

# Solving PDEs using Spectral Methods in the Chebyshev basis by example of the Heat Equation

Special Topic on [APPROXIMATION OF FUNCTIONS](#)

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## Abstract

This work shall attempt to numerically solve the heat equation  $u_t = \alpha u_{xx}$  with Dirichlet boundary conditions over the domain  $[-1, 1] \times [0, T]$  by representing the spatial component as a *Chebfun* (Chebyshev series) and moving on in time by the Forward Euler numerical scheme.

**Our Goal:** Numerically obtain the solution  $u(x, T)$  of

$$\begin{cases} \frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2} & u : [-1, 1] \times [0, T] \mapsto \mathbb{R}, T \in \mathbb{R}^+, \alpha \in \mathbb{R}^+ \\ u(x_j, 0) = u_0(x_j) & \forall x_j \in X_N, N \in \mathbb{N}, N > 1, u_0 : [-1, 1] \mapsto \mathbb{R} \\ u(b, t) = u_0(b) & \forall b \in \{-1, 1\}, \forall t \in (0, T]. \end{cases}$$

The implementation, centered around what we will refer to as **Tscheb-Fun**, including three major algorithms **TschebFun::interpolantThrough()**, **TschebFun::evaluateOn()** and **TschebFun::derivative()**, is done manually in C++, extended to work as a Python module and for demonstration, even features a high-level graphical interface to play with. Finally, we will compare the numerical results with the output of *Chebfun*'s high-level **pde15s()**.



**Figure 1:** Screenshot of the graphical user interface. After entering an initial expression  $u_0(x)$ , depicted in grey, the simulation will run upon pressing 'Start'. The solution at time  $t$ , depicted in blue, is represented as a Chebyshev series of degree 29.

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## Contents

<b>1</b>	<b>Motivation</b>	<b>3</b>
1.1	The heat equation and its solution . . . . .	3
<b>2</b>	<b>Chebyshev Interpolation</b>	<b>3</b>
2.1	An Orthogonal Basis . . . . .	4
2.2	Well, and numerically? . . . . .	6
2.3	A Numerical Issue . . . . .	7
2.4	The Algorithm . . . . .	8
<b>3</b>	<b>The Spectral Method</b>	<b>9</b>
3.1	Forward Euler . . . . .	9
3.2	We need Differentiation . . . . .	9
3.3	Enforcing boundary conditions . . . . .	10
3.4	The Destination . . . . .	10
<b>4</b>	<b>Evaluation, Analysis, Comparison and Results</b>	<b>11</b>
4.1	Evaluation: The Clenshaw Algorithm . . . . .	11
4.2	Analysis: Interface to Python . . . . .	11
4.3	Comparison: ChebFun . . . . .	11
4.4	Results: On Target? . . . . .	12
<b>5</b>	<b>Outlook and Discussion</b>	<b>13</b>
<b>A</b>	<b>Title of Appendix</b>	<b>14</b>

# 1 Motivation

Partial differential equations are notoriously hard to solve. One more possible approach to make way in this important class of problems is by the technique of spectral methods, incidentally closely related to finite element methods. The key idea is to perform the problem solution by representation of the occurring functions in a certain basis. For non-periodic problem settings, CHEBYSHEV series are a fantastic choice.

## 1.1 The heat equation and its solution

Classical approaches to solving the heat equation are

- Using the separation Ansatz

$$u(x, t) = \underline{X}(x)\underline{T}(t), \quad X : [-1, 1] \mapsto \mathbb{R}, \quad \underline{T} : [0, T] \mapsto \mathbb{R}$$

- By Fourier-analysis (similar to below, or maybe even simpler: differentiation simply multiplies the Fourier coefficients by  $ik$ ).
- Numerically, a finite difference scheme such as

$$\frac{U_j^{(n+1)} - U_j^{(n)}}{\Delta t} = \alpha \frac{U_{j+1}^{(n)} - 2U_j^{(n)} + U_{j-1}^{(n)}}{(\Delta x)^2}$$

elaborate

# 2 Chebyshev Interpolation

Let  $\mathbb{N}$  denote the nonnegative integers, so  $0 \in \mathbb{N}$ .

