

# Introduction to quantum theory: Quantum states and quantum measurements

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# What is “Quantum theory”?

## Quantum theory

- Physics for microscopic phenomena, e.g., atoms, light.

## Why is quantum theory important ?

- Just because it's **reality**.
- Because it gives more **efficient** information processing, e.g., quantum factoring algorithm, quantum secret-key sharing, etc.

## On this course

We study **mathematical foundation** of quantum theory.

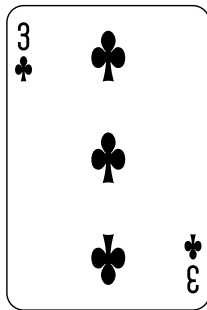
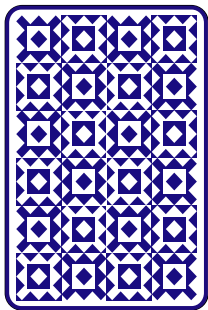
- Mathematical foundation of quantum physics
- Quantum algorithms
- Other quantum information processing, e.g, quantum communication, quantum error-correction.

Score

- Assignments: 100%

<https://github.com/QuantumComputationQuantumInformation/slides2020>

## Experimental facts



- ✗ The card was the three of club.
- ✓ When we flip the card, we see the three of club.
- State: The card
- Measurement: Flipping the card
- The number on the card before the flipping the card cannot be defined in quantum theory

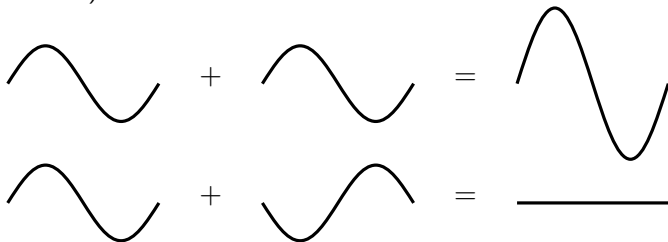
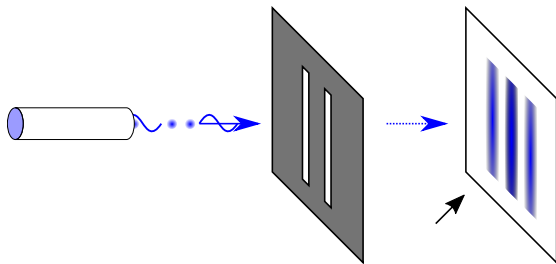
## States, measurements and distinguishability

- Two states are equivalent if the probability distributions of outcomes are equal for **arbitrary measurement**.
- Two measurements are equivalent if the probability distributions of outcomes are equal for **arbitrary states**.
- **We should not distinguish states (measurements) mathematically if we cannot distinguish them physically.**

## Light is wave (1801)



Thomas Young  
(1773 – 1829)

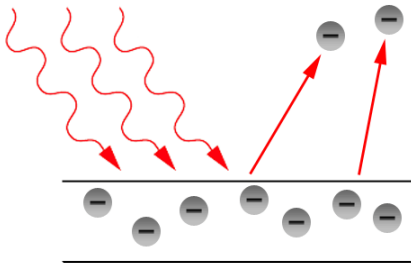


[https://en.wikipedia.org/wiki/Young's\\_interference\\_experiment](https://en.wikipedia.org/wiki/Young's_interference_experiment)

## Light is particle (1905)

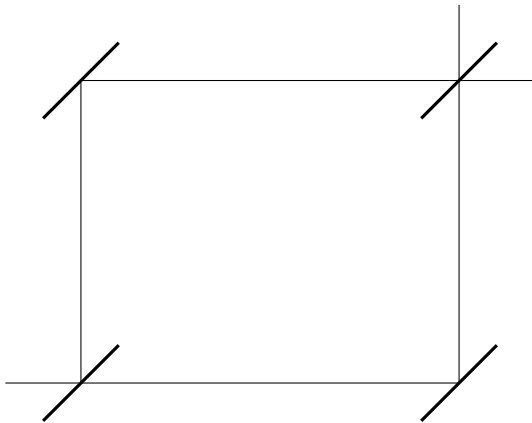


Albert Einstein  
(1879 – 1955)



