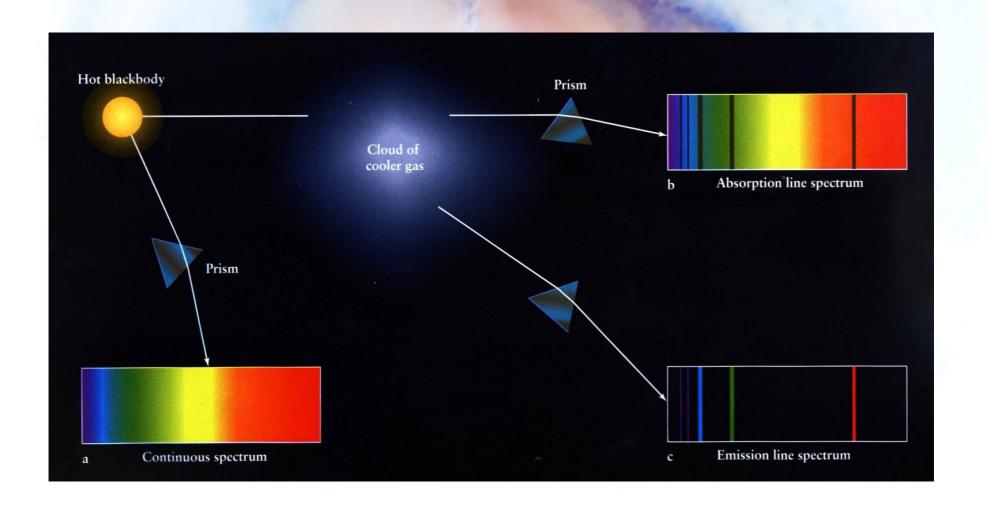
# Interstellar medium for dumnies

Christophe Morisset Instituto de Astronomía, UNAM Ensenada, Mexico

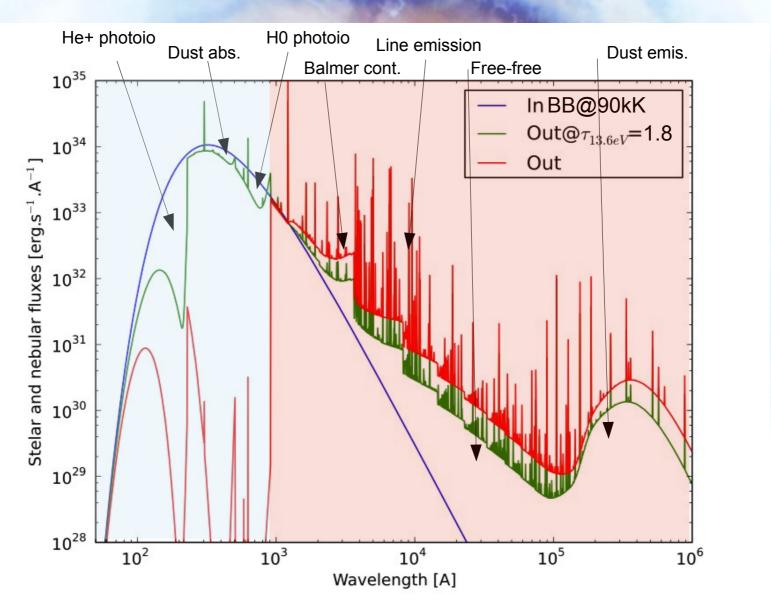
# Summary

- Introduction
- Emission processes
- Line emissivities
- PyNeb
- Atomic data

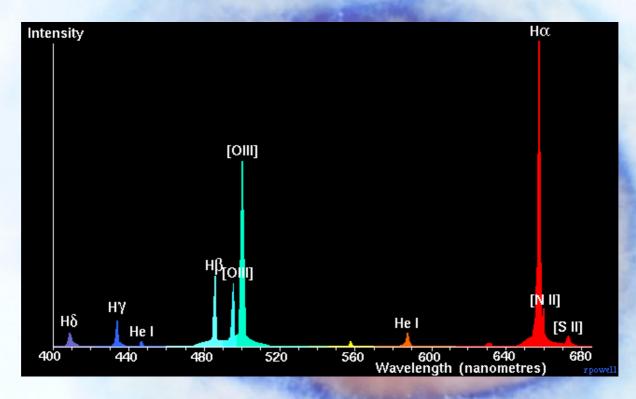
# Kirchhoff 1860



# Ionized ISM is an active filter to the ionizing photons

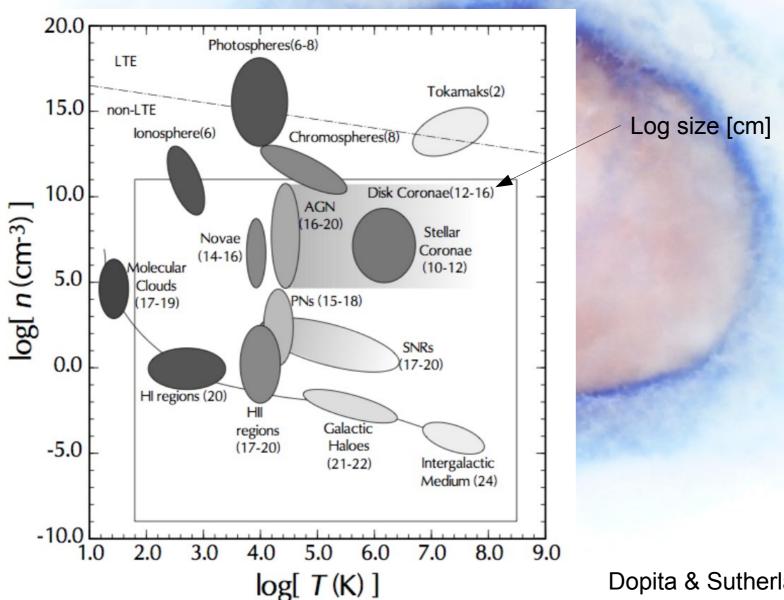


#### **Emission lines**



- Easy to detect and measure on a faint continuum. Redshifts
- Trace gas
- Close to ionizing source :