## 2.1 Definition: Chebyshev polynomial

Chebyshev<sup>1</sup> polynomials  $T_k : \mathbb{R} \mapsto \mathbb{R}$  are functions satisfying

$$T_k(x) = T_k(\cos \theta) := \cos(k\theta) = \frac{1}{2}(z^k + z^{-k})$$

$$z := e^{i\theta}, \quad x := \Re(z) = \cos(\theta) = \frac{1}{2}(z + z^{-1})$$

for degree  $k \in \mathbb{N}$ . Then,  $T_0(x) = 1$ ,  $T_1(x) = x$ ,  $T_2(x) = 2x^2 - 1$ , and so on.

<sup>1</sup>named after Pafnuty Lvovich CHEBYSHEV, alternatively transliterated as Tchebycheff, Tchebyshev (French) or TSCHEBYSCHOW (German)

These relations between  $x$ ,  $z$  and  $\theta$  reveal fundamental connections between three famous basis sets: CHEBYSHEV, LAURENT and FOURIER. [elaborate](#)

### 2.1 Theorem: Chebyshev Recursion Formula

The Chebyshev polynomials satisfy the three-term recurrence relation

$$T_{k+1}(x) = 2xT_k(x) - T_{k-1}(x).$$

*Proof.* Theorem 2.1 For  $k > 1$ ,

$$\begin{aligned} 2xT_k(x) - T_{k-1}(x) &= 2x \frac{1}{2}(z^k + z^{-k}) - \frac{1}{2}(z^{k-1} + z^{-(k-1)}) \\ &= 2 \frac{1}{2}(z + z^{-1}) \frac{1}{2}(z^k + z^{-k}) - \frac{1}{2}(z^{k-1} + z^{-k+1}) \\ &= \frac{1}{2}(z^{k+1} + z^{k-1} + z^{-k+1} + z^{-k-1}) - \frac{1}{2}(z^{k-1} + z^{-k+1}) \\ &= \frac{1}{2}(z^{k+1} + z^{-(k+1)}) = T_{k+1}(x) \end{aligned}$$

□

### 2.1 An Orthogonal Basis

The Chebyshev polynomials also satisfy an *orthogonality relation*,

$$\langle T_m, T_n \rangle := \int_{-1}^1 T_m(x) T_n(x) \frac{1}{\sqrt{1-x^2}} dx = \int_{\pi}^0 \cos(m\theta) \cos(n\theta) \frac{-\sin(\theta)}{\sqrt{1-\cos^2(\theta)}} d\theta,$$

which becomes, with the fitting substitution  $x = \cos(\theta)$  and  $dx = -\sin(\theta)d\theta$ ,

$$\begin{aligned} \langle T_m, T_n \rangle &= \int_0^{\pi} T_m(\cos \theta) T_n(\cos \theta) \frac{\sin \theta}{\sin \theta} d\theta = \int_0^{\pi} \cos(m\theta) \cos(n\theta) d\theta \\ &= \frac{1}{2} \int_0^{\pi} \left( \underbrace{\cos((m+n)\theta)}_{=\cos(2m\theta) \text{ for } m=n} + \underbrace{\cos((m-n)\theta)}_{=1 \text{ for } m=n} \right) d\theta \end{aligned}$$

along with the knowledge that  $\int_0^{\pi} \cos(k\theta) d\theta = k^{-1} [\sin(k\theta)]_0^{\pi} = 0$  for  $k \in \mathbb{Z} \setminus \{0\}$ ,

$$\langle T_m, T_n \rangle = \int_0^{\pi} T_m(\cos \theta) T_n(\cos \theta) d\theta = \begin{cases} 0 & \text{for } m \neq n \\ \pi/2 & \text{for } m = n \neq 0 \\ \pi & \text{for } m = n = 0 \end{cases}$$

which can be effectively utilised to define a function space  $(\mathcal{T}, +, \cdot)$  in the *orthogonal* basis of Chebyshev polynomials  $\mathcal{T} := \{T_k\}_{k \in \mathbb{N}}$ . Note that the operation  $\langle \cdot, \cdot \rangle$  satisfies

all axioms of an authentic inner product (linearity, etc.) over a function space due to the linearity of the integral.