[https://en.wikipedia.org/wiki/File:Photoelectric\\_effect.png](https://en.wikipedia.org/wiki/File:Photoelectric_effect.png)

## Mach-Zehnder interferometer





## Quantum states and quantum measurements

A single photon  $\Rightarrow$  BS1  $\Rightarrow$  BS2  $\Rightarrow$  detection

State	Measurement
A single photon $\Rightarrow$ BS1 $\Rightarrow$ BS2	detection
A single photon $\Rightarrow$ BS1	BS2 $\Rightarrow$ detection
A single photon	BS1 $\Rightarrow$ BS2 $\Rightarrow$ detection

All understandings are valid

## Mathematical representations of states and measurements

How “States” and “Measurements” are represented mathematically ?

A table of probabilities of outcome 'YES' for each binary measurement on each state

	Measurement 1	Measurement 2	...
State A	$p_{A1}$	$p_{A2}$	...
State B	$p_{B1}$	$p_{B2}$	...
⋮			

\* The number of states and measurements are not necessarily countable.

## Linear space

	Measurement 1	Measurement 2	...
State A	$p_{A1}$	$p_{A2}$	...
State B	$p_{B1}$	$p_{B2}$	...
State C	$0.7p_{A1} + 0.3p_{B1}$	$0.7p_{A2} + 0.3p_{B2}$	

### Assumption

- Probabilistic mixture of states is also state.
- Probabilistic mixture of binary measurement is also binary measurement.

States and measurements can be represented by vectors!

## Classical theory

States: 0, 1

Binary measurements: 0?, 1?

$$\underline{0}?(0) = 1,$$

$$\underline{0}?(1) = 0,$$

$$\underline{1}?(0) = 0,$$

$$\underline{1}?(1) = 1$$

	<u>0</u> ?	<u>1</u> ?
<u>0</u>	1	0
<u>1</u>	0	1

## State and measurement

	<u>0</u> ?	<u>1</u> ?
<u>0</u>	1	0
<u>1</u>	0	1

$S := \underline{0}$  with probability  $p$ ,  $\underline{1}$  with probability  $1 - p$ .

$S$  is also **regarded as a state**.

$$\underline{0}?(S) = p,$$

$$\underline{1}?(S) = 1 - p.$$

Similarly,

$E_1 := \underline{0}?$  with probability  $p$ ,  $\underline{1}?$  with probability  $1 - p$ .

$E_2 := (\underline{0} \text{ or } \underline{1})?$ .

$E_1$  and  $E_2$  are also **regarded as a binary measurement**.

## Linear space

$$\omega_{\underline{0}} := \begin{bmatrix} 1 \\ 0 \end{bmatrix},$$

$$e_{\underline{0}} := \begin{bmatrix} 1 \\ 0 \end{bmatrix},$$

$$\omega_{\underline{1}} := \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

$$e_{\underline{1}} := \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

$$\underline{0}?(0) = \langle e_{\underline{0}}, \omega_{\underline{0}} \rangle,$$

$$\underline{0}?(1) = \langle e_{\underline{0}}, \omega_{\underline{1}} \rangle$$

$$\underline{1}?(0) = \langle e_{\underline{1}}, \omega_{\underline{0}} \rangle,$$

$$\underline{1}?(1) = \langle e_{\underline{1}}, \omega_{\underline{1}} \rangle$$

$S := \underline{0}$  with probability  $p$ ,  $\underline{1}$  with probability  $1 - p$

$$\omega_S = p\omega_{\underline{0}} + (1 - p)\omega_{\underline{1}} = \begin{bmatrix} p \\ 1 - p \end{bmatrix}.$$

$$\underline{0}?(S) = p = \langle e_{\underline{0}}, \omega_S \rangle.$$

## States and measurements in a linear space

$$\omega_{\underline{0}} := \begin{bmatrix} 1 \\ 0 \end{bmatrix},$$

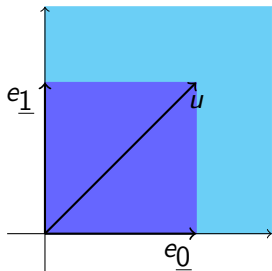
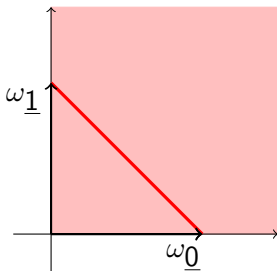
$$\omega_{\underline{1}} := \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

