\vec{d}) = \tilde{y}_i w(\vec{d}) \geq -\frac{\tilde{y}_{min}\tilde{y}_{max}}{|\tilde{y}|^2} \geq -\frac{1}{2}$$

implying $w_i(\vec{d}) \geq \frac{1}{n} - \frac{1}{2}$. Similarly,

$$\alpha_i w(\vec{d}) \leq \max\left(\frac{\tilde{y}_{min}^2}{|\tilde{y}|^2}, \frac{\tilde{y}_{max}^2}{|\tilde{y}|^2}\right) \leq \frac{n-1}{n}$$

implying $w_i(\vec{d}) \leq \frac{n-1}{n} + \frac{1}{n} = 1$.

Perhaps more clever choice can guarantee positivity of w_i s as well.

Cauchy distribution, mostly $n = 2$.

We illustrate this in the case when the sampling distribution is Cauchy: $p(y|\mu) = \frac{1}{\pi} \frac{1}{1+(y-\mu)^2}$.

For general n and the sample y_1, \dots, y_n the likelihood is $\prod_{i=1}^n \frac{1}{\pi} \frac{1}{1+(y_i-\mu)^2}$, and log likelihood is up to a constant $L_{\vec{y}}(\mu) = \sum_i -\ln(1+(y_i-\mu)^2)$. The extremality condition is $\frac{d}{d\mu} L_{\vec{y}}(\mu) = 0$ i.e. $\sum_i \frac{y_i-\mu}{1+(y_i-\mu)^2} = 0$. This is in general equivalent to a degree $2n-1$ polynomial equation in μ - there are many local optima for the likelihood.

Consider now the case $n = 2$.

As before, we write $y = \frac{y_1+y_2}{2}$ consider the problem shifted by $\delta = \frac{y_1+y_2}{2}$.

The optimality equation becomes $\frac{y-\hat{\mu}}{1+(y-\hat{\mu})^2} - \frac{y+\hat{\mu}}{1+(y+\hat{\mu})^2} = 0$ so $f(x) = \frac{x}{1+x^2} = \frac{1}{x+\frac{1}{x}}$ has $f(y-\hat{\mu}) = f(y+\hat{\mu})$. Either $y-\hat{\mu} = y+\hat{\mu}$, i.e. $\hat{\mu} = 0$ or $(y-\hat{\mu})(y+\hat{\mu}) = 1$, $\hat{\mu} = \pm\sqrt{y^2-1}$. This last pair of solution is real only if $|y| > 1$.

Suppose that in fact $|y| > 1$. Then one over the likelihood is a positive fourth degree polynomial which we now know has three local extrema - $-\sqrt{y^2-1}, 0, \sqrt{y^2-1}$. Therefore these extrema must be non-degenerate and be min, max, min. Correspondingly, the (log)likelihood extrema must be max, min, max.

The MLE estimate is thus indifferent between

$$\hat{\mu}_1 = \frac{y_1+y_2}{2} + \sqrt{\left(\frac{y_1-y_2}{2}\right)^2 - 1}$$

and

$$\hat{\mu}_2 = \frac{y_1 + y_2}{2} - \sqrt{\left(\frac{y_1 - y_2}{2}\right)^2 - 1}.$$

Let's see how this can be written in the form $\hat{\mu} = y_1 w_1(y_1 - y_2) + y_2 w_2(y_1 - y_2)$.
We treat $\hat{\mu}_1$

As prescribed,

$$w(y) = \sqrt{y^2 - 1} = y \cdot \text{sgn}(y) \sqrt{\frac{1}{4} - \frac{1}{y^2}}$$

so

$$w_1(y) = \frac{1}{2} + \text{sgn}(y) \sqrt{\frac{1}{4} - \frac{1}{y^2}}$$

and

$$w_2(y) = \frac{1}{2} - \text{sgn}(y) \sqrt{\frac{1}{4} - \frac{1}{y^2}}$$

and indeed,

$$y_1 w_1(y_1 - y_2) + y_2 w_2(y_1 - y_2) =$$

$$\frac{y_1 + y_2}{2} + (y_1 - y_2) \cdot \text{sgn}(y_1 - y_2) \sqrt{\frac{1}{4} - \frac{1}{(y_1 - y_2)^2}} =$$