  - Hot stars (Hot == OB == Young, CSPN == Old)
  - AGN
  - (shocks)

# Astrophysical plasmas



Dopita & Sutherland, 2003

# ISM: the two equilibria

Ionization equilibrium :

ionization <==> recombination

Photoionization
Collisions (micro)
Charge exchange

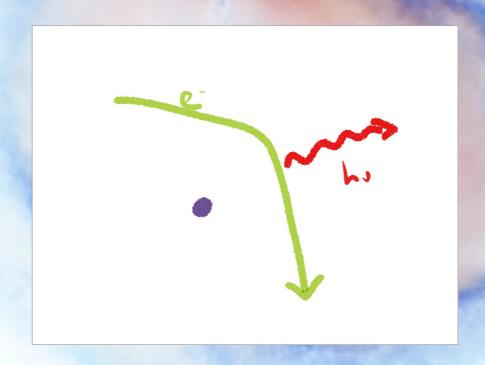
Radiative recombination
Dielectronic recombination
Charge exchange

Thermal equilibrium :

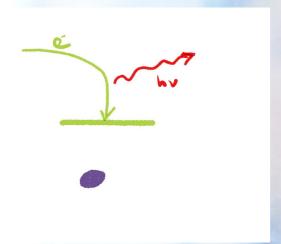
heating <==> cooling

Photoionization
Collisions (macro)

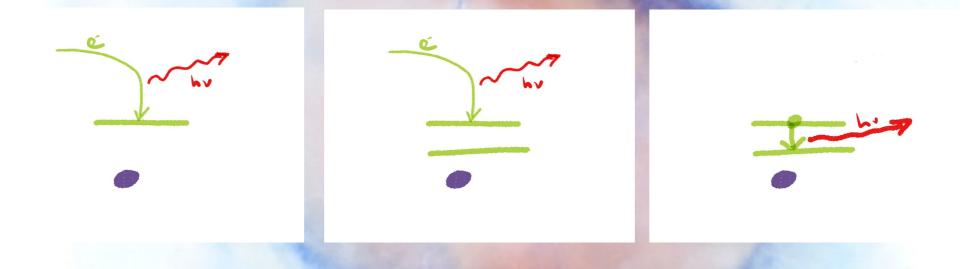
Free-free radiation
Free-bound radiation
Bound-bound radiation