$$e_{\underline{0}} := \begin{bmatrix} 1 \\ 0 \end{bmatrix},$$

$$e_{\underline{1}} := \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

$$\text{Set of states} = \left\{ \begin{bmatrix} x \\ y \end{bmatrix} \mid x \geq 0, y \geq 0, x + y = 1 \right\}.$$

$$\text{Set of binary measurements} = \left\{ \begin{bmatrix} x \\ y \end{bmatrix} \mid x \geq 0, y \geq 0, x \leq 1, y \leq 1 \right\}.$$



## State and measurement in a linear space

$$\text{Set of states} = \left\{ \begin{bmatrix} x \\ y \end{bmatrix} \mid x \geq 0, y \geq 0, x + y = 1 \right\}.$$

$$\text{Set of binary measurements} = \left\{ \begin{bmatrix} x \\ y \end{bmatrix} \mid x \geq 0, y \geq 0, x \leq 1, y \leq 1 \right\}.$$

$$\text{Let } C_{\geq 0} \text{ be the set of nonnegative vectors and } u := \begin{bmatrix} 1 \\ 1 \end{bmatrix}.$$

$$\text{Set of states} = \{ \omega \in \mathbb{R}^2 \mid \omega \in C_{\geq 0}, \langle u, \omega \rangle = 1 \}.$$

$$\text{Set of binary measurements} = \{ e \in \mathbb{R}^2 \mid e \in C_{\geq 0}, u - e \in C_{\geq 0} \}.$$

$$\text{Set of measurements} = \{ (e_1, \dots, e_k) \mid e_1 + \dots + e_k = u, e_i \in C_{\geq 0} \\ i = 1, 2, \dots, k, k = 1, 2, \dots \}$$

Outcome of the measurement  $M = (e_1, \dots, e_k)$  on  $\omega$  is  $i$  with probability  $\langle e_i, \omega \rangle$ .



## Quantum theory

$C_{\geq 0} \subseteq \mathbb{R}^2$  : the set of nonnegative vectors,  $u := \begin{bmatrix} 1 \\ 1 \end{bmatrix}$ .

Set of states =  $\{\omega \in \mathbb{R}^2 \mid \omega \in C_{\geq 0}, \langle u, \omega \rangle = 1\}$ .

Set of binary measurements =  $\{e \in \mathbb{R}^2 \mid e \in C_{\geq 0}, u - e \in C_{\geq 0}\}$ .

Set of measurements =  $\{(e_1, \dots, e_k) \mid e_1 + \dots + e_k = u, e_i \in C_{\geq 0}$   
 $i = 1, 2, \dots, k, k = 1, 2, \dots, \}$

$V$  : the linear space on  $\mathbb{R}$  spanned by  $2 \times 2$  Hermitian matrices.  
 $\langle e, \omega \rangle := \text{Tr}(e\omega)$  for  $\omega, e \in V$  (Hilbert-Schmidt inner product).

$C_{\succeq 0} \subseteq V$  : the set of positive semidefinite matrices,  $u := \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$ .

Set of states =  $\{\omega \in V \mid \omega \in C_{\succeq 0}, \langle u, \omega \rangle = 1\}$ .

Set of binary measurements =  $\{e \in V \mid e \in C_{\succeq 0}, u - e \in C_{\succeq 0}\}$ .

## Linear space spanned by 2x2 Hermitian matrices

Basis

$$A := \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}, B := \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}, C := \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, D := \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$$

Another choice of basis

$$I := \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, X := \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, Y := \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, Z := \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

Both are orthogonal basis.

The second basis ( $I$  and Pauli matrices  $X$ ,  $Y$  and  $Z$ ) has nice properties.

- 1  $\text{Tr}(I) = 2$ .  $\text{Tr}(X) = \text{Tr}(Y) = \text{Tr}(Z) = 0$ .
- 2  $X^2 = Y^2 = Z^2 = I$  ( $X$ ,  $Y$  and  $Z$  have eigenvalues  $\pm 1$ ).
- 3  $XY = -YX$ ,  $YZ = -ZY$ ,  $ZX = -XZ$ .