In the following proceedings, we will restrict our view on functions over the interval  $[-1, 1] \subset \mathbb{R}$ . Any (real) Lipschitz-continuous function  $f \in \mathcal{C}_L$ , where  $\mathcal{C}_L := \{g : [-1, 1] \mapsto \mathbb{R} \mid \exists L \text{ s.t. } \forall x_1, x_2 \in \mathbb{R}, |g(x_1) - g(x_2)| \leq L \cdot |x_1 - x_2|\}$  can be represented in the Chebyshev basis  $\mathcal{T}$ , as Lipschitz continuity is a sufficient condition for absolute and uniform convergence of the corresponding series representation

$$f(x) = \sum_{k=0}^{\infty} a_k T_k(x), \quad a_k \in \mathbb{R}, \quad k \in \mathbb{N}.$$

Utilising orthogonality, for any  $f \in \mathcal{C}_L$ , we find coefficients  $a_l \in \mathbb{R}$  by 'right-multiplying' the equation  $f = \sum_{k=0}^{\infty} a_k T_k$  with any one of the Chebyshev polynomials  $T_l$ .

$$\begin{aligned} \langle f, T_l \rangle &= \left\langle \sum_{k=0}^{\infty} a_k T_k, T_l \right\rangle = \int_0^{\pi} a_k T_k(\cos \theta) T_l(\cos \theta) d\theta \\ &= \sum_{k=0}^{\infty} a_k \langle T_k, T_l \rangle \quad \text{by linearity} \\ &= \begin{cases} a_0 \pi & \text{for } l = 0 \\ a_l \pi / 2 & \text{for } l \neq 0 \end{cases} \end{aligned}$$

which can easily be rearranged to give explicit relations for  $a_0$  and  $a_k$ , summarised in the below theorem.

## 2.2 Theorem: Chebyshev series coefficient formula

For any  $f \in \mathcal{C}_L$ , one can obtain the Chebyshev series coefficients  $a_k$ ,  $k \in \mathbb{N}$  as

$$\begin{aligned} a_0 &= \frac{1}{\pi} \langle f, T_0 \rangle = \frac{1}{\pi} \int_0^{\pi} f(\cos \theta) d\theta \\ a_k &= \frac{2}{\pi} \langle f, T_k \rangle = \frac{2}{\pi} \int_0^{\pi} f(\cos \theta) \cos(k\theta) d\theta, \quad k \neq 0. \end{aligned}$$

*Proof.* As given in the discussion above. A different approach for the derivation of the explicit coefficient integrals can be found in [Trefethen 2019](#) along with a complex analysis styled proof.  $\square$

Dealing with a numerical problem, we shall then approximate the above two integrals by the rectangular integral rule.

## 2.2 Well, and numerically?

As computers rarely allow us to store infinitely many coefficients  $a_k$ , we will work with the truncated Chebyshev series

$$f_N(x) = \sum_{k=0}^{N-1} a_k T_k(x), \quad k \in \{0, \dots, N\}, \quad N \in \mathbb{N}, N > 1$$

which approximates the function  $f$  with a degree  $N - 1$  polynomial up to an error

$$f(x) - f_N(x) = \sum_{k=0}^{\infty} a_k T_k(x) - \sum_{k=0}^{N-1} a_k T_k(x) = \sum_{k=N}^{\infty} a_k T_k(x).$$

### Method One: Integral Approximation Text

#### 2.3 Theorem: Rectangular integral rule

$$\int_a^b f(x) dx = \lim_{N \rightarrow \infty} \frac{b-a}{N} \sum_{k=0}^{N-1} f(x_k), \quad x_k := a + \frac{b-a}{N} k$$

Bonthuis and Schachinger 2021, p. 128 [elaborate](#)

Most importantly, this quadrature-style integral approximation is only one way of numerically determining the coefficients  $a_k$ . Another is to recognise the structure of the above integral for  $k \neq 0$  as a cosine transform of the function  $(f \circ \cos)$ .

### Method Two: Relation to the DCT Text

#### 2.2 Definition: Cosine Transform

$$\hat{f}(\omega) := \int_0^\pi f(t) \cos(\omega t) dt$$

#### 2.3 Definition: DCT-II: Discrete Cosine Transform - type II

$$\hat{y}_k := \sum_{n=0}^{N-1} y_n \cos \left[ \frac{\pi}{N} \left( n + \frac{1}{2} \right) k \right] \quad \text{for } k = 0, \dots, N-1.$$

Most significantly, this approach via the Discrete Cosine Transform can be sped up by means of the *Fast Fourier Transform* (Cooley and Tukey 1965).

The above two methods demonstrate why it is sensible to choose to sample a function in *chebpoints* instead of equispaced points, simply because the mathematics works

out nicely in this basis. A much more fundamental numerical insight relating to the choice of points is the *Runge phenomenon* (Runge 1901).