$$\frac{y_1 + y_2}{2} + \sqrt{\left(\frac{y_1 - y_2}{2}\right)^2 - 1} =$$

$$y_1 w_1(y_1 - y_2) + y_2 w_2(y_1 - y_2) = \hat{\mu}_1$$

Observe that $w_1(y) + w_2(y) = 1$ as expected. Observe also that for $\hat{\mu}_2$ we would obtain $\tilde{w}_1 = w_2$ and $\tilde{w}_2 = w_1$ - the individual μ_i and w_i are not symmetric in exchanging y_i s, but the sets of μ s and w s are.

Exercise 7.4

We are minimizing $\vec{w}^T C \vec{w}$. Let's minimize over the hyperplane $\sum w_i = 1$. Since C is positive definite on \mathbb{R}^n , at infinity the values are large and positive. So the minimum is achieved at finite distance and must satisfy the Lagrange multiplier equation is $C\vec{w} = \lambda \vec{1}$, so $\vec{w} = \lambda C^{-1} \vec{1}$ and $\sum w_i = 1$ gives, with $C^{-1} = K$

the answer $w_i = \sum_j K_{ij} / \sum_{i,j} K_{ij}$, as wanted (note that the denominator is $\vec{1}^T K \vec{1} > 0$).

Corresponding value is

$$\vec{w}^T C \vec{w} = \lambda^2 (C^{-1} \vec{1})^T C (C^{-1} \vec{1}) = \lambda = \left(\sum_{ij} K_{ij} \right)^{-1}$$

HOWEVER this answer satisfies does not always satisfy the constraints $w_i \geq 0$: consider $C = \begin{pmatrix} 1 & 4 \\ 4 & 17 \end{pmatrix}$ so that $K = \begin{pmatrix} 17 & -4 \\ -4 & 1 \end{pmatrix}$; then $w = (1.3, -0.3)^T$.

Exercise 7.5

See <https://www.cs.toronto.edu/~yuvalf/CLT.pdf>

7.84

$$\exp\{xa - \frac{a^2}{2}\} = \exp\{\frac{x^2}{2}\} \exp\{-\frac{(x-a)^2}{2}\}$$

so

$$\frac{d^n}{da^n} \exp\{xa - \frac{a^2}{2}\} = \exp\{\frac{x^2}{2}\} \frac{d^n}{da^n} \left(\exp\{-\frac{(x-a)^2}{2}\} \right)$$

7.85

$$\frac{\phi(x-a)\phi(x-b)}{\phi(x)} = \phi(x) \left(\sum_n R_n(x) \frac{a^n}{n!} \right) \left(\sum_m R_m(x) \frac{b^m}{m!} \right)$$

LHS: Observe

$$\phi(x-a)\phi(x-b) = \phi(x)\phi(x-(a+b))\exp\{ab\}$$

Then

$$\int \frac{\phi(x-a)\phi(x-b)}{\phi(x)} dx = \int \phi(x-(a+b)) \exp\{ab\} dx = \exp\{ab\}$$

As a power series in ab it is $\sum_i \frac{a^i b^i}{i!}$.

RHS:

$$\sum_{n,m} (\int \phi(x) R_n(x) R_m(x) dx) \frac{a^n b^m}{n! m!}$$

Equating coefficients we get 7.85.

7.86

From 7.83 $\phi(x-y) = \phi(x) \sum_m R_m(x) \frac{y^m}{m!}$ so

$$\phi(x-y) R_n(x) = R_n(x) \phi(x) \sum_m R_m(x) \frac{y^m}{m!}$$

and integrating and using 7.85 we get

$$\int \phi(x-y) R_n(x) dx = y^n.$$

7.89

$$\exp\{xa\} \exp\{-a^2/2\} = \sum_k \frac{x^k a^k}{k!} \sum_m \frac{a^{2m}}{(-2)^m m!}$$

Isolating the term in front of $n = k + 2m$ gives 7.89