## Positive semidefinite cone

$$I := \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, X := \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, Y := \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, Z := \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

$$\rho = \frac{1}{\sqrt{2}} (a_I I + a_X X + a_Y Y + a_Z Z)$$

$$\lambda_1 \geq 0, \lambda_2 \geq 0 \iff \lambda_1 + \lambda_2 \geq 0, \lambda_1 \lambda_2 \geq 0$$

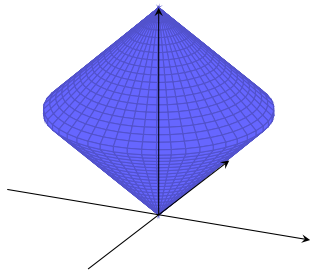
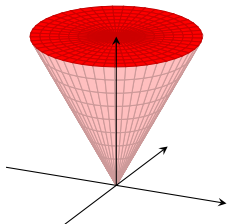
$$\iff \text{Tr}(\rho) \geq 0, \text{Tr}(\rho)^2 - \text{Tr}(\rho^2) \geq 0$$

$$\iff a_I \geq 0, 2a_I^2 - (a_I^2 + a_X^2 + a_Y^2 + a_Z^2) \geq 0$$

$$\iff a_I \geq 0, a_I^2 \geq a_X^2 + a_Y^2 + a_Z^2$$

$$\text{Tr}(\rho) = 1 \iff a_I = \frac{1}{\sqrt{2}}$$

# Geometry of quantum states and effects



## Convex cone and dual cone

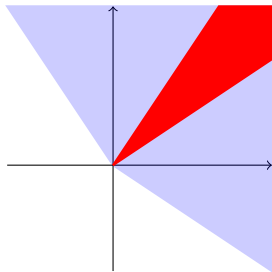
$$C \subseteq V \text{ is a convex cone} \iff x + y \in C, \lambda x \in C, \\ \forall x \in C, y \in C, \lambda \geq 0$$

Proper cone: closed, not  $V$ , full-dimensional.

$C^* \subseteq V$  is a dual cone of  $C$

$$\iff C^* := \{x \in V \mid \langle x, y \rangle \geq 0, \forall y \in C\}$$

$C_{\geq 0}$  and  $C_{\leq 0}$  are self-dual cones.



## Generalized probabilistic theories

$C$  : convex cone.

$u \in \text{interior of } C^*$ .

Set of states =  $\{\omega \in V \mid \omega \in C, \langle u, \omega \rangle = 1\}$ .

Set of effects =  $\{e \in V \mid e \in C^*, u - e \in C^*\}$ .

Set of measurements =  $\{(e_1, \dots, e_k) \mid e_1 + \dots + e_k = u, e_i \in C^*,$   
 $i = 1, 2, \dots, k, k = 1, 2, \dots\}$

Classical theory

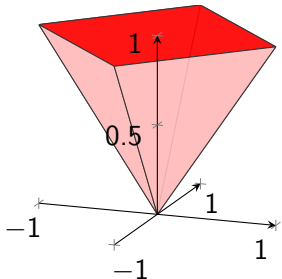
$V = \mathbb{R}^n$ ,  $C = C_{\geq 0}$ ,  $u$  = the all-1 vector.

Quantum theory

$V$  = A set of  $n \times n$  Hermitian matrices,  $C = C_{\geq 0}$ ,  $u = I$ .

Other theories ?

## Toy theory: Gbit



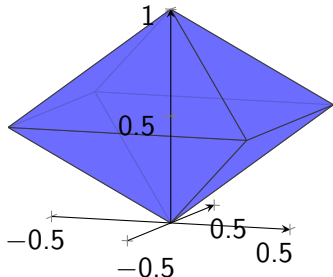
$$\omega_0 = [1 \ 0 \ 1],$$

$$\omega_2 = [-1 \ 0 \ 1],$$

$$e_0 = \frac{1}{2} [1 \ 1 \ 1],$$

$$e_2 = \frac{1}{2} [1 \ -1 \ 1],$$

$$u = [0 \ 0 \ 1].$$



$$\omega_1 = [0 \ 1 \ 1],$$

$$\omega_3 = [0 \ -1 \ 1].$$

$$e_1 = \frac{1}{2} [-1 \ 1 \ 1],$$

$$e_3 = \frac{1}{2} [-1 \ -1 \ 1],$$

## Nonlocality

OK, generalized probabilistic theory is quite simple and easy to understand.