**Method Three: Barycentric Interpolation Formula** Numerically speaking, a significant improvement to these two approaches can be made by using the *Barycentric interpolation formula in Chebyshev points* (Trefethen 2019). Given more time, one should implement this feature in TschebFun as well.

Continuing with integral approximation ... Text

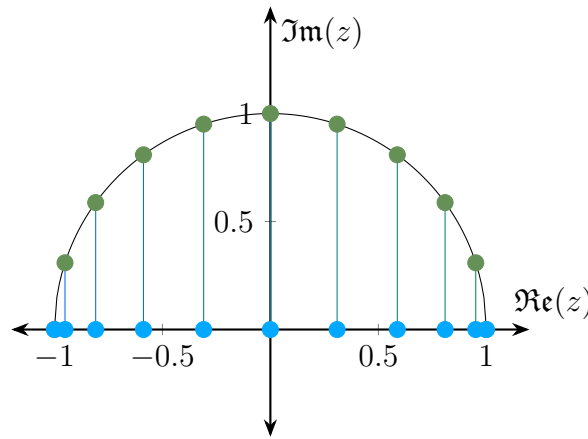
#### 2.4 Definition: Chebyshev points

From the equispaced points

$$\Theta_N := \left\{ \theta_j := \frac{j\pi}{N-1} \mid j = 0, \dots, N-1 \right\},$$

we can further define the Chebyshev points as the corresponding  $\cos(\theta_j)$ ,

$$X_N := \{x_j := \cos(\theta_j) \mid \theta_j \in \Theta_N\}.$$



**Figure 2:** The Chebyshev points  $\{x_j = \cos(\theta_j)\}$  are projections of the equispaced points  $\{\theta_j\}$  on the unit circle onto the x-axis.

### 2.3 A Numerical Issue

Wanting to interpolate through  $N$  points, namely through  $\{(x_j, f(x_j)) \mid x_j \in X_N\}$ , the corresponding truncated Chebyshev series will have  $N$  coefficients. To approximate the coefficient integral, we use Theorem 2.3 without the limit.

Even approximately speaking, the rectangular integral rule in application to Theorem 2.2 does, in some situations, not approach the correct coefficients. We now consider three different variations of the rectangular rule integral approximation, see why none of them are optimal in this situation and look at a possible resolution to this numerical issue, inspired by the DCT (Definition 2.3).

Intuitively, the problem is related to the endpoints of the interval  $[0, \pi]$ , the corresponding areas under the curve we aim to integrate and symmetry. In the following, consider the (near-) simplest case  $N = 2$  and  $f(x) = x$ . The analytical solution is  $a_0 = 0$  and  $a_1 = 1$ .

- With  $\Theta_N$  as in Definition 2.4, consider the issue that

$$a_1 = \frac{2}{\pi} \int_0^\pi \cos^2(\theta) d\theta \approx \frac{2}{N} \sum_{j=0}^{N-1} \cos^2\left(\frac{j\pi}{N-1}\right) = \frac{2}{2} (\cos^2(0) + \cos^2(\pi)) = 2 \neq 1.$$

- Instead sampling at  $\{\cos\left(\frac{j\pi}{N}\right) \mid j = 0, \dots, N-1\}$ ,

$$a_0 = \frac{1}{\pi} \int_0^\pi \cos(\theta) d\theta \approx \frac{1}{N} \sum_{j=0}^{N-1} \cos\left(\frac{j\pi}{N}\right) = \frac{1}{2} (\cos(0) + \cos(\pi/2)) = \frac{1}{2} \neq 0.$$

- Even for a creative  $\{\cos\left(\frac{j\pi}{N+1}\right) \mid j = 1, \dots, N\}$ ,

$$a_1 \approx \frac{2}{N} \sum_{j=1}^N \cos^2\left(\frac{j\pi}{N+1}\right) = \frac{2}{2} (\cos^2(\pi/3) + \cos^2(2\pi/3)) = \frac{1}{2} \neq 1.$$

So we cannot use normal chebpoints in this interpolation method.

elaborate

### 2.5 Definition: Modified Chebyshev points

$$\Theta_N := \left\{ \theta_j := \frac{(j + \frac{1}{2})\pi}{N-1} \mid j = 0, \dots, N-1 \right\},$$

$$X_N := \{x_j := \cos(\theta_j) \mid \theta_j \in \Theta_N\}.$$

Benjamin et al. 2010.