Bremstrallung = free-free continuous emission

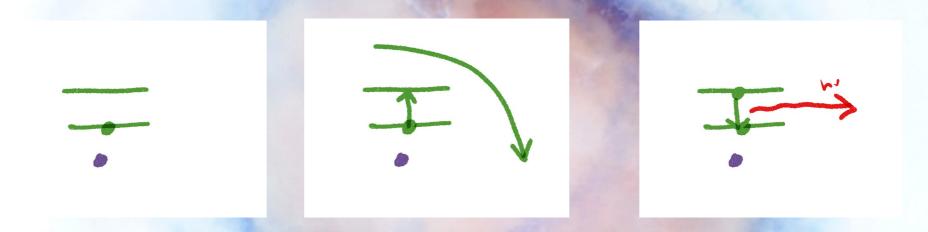


Recombination of an electron to ion:
continuous emission



Recombination followed by transition between 2 levels:

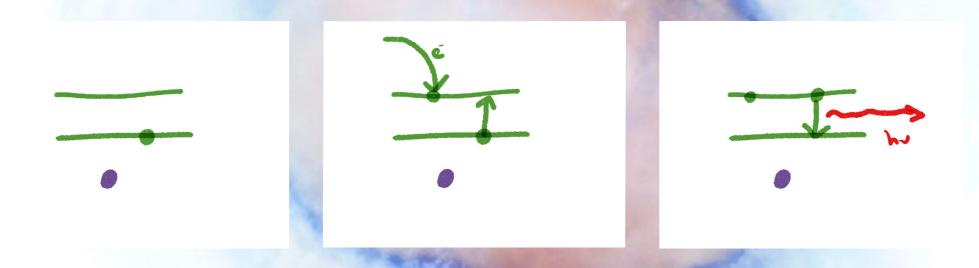
emission line



Collisional excitation (an electron gives part of its kinetic energy)

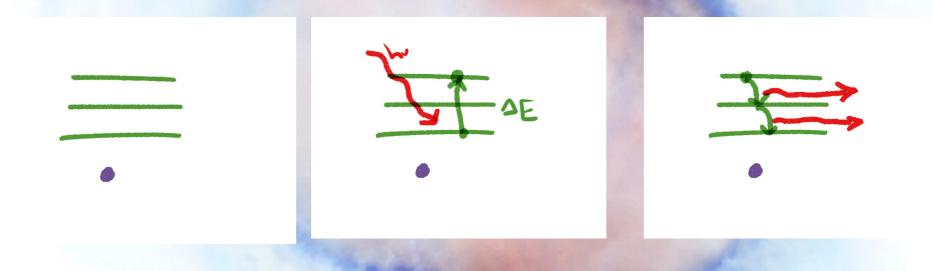
followed by transition:

emission line



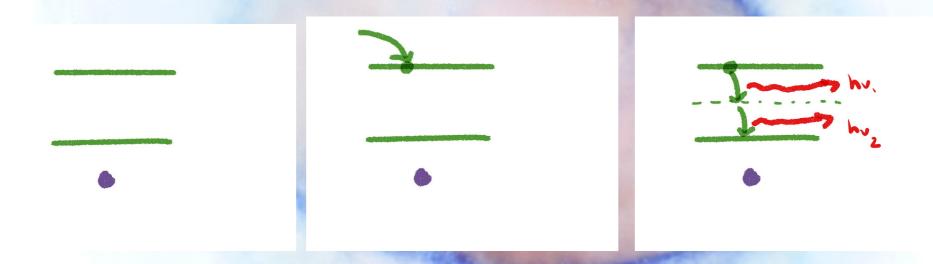
Dielectronic recombination : recombining electron excites inner electron.

emission line



Fluorescence: a photon is absorbed, followed by one or more decay(s)

emission line(s)



2 photons recombination (H) continuous emission

- Free-free, bound-free and bound-bound processes.
- Recombination works better at low temperature.
- Collision excitation needs enough energetic electrons (high Te if high energy level).