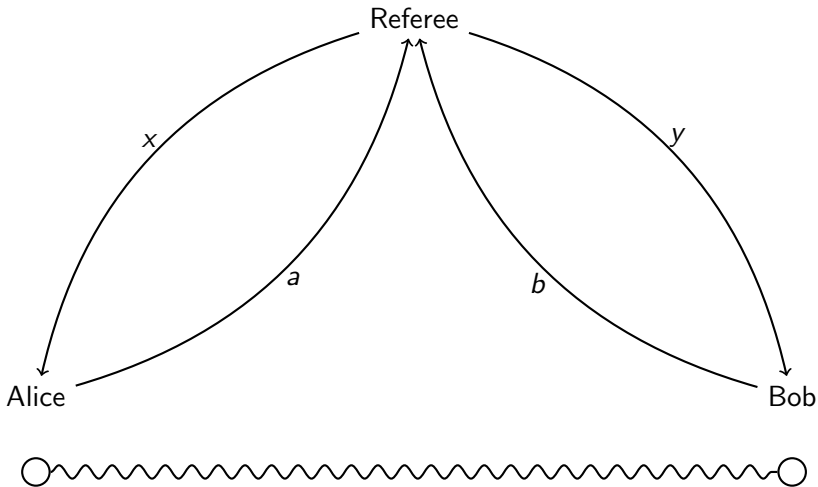
But, what is essential difference between classical theory and quantum theory ?

Can classical theory “simulates” or “explains” quantum theory ?

Are Hermitian positive-semidefinite matrices really needed for explaining reality ?



## Bell test: CHSH game (1964, 1969)



Alice and Bob win iff  $a \oplus b = x \wedge y$ .

## Bell inequality

$a_x$ : Output of Alice for given  $x$ .

$b_y$ : Output of Bbob for given  $y$ .

$$a_0 \oplus b_0 = 0$$

$$a_1 \oplus b_0 = 0$$

$$a_0 \oplus b_1 = 0$$

$$a_1 \oplus b_1 = 1$$

By adding all equations, we get  $0 = 1$ , which means there is no solution. Hence, the winning probability 1 cannot be achieved.

Three equalities can be satisfied, so that the largest winning probability is  $3/4$  (Bell inequality or CHSH inequality).

If Alice and Bob share quantum states, then the largest winning probability is  $(2 + \sqrt{2})/4 \approx 0.854$  (Violation of Bell/CHSH inequality)

## Locality (Hidden variable model)

Joint preparation and independent measurements.

Probability distribution  $P(a, b \mid x, y)$  is said to be **local** if

$$P(a, b \mid x, y) = \sum_{\lambda} P(\lambda) P(a \mid x, \lambda) P(b \mid y, \lambda).$$

Quantum physics allow **nonlocal** behaviors.

## Summary

- Let  $V$  be the linear space spanned by  $n \times n$  Hermitian matrices. Let  $C_{\succeq 0}$  be the set of PSD matrices. Let  $I$  be the  $n \times n$  identity matrix. Let  $\langle e, \omega \rangle := \text{Tr}(e\omega)$  for  $\omega, e \in V$  (Hilbert-Schmidt inner product).

$$\text{Set of states} = \{\omega \in V \mid \omega \in C_{\succeq 0}, \langle u, \omega \rangle = 1\}.$$

$$\text{Set of binary measurements} = \{e \in V \mid e \in C_{\succeq 0}, u - e \in C_{\succeq 0}\}.$$

- Classical theory and quantum theory are special cases of generalized probabilistic theories.
- Violation of Bell (CHSH) inequality show that quantum theory is essentially different from classical theory (Arguments on joint system is needed).

## Assignments

- 1 Show the dimension and one of the basis of the real linear space spanned by  $n \times n$  Hermitian matrices.
- 2 Show that  $XY = -YX$ ,  $YZ = -ZY$  and  $ZX = -XZ$ .
- 3 Show that the Hilbert–Schmidt inner product satisfies the axioms of inner product.
- 4 [Advanced] Show that  $C_{\succeq 0}$  is a self-dual cone.