

Seeded QED cascades in counter-propagating laser pulses

T. Grismayer, M. Vranic, J. L. Martins, R. A. Fonseca, and L. O. Silva, Phys Rev E 95, 023210 (2017)

Notebook: Óscar Amaro, November 2022 @ [GoLP-EPP](#)

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Introduction

In this notebook we reproduce some results from the paper.

Setups A,B,C field components

```
In[437]:= (* setup 1 (lp-lp) *)
Clear[a0, k0, ω0, x, t, ap, am]
ap = {0, a0 Cos[ω0 t + k0 x], 0};
am = {0, a0 Cos[ω0 t - k0 x], 0};

(* E: only y component. pre-
   factor ω0 is due to choice of reduced units/normalization *)
E = -D[ap + am, t] // Simplify
(* B: only z component. pre-
   factor k0 is due to choice of reduced units/normalization *)
B = Curl[ap + am, {x, y, z}] // Simplify

Out[440]= {0, 2 a0 ω0 Cos[k0 x] Sin[t ω0], 0}

Out[441]= {0, 0, -2 a0 k0 Cos[t ω0] Sin[k0 x]}
```

```

In[442]:= (* setup 2 (cw-cw) *)
(*"...results in a helical field structure
growing or shrinking uniformly in space..."*)
Clear[a0, k0, ω0, x, t, ap, am]
ap = {0, a0 Cos[ω0 t + k0 x], a0 Sin[ω0 t + k0 x]};
am = {0, a0 Cos[ω0 t - k0 x], -a0 Sin[ω0 t - k0 x]};

(* E: y+z components *)
E = -D[ap + am, t] // Simplify
(* B: y+z components *)
B = Curl[ap + am, {x, y, z}] // Simplify
Out[445]= {0, 2 a0 ω0 Cos[k0 x] Sin[t ω0], 2 a0 ω0 Sin[k0 x] Sin[t ω0]}
Out[446]= {0, -2 a0 k0 Cos[k0 x] Cos[t ω0], -2 a0 k0 Cos[t ω0] Sin[k0 x]}

In[447]:= (* setup 3 (cw-cp) *)
(*"...fixed planar beating pattern that rotates around the laser propagation
axis. This setup consists in a rotating field structure..."*)
Clear[a0, k0, ω0, x, t, ap, am]
ap = {0, a0 Cos[ω0 t + k0 x], -a0 Sin[ω0 t + k0 x]};
am = {0, a0 Cos[ω0 t - k0 x], -a0 Sin[ω0 t - k0 x]};

(* E: y+z components *)
E = -D[ap + am, t] // Simplify
(* B: y+z components *)
B = Curl[ap + am, {x, y, z}] // Simplify
Out[450]= {0, 2 a0 ω0 Cos[k0 x] Sin[t ω0], 2 a0 ω0 Cos[k0 x] Cos[t ω0]}
Out[451]= {0, -2 a0 k0 Sin[k0 x] Sin[t ω0], -2 a0 k0 Cos[t ω0] Sin[k0 x]}

```

A. Ideal model (check some of the calculations of the paper)

```
In[452]:= Clear[n0, s, t, Wγ, Wp, np, nppt, tt, intt, dnpdt]
```

```
(* invert Laplace transform of equation 6 *)
```

```
np = InverseLaplaceTransform[ $\frac{n0}{s - \frac{2 W_\gamma W_p}{s + W_p}}$ , s, t] // Simplify
```

```
(* LHS of eq 5 *)
```

```
dnpdt = D[np, t] // Simplify
```

```
(* change of variable *)
```

```
nppt = np /. {t → tt} // Simplify;
```

```
(* RHS of eq 5 *)
```

```
intt = 2 Integrate[nppt Wγ Wp Exp[-Wp (t - tt)], {tt, 0, t}] // Simplify
```

```
(* check validity of eq 5 using Laplace transform of eq 6*)
```

```
Refine[intt - dnpdt // FullSimplify, {Wp > 0, Wγ > 0}] // Simplify
```

$$\text{Out[453]} = \frac{e^{-\frac{1}{2} t (W_p + \sqrt{W_p} \sqrt{W_p + 8 W_\gamma})} n_0 \left((-1 + e^{t \sqrt{W_p} \sqrt{W_p + 8 W_\gamma}}) \sqrt{W_p} + (1 + e^{t \sqrt{W_p} \sqrt{W_p + 8 W_\gamma}}) \sqrt{W_p + 8 W_\gamma} \right)}{2 \sqrt{W_p + 8 W_\gamma}}$$