- Forbidden transition: not allowed (!)
- In earth laboratory, not in ISM → 5007 is not emitted by Nebulium

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#### THE ORIGIN OF THE CHIEF NEBULAR LINES By I. S. Bowen

Several of the strongest lines in the spectra of the gaseous nebulae have not been observed in terrestrial sources. Since the spectra of the light elements, which are thought to form the chief constituents of nebulae, have been thoroughly studied, this leads to the conclusion that some cause, such as low density, must be operating in the nebulae to bring out lines in addition to those found in laboratory sources.

- Permitted lines vs. forbidden lines: a matter of transition probability of the upper level. Einstein coefficients As.
- If another collision occurs while the electron is in the upper level: collisional desexcitation, no emission.
- The criterium to determine which desexcitation (radiative or collisional) dominates is the density (critical density).

- Forbidden vs. permitted lines: related to the decay (upper level lifetime).
- Collisionally excited or recombination : related to the upper level population process.
- Not systematically related: there is a contribution from recombination to forbidden lines, and there is collisional excitation of permitted lines.

Kinetic equilibrium -> electron temperature :

$$E = \frac{1}{2}.m.v^{2} = \frac{3}{2}.k_{B}.T_{e}$$
$$E(eV) = T_{e}/7736K$$

Energy [eV] to ionize H<sup>0</sup> into H<sup>+</sup> + e-?

Corresponding wavelength?

Corresponding Te?

Kinetic equilibrium -> electron temperature :

$$E = \frac{1}{2}.m.v^2 = \frac{3}{2}.k_B.T_e$$

$$E(eV) = T_e/7736K$$

Energy [eV] to ionize Ho into H+ e-: 13.6 eV

Corresponding wavelength: 912 A

Corresponding Te: 105,000 K

Collisional ionization

#### Planck function

$$B(\lambda, T) = \frac{2hc^2}{\lambda^5} \frac{1}{e^{\frac{hc}{\lambda k_B T}} - 1}$$

Peak at T / 3030K (eV)

Te corresponding to peak at 13.6 eV?

Stellar type?

#### Planck function

$$B(\lambda, T) = \frac{2hc^2}{\lambda^5} \frac{1}{e^{\frac{hc}{\lambda k_B T}} - 1}$$

Peak at T / 3030K (eV)

Te corresponding to peak at 13.6 eV: 45,000 K

Stellar type: **O2** 

#### Planck function

$$B(\lambda, T) = \frac{2hc^2}{\lambda^5} \frac{1}{e^{\frac{hc}{\lambda k_B T}} - 1}$$

Any star with Teff > 24,000 K (B2) emits more than 10 % of its radiation with  $\lambda$  < 912 Å.

Relation of the ionizing photons emitting rate and the volume of ionized gas :  $Q_H = \int_V n_H^2 . \alpha_B(H) . dv$ 

In the case of constant density filled sphere of Strömgren radius R<sub>s</sub>:

$$Q_H = ff.\frac{4}{3}.\pi.R_S^3.n_H^2.\alpha_B(H)$$

In the case of decreasing density  $n_H(r) = \bar{n}_H.r^{-a}$ 

$$Q_H = ff.4./(3-2a).\pi.\left[R_{rec}^{(3-2a)} - R_{in}^{(3-2a)}\right].\bar{n}_H^2.\alpha_B(H)$$