## 2.4 The Algorithm

Pseudo-Vandermonde bla



```

1  TschebFun TschebFun::interpolantThrough(Vector y) {
2      int order = y.size(), degree = order - 1;
3      Vector j = (xt::linspace((double)degree, 0.0, order) + 0.5) * (pi /
↪   order);
4      Vector coeffs = xt::zeros_like(y); // as many coefficients as data
↪   points
5      coeffs[0] = xt::sum(y)() / order;
6      for (size_t k = 1; k < order; k++)
7          coeffs[k] = (2.0 / order) * xt::sum(y * xt::cos(j * k))();
8      return TschebFun(coeffs);
9  }

```

## 3 The Spectral Method

### 3.1 Forward Euler

$$\frac{\partial u}{\partial t} \approx \frac{U_{n+1} - U_n}{\Delta t}$$

$$U_{n+1} \approx U_n + \Delta t \frac{\partial u}{\partial t} = U_n + \alpha \Delta t \frac{\partial^2 u}{\partial x^2}$$

$$\mathbf{a}^{(t+\Delta t)} = \mathbf{a}^{(t)} - \alpha \Delta t \cdot D_N^2 \mathbf{a}^{(t)}$$

elaborate

### 3.2 We need Differentiation

```

1  TschebFun TschebFun::derivative() {
2      int n = coefficients.size();
3      n = n - 1; // differentiation reduces the degree (order) by 1
4      Vector coeffs = coefficients; // make a copy
5      Vector derivative = xt::zeros<double>({n});
6      for (size_t j = n; j > 2; j--) {
7          derivative[j - 1] = (2 * j) * coeffs[j];
8          coeffs[j - 2] += (j * coeffs[j]) / (j - 2);
9      }
10     if (n > 1)
11         derivative[1] = 4 * coeffs[2];

```

```

12 derivative[0] = coeffs[1];
13 return TschebFun(derivative);
14 }

```

Differentiation matrix  $D_N$  according to [Trefethen 2000](#). [elaborate](#)

### 3.3 Enforcing boundary conditions

One way of forcing the boundary conditions, at least the first that came to my mind when thinking of this issue, is to pin down the two highest-order coefficients in the series representation after the iteration.

Let  $l := u_0(-1)$ ,  $r := u_0(1)$ . Recognise that

$$\begin{aligned} T_k(-1) &= T_k(\cos \pi) = \cos(k\pi) = (-1)^k \\ T_k(1) &= T_k(\cos 0) = \cos(k0) = 1 \end{aligned}$$

which leads to

$$\begin{aligned} u(-1, t) &= \sum_{k=0}^{N-1} a_k^{(t)} T_k(-1) = \sum_{k=0}^{N-3} \overbrace{a_k^{(t)} (-1)^k}^{:=\Sigma_1} + (-1)^{N-2} a_{N-2} + (-1)^{N-1} a_{N-1} = l \\ u(1, t) &= \sum_{k=0}^{N-1} a_k^{(t)} T_k(1) = \underbrace{\sum_{k=0}^{N-3} a_k^{(t)}}_{:=\Sigma_2} + a_{N-2} + a_{N-1} = r \end{aligned}$$

By adding up the above two equations, one obtains

$$\Sigma_1 + \Sigma_2 + \underbrace{\left((-1)^{N-2} + 1\right)}_{\in\{0,2\}} a_{N-2} + \underbrace{\left((-1)^{N-1} + 1\right)}_{\in\{0,2\}} a_{N-1} = l + r \quad (1)$$

For  $N$  even: Equation 1 has one unknown  $a_{N-2} = \frac{l+r-\Sigma_1-\Sigma_2}{2}$ ,  $a_{N-1} = r - a_{N-2} - \Sigma_2$ .

For  $N$  odd: Equation 1 has one unknown  $a_{N-1} = \frac{l+r-\Sigma_1-\Sigma_2}{2}$ ,  $a_{N-2} = r - a_{N-1} - \Sigma_2$ .