$$\text{Out[454]} = \frac{2 e^{-\frac{1}{2} t (W_p + \sqrt{W_p} \sqrt{W_p + 8 W_\gamma})} (-1 + e^{t \sqrt{W_p} \sqrt{W_p + 8 W_\gamma}}) n_0 \sqrt{W_p} W_\gamma}{\sqrt{W_p + 8 W_\gamma}}$$

$$\text{Out[456]} = \frac{4 e^{-\frac{t W_p}{2}} n_0 \sqrt{W_p} W_\gamma \sinh\left[\frac{1}{2} t \sqrt{W_p} \sqrt{W_p + 8 W_\gamma}\right]}{\sqrt{W_p + 8 W_\gamma}}$$

```
Out[457]= 0
```

```

In[458]:= (* eqs 7 and 8 *)
Clear[sp, sm, s, Wp, Wγ, sol, r]
sol = Refine[Solve[s2 + Wp s - 2 Wγ Wp == 0, s], {Wp > 0, Wγ > 0}] // FullSimplify
sp = sol[[2, 1, 2]]

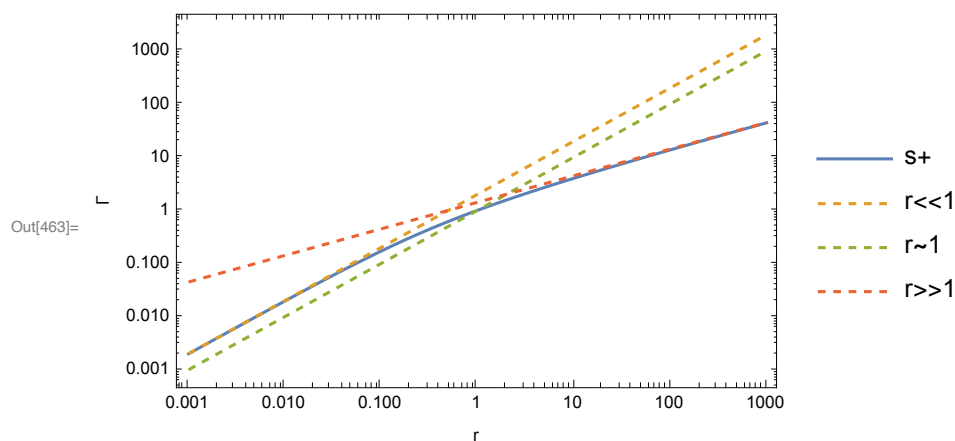
(* rescale *)
Wγ = r Wp;
Wp = 1;

LogLogPlot[{sp, 2 Wγ, Wγ, Sqrt[2 Wγ Wp]},
  {r, 10-3, 103}, PlotStyle → {Default, Dashed, Dashed, Dashed},
  PlotLegends → {"s+", "r<<1", "r~1", "r>>1"},
  Frame → True, FrameLabel → {"r", "τ"}]

```

Out[459]= $\left\{ \left\{ s \rightarrow -\frac{Wp}{2} - \frac{1}{2} \sqrt{Wp (Wp + 8 W\gamma)} \right\}, \left\{ s \rightarrow \frac{1}{2} (-Wp + \sqrt{Wp (Wp + 8 W\gamma)}) \right\} \right\}$