Hß luminosity of a nebula:

$$L_{H\beta} = n_H^2 . ff. V. \epsilon(H\beta)$$

Absorption of ionizing photons:

$$Q_{H,abs} = n_H^2.ff.V.\alpha_b(H)$$

Ionization-bounded case:

$$Q_{H,abs} = Q_H \to L_{H\beta} = Q_H.\epsilon(H\beta)/\alpha_B(H)$$

Density-bounded case:

$$L_{H\beta} = n_H.M_{neb}.\epsilon(H\beta)/m_H$$

$$U(r) = \frac{Q_H}{4.\pi \cdot r^2 \cdot n_H \cdot c}$$

#### For a Strömgren sphere:

$$Q_H = 4/3.\pi . R_S^3 . n_H^2 . ff.\alpha_B$$

$$U(R_S) = \frac{Q_H}{4.\pi . R_S^2 . n_H . c} = \frac{\Phi_H}{n_H . c}$$

$$< U > = \int_{V} U dv = \frac{3.Q_{H}}{4.\pi R_{S}^{2}.n_{H}.c}$$

$$< U> = A.(Q_H.n_H.ff^2)^{1/3}$$

$$where A = \left[\frac{3.\alpha_B^2}{4.\pi.c^2}\right]^{1/3}$$

#### Surface brightness:

$$S(H\beta) = F(H\beta)/\Theta^2 = L(H\beta)/(4.\pi \cdot d^2 \cdot \Theta^2) = L(H\beta)/(4.\pi \cdot R_S^2)$$

#### For a radiation-bounded nebula:

$$L(H\beta) \propto Q_H$$

$$< U > \propto S(H\beta)/n_H$$

- Line emissivities are computed using atomic data.
   They can strongly influence the result.
- Energy emitted by a line from level 2 to level 1:

$$I_{2,1} = n_2.A_{2,1}.h.\nu_{2,1}$$

Statistical equilibrium :

$$n_1.n_e.q_{1,2} = n_2.n_e.q_{2,1} + n_2.A_{2,1}$$

A<sub>i,j</sub>: Einstein coefficient = transition probability = s<sup>-1</sup>

q : (de)excitation coefficient = collision rates = cm<sup>-3</sup>.s<sup>-1</sup>

Effective collision strengths:

$$\Upsilon_{2,1}(T) = \int_{0}^{\infty} \Omega_{2,1}(E) \cdot e^{-\frac{E}{k \cdot T}} \cdot d\frac{E}{k \cdot T}$$

Collision coefficients:

$$q_{2,1} = \frac{8.629.10^{-6}}{T^{1/2}} \frac{\Upsilon_{2,1}(T)}{g_2}$$

$$q_{1,2} = \frac{g_2}{g_1} \cdot q_{2,1} \cdot e^{-\frac{h \cdot \nu_{2,1}}{k \cdot T}}$$

#### General 2-levels emission:

$$I_{2,1} = \frac{n_1 \cdot \frac{g_2}{g_1} \cdot e^{\frac{-h \cdot \nu_{2,1}}{k \cdot T}} \cdot A_{2,1} \cdot h \cdot \nu_{2,1}}{1 + \frac{A_{2,1}}{n_e \cdot q_{2,1}}}$$

#### Low density limit:

$$I_{2,1} = n_1.n_2.q_{1,2}.h.\nu_{2,1} \propto n^2$$

#### High density limit:

$$I_{2,1} = n_2.A_{2,1}.h.\nu_{2,1} \propto n$$

- Recombination lines: needs for effective radiative recombination coefficients
- H, He+ and He++
- But also for metals, hard to compute : ADF problem...

 As a summary, the emissivities of recombination lines go like:

$$\epsilon \propto rac{1}{T_e}$$

 While the emissivities of collisionally excited lines go like :

$$\epsilon \propto \frac{e^{-\frac{h.\nu}{k.Te}}}{\sqrt{Te}}$$

# pyStuff

#### PyNeb :

- Luridiana, Morisset, Shaw 2012
- Python « modern » version of FIVEL and nebular packages.
   Now extends to much more facilities:
  - More levels, recombination lines, continuum, Balmer decrement, plotting facilities.
- Easy manage atomic data
- Easy install : pip install pyneb
- Github: https://github.com/Morisset/PyNeb\_devel
- Google discussion group : https://groups.google.com/forum/#!forum/pyneb