### Adaptive time-steps

### 3.4 The Destination

**Algorithm 1:** Heat Equation Forward-Euler

---

```

1 currentU = currentU -  $\alpha \Delta t$ 

```

---

## 4 Evaluation, Analysis, Comparison and Results

### 4.1 Evaluation: The Clenshaw Algorithm

#### 4.1 Theorem: Clenshaw recurrence relation

(Press et al. 1987, pp. 172–178).

```

1 Vector TschebFun::evaluateOn(Vector x) {
2     Vector U_kp2;
3     Vector U_kp1 = xt::zeros_like(x);
4     Vector U_k = xt::ones_like(x) * coefficients[coefficients.size() -
    ↪ 1];
5     for (int k = coefficients.size() - 2; k >= 0; k--) {
6         U_kp2 = U_kp1;
7         U_kp1 = U_k;
8         U_k = 2.0 * x * U_kp1 - U_kp2 + coefficients[k];
9     }
10    return (U_k - U_kp2 + coefficients[0]) / 2.0;
11 }

```

### 4.2 Analysis: Interface to Python

Extension to a Python module, using pybind11.

```

1 import heatfun, numpy as np
2 u0 = lambda x: np.exp(-12 * x**2)
3 x_of_interest = np.linspace(-1.0, 1.0, 500)
4 cheb_x = heatfun.modifiedChebpoints(30)
5 solution = heatfun.solve(u0(cheb_x), 0.01, x_of_interest)