Out[460]= $\frac{1}{2} (-Wp + \sqrt{Wp (Wp + 8 W\gamma)})$



III. CASCADE MODELS - Weak-field limit

```

In[464]:= Clear[Wp, d2Pdt dχγ, α, τc, χγ, εγ, χeavg, γavg, δ, f]
Clear[ħ, m, c, e]
Clear[sols, argexp, dargexp, d2argexp]
Clear[fx0, d2fx0, hx0, int]

(* to get eq 11, just assume s+
   Wp ~ s in equation 10 and solve the quadratic equation s^2-2 ∫ ... = 0 *)

(* expressions in text after eq 10 *)
Wp =  $\frac{3 \pi}{50} \frac{\alpha}{\tau c} \text{Exp}\left[-\frac{8}{3 \chi \gamma}\right] \frac{\chi \gamma}{\epsilon \gamma};$ 
d2Pdt dχγ =  $\text{Sqrt}\left[\frac{2}{3 \pi}\right] \frac{\alpha}{\tau c} \frac{\text{Exp}[-\delta]}{\delta^{1/2} \chi \text{eavg} \gamma \text{avg}};$ 
δ =  $\frac{2 \chi \gamma}{3 \chi \text{eavg} (\chi \text{eavg} - \chi \gamma)};$ 

(* integrand of equation 11 *)
d2Pdt dχγ Wp // Simplify

(* "...the argument of the exponential possesses a unique maximum..."*)
argexp =  $-\frac{8}{3 \chi \gamma} - \frac{2 \chi \gamma}{3 \chi \text{eavg}^2 - 3 \chi \text{eavg} \chi \gamma};$  (* argument of exponential *)
dargexp = D[argexp, χγ]; (* derivative of this argument *)
d2argexp = D[argexp, {χγ, 2}] // Simplify;
(* second derivative of this argument *)
sols = Solve[dargexp == 0, χγ] (* get stationary points *)

d2argexp /. {χγ → sols[[1, 1, 2]]} (* 2 χeavg/3 is a maximum *)
d2argexp /. {χγ → sols[[2, 1, 2]]} (* 2 χeavg is a minimum *)

Out[471]= 
$$\frac{3 e^{-\frac{8}{3 \chi \gamma} - \frac{2 \chi \gamma}{3 \chi \text{eavg}^2 - 3 \chi \text{eavg} \chi \gamma}} \alpha^2 \chi \gamma}{50 \gamma \text{avg} \epsilon \gamma \tau c^2 \chi \text{eavg} \sqrt{\frac{\chi \gamma}{\pi \chi \text{eavg}^2 - \pi \chi \text{eavg} \chi \gamma}}}$$


Out[475]= 
$$\left\{ \left\{ \chi \gamma \rightarrow \frac{2 \chi \text{eavg}}{3} \right\}, \left\{ \chi \gamma \rightarrow 2 \chi \text{eavg} \right\} \right\}$$


Out[476]= 
$$-\frac{54}{\chi \text{eavg}^3}$$


Out[477]= 
$$\frac{2}{3 \chi \text{eavg}^3}$$


```

```

In[478]:= (* physical derived constants *)
(*  $\tau c = \frac{\hbar}{m c^2}$ ;
 $\alpha = \frac{e^2}{\hbar c}$ ; *)

(* text before eq 12 *)
f = - $\delta - \frac{8}{3 \chi \gamma}$  // Simplify
fx0 = f /. { $\chi \gamma \rightarrow \frac{2 \chi_{\text{eavg}}}{3}$ } // Simplify
d2fx0 = (D[f, { $\chi \gamma$ , 2}] // Simplify) /. { $\chi \gamma \rightarrow \frac{2 \chi_{\text{eavg}}}{3}$ }
hx0 = Refine[ $\frac{1}{\text{Sqrt}[\delta]}$  /. { $\chi \gamma \rightarrow \frac{2 \chi_{\text{eavg}}}{3}$ }, { $\chi_{\text{eavg}} > 0$ }]

(* apply Laplace method following text before eq 12 *)
int = Sqrt[ $\frac{2 \pi}{\text{Abs}[d2fx0]}$ ] hx0 Exp[fx0];

(* equation 12 for growth rate in weak-field limit. however,
in the analysis in the text the pre-
factors are removed from the Laplace method. We have to put them back *)
(* attention:  $\epsilon \gamma = \gamma_{\text{avg}} \chi \gamma / \chi_{\text{eavg}}$  according to text before eq 10. *)
(* this then leads to eq 12 *)
Refine[Sqrt[ $\frac{3 \pi}{50} \frac{\alpha}{\tau c} \text{Sqrt}[\frac{2}{3 \pi}] \frac{\alpha}{\tau c} \frac{1}{\gamma_{\text{avg}}^2} 2 \text{int}$ ] // Simplify,
{ $\chi_{\text{eavg}} > 0$ ,  $\alpha > 0$ ,  $\tau c > 0$ ,  $\gamma_{\text{avg}} > 0$ }]

Out[478]= 
$$\frac{2 (-2 \chi_{\text{eavg}} + \chi \gamma)^2}{3 \chi_{\text{eavg}} \chi \gamma (-\chi_{\text{eavg}} + \chi \gamma)}$$


Out[479]= 
$$-\frac{16}{3 \chi_{\text{eavg}}}$$


Out[480]= 
$$-\frac{54}{\chi_{\text{eavg}}^3}$$


Out[481]= 
$$\frac{\sqrt{3} \sqrt{\chi_{\text{eavg}}}}{2}$$


Out[483]= 
$$\frac{e^{-\frac{8}{3} / \chi_{\text{eavg}}} \sqrt{\pi} \alpha \chi_{\text{eavg}}}{5 \times 6^{1/4} \gamma_{\text{avg}} \tau c}$$


```