```

### 4.3 Comparison: ChebFun

```

1 % example adapted from 'Exploring ODEs', page 282
2 pdefun = @(t, x, u) diff(u, 2);
3 bc.left = @(t, u) u;
4 bc.right = @(t, u) u;

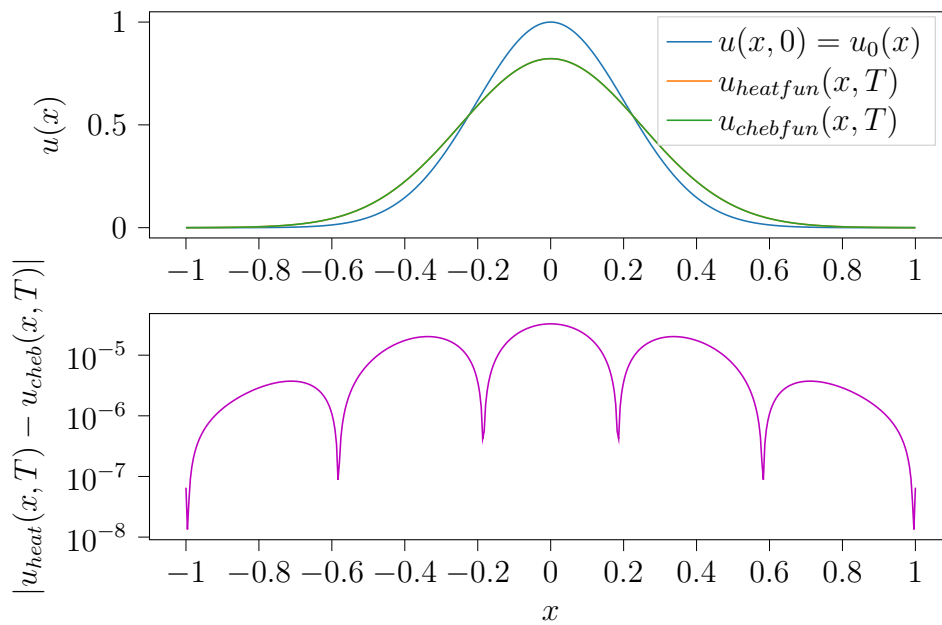
```

```

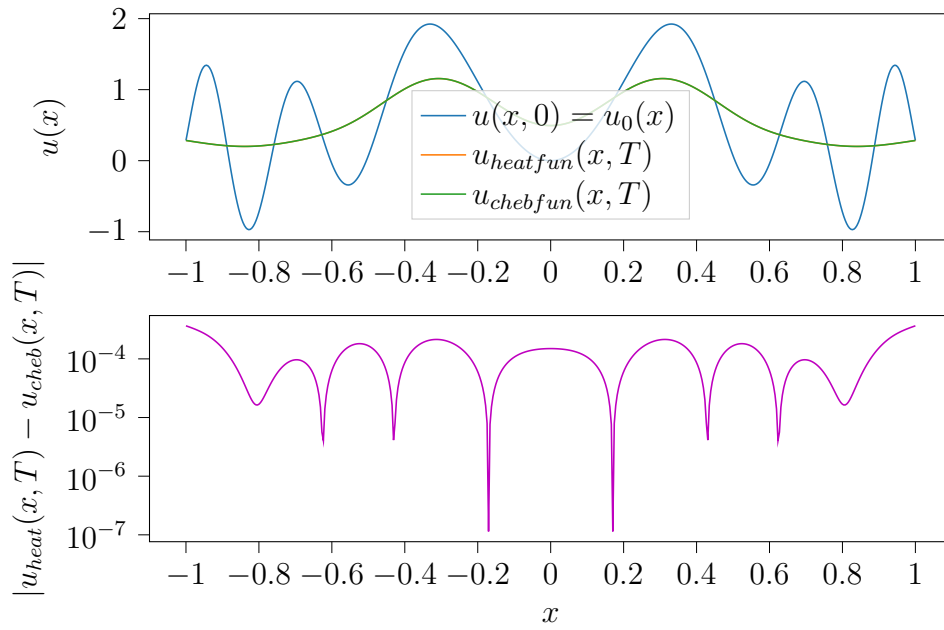
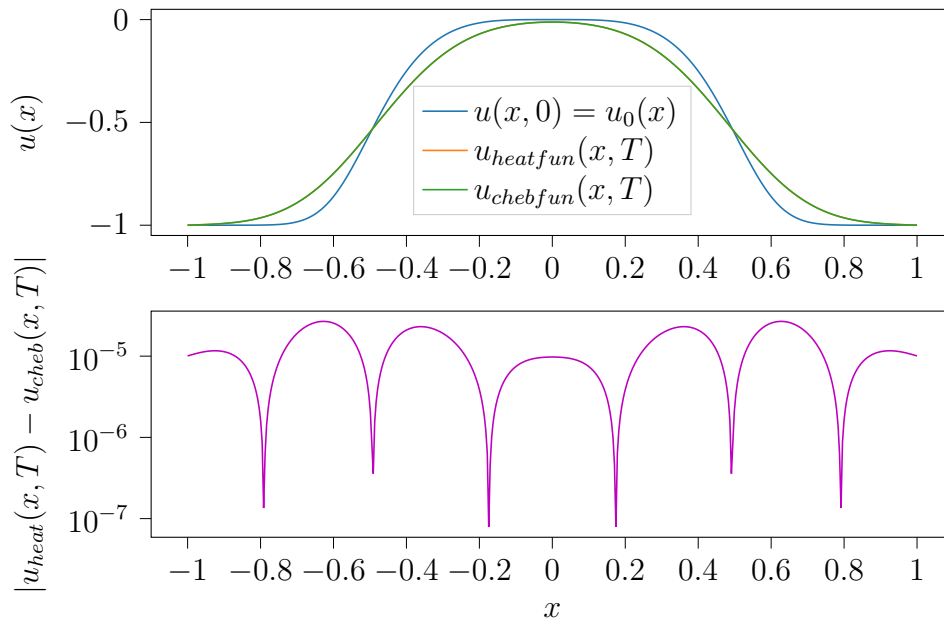
5 opts = pdeset('plot', 'off');
6 [t, u] = pde15s(pdefun, [0 0.005 0.010], u0, bc, opts);
7
8 x = linspace(-1.0, 1.0, 500).';
9 all_outputs = u(x);
10 output = all_outputs(:, end);
11 dlmwrite(filename, output, 'precision', '%.16f')

```

#### 4.4 Results: On Target?



**Figure 3:** Comparison of heatfun and chebfun

**Figure 4:** Comparison of heatfun and chebfun**Figure 5:** Comparison of heatfun and chebfun

## 5 Outlook and Discussion

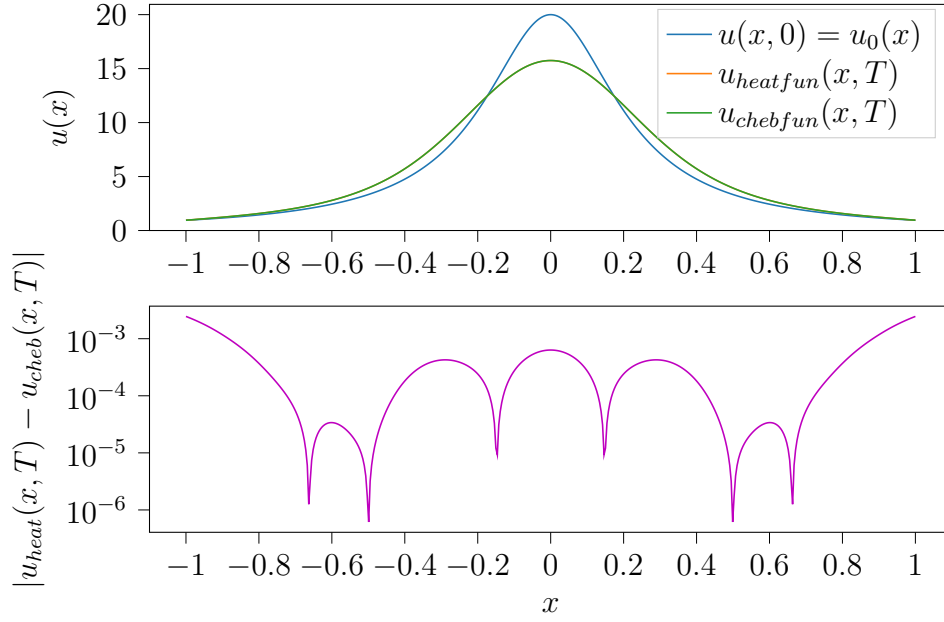
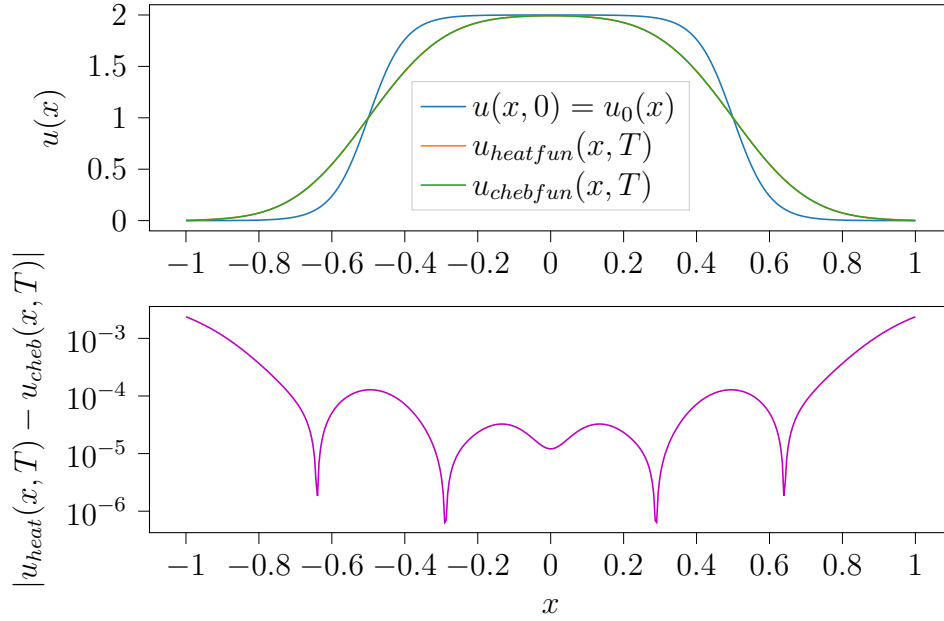
Given more time, one should.

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## A Title of Appendix

Appendices are definitely not necessary and assessors are not obliged to read them so only use them for non-vital text, figures or calculations.

**Figure 6:** Comparison of heatfun and chebfun**Figure 7:** Comparison of heatfun and chebfun

```

1  #pragma once
2
3  #include <xtensor/xarray.hpp>
4  #include <xtensor/xindex_view.hpp>
5  #include <xtensor/xview.hpp>
6

```

```
7 typedef xt::xarray<double> Vector;
8
9 class TschebFun {
10 public:
11     xt::xarray<double> coefficients;
12
13 public:
14     TschebFun(Vector coeffs);
15     size_t order() { return coefficients.size(); };
16     size_t degree() { return coefficients.size() - 1; };
17     Vector evaluateOn(Vector x);
18     TschebFun derivative();
19     TschebFun operator+(const TschebFun &other);
20     TschebFun operator-(const TschebFun &other);
21     TschebFun operator*(const double &factor);
22     static TschebFun interpolantThrough(Vector y);
23
24     static Vector equipoints(size_t N);
25     static Vector chebpoints(size_t N);
26     static Vector modifiedEquipoints(size_t N);
27     static Vector modifiedChebpoints(size_t N);
28 };
```